## Systematic Study of the Boundary Composition in Poisson Boltzmann Calculations

Parimal Kar<sup>1</sup>, Yanj

Yanjie Wei<sup>1</sup>, Ulrich H. E. Hansmann<sup>1,2</sup>

Siegfried Höfinger<sup>1\*</sup>

<sup>1</sup>Department of Physics, Michigan Technological University, 1400 Townsend Drive, Houghton, MI, 49331 - 1295, USA <sup>2</sup>John v. Noumann Institute for Computing Forschungszontrum Jülich, 52425 Jülich, Gor

<sup>2</sup>John v. Neumann Institute for Computing, Forschungszentrum Jülich, 52425 Jülich, Germany

\*Corresponding Author:

Siegfried Höfinger Michigan Technological University Department of Physics 1400 Townsend Drive, Houghton, Michigan, 49931 - 1295, USA Phone: 1-906-487-1496 Fax: 1-906-487-2933 Email: shoefing@mtu.edu

**Keywords:** Solvation, Continuum Electrostatics, Polarizable Continuum Model, Poisson Boltzmann, Implicit Solvation Models

Abbreviations: DNA, desoxyribosenucleic acid; RNA, ribonucleic acid; PB, Poisson Boltzmann; GB, Generalized Born; FDPB, finite difference Poisson Boltzmann; PB/BEM, boundary element method Poisson Boltzmann; SASA, solvent accessible surface area; PCM, polarizable continuum model; PDB, protein data bank; PDB-REPRDB, representative protein chains from PDB; MSROLL, molecular surface program by M. Connolly; SIMS, molecular surface program by Y. Vorobjev; MSMS, molecular surface program by M. Sanner; MOLDEN, a molecular graphics program; AMBER, a molecular modelling package for biomolecular simulation; DFT, density functional theory; GAUSSIAN-03, a quantum chemical software package; GG,AA,VV, di-glycine, di-alanine, di-valine;

#### Abstract

We describe a three-stage procedure to analyze the dependence of Poisson Boltzmann calculations on the shape, size and geometry of the boundary between solute and solvent. our study is carried out within the boundary element formalism, but our results are also of interest to finite difference techniques of Poisson Boltzmann calculations. At first, we identify the critical size of the geometrical elements for discretizing the boundary, and thus the necessary resolution required to establish numerical convergence. In the following two steps we perform reference calculations on a set of dipeptides in different conformations using the Polarizable Continuum Model and a high-level Density Functional as well as a high-quality basis set. Afterwards, we propose a mechanism for defining appropriate boundary geometries. Finally, we compare the classic Poisson Boltzmann description with the Quantum Chemical description, and aim at finding appropriate fitting parameters to get a close match to the reference data. Surprisingly, when using default AMBER partial charges and the rigorous geometric parameters derived in the initial two stages, no scaling of the partial charges is necessary and the best fit against the reference set is obtained automatically.

#### 1 Introduction

A common way of describing solvation effects to biomolecular structure is to treat the solvent as a continuum of characteristic dielectric constant. The biomolecule of interest, i.e. a protein, DNA, RNA, glycolipid, etc. is considered in full atomic detail, while the surrounding medium is represented as structureless continuum interacting primarily via polarization, dispersion, repulsion and cavitation effects (1, 2, 3, 4, 5). The underlying physics concerned with polarization is then often expressed in terms of solutions to the Poisson-Boltzmann equation (PB) (6, 7, 8, 9, 10, 11, 12, 13). Approximations to the PB — motivated by simplified computational protocols — are standard practice e.g. the Generalized Born model (GB) (14, 15). However, PB and GB are dealing with the polarization term only, and the other above mentioned interactions are usually treated by either first-principle (16) or semiempirical (17) character.

Solutions to the PB are computed either by the finite difference method (FDPB) (6, 7, 8, 9) or by the boundary element method (PB/BEM) (12, 13). The latter is particularly intriguing since it reduces a three-dimensional integral over the entire volume to a two-dimensional surface integral, leading to considerable savings in computational time. Both approaches depend fundamentally on the exact definition of the boundary between solute and solvent. All definitions are based on the area of the atoms exposed to the solvent, for instance the solvent accessible surface area (SASA), the solvent excluded volume, or the molecular surface (18), which all depend on a chosen set of van der Waals radii (19, 20, 21, 22) assigned to the center of the atoms.

Given the dependence on the exact geometry and quality of the boundary it appears necessary to study the geometric factors that influence the outcome of PB calculations in greater detail. This is particularly appropriate for semi-quantitative approaches (23) where the demand on accuracy is a very sensitive issue (24). Particular attention has to be drawn to factors such as i) surface type and surface resolution, ii) dependence on atomic model parameters, i.e. van der Waals radii, iii) generality and physicochemical significance. In this present work we provide such an analysis by focussing on each of these three points separately. At first, we employ different surface generation algorithms to a subset of randomly chosen protein structures of variable size and shape. PB/BEM calculations are carried out with increasing resolution of the boundary. Optimal surface resolution and surface generation parameters that guarantee numerical convergence and methodic stability are derived. Next, we use these optimized parameters for a set of model peptides and vary the van der Waals radii in a systematic way. The reference set of model peptides is considered at a high level of quantum chemical theory, i.e. PCM (3) using the Becke-98 density functional (25) and the basis set of Sadlej (26). The aim of this second step is to identify optimal van der Waals radii within the PB/BEM approach that will lead to boundaries and solute geometries of similar size and shape as those used in the high-level PCM calculations. Finally, with the optimized parameters determined in the initial two stages we compute actual PB/BEM polarization energies in order to obtain a close match with the quantum chemical results obtained from the reference set.

#### 2 Methods

## 2.1 Sample Selection, Preparation and Set Up of Structures and Computation of Molecular Surfaces with Different Programs

A set of different protein structures is randomly selected from the Protein Data Bank (27). The actual download site used is the repository PDB-REPRDB (28). Default options are applied with the following exceptions: i) Number of residues less than 40 excluded – NO, ii) Include MUTANT – NO, iii) Exclude COMPLEX, iv) Exclude FRAG-MENT, v) Include NMR – NO, vi) Include Membrane Proteins – NO. A total of 28 structures of different protein sizes and shapes (see table I) are chosen. The PDB codes of the samples are, 2ERL, 1P9GA, 1FD3A, 1N13E, 1BRF, 1PARB, 1K6U, 1AVOA, 1SCMA, 1OTFA, 1DJTA, 1KU5, 1K3BC, 1R2M, 1CC8, 1L9LA, 1ZXTD, 1GYJA, 1T8K, 1XMK,

#### 1YNRB, 1EZGA, 1C5E, 1SAU, 1WN2, 1JBE, 1C7K and 1WKR.

Four different programs to calculate molecular surfaces have been employed: the Connolly program **MSROLL** (18), the **SIMS** program (29), the **MSMS** program by Sanner (30) and a self-written program based on an estimation of the SASA (31). However, we detected already in early stages of this investigations problems with varying van der Waals radii and found indications that the SASA is not an appropriate choice for this kind of application. Hence, we have dropped the latter two programs, and most of the analysis is done with programs **MSROLL** and **SIMS**.

Downloaded PDB structures are cleaned from multichain entries, HETATM lines CONNECT lines, ANISOU lines, counter ions, water molecules and the footer section. Program MOLDEN (32) is used to visualize the downloaded PDB structures after cleaning and the force field **Tinker Amber** is selected before a new PDB file is written out from within MOLDEN using option 'Write\_With\_Hydrogens'. Since MOLDEN always uses the default HIP-type in AMBER jargon, HIS residues need to be converted to HIP types, as well as CYS residues engaged in disulfide bonds need to be converted to CYX-type residues. Occasional cases with PRO being the initial residue are manually edited and initial PROs removed. AMBER non-bonded parameters (33), i.e. charges and van der Waals radii are assigned to all the atoms in the protein structures. In this first part of the study, the vdW-radii are increased by a factor of 1.12 and atomic partial charges are scaled down by another factor of 0.9 (34).

The MSROLL program is used with varying choices of the fineness value (the **-f** command line argument) which defines the resolution of the surface. With smaller values the resolution of the surface becomes better but computational cost will increase. The probe radius (the **-p** command line argument) is set to 1.5 Å. Analytically calculated SASA and molecular volumes are recorded, and the data file containing triangulation details is translated into a human readable format, and critical items (for example almost coinciding triangles) removed.

The SIMS program is used with identical arguments to those employed in MSROLL.

Similarly, varying the resolution of the surface triangulation into small sized triangles means adjusting the **dot-density** parameter in SIMS. Higher values for this parameter will yield higher surface resolutions but also increase the computational demand. We record the number of BE, number of iterations, SASA and volume for comparison.

# 2.2 Computation of Polarization Free Energies, $\Delta G^{Pol}$ , Based on solutions to the Poisson Boltzmann Equation

Inner/outer dielectric constants at the molecular boundary are set to 1.0 and 80.0 respectively. The serial version of the PB/BEM program **POLCH** (35) is used. Critical cases with additional secondary cavities located in the interior part of the proteins are excluded. AMBER van der Waals radii and partial charges (33) are applied. Using our own tool chain for the assignment allows us to conveniently scale these data, as well as to write out in the same instance the corresponding parameter files required by the molecular surface programs.

### 2.3 Density Functional Theory Calculations Including the Polarizable Continuum Model on a Reference Set of Dipeptides

The most prominent combinations of peptidic  $\Phi$ ,  $\Psi$ -angles (36) are used to construct different conformations of dipeptides. Only homodimers are considered. All 20 types of different amino acids are used for this combinatorial approach. Zwitter-ionic forms are built and 9 conformations per class of amino acid are taken into account leading to all in all 180 structures. Program "protein.x" from the TINKER package version 4.2 is employed (37). Each of these reference structures is subjected to PCM (3) calculations at the Becke-98 (25) level of density functional theory (DFT) using the high-quality basis set of Sadlej (26) within the Gaussian-03 suite of programs (38). Geometric properties, i.e. the molecular volume and the molecular surface area, as well as polarization free energies are extracted from the reference calculations and used as a base line when comparing to PB/BEM data. The computational demand of these reference calculations is significant. For example, WW- conformations require on the order of 6 weeks (and beyond) single-processor time on modern computing architectures.

#### 3 Results

## 3.1 Stage I: Rather small-sized BEs are Needed to Obtain Consistently Convergent Polarization Free Energies $\Delta G^{Pol}$

We start with PB/BEMA calculations for a set of protein structures (PDB codes summarized in Table I). The boundary discretization is achieved with two independent programs, MSROLL (18) and SIMS (29). Boundary resolution into BEs is steadily increased with either program and independent PB/BEM results are computed for each particular boundary decomposition. A typical plot of the trend of  $\Delta G^{Pol}$  as a function of number of BEs is shown in Figure 1 for the protein structure with PDB code 1C5E. Similar plots for the other examples in Table I are provided as supplementary material. Both approaches converge to identical results in the limit of large numbers of BEs. The importance of well-resolved boundaries becomes clear from Figure 1. Errors on the order of  $\pm 40 \frac{kcal}{mol}$  are easily introduced when working in the non-converged domain. Connolly's MSROLL program (red triangles in Figure 1) reaches a plateau value in a continuous manner, while the SIMS program (blue spheres in Figure 1) finds its limit value within an alternating sequence. The SIMS program reaches convergence much faster than the Connolly program. The quality of the computed molecular boundaries is comparable, see, for instance, the values of molecular surfaces and volumes (final two columns in Table I) obtained with either program. SIMS seems to overestimate the volume by a small margin of roughly 1%. The recommended average size of BEs for converged results using MSROLL is on the order of 0.11 Å<sup>2</sup> while SIMS would require an average size of 0.31 Å<sup>2</sup>. Both numbers are close to the value of 0.4 Å<sup>2</sup> advocated in Quantum Chemistry (39).

## 3.2 Stage II: Systematic Geometric Comparison to High Level Quantum Chemistry Calculations Suggests a Uniform Scaling of AMBER van der Waals Radii by a Factor of 1.07

A reference set of dipeptides in different conformations (9 per species) is constructed. Only homodipeptides comprising all 20 types of naturally occurring amino acids are considered. Thus a total number of 180 dipeptidic reference structures is set up. The zwitterionic form is used throughout. Each of these structures is computed at the Becke-98 level of theory (25) using the basis set of Sadlej (26) and the PCM model (3) for solvation free energies. Geometric properties such as the cavity volume and the cavity surface area are extracted from each of the reference calculations. All 180 structures are also computed within the PB/BEM approach using optimized parameters for the boundary resolution determined in Stage I of this study. However, only the SIMS program is used. We define a global deviation from the reference data by

$$\Delta^{Surf} = \frac{1}{20} \sum_{i=1}^{20} \frac{1}{9} \sum_{j=1}^{9} \sqrt{(Surf_{i,j}^{PCM} - Surf_{i,j,\alpha}^{PB/BEM})^2}$$
(1)

where j runs over the conformations and i over the different types of homodipeptides, i.e. GG, AA, VV, etc. The parameter  $\alpha$  refers to a specific scaling factor used when constructing the boundaries within the PB/BEM approach. In particular this scaling makes the van der Waals radii larger or smaller by a certain fraction. The AMBER default set of van der Waals radii is used (33). A similar criterion is used for comparing molecular volumes,

$$\Delta^{Vol} = \frac{1}{20} \sum_{i=1}^{20} \frac{1}{9} \sum_{j=1}^{9} \sqrt{(Vol_{i,j}^{PCM} - Vol_{i,j,\alpha}^{PB/BEM})^2}$$
(2)

and the dependence on the scaling factor  $\alpha$  is shown in Figures 2 and 3. As becomes clear from Figures 2 and 3 the best match to the reference data is obtained when scaling the AMBER van der Waals radii by a factor of 1.07. Detailed data with respect to conformational averages per type of dipeptide are shown in Tables II and III.

#### 3.3 Stage III: Charge Scaling is Not Required

Using the optimized parameters obtained in the previous two stages leads us to the final step of directly comparing polarization free energies  $\Delta G^{Pol}$  computed within the PB/BEM approximation and at the PCM level of theory. The idea is to identify another uniform scaling factor  $\beta$  which applied to the AMBER default charges would result in an optimal match to the reference polarization free energies. Thus another deviation criterion is introduced,

$$\Delta^{\Delta G^{Pol}} = \frac{1}{20} \sum_{i=1}^{20} \frac{1}{9} \sum_{j=1}^{9} \sqrt{(\Delta G_{i,j}^{Pol,PCM} - \Delta G_{i,j,\beta}^{Pol,PB/BEM})^2}$$
(3)

that allows to identify the optimal value of  $\beta$ . The dependence of the PB/BEM polarization free energies on the charge scaling factor  $\beta$  is shown in Figure 4. The trend shown in Figure 4 suggests an optimal value of  $\beta$  very close to 1.0, hence no charge scaling is required. This result i) emphasizes the broad applicability of AMBER partial charges and ii) circumvents conceptual difficulties that would arise when charges had to be scaled, i.e. modified net charges in proteins, non-neutral forms, etc. A detailed analysis with respect to the magnitude of the average deviation of each particular type of dipeptide studied is shown in Table IV.

#### 4 Discussion

Motivated by our recent high-performance solution to Poisson Boltzmann calculations (35) we have now tested the influence of the many critical parameters involved. One obvious issue is the exact choice and composition of the boundary between solute and solvent. At first, we have to ensure the numerical stability within the selected level of approximation. In order to address this problem we have carried out PB/BEM calculations on a large sample of different proteins. When using different programs to create the boundary surface and increasing systematically the resolution of these surfaces into small-sized boundary elements, a recommended threshold size of about 0.31 Å<sup>2</sup> for the average BE is identified when using program SIMS (29) which showed faster convergence than the well-known Connolly program (18). Although giving rise to very fine-resolved boundary surfaces, hence large numbers of

BEs, this value is close to the corresponding value of  $0.4 \text{ Å}^2$  frequently advised in Quantum Chemical models (39). As a consequence, even proteins of modest size thus require consideration of vast numbers of BEs (see for example Table I), and the importance of efficient means of solving the computational problem is underlined again.

After having established the necessary degree of boundary partitioning in the first stage, we performed a systematic comparison against a reference set of dipeptides computed at a high level of Quantum Chemical theory. Consideration of geometric factors revealed that when applying a scaling factor of about 1.07 to AMBER default van der Waals radii, rather good agreement can be reached between the reference geometries and the geometries in the PB/BEM approach. The recommended value of 1.07 is somewhat smaller than a factor found previously (1.12 of ref (34)) and reflects the much finer resolved boundary surfaces used in this present work.

The final step was to compare actual calculations of the polarization free energies to each other. Following previous attempts, we wanted to derive another scaling factor that, when applied to AMBER partial charges, would yield a close match to the reference polarization free energies. The trend visible in Figure 4 indicates that no scaling of the charges is necessary: they are already close to optimal. This is an unexpected — but very welcome — result, as it eliminates potential secondary problems that would emerge with modifying charges. Again, this is another consequence of the much finer resolved boundary surfaces in this present work as opposed to previous results (34) where a scaling factor of 0.9 had been found.

#### 5 Conclusion

Combined employment of small-sized BEs ( $\approx 0.3 \text{ Å}^2$  on average), slightly increased AMBER van der Waals radii (by a factor of 1.07), and default AMBER partial charges leads to good quality estimates of the polarization free energy,  $\Delta G^{Pol}$ , for proteins within the PB/BEM framework. **Acknowledgements** This work was supported in part by the National Institutes of Health Grant GM62838 and by a grant of the National Science Foundation (CHE-0313618). Dr. Michael Connolly is acknowledged for providing a test version of his molecular surface program.

#### References

- 1. Tomasi, J.; Persico, M. Chem Rev 1994, 94, 2027.
- 2. Cramer, C. J.; Truhlar, D. G. Chem Rev 1999, 99, 2161.
- 3. Tomasi, J.; Mennucci, B.; Cammi, R. Chem Rev 2005, 105, 2999.
- 4. Roux, B.; Simonson, T. Biophys Chem 1999, 78, 1.
- 5. Rinaldi, D.; Ruiz-López, M. F., Rivail, J. L. J Chem Phys 1983, 78, 834.
- 6. Warwicker, J.; Watson, H. C. J Mol Biol 1982, 157, 671.
- 7. Honig, B.; Nicholls, A. Science 1995, 268, 1144.
- 8. Luo, R.; David, L.; Gilson, M. K. J Comput Chem 2002, 23, 1244.
- Baker, N. A.; Sept, D.; Joseph; S.; Holst, M. J.; McCammon, J. A. Proc Natl Acad Sci USA 2001, 98, 10037.
- 10. Bashford, D.; Karplus, M. Biochemistry-US 1990, 29, 10219.
- 11. Tironi, I.; Sperb, R.; Smith, P. E.; van Gunsteren, W. F. J Chem Phys 1995, 102, 5451.
- 12. Zauhar, R. J.; Morgan, R. S. J Mol Biol 1985, 186, 815.
- Juffer, A. H.; Botta, E. F. F.; van Keulen, B. A. M.; van der Ploeg, A.; Berendsen, H. J. C. J Comput Phys 1991, 97, 144.
- Still, W. C.; Tempczyk, A.; Hawley, R. C.; Hendrickson, T. J Am Chem Soc 1990, 112, 6127.
- 15. Onufriev, A.; Bashford, D.; Case, D. A. J Phys Chem B 2000, 104, 3712.
- Curutchet, C.; Orozco, M.; Luque, J. F.; Mennucci, B.; Tomasi, J. J Comput Chem 2006, 27, 1769.

- 17. Carlson, H. A.; Jorgensen, W. L. J Phys Chem 1995, 99, 10667.
- Connolly, M. L. J Am Chem Soc 1985, 107, 1118. http://www.netsci.org/Science/Compchem/feature14e.html
- 19. Lii, J.-H.; Allinger, N. L. J Am Chem Soc 1989, 111, 8576.
- 20. Aqvist, J. J Phys Chem 1990, 94, 8021.
- Weiner, S. J.; Kollman, P. A.; Case, D. A.; Singh, U. C.; Ghio, C.; Alagona, G.; Profeta, S.; Weiner, P. J Am Chem Soc 1984, 196, 765.
- 22. Halgren, T. J Am Chem Soc 1992, 114, 7827.
- Kollman, P. A.; Massova, I.; Reyes, C.; Kuhn, B.; Huo, S.; Chong, L.; Lee, M.; Lee, T.; Duan, Y.; Wang, W.; Donini, Q.; Cieplak, P.; Srinivasan, J.; Case, D. A.; Cheatham III, T. E. Acc Chem Res 2000, 33, 889.
- 24. Page, C. S.; Bates, P. A. J Comput Chem 2006, 27, 1990.
- 25. Becke, A. D. J Chem Phys 1997, 107, 8554.
- 26. Sadlej, A. J. Theor Chim Acta 1991, 79, 123.
- Berman, H. M.; Westbrook, J.; Feng, Z.; Gilliland, G.; Bhat, T. N.; Weissig, H.; Shindyalov, I. N.; Bourne; P. E. Nucleic Acids Res 2000, 28, 235.
- Akiyama, Y.; Onizuka, K.; Noguchi, T.; Ando, M. Proc. 9th Genome Informatics Workshop (GIW'98), Universal Academy Press, 1998, ISBN 4-946443-52-5, 131-140 http://mbs.cbrc.jp/pdbreprdb-cgi/reprdb\_menu.pl
- 29. Vorobjev, Y. N.; Hermans, J. Biophys J 1997, 73, 722.
- Sanner, M. F.; Olson, A. J.; Spehner, J.-C. Proc. 11th ACM Symp Comp Geom, 1995, C6-C7.

- Hayryan, S.; Hu, C.-K.; Skřivánek, J.; Hayryan, E.; Pokorný, I. J Comput Chem 2005, 26, 334.
- 32. Schaftenaar, G.; Noordik, J. H. J Comput Aid Mol Des 2000, 14, 123.
- Cornell, W. D.; Cieplak, P.; Bayly, C. I.; Gould, I. R.; Merz, K. M.; Ferguson, D. M.; Spellmeyer, D.; Fox, T.; Caldwell, J. W.; Kollman, P. A. J Am Chem Soc 1995, 117, 5179.
- Höfinger, S. Modelling Molecular Structure and Reactivity in Biological Systems, RSC Publishing, 2006, ISBN 0-85404-668-2, 151.
- 35. Höfinger, S. J Comput Chem 2005, 26, 1148.
- 36. Marshall, N. J.; Grail, B. M. J Pept Sci, 2000, 6, 186.
- 37. Ren, P.; Ponder, J. W. J Phys Chem B 2003, 107, 5933.
- Gaussian 03, Revision B.05, Frisch, M. J.; Trucks, G. W.; Schlegel, H. B.; Scuseria, G. E.; Robb, M. A.; Cheeseman, J. R.; Montgomery Jr., J. A.; Vreven, T.; Kudin, K. N.; Burant, J. C.; Millam, J. M.; Iyengar, S. S.; Tomasi, J.; Barone, V.; Mennucci, B.; Cossi, M.; Scalmani, G.; Rega, N.; Petersson, G. A.; Nakatsuji, H.; Hada, M.; Ehara, M.; Toyota, K.; Fukuda, R.; Hasegawa, J.; Ishida, M.; Nakajima, T.; Honda, Y.; Kitao, O.; Nakai, H.; Klene, M.; Li, X.; Knox, J.E.; Hratchian, H. P.; Cross, J. B.; Bakken, V.; Adamo, C.; Jaramillo, J.; Gomperts, R.; Stratmann, R. E.; Yazyev, O.; Austin, A. J.; Cammi, R.; Pomelli, C.; Ochterski, J. W.; Ayala, P. Y.; Morokuma, K.; Voth, G. A.; Salvador, P.; Dannenberg, J. J.; Zakrzewski, V. G.; Dapprich, S.; Daniels, A. D.; Strain, M. C.; Farkas, O.; Malick, D. K.; Rabuck, A. D.; Raghavachari, K.; Foresman, J. B.; Ortiz, J. V.; Cui, Q.; Baboul, A. G.; Clifford, S.; Cioslowski, J.; Stefanov, B. B.; Liu, G.; Liashenko, A.; Piskorz, P.; Komaromi, I.; Martin, R. L.; Fox, D. J.; Keith, T.; Al-Laham, M. A.; Peng, C. Y.; Nanayakkara, A.; Challacombe, M.; Gill, P. M. W.; Johnson, B.; Chen, W.; Wong, M. W.; Gonzalez, C.; Pople, J. A. Gaussian, Inc., Wallingford CT, 2004.

 European Summerschool in Quantum Chemistry, Book II, Björn O. Roos and Per-Olof Widmark, Eds, 2000, 561.

## List of Figures

1	PB/BEM derived polarization free energies $\Delta G^{Pol}$ as a function of boundary	
	resolution obtained from two independent programs MSROLL $\left(18\right)$ and SIMS	
	(29). The example represents results for PDB structure 1C5E	17
2	Comparison of employed molecular surfaces in the PB/BEM series based on	
	scaling the AMBER default van der Waals radii by a factor $\alpha$ to the reference	
	data obtained from PCM calculations at the Becke-98/Sadlej level. $\ . \ . \ .$	18
3	Comparison of employed molecular volumes in the PB/BEM series based on	
	scaling the AMBER default van der Waals radii by a factor $\alpha$ to the reference	
	data obtained from PCM calculations at the Becke-98/Sadlej level. $\ . \ . \ .$	19
4	Comparison of PB/BEM polarization free energies $\Delta G^{Pol}$ based on scaling the	
	AMBER default charges by a factor $\beta$ to the reference data obtained from	
	PCM calculations at the Becke-98/Sadlej level.	20



Figure 1:





Scaling Factor  $\alpha$  for vdW Radii

Figure 2:



Scaling Factor  $\alpha$  for vdW Radii

Figure 3:



Scaling Factor  $\beta$  for AMBER Partial Charges

Figure 4:

Table I: PDB codes of studied structures and corresponding number of residues. Columns three and four show the number of BEs needed to reach converged PB/BEM results using molecular surface algorithms MSROLL and SIMS respectively. Final two columns show average molecular surface areas and average molecular volumes derived when using data of both programs, MSROLL and SIMS. The differences obtained when subtracting the SIMS results from the MSROLL results are given in parentheses.

PDB Code	Number of	Number of BEs	Number of BEs	Molecular	Solvent
	Residues	at Convergence	at Convergence	Surface Area	Excluded Volume
		Using MSROLL	Using SIMS	(Difference)	(Difference)
		(18)	(29)	$[Å^2]$	$[Å^3]$
2ERL	40	15661	9807	2370 (+1)	5653 (-43)
1P9GA	41	22302	5751	2091 (-5)	5055 (-72)
1FD3A	44	25865	6699	2408 (+7)	5819 (-56)
1N13E	52	18419	10353	3750 (+11)	6542 (-69)
1BRF	53	33879	11810	2796 (-7)	7734 (-77)
1PARB	53	42336	11006	3968 (-8)	8509 (-108)
$1 \mathrm{K6U}$	58	24220	13406	3195 (+12)	8603 (-65)
1AVOA	60	43916	13335	4777 (-3)	9325 (-186)
1SCMA	60	54464	14603	5131 (-13)	$10601 \ (-179)$
10TFA	62	40128	10610	3767 (+6)	8942 (-86)
1DJTA	64	35828	9134	3331 (+26)	9422 (-43)
$1 \mathrm{KU5}$	66	46390	12208	4310 (+3)	10153 (-133)
1K3BC	69	61667	16297	5768 (+19)	18193 (-165)
1R2M	71	39316	13659	3244 (-1)	9596 (-74)
1CC8	73	27668	15091	3644 (+2)	11094 (-65)
1L9LA	74	20278	11636	4182 (+5)	11728 (-112)
1ZXTD	76	37335	11259	4089 (+8)	10809 (-93)
1GYJA	76	44770	13665	4885 (-3)	11464 (-118)
1T8K	77	35978	13846	3925 (+4)	11410 (-119)
1XMK	79	56033	16468	4294 (+9)	12288 (-98)
1YNRB	80	31529	12630	4417 (+16)	11911 (-157)
1EZGA	84	34628	9122	3258 (+3)	10103 (-95)
1C5E	95	48306	19880	4480(0)	13285 (-110)
1SAU	115	47613	17765	5197 (+25)	17897 (-116)
1WN2	121	51325	21555	5614 (+15)	17836 (-118)
1JBE	128	58119	16729	5409(+22)	18905 (-188)
1C7K	132	54104	16675	5389~(0)	18858 (-182)
1WKR	340	74167	55378	11008 (-41)	47105 (-299)

Table II: Comparison of average molecular surfaces based on unscaled and scaled AMBER van der Waals radii with data from PCM reference calculations. The average comprises 9 different conformations per type of dipeptide considered. The value in parentheses monitors that conformation that deviates most severely from the mean.

Dipeptide	Mean Surface	Mean Surface	Mean Surface	
Type	AMBER Unscaled	AMBER Scaled	PCM Reference	
	$[Å^2]$	$[Å^2]$	$[Å^2]$	
AA	191.764(5.205)	204.388(3.697)	214.747(4.955)	
$\mathbf{C}\mathbf{C}$	214.592(4.622)	229.274(5.461)	232.910(6.522)	
DD	228.687(6.032)	242.724(5.972)	240.839(6.776)	
EE	273.402(5.835)	287.557(6.951)	283.025(5.766)	
GG	150.255 (4.460)	161.637(4.936)	167.764(3.249)	
II	279.535(10.274)	294.402 (11.331)	302.173 (12.233)	
KK	314.712(6.273)	332.593(7.677)	340.545(7.599)	
LL	276.435(10.012)	290.497(12.440)	293.172(10.465)	
MM	301.384(6.691)	318.033(8.261)	329.815(8.374)	
NN	232.466(6.072)	247.553(7.325)	247.040(7.310)	
QQ	276.939(5.957)	293.663(7.531)	292.648(6.582)	
RR	354.636(6.925)	377.310 (7.568)	380.408(6.739)	
SS	196.528(4.982)	207.904(4.610)	212.107(4.695)	
TT	$224.181 \ (8.047)$	238.570(8.359)	239.112 (10.124)	
VV	251.913(8.574)	265.008(8.296)	276.423(8.140)	
YY	340.272 (17.100)	356.042(17.293)	346.378 (14.704)	
$\mathbf{FF}$	329.058(17.123)	343.245 (17.402)	326.947(15.489)	
WW	355.790(26.425)	377.209(27.865)	361.573(25.182)	
HH	282.802 (13.007)	299.235(12.829)	296.885(11.461)	
PP	224.999 (10.768)	237.525 (10.500)	233.118 (9.900)	

Table III: Comparison of average molecular volumes based on unscaled and scaled AMBER van der Waals radii with data from PCM reference calculations. The average comprises 9 different conformations per type of dipeptide considered. The value in parentheses monitors that conformation that deviates most severely from the mean.

Dipeptide	Mean Volume	Mean Volume	Mean Volume	
Type	AMBER Unscaled	AMBER Scaled	PCM Reference	
	$[Å^3]$	$[Å^3]$	$[Å^3]$	
AA	191.804 (4.553)	215.661 (3.355)	228.591(2.885)	
$\mathbf{C}\mathbf{C}$	223.713(3.799)	252.135(4.069)	255.902(3.594)	
DD	242.406(4.242)	270.928(5.037)	260.904(4.723)	
EE	296.037 (4.888)	327.303(3.864)	311.645(3.810)	
$\mathbf{FF}$	381.346(5.875)	417.696(7.992)	388.363(6.223)	
GG	136.315 (3.565)	154.314(2.631)	164.287(2.048)	
HH	317.309(3.605)	353.465(5.228)	340.591(5.324)	
II	323.590(6.689)	357.080(7.564)	367.182(6.520)	
KK	343.720(3.747)	382.391(4.384)	388.641 (3.903)	
LL	313.537(5.063)	346.507(6.450)	344.950(7.458)	
MM	325.073(5.069)	363.563(5.484)	377.789(3.905)	
NN	248.729(4.628)	278.449(4.968)	271.305(5.415)	
PP	$242.861 \ (8.154)$	$269.561 \ (9.116)$	263.128(9.957)	
QQ	301.279(3.985)	336.468(5.590)	325.660(5.591)	
RR	384.833(4.089)	431.811 (4.444)	424.874(3.559)	
SS	$198.981 \ (3.533)$	221.590(3.162)	225.193(2.964)	
TT	244.397(5.673)	272.848(7.051)	273.546(6.533)	
VV	282.296(6.114)	311.383(6.895)	330.378(8.172)	
WW	431.248 (14.910)	480.974 (17.118)	447.693 (15.019)	
YY	393.622 (5.433)	433.845 (7.433)	407.695 (5.466)	

Table IV: Comparison of average PB/BEM polarization free energies  $\Delta G^{Pol}$  using AMBER default charges to corresponding data obtained from PCM reference calculations. The average comprises a variable number of conformations for each different type of dipeptide considered depending on how many PCM calculations terminated faithfully. Numbers in parentheses reflect maximum deviation from the mean.

Dipeptide	Mean $\Delta G^{Pol,PB/BEM}$	Mean $\Delta G^{Pol,PCM}$	Mean $\Delta\Delta G^{Pol}$	Number of
Type	AMBER Default Charges	PCM Reference	Deviation	References
	$[\rm kcal/mol]$	$[\rm kcal/mol]$	$[\rm kcal/mol]$	
AA	-91.36 ( 8.56 )	-83.89 (10.12)	7.47	9
CC	-115.11 (11.02)	-96.80 (12.83)	18.31	9
DD	-296.25 (17.08)	-285.27 (18.24)	10.98	9
EE	-266.54 (14.09)	-259.29(13.76)	7.25	9
GG	-96.52(10.09)	-89.41 (11.36)	7.11	9
II	-82.72 (7.49)	-75.97 (8.70)	6.77	9
KK	-249.63 (16.51 )	-236.37 (19.64)	13.26	9
LL	-85.54 (7.20)	-64.51 (8.64)	21.03	9
MM	-88.82 (7.55)	-82.10 ( 9.42 )	6.72	9
NN	-105.11 (8.19)	-101.80 (12.20)	4.25	9
QQ	-119.08 (10.88)	-115.33 (12.60)	3.89	9
RR	-235.39 (17.79)	-228.71 (21.45)	6.68	6
SS	-112.78 (13.38)	-105.47(13.90)	7.32	9
TT	-106.87 (12.06)	-100.61 (12.88)	6.55	9
VV	-85.17 (7.44)	-77.46 (8.73)	7.70	9
YY	-93.35 ( 4.36 )	-90.11 ( 8.48 )	3.59	5
$\mathbf{FF}$	-89.92(10.55)	-82.51 (13.94)	7.41	6
HH	-237.74 (19.08)	-236.06 (22.15)	3.66	9
PP	-79.15 ( 5.73 )	-82.71 (7.50)	3.56	9
WW	-100.50 (4.27)	-88.11 (12.93)	12.39	2

List of Figures