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# A Tunable 3-D Statistical Channel Model for Spatio-temporal Characteristics of Wireless Communication Networks 

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#### Abstract

A new three-dimensional (3-D) geometric channel model is proposed with multiple degrees of freedom in its scattering volume geometry, which enables accurate modeling of realistic propagation scenarios for emerging Machine-to-Machine (M2M) wireless networks. Analytical expressions for the joint and marginal statistics in delay and angular domains are derived and we present the impact of various physical channel parameters such as link distance, antenna heights, size and position of scattering region, as well as shape and orientation of scattering region on delay and angular spread of the proposed model. The derived analytical results are validated with simulations and are also shown to provide a good fit to published empirical data-sets. Relative to existing models, the proposed model is shown to more accurately adapt to different propagation environments envisioned for emerging communication scenarios such as in indoor pico-cellular and Vehicle-to-Vehicle (V2V) communication applications. Copyright © 2016 John Wiley \& Sons, Ltd.

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## 1. INTRODUCTION

Significant research interest has been recently devoted to emerging paradigms of wireless networks encountered in Machine-to-Machine (M2M) communications; these include application scenarios such as intelligent transport systems, smart metering, and remote health diagnostics [1, 2, 3]. A real-world performance evaluation of M2M communications necessitates an appropriate understanding of the relevant propagation scenario in terms of the Angle-of-Arrival (AoA) and Time-of-Arrival (ToA) information of multipath components of the wireless propagation channel. For this purpose geometry-based stochastic
channel models have been extensively studied in the literature for conventional cellular systems in which a fixed Base-Station (BS) communicates with a MobileStation (MS) [4, 5, 6]. Since geometric channel models inherently establish probabilistic relations between spatial locations of the Transmit (Tx), Receive (Rx), and scatterers to model wireless propagation, it becomes imperative to investigate the distributions of the multipath ToA, AoA, and Angle-of-Departure (AoD) to gain insights into the physics of any novel geometric channel model. For this reason the problem of determining the multipath delay and angular distributions has received significant attention in the literature, see for example $[7,8,9,10]$ and
references therein. Angular and temporal characteristics of a radio propagation channel are necessary for a realistic performance evaluation and designing of a mobile radio communication system. The spread of a channel in temporal and angular domains helps in determining frequency and space selective nature of the channel, respectively. The angular spread quantified in terms of multipath shape factors (i.e., angular Spread, angular constriction and direction of maximum fading) helps in determining second order fading statistics of the channel; the second order fading statistics include, level crossing rate, average fade duration, coherence distance, and spatial correlation [11, 12]. Such characterization of fading statistics assists in devising efficient schemes at Tx and Rx nodes for interleaving, coding, and equalization algorithms.

Unlike conventional cellular systems that have BS antenna heights above or at rooftop levels, the antenna heights of M2M nodes are expected to be typically much lower than the surrounding scatterers such as buildings, trees, and other objects. For example in Vehicle-to-Vehicle (V2V) channels, which constitute a special case of M2M communications, this differing propagation environment has led to many investigations into appropriate twodimensional (2-D) and three-dimensional (3-D) models [13, 14, 15, 16, 17, 18]. In [13], a uniform distribution of scatterers in circular regions is assumed around each MS and the ToA and AoA statistics are investigated. In [14], a scattering model for V2V scenarios is presented in which the scatterers are confined to elliptical regions and closedform expressions for the distribution of AoA and ToA are derived. In [15], a more general 2-D elliptical model is proposed for V 2 V channels by assuming the scatterers to be uniformly distributed within adjustable hollow ellipses around each MS and the distribution of the AoA at each MS is then derived.

Subsequent investigations have revealed that low antenna heights of the communicating nodes result in signal reflections from nearby scatterers that are more elevated than the communicating antennas. This makes it necessary to include the elevation plane for an accurate description of the relevant propagation channels. In V2V channel modeling efforts, various 3-D channel models have been previously proposed for a more precise spatiotemporal description of the propagation scenario, see [16, 17, 18] and references therein. In [16], a 3-D
multiple-input-multiple-output model for V2V channels is proposed and the authors derive a closed-form joint space-time correlation function for the considered 3D scattering environment; the effect of space-time correlation on outage capacity of different antenna array configurations is then studied. In [17], a 3-D semiellipsoid model is proposed to effectively model V2V channels for urban street canyons. The scatterers are assumed to be uniformly distributed within semi-ellipsoids around each MS and the authors derive closed-form expressions for the AoA jointly in azimuth and elevation. Furthermore, for various street orientations and different elevations of the scatterers the marginal probability density functions (PDFs) of the azimuth and elevation angles are also presented. In [18], the temporal characteristics of a 3-D semi-ellipsoid geometrical channel model for V2V communications is presented. Assuming a uniform scatterer distribution around the MSs, the authors derive closed-form expressions for the joint PDF of the ToA and AoA in azimuth and elevation planes.

The present work is a generalization of the 3-D geometric channel modeling approach of $[17,18]$ in that a more flexible placement of communicating nodes, relative to their surrounding scatterers, is proposed to admit M2M communication scenarios. Specifically, the novel contributions of this work are:

1. The maximum propagation delay is incorporated into the model through its associated maximum-delay-ellipsoid. The overlap of this ellipsoid with the scattering ellipsoids at Tx and Rx determine the effective scattering volumes at respective link ends.
2. The scattering ellipsoids at Tx and Rx side are independently rotatable in the horizontal plane, about their respective origins. Moreover, the origin of each of these two ellipsoids can be arbitrarily shifted along the vertical axis of the coordinate system.
3. The antennas at Tx and $R x$ can have a vertical offset relative to the center of their respective scattering regions.
4. Closed-form expressions are derived for the delay distribution, joint and marginal distributions of the AoA in azimuth and elevation of the multipath components of the modeled channel.

Table I: Proposed Model's Input Parameters

| Symbol | Definition |
| :---: | :--- |
| $h_{m}$ | Elevation of antenna employed at node $m$. |
| $z_{m}$ | Vertical displacement of the origin of scattering volume (ellipsoid) around Node $m$ with <br> reference to the ground level. |
| $d$ | Horizontal separation between $N_{1}$ and $N_{2}$. |
| $a_{m}, b_{m}$, and $c_{m}$ | Major, intermediate, and minor axes of the bounding ellipsoid at node $m$. |
| $\theta_{m}$ | Rotation angle (in horizontal plane) in the orientation of bounding ellipsoid around node $m$. |
| $\tau_{o}$ and $\tau_{\max }$ | Delay associated with LoS and longest propagation path (Minimum and maximum delay). |

The remainder of this paper is organized as follows. Sec. II describes the system model of the proposed 3D geometrical channel model. Furthermore, closed-form expressions for the distributions of delay and AoA are also derived in this section. In Sec. III, a validation of the theoretical results with simulations and published measurement data is presented with relevant discussions. Finally, Sec. IV concludes this work.

## 2. SYSTEM MODEL

The proposed model provides a high order of Degrees-of-Freedom (DoF) in it's parameters, which enables adjustability of the model according to the targeted propagation environment to obtain realistic channel characteristics. The geometrical composition of the proposed model is illustrated in Fig. 1. Scattering objects in the vicinity of communicating nodes $\mathrm{N}_{m}$ are modeled as bounded within (triaxial) ellipsoidal shaped regions denoted by $\mathrm{E}_{m}$, where, the subscript $m$ takes value as, $m=1$ and $m=2$ to represent the nodes 1 and 2 , respectively. The independently adjustable major, intermediate, and minor dimensions of these ellipsoids are represented by $a_{m}, b_{m}$, and $c_{m}$, respectively. Moreover, these scattering ellipsoids ( $\mathrm{E}_{m}$ ) are modeled as independently rotatable about their origin in $\mathrm{x}-\mathrm{y}$ plane with angles $\theta_{m}$, where $-\pi / 2 \leq \theta_{m} \leq \pi / 2$. The horizontal position of the origin of these scattering ellipsoids $\left(\mathrm{E}_{m}\right)$ is fixed at the horizontal position of communication nodes $\left(\mathrm{N}_{m}\right)$; however, their origin in vertical plane is designed as shiftable with parameter $z_{m}$. The elevation of antennas is denoted by $h_{m}$. Difference between the elevation of antennas employed at $\mathrm{N}_{1}$ and $\mathrm{N}_{2}$ can thus be found as, $h_{e}=h_{1}-h_{2}$. The angles observed at a certain node
$\mathrm{N}_{m}$ along the arriving multipath component corresponded from a certain scattering point $\mathrm{s}_{\mathrm{p}}$, are denoted by $\phi_{m}$ and $\beta_{m}$ in azimuth and elevation planes, respectively. The range of azimuth and elevation AoA parameters is, $-\pi \leq$ $\phi_{m}<\pi$ and $0 \leq \beta_{m} \leq \pi / 2$, respectively. The input geometric parameters are enlisted in Table I. Distance of a certain scattering point $s_{p}$ from a communication node $\mathrm{N}_{m}$ is denoted by $d_{m}$. Length of the propagation path corresponded from $\mathrm{p}^{\text {th }}$ scattering point can thus be found as, $l_{p}=d_{1}+d_{2}$; subsequently, propagation delay of this path can be obtained as $\tau_{p}=l_{p} / c$. The parameter $c$ represents the velocity of propagation of electromagnetic waves. The x -axis of the coordinate system is fixed along the azimuth LoS direction (i.e., azimuth LoS angle is fixed at $0^{\circ}$ ). The elevation LoS angle can be obtained from the following equation,

$$
\begin{equation*}
\beta_{\mathrm{LoS}}=\arctan \left(\frac{h_{e}}{d}\right) \tag{1}
\end{equation*}
$$

The horizontal and Line-of-Sight (LoS) distance from $\mathrm{N}_{1}$ to $\mathrm{N}_{2}$ is $d$ and $d_{\text {LoS }}$, respectively.

$$
\begin{equation*}
d_{\mathrm{LoS}}=\sqrt{d^{2}+h_{e}^{2}} \tag{2}
\end{equation*}
$$

One of the unique properties of an (oblate) ellipsoid is that the sum of distances from its (two) foci points to any point on its surface is the same. Exploiting this property, an ellipsoid $\left(\xi_{\tau}\right)$ can be drawn corresponding to a certain propagation delay $\tau$ with its foci points at $\mathrm{N}_{1}$ (Tx) and $\mathrm{N}_{2}$ (Rx). The effective scattering region (ESR) for delay $\tau$ is thus the overlapping portion of the scattering ellipsoids ( $\mathrm{E}_{m}$ ) with the surface of virtual bounding (oblate) ellipsoid $\left(\xi_{\tau}\right)$. Similarly, the overall boundary of ESR can be defined by an ellipsoid corresponding to longest propagation delay $\tau_{\text {max }}$. The ESR is thus the union of $\left.\mathrm{R}_{\tau 1}\right|_{\tau=\tau_{\text {max }}}$ and $\left.\mathrm{R}_{\tau 2}\right|_{\tau=\tau_{\text {max }}}$, truncated by the ground plane. The major and


Figure 1: Geometrical composition of proposed 3-D scattering model.
minor dimensions of the bounding (oblate) ellipsoid $\left(\xi_{\tau}\right)$ can be calculated from the given delay $\tau$, as given under,

$$
\begin{gather*}
a_{\tau}=\frac{c \tau}{2}  \tag{3}\\
b_{\tau}=\frac{1}{2} \sqrt{c^{2} \tau^{2}-d_{\mathrm{LoS}}^{2}} \tag{4}
\end{gather*}
$$

The origin of coordinate system is assumed at the base of communication node $\mathrm{N}_{1}$. The position of $\mathrm{N}_{1}$ and $\mathrm{N}_{2}$ can thus be represented in cartesian coordinates as $\left(0,0, h_{1}\right)$ and $\left(d, 0, h_{2}\right)$, respectively. The scattering ellipsoids $\mathrm{E}_{1}$ can be expressed as,

$$
\begin{align*}
& \frac{\left(x \cos \theta_{1}-y \sin \theta_{1}\right)^{2}}{a_{1}^{2}}+\frac{\left(y \cos \theta_{1}+x \sin \theta_{1}\right)^{2}}{b_{1}^{2}} \\
& +\frac{\left(z-z_{1}\right)^{2}}{c_{1}^{2}}=1 \tag{5}
\end{align*}
$$

The equation for scattering ellipsoid $\mathrm{E}_{2}$ is given as below,

$$
\begin{align*}
& \frac{\left((x-d) \cos \theta_{2}-y \sin \theta_{2}\right)^{2}}{a_{2}^{2}}+\frac{\left(y \cos \theta_{2}+(x-d) \sin \theta_{2}\right)^{2}}{b_{2}^{2}} \\
& +\frac{\left(z-z_{2}\right)^{2}}{c_{2}^{2}}=1 . \tag{6}
\end{align*}
$$

The spherical coordinates system based representation can be obtained by using the transformation relationships, $x=\rho \cos \beta_{m} \cos \phi_{m}, \quad y=\rho \cos \beta_{m} \sin \phi_{m}, \quad$ and $z=$ $\rho \sin \beta_{m}$, where, $\rho$ is the radial coordinate. The virtual bounding ellipsoid $\xi_{\tau}$ can be expressed in cartesian and
spherical coordinates as under,

$$
\begin{align*}
& \frac{\left(\left(x-k_{x}\right) \cos \beta_{\mathrm{LoS}}-\left(z-k_{z}\right) \sin \beta_{\mathrm{LoS}}\right)^{2}}{a_{\tau}^{2}}+\frac{y^{2}}{b_{\tau}^{2}} \\
& +\frac{\left(\left(z-k_{z}\right) \cos \beta_{\mathrm{LoS}}+\left(x-k_{x}\right) \sin \beta_{\mathrm{LoS}}\right)^{2}}{b_{\tau}^{2}}=1 \tag{7}
\end{align*}
$$

where the parameters $k_{x}$ and $k_{z}$ represent the shift from the origin (point k ) of the ellipsoid $\xi_{\tau}$ to the origin (point $o$ ) of the coordinates system along x -axis and z -axis, respectively. These shifts can be expressed as,

$$
\begin{gather*}
k_{x}=-\frac{d_{\mathrm{LoS}}}{2} \cos \beta_{\mathrm{LoS}},  \tag{8}\\
k_{z}=-h_{1}+\frac{d_{\mathrm{LoS}}}{2} \sin \beta_{\mathrm{LoS}} . \tag{9}
\end{gather*}
$$

The ESR is union of regions $\mathrm{R}_{\tau 1}$ and $\mathrm{R}_{\tau 2}$, formed with the overlapping of scattering ellipsoids $\mathrm{E}_{1}$ and $\mathrm{E}_{2}$ with the virtual bounding ellipsoid $\xi_{\tau}$. The volume of region $\mathrm{R}_{\tau 1}$ can be obtained by adopting the approach in [19], i.e., adding the volume of intersection of $\mathrm{R}_{\tau 1}$ with infinite elementary pyramidal frustums obtained from four rays originated from $\mathrm{N}_{2}$ in the directions ( $\phi_{2}, \beta_{2}$ ), $\left(\phi_{2}+\Delta \phi_{2}, \beta_{2}\right),\left(\phi_{2}, \beta_{2}+\Delta \beta_{2}\right)$, and $\left(\phi_{2}+\Delta \phi_{2}, \beta_{2}+\right.$ $\Delta \beta_{2}$ ); where $\phi_{2} \ll \Delta \phi_{2}$ and $\beta_{2} \ll \Delta \beta_{2}$. The volume of $\mathrm{R}_{\tau 2}$ can be obtained with the sum of volume of infinite elementary pyramids originated from $\mathrm{N}_{2}$. The volume $V$ of illuminated scattering region, for a given delay $\tau$, can thus be obtained by adding the volume of regions $\mathrm{R}_{\tau 1}$ and
$\mathrm{R}_{\tau 2}$, which can be expressed as follows,

$$
\begin{align*}
& V_{\tau}(\tau)=\frac{1}{3} \\
& \int_{\phi_{t 2}^{-}}^{\phi_{t 1}^{+}} \int_{\beta_{l}}^{\beta_{t 1}^{+}}\left\{\left(\rho_{\beta 1}^{+}\right)^{3}-\left(\min \left(\rho_{\beta 2}^{-}, \rho_{\tau}, \rho_{e}\right)\right)^{3}\right\} \sin \beta_{2} d \beta_{2} d \phi_{2} \\
& +\frac{1}{3} \int_{-\pi}^{+\pi} \int_{-\frac{\pi}{2}}^{\beta_{t 4}}\left(\min \left(\rho_{2}, \rho_{\tau}, \rho_{e}\right)\right)^{3} \sin \beta_{2} d \beta_{2} d \phi_{2}
\end{align*}
$$

where $\beta_{l}=\min \left(\beta_{t 1}^{-}, \beta_{t 3}\right)$. Similarly, the volume of ESR, $V_{\tau_{\max }}$, for a known longest propagation path's delay, $\tau_{\max }$, can be obtained by substituting $\tau_{\text {max }}$ instead of $\tau$ in (10). The total accumulated volume of scattering solids can be computed as,

$$
\begin{gather*}
V=\max \left(V_{\tau_{\max }}\right)=\frac{4 \pi}{3}\left(a_{1} b_{1} c_{1}+a_{2} b_{2} c_{2}\right)- \\
\frac{\pi}{3}\left\{\frac{a_{1} b_{1}\left(c_{1}-z_{1}\right)^{2}\left(2 c_{1}+z_{1}\right)}{c_{1}^{2}}\right.  \tag{11}\\
\\
\left.\quad+\frac{a_{2} b_{2}\left(c_{2}-z_{2}\right)^{2}\left(2 c_{2}+z_{2}\right)}{c_{2}^{2}}\right\} .
\end{gather*}
$$

### 2.1. Derivation of Delay Statistics

The scatterers responsible for contribution towards a certain delay $\tau$ are located on the surface of ellipsoid $\xi_{\tau}$. The commutative distribution function (CDF) is defined as the probability that a scatterer is placed inside the virtual bounding ellipsoid $\xi_{\tau}$ corresponding to a certain delay $\tau$. For uniformly distributed scatterers, the CDF of ToA can be calculated from the ratio of, the volume of overlap scattering regions with virtual bounding ellipsoid $V_{\tau}$, to the composite volume of whole scattering region $V$, [14]. The ToA CDF can thus be expressed as,

$$
\begin{equation*}
F_{\tau}(\tau)=\frac{V_{\tau}(\tau)}{V_{\tau_{\max }}} \tag{12}
\end{equation*}
$$

The PDF of ToA can thus be calculated by taking differential of $F_{\tau}(\tau)$ with respect to delay $\tau$, which is given as below,

$$
\begin{equation*}
p(\tau)=\frac{\mathrm{d}}{\mathrm{~d} \tau} F_{\tau}(\tau) . \tag{13}
\end{equation*}
$$

The power delay profile (PDP) gives the intensity of a signal received through a multipath channel as a function
of time delay. The PDP can be calculated from PDF of ToA as below,

$$
\begin{equation*}
P_{r}(\tau)=p(\tau) V_{\tau_{\max }} P_{o}\left(\frac{c \tau}{d_{\mathrm{LoS}}}\right)^{-n} \tag{14}
\end{equation*}
$$

where, $P_{o}$ is the power received by the LoS path. The Root Mean Square (RMS) Delay Spread (DS), ( $\sigma_{\mathrm{rms}}$ ), is the square root of the second central moment of the PDP, $P_{r}(\tau)$, which is given below,

$$
\begin{equation*}
\sigma_{\mathrm{rms}}=\sqrt{\frac{\sum P_{r}(\tau) \tau^{2}}{\sum P_{r}(\tau)}-\left(\frac{\sum P_{r}(\tau) \tau}{\sum P_{r}(\tau)}\right)^{2}} \tag{15}
\end{equation*}
$$

### 2.2. Derivation of Angular Statistics

The expressions for the joint and marginal PDF of AoA for the proposed generalized 3-D channel model are derived below. The Joint density function as a function of angles seen at $N_{2}$ can be found as,

$$
p\left(\rho, \phi_{2}, \beta_{2}\right)=\left.\frac{f\left(\mathrm{x}_{2}, \mathrm{y}_{2}, \mathrm{z}_{2}\right)}{\left|J\left(\mathrm{x}_{2}, \mathrm{y}_{2} \cdot \mathrm{z}_{2}\right)\right|}\right|_{\begin{array}{l}
\mathrm{x}_{2} \\
\mathrm{y}_{2}
\end{array}=\rho \cos \beta_{2} \sin \phi_{2}},
$$

where the Jacobean transformation $J\left(\mathrm{x}_{2}, \mathrm{y}_{2}, \mathrm{z}_{2}\right)$ can be found as,

$$
\begin{equation*}
J\left(\mathrm{x}_{2}, \mathrm{y}_{2}, \mathrm{z}_{2}\right)=\frac{1}{\rho^{2} \cos \beta_{2}} . \tag{17}
\end{equation*}
$$

When the scatterers are uniformly distributed with in the defined scattering region with volume $V_{\tau_{\text {max }}}$, the scatterer density function can be written as,

$$
f\left(\mathrm{x}_{2}, \mathrm{y}_{2}, \mathrm{z}_{2}\right)=\left\{\begin{array}{lr}
\frac{1}{V_{\tau_{\max }}} & ;  \tag{18}\\
0 \quad ; \quad\left(\mathrm{x}_{2}, \mathrm{y}_{2}, \mathrm{z}_{2}\right) \in \mathrm{ESR} \\
0 \quad & \text { otherwise }
\end{array}\right.
$$

Substituting (17) and (18) in (16), the joint density function can be rewritten as,

$$
\begin{equation*}
p\left(\rho, \phi_{2}, \beta_{2}\right)=\frac{\rho^{2} \cos \beta_{2}}{V_{\tau_{\max }}} . \tag{19}
\end{equation*}
$$

The joint PDF of AoA for azimuth and elevation angles can be found by integrating (19) over $\rho$ for appropriate limits. The closed-form expression of joint PDF of AoA for different partitions can be written in simplified form as
below,

$$
p\left(\phi_{2}, \beta_{2}\right)=\left\{\begin{array}{l}
\frac{\left(\min \left(\rho_{2}, \rho_{\tau}\right)\right)^{3} \cos \beta_{2}}{3 V_{\tau_{\max }}}  \tag{20}\\
; \pi / 2>\beta_{2}>\beta_{t 1} \text { or } \beta_{t 3}>\beta_{2}>\beta_{t 4} \\
\frac{\left(\min \left(\rho_{2}, \rho_{\tau}\right)+\min \left(\rho_{\beta_{2}}^{+}, \rho_{\tau}\right)-\rho_{\beta_{1}}^{+}\right)^{3} \cos \beta_{2}}{3 V_{\tau_{\max }}} \\
; \beta_{t 1}>\beta_{2}>\beta_{t 2} \\
\frac{\left(\min \left(\rho_{2}, \rho_{\tau}\right)+\min \left(\rho_{e}, \rho_{\tau}\right)-\rho_{\beta_{1}}^{+}\right)^{3} \cos \beta_{2}}{3 V_{\tau_{\max }}} \\
; \beta_{t 2}>\beta_{2}>\beta_{t 3} \\
\frac{\left(\min \left(\rho_{2}, \rho_{\tau}, \rho_{e}\right)\right)^{3} \cos \beta_{2}}{3 V_{\tau_{\max }}} \\
; \beta_{t 4}>\beta_{2}>-\pi / 2 .
\end{array}\right.
$$

The marginal PDF of elevation AoA is found by integrating (20) over $\phi_{2}$ for appropriate limits, as given below,

$$
\begin{equation*}
p\left(\beta_{2}\right)=\int_{\phi_{t 1}}^{\phi_{t 2}} p\left(\phi_{2}, \beta_{2}\right) \mathrm{d} \phi_{2} . \tag{21}
\end{equation*}
$$

Similarly, the marginal PDF of azimuth AoA is found by integrating (20) over $\beta_{2}$ for appropriate limits, as given belwo,

$$
\begin{equation*}
p\left(\phi_{2}\right)=\int_{\beta_{\min }}^{\beta_{\max }} p\left(\phi_{2}, \beta_{2}\right) \mathrm{d} \beta_{2} . \tag{22}
\end{equation*}
$$

The derivations of parameters $\rho_{e}, \rho_{2}, \rho_{\tau}, \rho_{\beta_{1}}^{+}, \rho_{\beta_{2}}^{+}$and azimuth and elevation threshold angles $\phi_{t 1}, \phi_{t 2}, \beta_{t 1} \ldots \beta t 4$ can be seen in next sub-section.

The multipath shape factors (SFs) proposed in [11] and [12] are notable quantifiers to measure the spread of 2D and 3-D angular data, respectively. Fourier coefficients are used to compute SFs for analyzing the PDF of AoA. The $\kappa^{\text {th }}$ complex trigonometric moment of any angular distribution is $\bar{R}_{\kappa}$. For example, for $p(\alpha)$, (i.e., where $\alpha$ may be $\phi$ or $\beta$ for azimuthal and elevational AoA), whose total power is equal to $P_{o}=\int_{\alpha_{\min }}^{\alpha_{\max }} p(\alpha) \mathrm{d} \alpha, \bar{R}_{\kappa}$ is defined as,

$$
\begin{equation*}
\bar{R}_{\kappa}=\frac{1}{P_{o}} \int_{\alpha_{\min }}^{\alpha_{\max }} p(\alpha) \exp \left(j \alpha_{\kappa}\right) \mathrm{d} \alpha . \tag{23}
\end{equation*}
$$

Only first and second moments are used in characterization of the SFs. $\left|\bar{R}_{1}\right|$ is the magnitude of first trigonometric moment. The SF angular spread, $\Lambda_{\alpha}$, is given as below,

$$
\begin{equation*}
\Lambda_{\alpha}=\sqrt{1-\left|\bar{R}_{1}\right|^{2}} \tag{24}
\end{equation*}
$$

The angular spread SF measures the spread with a value between 0 and 1 . A value close to 0 indicate that signals from a wider range of angles are received and a value close to 1 indicate a sharp angular spread. The other important SF angular constriction, $\gamma_{\alpha}$, is defined as below,

$$
\begin{equation*}
\gamma_{\alpha}=\frac{\left|\bar{R}_{2}-\bar{R}_{1}^{2}\right|}{1-\left|\bar{R}_{1}\right|^{2}} . \tag{25}
\end{equation*}
$$

The angular constriction shape factor measures the bias of angular data towards exactly two physical directions. A value 1 indicates the extreme case of signals arriving exactly from two physical directions, and a value 0 indicate no bias towards two directions.

### 2.3. Derivation of Angular Partitions

A cross-section view of the scattering region for a certain azimuth and elevation angle of observation at node $\mathrm{N}_{2}$ is shown in Fig. 2. The vertical distance from the origin of $\mathrm{E}_{1}$ to the antenna of $\mathrm{N}_{1}$ is denoted by $h_{e 1}=h_{1}-z_{1}$. Similarly, the vertical distance from the origin of $E_{2}$ to the antenna of $\mathrm{N}_{2}$ is shown by $h_{e 2}=h_{2}-z_{2}$. The difference between the elevation of the origin of $E_{1}$ and antenna of $\mathrm{N}_{2}$ is shown by $h_{e 3}=h_{2}-z_{1}$. Similarly, the difference between the elevation of the origin of $\mathrm{E}_{2}$ and antenna of $\mathrm{N}_{1}$ is shown by $h_{e 4}=h_{1}-z_{2}$. A line originated from the node $\mathrm{N}_{2}$ in the direction of observation (given, $\phi_{2}$ and $\beta_{2}$ ) intersects $\mathrm{E}_{2}$ at point $\varsigma_{2}, \mathrm{E}_{1}$ at points $\varsigma_{1}$ and $\varsigma_{1}^{\prime}$, and $\xi_{\tau}$ at point $\varsigma_{\tau}$; the distance from $\mathrm{N}_{2}$ to these intersection points is denoted by $\rho_{2}, \rho_{\beta 1}^{+}$and $\rho_{\beta 2}^{-}$, and $\rho_{\tau}$, respectively. Mathematical expressions for these distances can be obtained by substituting the equation of the line into the equation of respective ellipsoid ( $\mathrm{E}_{1}, \mathrm{E}_{2}$, or $\xi_{\tau}$ ) and solving for the radial distance. After performing tedious mathematical simplifications, the expression for $\rho_{2}, \rho_{\beta 1}^{+}$, $\rho_{\beta 2}^{-}$, and $\rho_{\tau}$ are derived as given below,

$$
\begin{align*}
\rho_{2}= & \frac{1}{2 a_{2}^{2} b_{2}^{2}+c_{2}^{2}\left(a_{2}^{2}+b_{2}^{2}\right)-c_{2}^{2} \Psi-\left(2 a_{2}^{2} b_{2}^{2}-c_{2}^{2}\left(a_{2}^{2}+b_{2}^{2}\right)\right) \cos \left(2 \beta_{2}\right)} \times \\
& \left\{4 a_{2}^{2} b_{2}^{2} h_{e 2} \sin \beta_{2}-2\left(a _ { 2 } ^ { 2 } b _ { 2 } ^ { 2 } c _ { 2 } ^ { 2 } \left(b_{2}^{2}\left(c_{2}^{2}-h_{e 2}^{2}\right)\right.\right.\right. \\
& +a_{2}^{2}\left(2 b_{2}^{2}+c_{2}^{2}-h_{e 2}^{2}\right)-\Psi\left(c_{2}^{2}-h_{e 2}^{2}\right)+\left(b_{2}^{2}\left(c_{2}^{2}-h_{e 2}^{2}\right)\right. \\
& \left.\left.\left.\left.+a_{2}^{2}\left(-2 b_{2}^{2}+c_{2}^{2}-h_{e 2}^{2}\right)\right) \cos \left(2 \beta_{2}\right)\right)\right)^{1 / 2}\right\}, \tag{26}
\end{align*}
$$

$$
\begin{align*}
& \left.\begin{array}{l}
\rho_{\beta 1}^{+} \\
\rho_{\beta 2}^{-}
\end{array}\right\}=\frac{-1}{a_{\phi}^{2}+c_{\phi}^{2}+\left(c_{\phi}^{2}-a_{\phi}^{2}\right) \cos \left(2 \beta_{2}\right)} \times \\
& \left\{-2 c_{\phi}^{2} \rho_{e 2} \cos \beta_{2}-2 h_{e 3} a_{\phi}^{2} \sin \beta_{2} \pm\left(2 a_{\phi}^{2} c_{\phi}^{2} \times\right.\right. \\
& \left(a_{\phi}^{2}+c_{\phi}^{2}-h_{e 3}^{2}-\rho_{e 2}^{2}-\left(a_{\phi}^{2}-c_{\phi}^{2}+h_{e 3}^{2}-\rho_{e 2}^{2}\right)\right. \\
& \left.\left.\left.\cos \left(2 \beta_{2}\right)+2 h_{e 3} \rho_{e 2} \sin \left(2 \beta_{2}\right)\right)\right)^{1 / 2}\right\},  \tag{27}\\
& \rho_{\tau}=\left\{-5 a_{\tau}^{2}-3 b_{\tau}^{2}+\left(a_{\tau}^{2}-b_{\tau}^{2}\right)\left(-\cos \left(2 \beta_{\mathrm{LoS}}\right)\right.\right. \\
& +\cos \left(2 \beta_{2}\right)\left(-1+3 \cos \left(2 \beta_{\mathrm{LoS}}\right)\right) \\
& +4 \cos ^{2} \beta_{2} \cos ^{2} \beta_{\mathrm{LoS}} \cos \left(2 \phi_{2}\right) \\
& \left.\left.-4 \cos \phi_{2} \sin \left(2 \beta_{2}\right) \sin \left(2 \beta_{\mathrm{LoS}}\right)\right)\right\}^{-1} \\
& \times\left\{-4 b_{\tau}^{2} d_{\mathrm{LoS}} \cos \beta_{2} \cos \beta_{\mathrm{LoS}} \cos \phi_{2}\right. \\
& +4 b_{\tau}^{2} d_{\mathrm{LoS}} \sin \beta_{2} \sin \beta_{\mathrm{LoS}}-\left(2 a _ { \tau } ^ { 2 } b _ { \tau } ^ { 2 } \left(20 a_{\tau}^{2}\right.\right. \\
& +12 b_{\tau}^{2}-5 d_{\mathrm{LoS}}^{2}+\left(4 a_{\tau}^{2}-4 b_{\tau}^{2}-d_{\mathrm{LoS}}^{2}\right) \\
& \left(\cos \left(2 \beta_{2}\right)\left(1-3 \cos \left(2 \beta_{\mathrm{LoS}}\right)\right)+\cos \left(2 \beta_{\mathrm{LoS}}\right)\right. \\
& -4 \cos ^{2} \beta 2 \cos ^{2} \beta_{\mathrm{LoS}} \cos \left(2 \phi_{2}\right) \\
& \left.\left.\left.\left.+4 \cos \phi_{2} \sin \left(2 \beta_{2}\right) \sin \left(2 \beta_{\mathrm{LoS}}\right)\right)\right)\right)^{1 / 2}\right\}, \tag{28}
\end{align*}
$$

where the parameters $\Psi, a_{\phi}, c_{\phi}$, and $\rho_{e 2}$ are the geometric simplification parameters. $\Psi$ can be expressed as given below,

$$
\begin{equation*}
\Psi=2\left(a_{2}^{2}-b_{2}^{2}\right) \cos ^{2} \beta_{2} \cos \left(2 \theta_{2}-2 \phi_{2}\right) . \tag{29}
\end{equation*}
$$

The parameters $a_{\phi}$ and $c_{\phi}$ are the major and minor dimensions of a vertical 2-D ellipse $\varepsilon_{\phi}$ formed at the vertical cross-section of ellipsoid $\mathrm{E}_{1}$ for a certain given azimuth angle $\left(\phi_{2}\right)$ observed at $\mathrm{N}_{2}$.

$$
\begin{align*}
a_{\phi}= & \frac{\sec \phi_{2}}{a_{1}^{2}+b_{1}^{2}+\left(b_{1}^{2}-a_{1}^{2}\right) \cos \left(2 \theta_{1}-2 \phi_{2}\right)} \times \\
& \left\{\begin{array}{l}
-2 a_{1}^{2} b_{1}^{2} \cos ^{2} \phi_{2}\left(d^{2}-a_{1}^{2}-b_{1}^{2}\right.
\end{array}\right. \\
& \left.\left.+\left(a_{1}^{2}-b_{1}^{2}\right) \cos \left(2 \theta_{1}-2 \phi_{2}\right)-d^{2} \cos \left(2 \phi_{2}\right)\right)\right\}^{1 / 2} \tag{30}
\end{align*}
$$

$$
\begin{align*}
& c_{\phi}=c_{1} \times  \tag{31}\\
& \sqrt{\frac{-b_{1}^{2} \rho_{e 1}^{2}+a_{1}^{2}\left(2 b_{1}^{2}-\rho_{e 1}^{2}\right)+\left(a_{1}^{2}-b_{1}^{2}\right) \rho_{e 1}^{2} \cos \left(2\left(\theta_{1}+\phi_{e}\right)\right)}{2 a_{1}^{2} b_{1}^{2}}},
\end{align*}
$$

where the parameter $\phi_{e}$ can be obtained as,

$$
\begin{equation*}
\phi_{e}=\arcsin \left(\frac{\rho_{e 2} \sin \phi_{2}}{\rho_{e 1}}\right) . \tag{32}
\end{equation*}
$$

The origin of the vertical 2-D ellipse $\varepsilon_{\phi}$ is labeled as $\mathrm{w}_{\phi}$ in Fig. 2. The horizontal distance from $\mathrm{N}_{1}$ and $\mathrm{N}_{2}$ to the point $\mathrm{w}_{\phi}$ is denoted by $\rho_{e 1}$ and $\rho_{e 2}$. These distances can be obtained as,

$$
\begin{gather*}
\rho_{e 1}=\sqrt{d^{2}+\rho_{e 2}^{2}-2 d \rho_{e 2} \cos \phi_{2}},  \tag{33}\\
\rho_{e 2}=-d \frac{\left(a_{1}^{2}-b_{1}^{2}\right) \cos \left(2 \theta_{1}-\phi_{2}\right)-\left(a_{1}^{2}+b_{1}^{2}\right) \cos \phi_{2}}{a_{1}^{2}+b_{1}^{2}+\left(b_{1}^{2}-a_{1}^{2}\right) \cos \left(2 \theta_{1}-2 \phi_{2}\right)} . \tag{34}
\end{gather*}
$$

The length of line drawn from $\mathrm{N}_{2}$ to the intersection points on ellipse $\varepsilon_{\phi}$ at an elevation of $z_{1}$ is denoted by $\rho_{\phi 1}^{+}$ and $\rho_{\phi 2}^{-}$. The simplified solution can be expressed as,

$$
\begin{align*}
& \left.\left.\rho_{\phi 1}^{+}\right\}=\frac{-1}{\rho_{\phi 2}^{-}}\right\}=\frac{a_{1}^{2}+b_{1}^{2}+\left(b_{1}^{2}-a_{1}^{2}\right) \cos \left(2 \theta_{1}-2 \phi_{2}\right)}{} \times \\
& \left\{\begin{array}{l}
\sec \phi_{2}\left(2 a_{1}^{2} b_{1}^{2} \cos ^{2} \phi_{2} \times\left(a_{1}^{2}+b_{1}^{2}-d^{2}\right.\right. \\
\left.\left.+\left(b_{1}^{2}-a_{1}^{2}\right) \cos \left(2 \theta 1-2 \phi_{2}\right)+d^{2} \cos \left(2 \phi_{2}\right)\right)\right)^{1 / 2} \\
-2 d\left(b_{1}^{2} \cos \theta_{1} \cos \left(\theta_{1}-\phi_{2}\right)\right. \\
\left.\left.+a_{1}^{2} \sin \theta_{1} \sin \left(\theta_{1}-\phi_{2}\right)\right)\right\}
\end{array} .\right.
\end{align*}
$$

Similarly, the length of line from $\mathrm{N}_{2}$ to the intersection points with $\varepsilon_{\phi}$ formed at any arbitrary elevation $z$ can be obtained by replacing the dimensions $a_{1}$ and $b_{1}$ with $a_{1, z}$ and $b_{1, z}$ in the above equation. The parameters $a_{1, z}$ and $b_{1, z}$ represent the major and minor dimensions of the horizontal cross-section of $\mathrm{E}_{1}$ formed at an height $z$; which can be expressed as,

$$
\begin{align*}
& a_{1, z}=\frac{a_{1} \sqrt{c_{1}^{2}-(z-z 1)^{2}}}{c_{1}},  \tag{36}\\
& b_{1, z}=\frac{b_{1} \sqrt{c_{1}^{2}-(z-z 1)^{2}}}{c_{1}} . \tag{37}
\end{align*}
$$

Azimuth threshold angles to differentiate among the angular regions of different geometrical composition of the ESR are illustrated in Fig. 3. The azimuth angles observed at $\mathrm{N}_{2}$ in the horizontal plane at an height $z_{1}$ (i.e., for maximum horizontal cross-section of $\mathrm{E}_{1}$ ) are labeled


Figure 2: Cross-section view of the proposed model for a fixed direction of observation in azimuth and elevation planes from node $\mathrm{N}_{2}$.


Figure 3: Top-view ( $x-y$ plane) of the proposed model with azimuth threshold angle observed from $\mathrm{N}_{2}$.
by $\phi_{t 1}^{+}$and $\phi_{t 2}^{-}$. These angles are subtended by the lines originated from $\mathrm{N}_{2}^{\prime}$ (projection of $\mathrm{N}_{2}$ at a certain elevation $z$ ) to the tangent points $\mathrm{T}_{\mathrm{z}}$ and $\mathrm{T}_{\mathrm{z}}^{\prime}$, see Fig. 3 at the surface of $\mathrm{E}_{1}$ for an elevation $z=z_{1}$. Simplified mathematical expression for these angles can be expressed as given
below,

$$
\left.\begin{array}{rl}
\phi_{t 1}^{+} \\
\phi_{t 2}^{-}
\end{array}\right\}=\arctan \left(\frac{1}{d^{2}-a_{1}^{2} \cos ^{2} \theta_{1}-b_{1}^{2} \sin ^{2} \theta_{1}} \times x .\left\{\left(b_{1}^{2}-a_{1}^{2}\right) \cos \theta_{1} \sin \theta_{1},\right.\right.
$$

Similarly, the azimuth threshold angles formed at any arbitrary elevation $z$ can be obtained by replacing the


Figure 4: Side-view (x-z plane) of the proposed model, indicating elevation threshold angles observed from $\mathrm{N}_{2}$ for a fixed azimuth angle of observation $\phi_{2}$.
dimensions $a_{1}$ and $b_{1}$ with $a_{1, z}$ and $b_{1, z}$ in the above equation. The dimensions of the 2-D ellipse ( $\varepsilon_{1}$ ) formed with the cross-section of scattering ellipsoids, $\mathrm{E}_{1}$, and the ground plane can also be obtained from (36) and (37) by setting $z=-z_{1}$. Subsequently, the azimuth threshold angles ( $\phi_{t 1, b}^{+}$and $\phi_{t 2, b}^{-}$, see Fig. 3) formed at the ground plane can be deduced from (38) by substituting $a_{1}=a_{1, z}$ and $b_{1}=b_{1, z}$ with $z=-z 1$.

Important elevation threshold angles, to differentiate among the angular regions based upon their geometrical composition, are illustrated in Fig. 4. The elevation threshold angles subtended by a line drawn at a certain fixed azimuth angle $\phi_{2}$ from $\mathrm{N}_{2}$ to the tangent points $T_{\phi 1}$ and $T_{\phi 1}^{\prime}$ at the vertical ellipse $\varepsilon_{\phi}$ are denoted by $\beta_{t 1}^{+}$and $\beta_{t 1}^{-}$; these angles can be expressed as,

$$
\left.\begin{array}{l}
\beta_{t 1}^{+}  \tag{39}\\
\beta_{t 1}^{-}
\end{array}\right\}=\left\{\begin{array}{c}
\arctan \left(\frac{-h_{e 3} \rho_{e 2} \pm \sqrt{a_{\phi}^{2} h_{e 3}^{2}+c_{\phi}^{2}\left(-a_{\phi}^{2}+\rho_{e 2}^{2}\right)}}{\rho_{e 2}^{2}-a_{\phi}^{2}}\right) \\
0 \quad ; \phi_{t 2}^{-} \leq \phi_{2} \leq \phi_{t 1}^{+} \\
0 \quad ; \text { otherwise }
\end{array}\right.
$$

The intersection points of vertical 2-D ellipse $\varepsilon_{\phi}$ and ground plane are labeled by $\mathrm{j}_{\phi}$ and $\mathrm{r}_{\phi}$, see Fig. 4. The angles formed with the line drawn from $\mathrm{N}_{2}$ to these intersection points are denoted by $\beta_{t 2}$ and $\beta_{t 3}$, which can
be expressed as given below,

$$
\begin{gather*}
\beta_{t 2}= \begin{cases}\arctan \left(\frac{h_{2}}{\rho_{\phi 2, b}^{-}}\right) & ; \phi_{t 2}^{-} \leq \phi_{2} \leq \phi_{t 1}^{+} \\
0 & ; \text { otherwise },\end{cases}  \tag{40}\\
\beta_{t 3}= \begin{cases}\arctan \left(\frac{h_{2}}{\rho_{+1, b}^{+}}\right) & ; \phi_{t 2}^{-} \leq \phi_{2} \leq \phi_{t 1}^{+} \\
0 & ; \text { otherwise. }\end{cases} \tag{41}
\end{gather*}
$$

Similarly, the elevation angle formed with the intersection point $s_{\phi}$ of the vertical slice of scattering region around $\mathrm{N}_{2}$ with the ground plane is denoted by $\beta_{t 4}$, which is given as under,

$$
\begin{equation*}
\beta_{t 4}=\arctan \left(\frac{h_{2}}{\rho_{2, b}}\right) \tag{42}
\end{equation*}
$$

where the parameter $\rho_{\phi 1, b}^{+}$and $\rho_{\phi 2, b}^{-}$represent the horizontal distance from $\mathrm{N}_{2}$ to the intersection points $\mathrm{r}_{\phi}$ and $j_{\phi}$, recall Fig. 4. These parameters can be obtained from (35) by substituting $a_{1}=a_{1, z}, b_{1}=b_{1, z}$, and $z=$ $-z_{1}$.

## 3. RESULTS AND DISCUSSIONS

This section presents the obtained analytical results on delay and angular characteristics of the channel. A thorough analysis on the impact of various physical channel


Figure 5: Trend of Delay Spread w.r.t. physical channel parameters. (a) Relationship between delay-span and RMS-DS for different values of path-loss exponent, $\left(a_{1}=\right.$ $a_{2}=500 \mathrm{~m}, b_{1}=b_{2}=250 \mathrm{~m}, c_{1}=c_{2}=175 \mathrm{~m}, h_{1}=h_{2}$ $=0 \mathrm{~m}, z_{1}=z_{2}=0 \mathrm{~m}, d=1000 \mathrm{~m}, \theta_{1}=0^{\circ}, \theta_{2}=0^{\circ}$ ). (b) Impact of the orientation of scattering volume and vertical spatial deviation of mobile node from the center of scattering volume, $\left(a_{1}=a_{2} 500 \mathrm{~m}, b_{1}=b_{2}=250 \mathrm{~m}, c_{1}=\right.$ $c_{2}=200 \mathrm{~m}, h_{2}=0 \mathrm{~m}, z_{2}=200 \mathrm{~m}, d=1000 \mathrm{~m}, \theta_{2}=0^{\circ}$, $\left.\tau_{\text {max }}=6.74 \mu \mathrm{sec}\right)$.
parameters on statistical characteristics of multipath signals in delay and angular domains is presented.

### 3.1. Delay Statistics

The spread of multipath signals in the delay domain leads to frequency selectivity of the channel. The measure of delay-span alone is not sufficient to provide realistic picture of multipath delay spread; therefore, it is necessary to consider both the delay-span and the peak(s) of the ToA distribution to quantify the delay spread. The definition of RMS-DS is the most favorable quantifier in the literature to gauge the spread of multipath waves in delay domain. The relationship between delay-span and RMS-DS for the proposed channel model is demonstrated in Fig. 5 (a), where the impact of path-loss exponent is also shown. The values of the parameters taken for the graphs are indicated in caption of the relevant figures. The RMS-DS of the multipath energy decreases nonlinearly with a decrease in delay-span, and it decreases sharply for the smaller values of the span. In Fig. 5 (b), the impact of vertical spatial deviation of mobile node from the center of scattering volume (i.e., $h_{e 1}$ ) and horizontal orientation of scattering volume (i.e., $\theta_{1}$ ) on the RMS-DS of the channel is shown.

Circular or elliptical shaped scattering region in the vicinity of a mobile node is usually considered better approximation of the physical reality to model


Figure 6: The impact of transformation of scattering volumes from spherical to ellipsoidal shape on the CDF and PDF of ToA, shown in (a) and (b), respectively, $\left(c_{1}=\right.$ $10 \mathrm{~m}, c_{2}=20 \mathrm{~m}, h_{1}=h_{2}=10 \mathrm{~m}, z_{1}=z_{2}=10 \mathrm{~m}, d=70 \mathrm{~m}$, $\left.\theta_{1}=\theta_{2}=0^{\circ}, \tau_{\max }=0.43 \mu \mathrm{sec}\right)$.


Figure 7: The impact of increase in the amount of scattering objects along the elevation axis in the vicinity of far-end communicating node on the CDF and PDF of ToA, shown in (a) and (b), respectively. ( $a_{1}=a_{2} 50 \mathrm{~m}, b_{1}$ $=25 \mathrm{~m}, b_{2}=30 \mathrm{~m}, c_{2}=20 \mathrm{~m}, h_{1}=h_{2}=10 \mathrm{~m}, z_{1}=z_{2}=$ $\left.10 \mathrm{~m}, d=70 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}, \tau_{\max }=0.43 \mu \mathrm{sec}\right)$.
sparsely populated small towns or dense urban streets, respectively. The shape of the scattering region changes in accordance with the terrain, clutter, and spatial position of scattering structures in the vicinity of communication nodes; therefore, fixing the geometric shape of the scattering volumes to be either spherical or ellipsoidal alone and fixing the position of antennas in the center of scattering volumes, does not allow full flexibility to the physical channel modeling. The proposed model is thus designed tunable in its geometric parameters to more realistically adapt the physical environment. The impact of transforming the scattering volumes around both the ends of the communication link from circular shape to elliptical shape (changing eccentricity of ellipsoids) in horizontal plane on ToA distribution characteristics are shown in Fig. 6. With a decrease in the ratio between minor


Figure 8: The impact of tuning the ratio between the dimensions of major axes of the scattering ellipsoids around both the ends of the communication link on the PDF of azimuth AoA and elevation AoA, shown in (a) and (b), respectively. ( $b_{1}=20 \mathrm{~m}, c_{1}=20 \mathrm{~m}, b_{2}=15 \mathrm{~m}, c_{2}=15 \mathrm{~m}$, $h_{1}=h_{2}=0 \mathrm{~m}, z_{1}=z_{2}=0 \mathrm{~m}, d=65 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}$, $\left.\tau_{\text {max }}=0.32 \mu \mathrm{sec}\right)$.
and major axes of scattering volumes, the probability of arrival of signal from the directions around LoS direction increases, which further increases the probability of arrival of signal with shorter delays. In the emerging nextgeneration wireless communication networks, the decrease in elevation of BS and size of cells in dense urban propagation environments, make it necessary to accurately model the spatial location of scattering objects in both elevation and azimuth plane. The effect of increasing the amount of scattering objects along elevation axis on distribution of ToA is shown in Fig. 7. Only the local scattering objects around a mobile node in rural areas (typical macro-cellular environment) correspond in the arrival of multipath signals; where the high elevation of BS makes the vicinity of BS a scattering free region. On the other hand, in densely populated and low BS antenna height outdoor and indoor environments (i.e., micro-cellular and pico-cellular), the effective scatterers are concentrated around both the ends of the link.

### 3.2. Angular Statistics

It is observed that the change in scattering volume around the observing-end changes the tails of AoA distribution and change in the scattering volume around observingend impacts the AoA distribution for the directions around $\operatorname{LoS}$ path. The impact of changing the relative volume of independent scattering regions around both the communicating nodes on the PDF of azimuth and elevation AoA is shown in Fig. 8. The impact of relative amount of scatterers in the elevation plane around both the ends


Figure 9: The impact of tuning the ratio between the dimensions of minor axes of the scattering ellipsoids around observing- and far-ends of the communication link on the PDF of azimuth AoA and elevation AoA, shown in (a) and (b), respectively. ( $a_{1}=a_{2}=30 \mathrm{~m}, b_{1}=b_{2}=20 \mathrm{~m}$, $h_{1}=h_{2}=0 \mathrm{~m}, z_{1}=z_{2}=0 \mathrm{~m}, d=65 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}$, $\left.\tau_{\text {max }}=0.32 \mu \mathrm{sec}\right)$.



Figure 10: The impact of transformation of scattering volumes from spherical to ellipsoidal shape on the PDF of azimuth AoA and elevation AoA shown in (a) and (b), respectively. ( $c_{1}=c_{2}=20 \mathrm{~m}, h_{1}=h_{2}=0 \mathrm{~m}, z_{1}=z_{2}=$ $\left.0 \mathrm{~m}, d=65 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}, \tau_{\text {max }}=0.32 \mu \mathrm{sec}\right)$.
of the link on the PDF of azimuth and elevation AoA is shown in Fig. 9. The increase in ratio of $c_{1} / c_{2}$ means an increase in the volume of scattering region around the far-end node with respect to the observing-end. It can be seen that values of the PDF of azimuth AoA increases along the LoS direction with an increase in ratio of $c_{1} / c_{2}$ and it shows converse behavior for tails of the distribution for the angles away from LoS direction. The elevational angular spread is observed to increase with a decrease in ratio $c_{1} / c_{2}$, which indicates that the local scattering objects in elevation plane around the observing-end has higher influence on the elevational AoA distribution compared to the scatterers in elevation plane around the far-end. The impact of transforming the scattering volumes around both the ends of the communication link from circular shape to elliptical shape (changing eccentricity of ellipsoids)


Figure 11: The impact of tuning vertical positions of the (origins of) scattering ellipsoids on the PDF of azimuth AoA and elevation AoA, shown in (a) and (b), respectively. $\left(a_{1}=a_{2}=30 \mathrm{~m}, b_{1}=b_{2}=20 \mathrm{~m} c_{1}=c_{2}=\right.$ $20 \mathrm{~m}, h_{1}=h_{2}=0 \mathrm{~m}, d=65 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}, \tau_{\max }=0.32$ $\mu \mathrm{sec})$.


Figure 12: The impact of vertical displacement of (origin of) scattering volume around observing-end w.r.t. the position of far-end antenna on the PDF of azimuth AoA and elevation AoA, shown in (a) and (b), respectively. ( $a_{1}$ $=30 \mathrm{~m}, b_{1}=30 \mathrm{~m}, c_{1}=30 \mathrm{~m}, a_{2}=25 \mathrm{~m}, b_{2}=15 \mathrm{~m}, c_{2}=$ $15 \mathrm{~m}, h_{2}=0 \mathrm{~m}, z_{1}=0 \mathrm{~m}, d=65 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}, \tau_{\text {max }}=$ $0.32 \mu \mathrm{sec})$.
in horizontal plane on AoA distribution characteristics are shown in Fig. 10. With a decrease in the ratio between minor and major axes of scattering volumes, the probability of arrival of signal from the directions around LoS direction increases, which further increases the probability of arrival of signal with shorter delays. The impact of change in the ratio between the height of origins of scattering volumes around both the ends of the communication link on the PDF of AoA in azimuth and elevation planes is shown in Fig. 11 (a) and (b), respectively. This graph is left non-normalized, in order to clearly present the effect of the ratio. It can be observed that with an increase in the ratio $z_{1} / z_{2}$, the PDF of azimuth AoA becomes sharper along the LoS direction and tails of the PDF also become prominent. On the other hand, with a change in $z_{1} / z_{2}$ the peak of elevation AoA PDF


Figure 13: The impact of change in delay span and horizontal distance between communicating nodes on the angular spread in azimuth and elevation plane shown in (a) and (b), respectively. $\left(a_{1}=30 \mathrm{~m}, b_{1}=20 \mathrm{~m}, c_{1}=20 \mathrm{~m}, a_{2}=\right.$ $25 \mathrm{~m}, b_{2}=15 \mathrm{~m}, c_{2}=15 \mathrm{~m}, h_{1}=30 \mathrm{~m}, h_{2}=0 \mathrm{~m}, z_{1}=5 \mathrm{~m}$, $\left.z_{2}=0 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}\right)$.


Figure 14: The impact of change in horizontal distance between communicating nodes and vertical displacement of (the origin of) observing-end's scattering ellipsoid w.r.t. the height of far-end's antenna on angular spread observed in azimuth and elevation planes, shown in (a) and (b), respectively. $\left(a_{1}=30 \mathrm{~m}, b_{1}=20 \mathrm{~m}, c_{1}=20 \mathrm{~m}, a_{2}=25 \mathrm{~m}\right.$, $\left.b_{2}=15 \mathrm{~m}, c_{2}=15 \mathrm{~m}, h_{2}=0 \mathrm{~m}, z_{2}=0 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}\right)$.
shifts along elevation axis however the shape of the PDF remains similar. This indicates the importance of inclusion of scatterers in the elevation plane while modeling the scattering region. The effect of vertical deviation of the scattering volume around the observing-end node (i.e., $\mathrm{N}_{2}$ ) with respect to the elevation of antenna of far-end node on the PDF of azimuth and elevation AoA is shown in Fig. 12 (a) and (b), respectively. With an increase in $h_{e 4}$, the PDF of azimuth AoA becomes prominent along LoS direction and angular spread decreases. In elevation plane, with a decrease in $h_{e 4}$, the tails of the PDF (for $\beta_{2}>20^{\circ}$ ) become significantly prominent and thus the angular decreases. The spatial location of the effective scattering objects in a particular propagation environment determine the angular characteristics of the multipaths arrived at the receiver. The angular spread of the mulipath


Figure 15: The impact of change in horizontal distance between communicating nodes and vertical displacement of (the origin of) far-end's scattering region w.r.t. the height of far-end's antenna on angular constriction observed in azimuth and elevation plane, shown in (a) and (b), respectively. ( $a_{1}=30 \mathrm{~m}, b_{1}=20 \mathrm{~m}, c_{1}=20 \mathrm{~m}, a_{2}=25 \mathrm{~m}$, $\left.b_{2}=15 \mathrm{~m}, c_{2}=15 \mathrm{~m}, h_{2}=0 \mathrm{~m}, z_{2}=0 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}\right)$.
waves further determine the space, time, and frequency selective characteristics of the channel. Therefore, it is of vital importance to study the behaviour of angular spread for characterization of radio propagation channels. Various quantifier to measure the dispersion of energy in angular domain are available in the literature. The multipath shape factors theory is considered most favorable to measure the dispersion of energy in angular domain in the form of angular spread $\Lambda$ and angular constriction $\gamma$. The impact of different physical channel parameters on angular spread is studied on the base of these shape factor quantifiers. The angular spread of the mutipath waves with respect to the horizontal link-distance and maximum considered excess delay is plotted in Fig. 13 (a) and (b) for azimuth $\Lambda_{\phi}$ and elevation $\Lambda_{\beta}$ planes, respectively. The azimuthal and elevational angular spread is observed to increase gradually with an increase in the horizontal link-distance $d$, since increasing the link distance enhances the tails and sharpens the peak along LoS direction. However, after a certain value of link-distance the influence of the peak of the PDF along LoS direction (caused by farend scattering volume) on the angular spread measure becomes leaner compared to the peaks in the tails of the PDF caused by local scattering volume; therefore the behaviour of angular spread reverses along link-distance axis. The trend of decrease in the angular spread for higher values of link-distance becomes sharper with a decrease in the maximum excess delay. The effect of change in vertical deviation $h_{e 1}$ of far-end scattering


Figure 16: Summary of the influence of the proposed model's geometric parameters on PDF of azimuth AoA.
volume $\mathrm{E}_{2}$ with reference to the antenna height of farend node $N_{1}$ on angular spread and angular constriction (in both azimuth and elevation planes) with respect to horizontal link distance $d$ is shown in Fig. 14 and Fig. 15, respectively. Both the azimuthal and elevational angular spread is observed to decrease with a decrease in linkdistance, however, it decreases more sharply for higher values of the vertical deviation $h_{e 1}$. On the other hand, the azimuthal angular constriction is observed to increase nonlinearly with a linear decrease in the link-distance and the rate of it's increase further increases with an increase in vertical deviation $h_{e 1}$. However, the elevational angular constriction shows a converse behaviour along both the parameters $d$ and $h_{e 1}$. Summary of the effect of the proposed model's input parameters on the PDF of azimuth AoA discussed above is illustrated in Fig. 16. The two curves show the contribution of observing-end and farend scattering objects in the PDF of azimuth AoA; the composite pdf of azimuth AoA can thus be obtained by adding these two.

### 3.3. Model Validation

### 3.3.1. Validation through Computer Simulations

In order to demonstrate the validity of the derived analytical expressions, computer based simulations are performed. A comparison of simulation results with analytical results obtained for the proposed model is provided in Fig. 17. Realization of effective scattering points in generated from uniform distribution limited within the defined scattering region. The results are taken for $10^{3}$ uniformly distributed scattering points, and averaged over $10^{4}$ monte carlo runs. A good match is observed between simulation and analytical results for all


Figure 17: Comparison of simulations and analytical results. (a) Spatial distribution of scattering points taken for simulation setup, (b) PDF of azimuth AoA, (c) PDF of elevation AoA, (d) PDF of ToA. ( $a_{1}=300 \mathrm{~m}, b_{1}=500 \mathrm{~m}, c_{1}=100 \mathrm{~m}, a_{2}=$ $500 \mathrm{~m}, b_{2}=250 \mathrm{~m}, c_{2}=175 \mathrm{~m}, d=1200 \mathrm{~m}, h_{1}=0 \mathrm{~m}, h_{2}=0 \mathrm{~m}, z_{1}=0 \mathrm{~m}, z_{2}=0 \mathrm{~m}, \theta_{1}=\theta_{2}=0^{\circ}$.


Figure 18: Curve-fitting analytical results after matching the ToA PDF obtained from the proposed model to the empirical data-sets for outdoor propagation environments in K. I. Pedersen et al. [20], K. I. Pedersen et al. [20], and M. D. Batariere et al. [21] shown in (a), (b), and (c), respectively.
azimuth AoA, elevation AoA, and ToA statistics of the channel.

### 3.3.2. Validation through Empirical Data

A detailed comparison of the proposed analytical results with empirical data-sets available in the literature is presented to demonstrate the validity of the proposed model in Fig. 18, 19, and 20 for ToA, outdoor AoA, and indoor AoA statistics, respectively. The discrete data samples of the empirical data-sets obtained for
different outdoor propagation environments provided in [20, 21, 22, 23, 24, 25] are extracted through "GetData Graph Digitizer" instead of human visual reading. See http://www.getdata-graph-digitizer.com/. However, the data samples for joint ToA and azimuth AoA (at $0^{\circ}$ elevation AoA ) indoor measurement campaign results in [26] are manually extracted from the provided contour maps, by manually employing a careful color decoding through image processing tools. These joint (3D) data samples are numerically integrated over delay $(\tau)$ to obtain


Figure 19: Curve-fitting proposed analytical results after matching the azimuth AoA PDF obtained from the proposed model to the empirical data-sets for outdoor propagation environments available in K. I. Pedersen et al. [20], B. H. Fleury et al. [22], P. A. Matthews et al. [23], C. Kloch et al. [24], K. I. Pedersen et al. [25], and K. I. Pedersen et al. [20] shown in (a), (b), (c), (d), (e), and (f), respectively.


Figure 20: Curve-fitting proposed analytical results after matching the azimuth AoA PDF obtained from the proposed model to the empirical data-sets for indoor propagation environments available in X. Yin et al. [26] in (a) LoS measurement scenario, (b) NLoS measurement scenario 1, and (c) NLoS measurement scenario 2.
marginal PDF of azimuth AoA. All the extracted empirical data-sets for outdoor and indoor environments are scaled along vertical axis to fix their $\operatorname{peak}(\mathrm{s})$ at 1 , which is performed by dividing all the values of a data-set by its peak value. The normalized data-sets and/or analytical
plots obtained from proposed model are shifted (circularly) and scaled along horizontal axis to align with each other.

There are various notable methods available in the literature, see e.g., [27, 28, 29], which can be used to adopt the input geometric parameters for modeling a specific propagation environment. Considering the available details
of the targeted measurement campaigns and the knowledge of the impact of input geometric parameters on the output curves (see the summary plotted in Fig. 16), a reasonable range can be defined for the those input geometric parameters for which no direct information is available in the measurement-campaign details. The strategy used for setting the input geometric parameters of the proposed model (for comparison with measurement results) can be defined in following three steps,

1. The first category of input parameters (with directly available values in measurement-campaign details) are substituted the exact same value as in measurement campaign (e.g., link distance $d$, height of nodes $h_{m}$, and delay span $\tau_{\text {span }}$ ).
2. For the second category of input parameters, which cannot be directly set from the measurement campaign details due to insufficient available information, a reasonable range (constraints) is defined based on the available (if any) information about the environment in measurement-campaign details and from the behaviour (shape of graphs) of the measurement results.
3. An optimal value for each parameter from second category of the input parameters is set by minimizing the error in fitting by searching within the defined constraints.

Since, the variance of spread in the distribution is more important than the shape of distribution curves [12]; therefore, the absolute difference of the true standard deviation (TSD) [12] and RMS-DS of the proposed analytical results with empirical data-sets is taken as goodness-of-fit metric for AoA and ToA results, respectively, to gauge the fitness in mapping. In Fig. 19(a), the main lobe in the measurement data is caused by the far-end scattering region. The angular span of the main-lob can help in determining an appropriate value for the axes of far-end scattering volume compared to the link distance, i.e., $d / a_{1}$. Its shape being similar to bellshape indicate an equal value for major and intermediate axes and a significantly higher value for minor axes of far-end scattering volume, i.e., $a_{1}=b_{1}$ and $c_{1} \gg a_{1}$. Furthermore, no (very minor) peaks in the tails of the curve indicate an equal and very small value for all the axes of observing-end scattering volume, i.e., $a_{1}=b_{1}=$ $c_{1} \approx 0$. In Fig. 19(b) and (e), minor lobs along the tails of
distribution indicate a higher value for the ratio between major and intermediate axes of observing-end scattering volume, i.e., more elliptical-shaped base by setting a high value for $a_{1} / b_{1}$. Decrease in angular span and increase in sharpness of the peaks in tails can be ensured by setting a higher value for minor axis compared to major axis of observing-end scattering volume, i.e., increasing $a_{1} / c_{1}$. In Fig. 19(c), a minor tilt in the main-lob can be adapted by setting the orientation of far-end scattering volume, i.e., $-\varrho \leq \theta_{1} \leq+\varrho$, where $\varrho$ is a small value proportional to the level of tilt. In Fig. 19(d), a wide angular span with two similar slanted peaks can be modeled by increasing the values for far-end major and intermediate axes compared to the link distance, such that the major-lob caused by far-end volume overlaps a large portion of minor-lobs caused by observing-end scattering volume. Minor slant in the peaks can be tweaked by adjusting the orientation of both far- and observing-end scattering volumes, i.e., adjusting $\theta_{m}$. In Fig. 19(c), the more prominent peaks in the tails compared to the peak along LoS direction indicate a higher volume of observing-end compared to the farend scattering region, i.e., $V_{2}>V_{1}$. Establishing a direct relationship between the input geometric parameter of the model with physical parameters from the propagation environment is of high significance. The optimal setting of input physical and geometric parameters, for which a good agreement (based on TSD and RMS-DS) between measurement and analytical results is presented in Table II, for all the comparative graphs presented in the Fig. 18,19 , and 20 . It can be witnessed that the proposed model can effectively model a dynamic range of outdoor and indoor radio propagation environments; which is due to high degree of flexibility in geometric parameters of the proposed model. The proposed model not only effectively models the peaks of distribution of energy in angular (specially) and temporal domains, but also the tails of distribution; which is unlike its counterparts in the literature. The absolute difference between TSD of analytical (Proposed) and empirical curves is observed to be $0.05^{\circ}, 2.707^{\circ}, 0.55^{\circ}, 4.93^{\circ}, 4.38^{\circ}$, and $1.42^{\circ}$ in Fig. 19 (a), (b), (c), (d), (e), and (f), respectively. Moreover, the proposed model provides an improvement of $0.01^{\circ}$, $0.15^{\circ}, 0.03^{\circ}, 3.563^{\circ}, 1.721^{\circ}$, and $4.722^{\circ}$ in the accuracy of model fitness over measurement results shown in Fig. 19 (a), (b), (c), (d), (e), and (f), respectively, when compared to a notable 3D analytical model in the literature [17,

18]. Moreover, for indoor propagation environments, the difference in fitness of analytical (proposed) results over measurement data-sets is observed to be $0.12^{\circ}, 3.55^{\circ}$, and $0.34^{\circ}$ in Fig. 20 (a), (b), and (c), respectively. This good fit of the proposed analytical results on empirical datasets establishes the validity and usability of the proposed model.

Acknowledging the importance of comparing the proposed analytical results with appropriate M2M environments' measurement results, authors have conducted an exhaustive search for availability of such empirical datasets in the open literature. Among the available measurement data-sets presented in the open literature, most are ineligible, presented in graphical three-dimensional plots or contour maps, from which no numerical data can be reliably extracted. Overall, for all the propagation environments, only a few readable data-sets can be located, among which mostly are for outdoor environments. This limitation has restricted the authors to provide a comparison with only three indoor measurement results. However, to prove the validity and worth of the proposed model, a detailed comparison with the available outdoor and indoor measurement data-sets is presented. Furthermore, for M2M communication scenarios, a comparison of the proposed model with a notable analytical model in the literature $[17,18]$ and performed computer simulations is presented, where a good fit can be observed. The proposed model provides high degree of accuracy in modeling a certain targeted propagation environment (for M2M) compared to its counterparts in the literature. This high degree of accuracy is achieved at the cost of model's complexity. There may be practical applications, where high degree of accuracy is highly desirable even at very high computational complexity cost. Demands for such high accuracy in channel modeling of emerging wireless communication networks are obvious due to expected increase in the amount of antennas at BSs (highly direction links), increase in carrier frequency (millimeter Waves), decrease in elevation of BSs and size of cells, and increase in density of nodes with direct device-to-device links, etc. The proposed theoretical results can stand as a good reference for conducting measurement campaigns for M2M communication scenarios. The proposed performance analysis on azimuthal and elevational angular spread of M2M communication scenarios is also helpful in enhancing the data rates of MIMO streams and/or designing sharp and precise antenna
beamwidths using planar antenna arrays in correspondence with vertical and horizontal spacing among antenna elements.

## 4. CONCLUSIONS

A 3-D statistical channel model with high degrees of flexibility in geometric parameters has been proposed for emerging M2M communication environments. The proposed model has been demonstrated to effectively model a dynamic range of radio propagation environments. Independently rotatable and scalable ellipsoids have been considered around both the nodes of a communication link and another bounding ellipsoid has been considered associated with maximum propagation delay. The effective illuminated scattering region has been designed with the overlap region of this ellipsoid with the scattering ellipsoids around Tx and Rx. Analytical expressions for joint and marginal PDFs of ToA and AoA have been derived. The impact of various physical channel parameters such as link distance, antenna heights, size and position of scattering region, as well as shape and orientation of scattering region on delay and angular spread has thoroughly been investigated. In order to establish validity of the proposed model, the obtained analytical results have been compared with simulation results and with various experimental datasets given for different propagation scenarios in the literature. A good fit of proposed analytical results on empirical data-sets has been observed, which establishes the validity and flexibility of the proposed model.

A mathematical relationship between input geometric parameters (like, major, intermediate, and minor axes of scattering volumes and their horizontal orientation) can be established with constellation of the real environment in vicinity of Tx and Rx nodes, by conducting an intensive measurement campaign for M2M propagation environments. The fitness error optimization based parameters selection strategy (for some of geometric parameters) has been followed rather constellation of propagation environment. There remains the scope, to conduct a thorough study to determine such relationship of input geometric parameters with a constellation of the real environment in the future. Moreover, another possible

Table II: Comparison of the proposed analytical model with empirical data-sets found in the literature.

| Environment | Channel Statistics | Figure | Measurement Data-set | Measurement campaign region | Antenna(s) <br> Elevation | Proposed model input Parameters' setting | Absolute difference in RMS-DS | Absolute difference in angular TSD |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | ToA | Fig. 18(a) | K. I. Pedersen et al. <br> [20] (Figure 11) | Downtown, Aarhus, Denmark. | Above Rooftop level. | $\begin{gathered} a_{1} / b_{1} / c_{1}=3 / 1.5 / 1, \\ a_{2} / b_{2} / c_{2}=1 / 0.83 / 0.71, \\ h_{1}-h_{2}=25 \mathrm{~m}, z_{1}-z_{2}= \\ 5 \mathrm{~m}, d / a_{1}=20 / 11, \\ \left\|\theta_{1}-\theta_{2}\right\|=0^{\circ}, \tau_{\text {Span }}= \\ 3.64 \mu \mathrm{sec} \end{gathered}$ | $0.087 \mu \mathrm{sec}$ | - |
|  |  | Fig. 18(b) | K. I. Pedersen et al. [20] (Figure 9) | Downtown, Aarhus, Denmark. | Above Rooftop level. | $\begin{gathered} a_{1} / b_{1} / c_{1}=3 / 2 / 1, \\ a_{2} / b_{2} / c_{2}=3.4 / 1.7 / 1, \\ h_{1}-h_{2}=25 \mathrm{~m}, z_{1}-z_{2}= \\ 4 \mathrm{~m}, d / a_{1}=3 / 1, \\ \left\|\theta_{1}-\theta_{2}\right\|=0^{o}, \tau_{\text {span }}= \\ 3.66 \mu \mathrm{sec} \end{gathered}$ | $0.135 \mu \mathrm{sec}$ | - |
|  |  | Fig. 18(c) | M. D. Batariere et al. [21] (Figure 6) | Chicago, Illinois. | Rooftop level. | $\begin{gathered} a_{1} / b_{1} / c_{1}=3 / 2 / 1, \\ a_{2} / b_{2} / c_{2}=3.2 / 1.5 / 1,15 \mathrm{~m}, \\ h_{1}-h_{2}=22 \mathrm{~m}, z_{1}-z_{2}= \\ 4 \mathrm{~m}, d / a_{1}=3 / 1, \\ \left\|\theta_{1}-\theta_{2}\right\|=0^{o}, \tau_{\text {span }}= \\ 1.5 \mu \sec \mu \text { sec } \end{gathered}$ | $0.174 \mu \mathrm{sec}$ | - |
|  | $\begin{aligned} & \text { Azimuth } \\ & \text { AoA } \end{aligned}$ | Fig. 19(a) | K. I. Pedersen et al. <br> [20] (Figure 10) | Downtown, Aarhus, Denmark. | Above Rooftop level. | $\begin{gathered} a_{1} / b_{1} / c_{1}=9 / 2 / 7, \\ a_{2} / b_{2} / c_{2}=1 / 1 / 1, h_{1}-h_{2} \\ =0 \mathrm{~m}, z_{1}-z_{2}=0 \mathrm{~m}, d / a_{1}= \\ 6 / 1,\left\|\theta_{1}-\theta_{2}\right\|=0^{\circ}, \tau_{\mathrm{span}} \\ =0.85 \mu \mathrm{sec} \end{gathered}$ | - | $0.05^{\circ}$ |
|  |  | Fig. 19(b) | B. H. Fleury et al. [22] (Figure 16) | Downtown Area, Aalborg, Denmark | unspecified | $\begin{gathered} a_{1} / b_{1} / c_{1}=1.5 / 5 / 1, \\ a_{2} / b_{2} / c_{2}=1 / 1.3 / 1, \\ h_{1}-h_{2}=5 \mathrm{~m}, z_{1}-z_{2}= \\ 4 \mathrm{~m}, d / a_{1}=5 / 1 \mathrm{~m}, \\ \left\|\theta_{1}-\theta_{2}\right\|=-17^{\circ}, \tau_{\text {span }} \\ =0.65 \mu \mathrm{sec} \end{gathered}$ | - | $2.707^{\circ}$ |
|  |  | Fig. 19(c) | P. A. Matthews $e t$ al. [23] (Figure 7) | Open site, Valley area, flats and few buildings, Leeds, UK | Above Rooftop level | $\begin{gathered} a_{1} / b_{1} / c_{1}=1 / 1.2 / 1 \\ a_{2} / b_{2} / c_{2}=1.5 / 3 / 1, \\ h_{1}-h_{2}=0 \mathrm{~m}, z_{1}-z_{2}= \\ 0 \mathrm{~m}, d / a_{1}=1.5 / 1 \mathrm{~m}, \\ \left\|\theta_{1}-\theta_{2}\right\|=45^{\circ}, \tau_{\text {span }}= \\ 0.56 \mu \mathrm{sec} \end{gathered}$ | - | $0.55^{\circ}$ |
|  |  | Fig. 19(d) | C. Kloch et al. [24] (Figure 6) | Small Campus, Aalborg University. | Low ( $\approx 4 \mathrm{~m}$ ) . | $\begin{gathered} a_{1} / b_{1} / c_{1}=3 / 1 / 14, \\ a_{2} / b_{2} / c_{2}=25 / 50 / 1, \\ h_{1}-h_{2}=6 \mathrm{~m}, z_{1}-z_{2}= \\ 5 \mathrm{~m}, d / a_{1}=1 / 1, \\ \left\|\theta_{1}-\theta_{2}\right\|=-5^{\circ}, \tau_{\text {span }}= \\ 0.68 \mu \mathrm{sec} \end{gathered}$ | - | $4.93{ }^{\circ}$ |
|  |  | Fig. 19(e) | K. I. Pedersen et al. [25] (Figure 5) | Downtown, Aarhus, Denmark. | Rooftop level. | $\begin{gathered} a_{1} / b_{1} / c_{1}=1 / 1 / 1.5, \\ a_{2} / b_{2} / c_{2}=1 / 20 / 11, \\ h_{1}-h_{2}=15 \mathrm{~m}, z_{1}-z_{2}= \\ 5 \mathrm{~m}, d / a_{1}=3 / 1, \\ \left\|\theta_{1}-\theta_{2}\right\|=0^{o}, \tau_{\text {span }}= \\ 1.1 \mu \mathrm{sec} \end{gathered}$ | - | $4.38{ }^{\circ}$ |
|  |  | Fig. 19(f) | K. I. Pedersen et al. [20] (Figure 14) | Riverside, <br> Stockholm, Sweden. | Rooftop level (21m). | $\begin{gathered} a_{1} / b_{1} / c_{1}=1 / 1.6 / 1.6, \\ a_{2} / b_{2} / c_{2}=1 / 7 / 1, h_{1}-h_{2} \\ =19 \mathrm{~m}, z_{1}-z_{2}=6 \mathrm{~m}, d / a_{1}= \\ 2.2 / 1,\left\|\theta_{1}-\theta_{2}\right\|=159^{\circ}, \\ \tau_{\text {span }}=1.2 \mu \mathrm{sec} \end{gathered}$ | - | $1.42{ }^{\circ}$ |
| $\begin{aligned} & \sum_{d} \\ & \sum_{0} \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & 0 \\ & \hline \end{aligned}$ |  | Fig. 20(a) | X. Yin et al. [26] <br> (Figure 8(b) ) | In-office (LoS Scenario) | 1.5 m | $\begin{gathered} a_{1} / b_{1} / c_{1}=4.5 / 22 / 1, \\ a_{2} / b_{2} / c_{2}=2.4 / 1.5 / 1, \\ h_{1}-h_{2}=0 \mathrm{~m}, z_{1}-z_{2}=0 \mathrm{~m} \\ , d / a_{1}=1 / 1, \\ \left\|\theta_{1}-\theta_{2}\right\|=8.5^{\circ}, \tau_{\text {span }} \\ =210 \mathrm{nsec} \\ \hline \end{gathered}$ | - | $0.12{ }^{\circ}$ |
|  |  | Fig. 20(b) | X. Yin et al. [26] <br> (Figure 8(e) ) | In-office (NLoS Scenarios) |  | $\begin{gathered} a_{1} / b_{1} / c_{1}=2 / 4 / 1, \\ a_{2} / b_{2} / c_{2}=2.2 / 1.5 / 1, \\ h_{1}-h_{2}=0 \mathrm{~m}, z_{1}-z_{2}=0 \mathrm{~m} \\ , d / a_{1}=6 / 1, \\ \left\|\theta_{1}-\theta_{2}\right\|=42^{\circ}, \tau_{\text {span }}= \\ 160 \mathrm{nsec} \end{gathered}$ | - | $3.55^{\circ}$ |
|  |  | Fig. 20(c) | X. Yin et al. [26] (Figure 8(h)) |  |  | $\begin{gathered} a_{1} / b_{1} / c_{1}=1 / 2 / 0.0, \\ a_{2} / b_{2} / c_{2}=1 / 1 / 1, h_{1}-h_{2} \\ =0 \mathrm{~m}, z_{1}-z_{2}=0 \mathrm{~m}, d / a_{1}= \\ 1.5 / 1,\left\|\theta_{1}-\theta_{2}\right\|=70^{\circ}, \\ \tau_{\text {span }}=0.85 \mu \mathrm{sec} \\ \hline \end{gathered}$ | - | $0.34{ }^{\circ}$ |

extension of the proposed model can be the consideration of multi-bounce scattering into the account.

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