Determinantal expressions of certain integrals on symmetric spaces

Salem Said¹ and Cyrus Mostajeran²

¹ CNRS, Laboratoire Jean Kuntzmann (UMR 5224)
 ² School of Physical and Mathematical Sciences, NTU Singapore

Abstract. The integral of a function f defined on a symmetric space $M \simeq G/K$ may be expressed in the form of a determinant (or Pfaffian), when f is K-invariant and, in a certain sense, a tensor power of a positive function of a single variable. The paper presents a few examples of this idea and discusses future extensions. Specifically, the examples involve symmetric cones, Grassmann manifolds, and classical domains.

Keywords: symmetric space · matrix factorisation · random matrices.

1 Introduction

Riemannian symmetric spaces were classified by É. Cartan, back in the 1920s. A comprehensive account of this classification may be found in the monograph [1]. In the 1960s, a classification of quantum symmetries led Dyson to introduce three kinds of random matrix ensembles, orthogonal, unitary, and symplectic [2]. These three kinds of ensembles are closely related to the symmetric spaces known as symmetric cones, and also to their compact duals, which provide for so-called circular ensembles. More recently, Dyson's classification of quantum symmetries has been extended to free fermionic systems. It turned out that this extended classification is in on-to-one correspondance with Cartan's old classification of symmetric spaces [3]. This correspondance has motivated the notion that the relationship between random matrices and symmetric spaces extends well beyond symmetric cones, and is of a general nature (for example [4] or [5,6]).

The present submission has a modest objective. It is to show how the integral of a function f, defined on a symmetric space $M \simeq G/K$, can be expressed in the form of a determinant or Pfaffan, when f is K-invariant and satisfies an additional hypothesis, formulated in Section 4 below. This is not carried out in a general setting, but through a non-exhaustive set of examples, including symmetric cones, Grassmann manifolds, classical domains, and their duals (for the case of compact Lie groups, yet another example of symmetric spaces, see [7]).

The determinantal expressions obtained here, although elementary, are an analytic pre-requisite to developing the *random matrix theory of Riemannian symmetric spaces*. This long-term goal is the motivation behind the present work.

Unfortunately, due to limited space, no proofs are provided for statements made in the following. These will be given in an upcoming extended version. 2 Salem Said and Cyrus Mostajeran

2 Integral formulas

Let M be a Riemannian symmetric space, given by the symmetric pair (G, K). Write $\mathfrak{g} = \mathfrak{k} + \mathfrak{p}$ the corresponding Cartan decomposition, and let \mathfrak{a} be a maximal abelian subspace of \mathfrak{p} . Then, denote by Δ a set of positive reduced roots on \mathfrak{a} [1].

Assume that $\mathfrak{g} = \mathfrak{z}(\mathfrak{g}) + \mathfrak{g}_{ss}$ where $\mathfrak{z}(\mathfrak{g})$ is the centre of \mathfrak{g} and \mathfrak{g}_{ss} is semisimple and non-compact (\mathfrak{g}_{ss} is a real Lie algebra). The Riemannian exponential Exp maps \mathfrak{a} isometrically onto a totally flat submanifold of M, and any $x \in M$ is of the form $x = k \cdot \operatorname{Exp}(a)$ where $k \in K$ and $a \in \mathfrak{a}$.

Let $f: M \to \mathbb{R}$ be a K-invariant function, $f(k \cdot x) = f(x)$ for $k \in K$ and $x \in M$. There is no ambiguity in writing f(x) = f(a) where $x = k \cdot \text{Exp}(a)$. With this notation, there exists a constant C_M such that [1]

$$\int_{M} f(x) \operatorname{vol}(dx) = C_{M} \int_{\mathfrak{a}} f(a) \prod_{\lambda \in \Delta} \sinh^{m_{\lambda}} |\lambda(a)| \, da \tag{1}$$

where da is the Lebesgue measure on \mathfrak{a} .

The dual \hat{M} of M is a symmetric space given by the symmetric pair (U, K), where U is a compact Lie group, with the Cartan decomposition $\mathfrak{u} = \mathfrak{k} + \mathfrak{i}\mathfrak{p}$ $(\mathfrak{i} = \sqrt{-1})$. Now, Exp maps $\mathfrak{i}\mathfrak{a}$ onto a torus T which is totally flat in \hat{M} , and any point $x \in \hat{M}$ is of the form $x = k \cdot \operatorname{Exp}(\mathfrak{i}a)$ where $k \in K$ and $a \in \mathfrak{a}$.

If $f: M \to \mathbb{R}$ is K-invariant, there is no ambiguity in writing f(x) = f(t)where $x = k \cdot t$, t = Exp(ia). In this notation [1],

$$\int_{\hat{M}} f(x) \operatorname{vol}(dx) = C_M \int_T f(t) \prod_{\lambda \in \Delta} \sin^{m_\lambda} |\lambda(t)| \, dt \tag{2}$$

where dt is the Haar measure on T. Here, $\sin|\lambda(t)| = \sin|\lambda(a)|$ where $t = \exp(ia)$, and this does not depend on the choice of a.

3 Determinantal expressions

Let μ be a positive measure on a real interval *I*. Consider the multiple integrals,

$$z_{\beta}(\mu) = \frac{1}{N!} \int_{I} \dots \int_{I} |V(u_1, \dots, u_N)|^{\beta} \ \mu(du_1) \dots \mu(du_N)$$
(3)

where V denotes the Vandermonde determinant and $\beta = 1, 2$ or 4. Consider also the following bilinear forms,

$$(h,g)_{(\mu,1)} = \int_I \int_I (h(u)\varepsilon(u-v)g(v))\,\mu(du)\mu(dv) \tag{4}$$

$$(h,g)_{(\mu,2)} = \int_{I} h(u)g(u)\,\mu(du) \tag{5}$$

$$(h,g)_{(\mu,4)} = \int_{I} (h(u)g'(u) - g(u)h'(u))\,\mu(du) \tag{6}$$

Here, ε denotes the unit step function and the prime denotes the derivative. In the following proposition, det denotes the determinant and pf the Pfaffian. **Proposition 1.** The following hold for any probability measure μ as above. (a) if N is even,

$$z_{1}(\mu) = \operatorname{pf}\left\{\left(u^{k}, u^{\ell}\right)_{(\mu, 1)}\right\}_{k, \ell=0}^{N-1}$$
(7)

(b) on the other hand, if N is odd,

$$z_{1}(\mu) = \operatorname{pf} \left\{ \begin{array}{c} \left(u^{k}, u^{\ell}\right)_{(\mu,1)} \left(1, u^{k}\right)_{(\mu,2)} \\ -\left(u^{\ell}, 1\right)_{(\mu,2)} & 0 \end{array} \right\}_{k,\ell=0}^{N-1}$$
(8)

(c) moreover,

$$z_{2}(\mu) = \det\left\{\left(u^{k}, u^{\ell}\right)_{(\mu, 2)}\right\}_{k, \ell=0}^{N-1}$$
(9)

(d) and, finally,

$$z_4(\mu) = \operatorname{pf}\left\{ \left(u^k, u^\ell \right)_{(\mu, 4)} \right\}_{k,\ell=0}^{2N-1}$$
(10)

On the other hand, if μ is a probability measure on the unit circle S^1 , and

$$z_{\beta}(\mu) = \frac{1}{N!} \int_{S^1} \dots \int_{S^1} |V(u_1, \dots, u_N)|^{\beta} \ \mu(du_1) \dots \mu(du_N)$$
(11)

consider the bilinear form

$$(h,g)_{(\mu,1)} = \int_0^{2\pi} \int_0^{2\pi} (h(e^{ix})\varepsilon(x-y)g(e^{iy}))\,\tilde{\mu}(dx)\tilde{\mu}(dy)$$
(12)

where $\tilde{\mu}$ is the pullback of the measure μ through the map that takes x to e^{ix} , and let $(h,g)_{(\mu,2)}$ and $(h,g)_{(\mu,4)}$ be given as in (5) and (6), with integrals over S^1 instead of I.

Proposition 2. The following hold for any probability measure μ on S^1 . (a) if N is even,

$$z_1(\mu) = (-i)^{N(N-1)/2} \times \operatorname{pf}\left\{ (g_k, g_\ell)_{(\mu, 1)} \right\}_{k,\ell=0}^{N-1}$$
(13)

where $g_k(u) = u^{k - (N-1)/2}$.

(b) on the other hand, if N is odd,

$$z_{1}(\mu) = (-i)^{N(N-1)/2} \times \text{pf} \begin{cases} (g_{k}, g_{\ell})_{(\mu,1)} (1, g_{k})_{(\mu,2)} \\ -(g_{\ell}, 1)_{(\mu,2)} & 0 \end{cases} \begin{cases} N^{-1} \\ k_{k,\ell=0} \end{cases}$$
(14)

with the same definition of $g_k(u)$.

4 Salem Said and Cyrus Mostajeran

(c) moreover,

$$z_{2}(\mu) = \det\left\{\left(u^{k}, u^{-\ell}\right)_{(\mu,2)}\right\}_{k,\ell=0}^{N-1}$$
(15)

(d) and, finally,

$$z_4(\mu) = \operatorname{pf}\left\{ (h_k, h_\ell)_{(\mu, 4)} \right\}_{k,\ell=0}^{2N-1}$$
(16)

where $h_k(u) = u^{k-(N-1)}$.

Both of the above Propositions 1 and 2 are directly based on [8].

4 Main idea

An additional hypothesis is made on the function f(a) (in (1)) or f(t) (in (2)): that there exists a natural orthonormal basis $(e_j; j = 1, ..., r)$ of \mathfrak{a} , such that

$$f(a) = \prod_{j=1}^{r} w(a_j) \qquad f(t) = \prod_{j=1}^{r} w(t_j)$$
(17)

where w is a positive function of a single variable, and a_j are the components of a in the basis $(e_j; j = 1, ..., r)$, while $t_j = \text{Exp}(ia_j e_j)$. In this sense, it may be said that f is the r-th tensor power of w.

What is meant by *natural* is that (17) will imply that the integral (1) or (2) can be transformed into a multiple integral of the form (3) or (11), respectively. Thus, in the case of (1), there exists a measure μ on an interval I, which satisfies

$$\int_{M} f(x) \operatorname{vol}(dx) = \tilde{C}_{M} \times z_{\beta}(\mu) \qquad (\tilde{C}_{M} \text{ is a new constant})$$

and, in the case of (2), there is a measure μ on S^1 , which yields a similar identity. It should be noted that this measure μ will depend on the function w from (17).

Then, Propositions 1 and 2 provide a determinantal (or Pfaffian) expression of the initial integral on the symmetric space M or \hat{M} .

At present, this is not a theorem, but a mere idea or observation, supported by the examples in the following section.

5 Examples

5.1 Symmetric cones

Consider the following Lie groups (in the usual notation, as found in [1]).

β	G_{eta}	U_{eta}	K_{β}
1	$GL_N(\mathbb{R})$	U(N)	O(N)
2	$GL_N(\mathbb{C})$	$U(N) \times U(N)$	U(N)
4	$GL_N(\mathbb{H})$	U(2N)	Sp(N)

Then, $M_{\beta} \simeq G_{\beta}/K_{\beta}$ is a Riemannian symmetric space, with dual $\hat{M}_{\beta} = U_{\beta}/K_{\beta}$. In fact, M_{β} is realised as a so-called symmetric cone : the cone of positive-definite real, complex, or quaternion matrices (according to the value of $\beta = 1, 2$ or 4).

Each $x \in M_{\beta}$ is of the form $k\lambda k^{\dagger}$ where $k \in K_{\beta}$ and λ is a positive diagonal matrix (\dagger denotes the transpose, conjugate-transpose, or quaternion conjugate-transpose). If $f : M_{\beta} \to \mathbb{R}$ is K_{β} -invariant, and can be written $f(x) = \prod w(\lambda_j)$,

$$\int_{M_{\beta}} f(x) \operatorname{vol}(dx) = \tilde{C}_{\beta} \times z_{\beta}(\mu)$$
(18)

where $\mu(du) = (w(u)u^{-N_{\beta}})du$, with $N_{\beta} = (\beta/2)(N-1) + 1$, on the interval $I = (0, \infty)$. The constant \tilde{C}_{β} is known explicitly, but this is irrelevant at present.

The dual \hat{M}_{β} can be realised as the space of symmetric unitary matrices $(\beta = 1)$, of unitary matrices $(\beta = 2)$, or of antisymmetric unitary matrices with double dimension 2N, $(\beta = 4)$.

If $\beta = 1, 2$, then $x \in \hat{M}_{\beta}$ is of the form $ke^{i\theta}k^{\dagger}$ where $k \in K_{\beta}$ and θ is real diagonal. However, if $\beta = 4$, there is a somewhat different matrix factorisation,

$$x = k \begin{pmatrix} -e^{i\theta} \\ e^{i\theta} \end{pmatrix} k^{\text{tr}} \qquad (\text{tr denotes the transpose})$$
(19)

where $k \in Sp(N)$ is considered as a $2N \times 2N$ complex matrix (rather than a $N \times N$ quaternion matrix). If $f : \hat{M}_{\beta} \to \mathbb{R}$ is K_{β} -invariant, $f(x) = \prod w(e^{i\theta_j})$,

$$\int_{\hat{M}_{\beta}} f(x) \operatorname{vol}(dx) = \tilde{C}_{\beta} \times z_{\beta}(\mu)$$
(20)

where $\mu(du) = w(u)|du|$ on the unit circle S^1 $(|du| = d\varphi$ if $u = e^{i\varphi})$.

Remark : in many textbooks, \hat{M}_1 is realised as the space of real structures on \mathbb{C}^N , and \hat{M}_4 as the space of quaternion structures on \mathbb{C}^{2N} . The alternative realisations proposed here seem less well-known, but more concrete, so to speak.

5.2 Grassmann manifolds

Consider the following Lie groups (again, for the notation, see [1]).

$$\begin{array}{lll} \beta & G_{\beta} & U_{\beta} & K_{\beta} \\ 1 & O(p,q) & O(p+q) & O(p) \times O(q) \\ 2 & U(p,q) & U(p+q) & U(p) \times U(q) \\ 4 & Sp(p,q) & Sp(p+q) & Sp(p) \times Sp(q) \end{array}$$

Then, $M_{\beta} \simeq G_{\beta}/K_{\beta}$ is a Riemannian symmetric space, with dual $\hat{M}_{\beta} = U_{\beta}/K_{\beta}$. The M_{β} may be realised as follows [9] ($\mathbb{K} = \mathbb{R}, \mathbb{C}$ or \mathbb{H} , according to β),

$$M_{\beta} = \{x : x \text{ is a } p \text{-dimensional and space-like subspace of } \mathbb{K}^{p+q}\}$$
 (21)

Here, x is space-like if $|\xi_p|^2 - |\xi_q|^2 > 0$ for all $\xi \in x$ with $\xi = (\xi_p, \xi_q)$, where $|\cdot|$ denotes the standard Euclidean norm on \mathbb{K}^p or \mathbb{K}^q . Moreover, for each $x \in M_\beta$, $x = k(x_\tau)$ where $k \in K_\beta$ and $x_\tau \in M_\beta$ is spanned by the vectors

$$\cosh(\tau_j)\xi_j + \sinh(\tau_j)\xi_{p+j} \qquad j = 1, \dots, p$$

with $(\xi_k; k = 1, \dots, p+q)$ the canonical basis of \mathbb{K}^{p+q} , and $(\tau_j; j = 1, \dots, p)$ real $(p \leq q \text{ throughout this paragraph}).$

If $f: M_{\beta} \to \mathbb{R}$ is K_{β} -invariant, $f(x) = f(\tau)$, the right-hand side of (1) reads (the positive reduced roots can be found in [10])

$$C_{\beta} \int_{\mathbb{R}^{p}} f(\tau) \prod_{j=1}^{p} \sinh^{\beta(q-p)} |\tau_{j}| \sinh^{\beta-1} |2\tau_{j}| \prod_{i< j} |\cosh(2\tau_{i}) - \cosh(2\tau_{j})|^{\beta} d\tau \quad (22)$$

and this can be transformed into the form (3), by introducing $u_j = \cosh(2\tau_j)$. This will reappear, with $\beta = 2$ and p = q, in the following paragraph.

Now, the duals \hat{M}_{β} are real, complex, or quaternion Grassmann manifolds,

$$\hat{M}_{\beta} = \{x : x \text{ is a } p \text{-dimensional subspace of } \mathbb{K}^{p+q}\}$$
 (23)

For each $x \in \hat{M}_{\beta}$, $x = k(x_{\theta})$ where $k \in K_{\beta}$ and x_{θ} is spanned by the vectors

$$\cos(\theta_j)\xi_j + \sin(\theta_j)\xi_{p+j} \qquad j = 1, \dots, p$$

with $(\theta_j; j = 1, \ldots, p)$ real.

If $f: \hat{M}_{\beta} \to \mathbb{R}$ is K_{β} -inariant, $f(x) = f(\theta)$, the right-hand side of (2) reads

$$C_{\beta} \int_{(0,\pi)^{p}} f(\theta) \prod_{j=1}^{p} \sin^{\beta(q-p)} |\theta_{j}| \sin^{\beta-1} |2\theta_{j}| \prod_{i< j} |\cos(2\theta_{i}) - \cos(2\theta_{j})|^{\beta} d\theta \quad (24)$$

which can be transformed into the form (11), by introducing $u_j = \cos(2\theta_j)$. In [4], this is used to recover the Jacobi ensembles of random matrix theory.

Remark : the angles θ_j may be taken in the interval $(-\pi/2, \pi/2)$ instead of $(0, \pi)$. In this case, $|\theta_j|$ are the principal angles between x_{θ} and the subspace x_o spanned by $(\xi_j; j = 1, \ldots, p)$. By analogy, it is natural to think of $|\tau_j|$ as the 'principal boosts' (using the language of special relativity) between x_{τ} and x_o .

5.3 Classical domains

Consider, finally, the following Lie groups (again, for the notation, see [1]).

$$\begin{array}{lll} \beta & G_{\beta} & U_{\beta} & K_{\beta} \\ 1 & Sp(N,\mathbb{R}) & Sp(N) & U(N) \\ 2 & U(N,N) & U(2N) & U(N) \times U(N) \\ 4 & O^{*}(4N) & O(4N) & U(2N) \end{array}$$

Then, $M_{\beta} \simeq G_{\beta}/K_{\beta}$ is a Riemannian symmetric space, with dual $\hat{M}_{\beta} = U_{\beta}/K_{\beta}$. The M_{β} are realised as classical domains, whose elements are $N \times N$ complex matrices (if $\beta = 1, 2$) or $2N \times 2N$ complex matrices (if $\beta = 4$), with operator norm < 1, and which are in addition symmetric ($\beta = 1$) or antisymmetric ($\beta = 4$).

If $\beta = 1, 2$, then any $x \in M_{\beta}$ may be written

$$x = k_1(\tanh(\lambda))k_2 \tag{25}$$

where k_1 and k_2 are unitary $(k_2 = k_1^{\text{tr}}, \text{ in case } \beta = 1)$, and λ is real diagonal. However, if $\beta = 4$,

$$x = k \begin{pmatrix} -\tanh(\lambda) \\ \tanh(\lambda) \end{pmatrix} k^{\mathrm{tr}}$$
(26)

where k is $2N \times 2N$ unitary. If $f: M_{\beta} \to \mathbb{R}$ is K_{β} -invariant, and $f(x) = \prod w(\lambda_j)$,

$$\int_{M_{\beta}} f(x) \operatorname{vol}(dx) = \tilde{C}_{\beta} \int_{\mathbb{R}^{N}} \prod_{j=1}^{N} w(\lambda_{j}) \sinh |2\lambda_{j}| \prod_{i < j} |\cosh(2\lambda_{i}) - \cosh(2\lambda_{j})|^{\beta} d\lambda_{j}$$

After introducing $u_j = \cosh(2\lambda_j)$, this immediately becomes

$$\int_{M_{\beta}} f(x) \operatorname{vol}(dx) = \tilde{C}_{\beta} \times z_{\beta}(\mu)$$
(27)

where $\mu(du) = w(\operatorname{acosh}(u)/2)du$ on the interval $I = (1, \infty)$.

Remark : the domain M_2 is sometimes called the Siegel disk. As an application of (27), consider a random $x \in M_2$ with a Gaussian probability density function

$$p(x|\bar{x},\sigma) = (Z(\sigma))^{-1} \exp\left[-\frac{d^2(x,\bar{x})}{2\sigma^2}\right]$$
(28)

with respect to vol(dx), where $d(x, \bar{x})$ denotes Riemannian distance and $\sigma > 0$. Then, following the arguments in [6], (27) can be used to obtain

$$Z(\sigma) = \tilde{C}_2 \times \det \{m_{k+\ell}(\sigma)\}_{k,\ell=0}^{N-1} \quad m_j(\sigma) = \int_1^\infty \exp\left(-\operatorname{acosh}^2(u)/8\sigma^2\right) u^j du$$

The integrals $m_j(\sigma)$ are quite easy to compute, and one is then left with a determinantal expression of $Z(\sigma)$. The starting point to the study of the random matrix x is the following observation. If x is written as in (25) and $u_j = \cosh(2\lambda_j)$, then the random subset $\{u_j; j = 1, \ldots, N\}$ of $I = (1, \infty)$ is a determinantal point process (see [11]). By writing down its kernel function, one may begin to investigate in detail many of its statistical properties, including asymptotic ones, such as the asymptotic density of the (u_j) , or the asymptotic distribution of their maximum, in the limit where $N \to \infty$ (of course, with suitable re-scaling).

8 Salem Said and Cyrus Mostajeran

6 Future directions

The present submission developed determinantal expressions for integrals on symmetric spaces on a case-by-case basis, only through a non-exhaustive set of examples. Future work should develop these expressions in a fully general way, by transforming (1) and (2) into (3) or (11), for any system of reduced roots.

The long-term goal is to understand the *random matrix theory of symmetric spaces*. One aspect of this is to understand the asymptotic properties of a joint probability density (in the notation of (1))

$$f(a)\prod_{\lambda\in\Delta}\sinh^{m_{\lambda}}|\lambda(a)|\,da$$

and analyse how these depend on the set of positive reduced roots Δ . It is worth mentioning that, in previous work [6], it was seen that a kind of universality holds, where different root systems lead to the same asymptotic properties.

Random matrix theory (in its classical realm of orthogonal, unitary, and symplectic ensembles) has so many connections to physics, combinatorics, and complex systems in general. A further important direction is to develop such connections for the random matrix theory of symmetric spaces.

References

- 1. Helgason, S.: Differential geometry and symmetric spaces. Academic Press, (1962)
- Dyson, F.: The threefold way. Algebraic structures of symmetry groups and ensembles in quantum mechanics. Journal of Mathematical Physics 3(6), 1199–1215 (1962)
- Zirnbauer, M.: Symmetry Classes. The Oxford Handbook of Random Matrix Theory (Editors, G. Akemann, J. Baik, P. Di Francesco), (2018)
- 4. Edelman, A., Jeong, S.: On the Cartan decomposition for classical random matrix ensembles. Journal of Mathematical Physics **63**(6), (2022)
- 5. Santilli, L., Tierz, M.: Riemannian Gaussian distributions, random matrix ensembles, and diffusion kernels. Nuclear Physics B **973**, (2021)
- Said, S., Heuveline, S., Mostajeran, C.: Riemannian statistics meets random matrix theory: towards learning from high-dimensional covariance matrices. IEEE Transactions on Information Theory 69(1), 472–481 (2023)
- Meckes, E.S.: The random matrix theory of the classical compact groups. Cambridge University Press, (2019)
- 8. Mehta, M.L.: Random Matrices (Third Edition). Elsevier, (2004)
- 9. Huang, Y.: A uniform description of Riemannian symmetric spaces as Grassmannians using magic square. PhD Thesis, The Chinese University of Hong Kong (2007)
- Sakai, T.: On cut loci of compact symmetric spaces. Hokkaido Mathematical Journal 6, 136–161 (1977)
- Johansson, K.: Random matrices and determinantal processes. arXiv:matchph/0510038 (2005)