# Distributed Quantum Programming

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#### Abstract

In this paper we explore the structure and applicability of the Distributed Measurement Calculus (DMC), an assembly language for distributed measurement-based quantum computations. We describe the formal language's syntax and semantics, both operational and denotational, and state several properties that are crucial to the practical usability of our language, such as equivalence of our semantics, as well as compositionality and context-freeness of DMC programs. We show how to put these properties to use by constructing a composite program that implements distributed controlled operations, in the knowledge that the semantics of this program does not change under the various composition operations. Our formal model is the basis of a quantum virtual machine construction for distributed quantum computations, which we elaborate upon in the latter part of this work. This virtual machine embodies the formal semantics of DMC such that programming execution no longer needs to be analysed by hand. Far from a literal translation, it requires a substantial concretisation of the formal model at the level of data structures, naming conventions and abstraction mechanisms. At the same time we provide automatisation techniques for program specification where possible to obtain an expressive and user-friendly programming environment.

**Keywords:** quantum computing, measurement-based quantum computing, distributed computing, formal models, virtual machines.

### 1 Introduction

During the last decennia, quantum information has managed to become a significant field of research in the exact and applied sciences. Although it is a relatively new discipline one can currently discern several sub-disciplines such as quantum cryptography, information theory, computability, error correction,

fault tolerance, computations and of course there is also the experimental research on the construction of quantum computers (Nielsen and Chuang, 2000). Nevertheless, the development of quantum information as a proper computational domain of computer science is lagging behind. Indeed there is no such thing as a quantum computational paradigm. By this we mean a framework in which quantum problems can be expressed and solved in terms of data structures, algorithms, techniques such as abstraction, all of these wrapped up in an associated programming language. Paradigm building has proved to be an extremely useful approach in computer science. It has led to theoretically equal but practically very different programming frameworks, such as functional, imperative, logic and object-oriented programming. For this reason we expect this approach to be crucial also in developing quantum programming paradigms.

The first step in paradigm building is to construct a low-level quantum programming language and determine its properties. By low-level we mean that we need to define syntactical expressions for each physically allowed quantum operation: preparation, unitary transformation, measurement, combined with classical control expressions if so desired. The syntax in its own is not the goal, but rather a means by which to facilitate investigations with computer science techniques. First, we need to determine the functionality of a quantum program, its semantics. The most obvious way to do this is the operational semantics, the most practical is the denotational semantics. While in the former a program's meaning is given as a sequence of state-changing operations, in the latter it is instead a mathematical object. Paramount is linking both, so that one can use whichever in future analyses. Through a programming language's semantics one can investigate notions such as composition and context-freeness. These properties are crucial when one wants to build more complex programs. Indeed, they allow the semantics of these larger programs to be built up from that of smaller components using rules for composition, rather than having to be determined from scratch. While these properties may seem obvious, computer science is littered with examples where they were mistakenly taken to be true, leading to problems in programming language development (see for example (Brock and Ackerman, 1981)).

Recent advances in quantum communication and cryptographic computations motivate the need for a programming paradigm centred on distributed quantum computations. In a distributed system one has concurrently acting agents in separated locations, operating locally on a quantum state, which may be entangled over agents, and coordinating their actions via communication. Formal frameworks for distributed quantum computation have only very recently begun to appear. Typically, these are a combination of classical process theory, which formalises notions of concurrency and communication, the quantum circuit model, i.e. local operations are unitary transformations of an agent's qubits, and given initial shared entanglement. First, there is the work in (Lalire and Jorrand, 2004; Jorrand and Lalire, 2005), which is directly built upon classical process calculi. While this model profits from being closely related to existing classical models, the disadvantage is that it is hard to focus on quantum behaviour. A different approach is advanced in the model known as communicating quantum processes or CQP (Gay and Nagarajan, 2004; Gay et al., 2005), where the typical communication primitives of process calculi are combined with computational primitives from QPL (Selinger, 2003). The basic model of CQP is more transparent. This model serves as a basis for the development of formal verification techniques, in particular for proving the security of larger scale quantum communication systems. In related work, a probabilistic model-checking tool built upon an existing automated verification component is developed (Gay et al., 2005). There is also the work in (Adão and Mateus, 2005), which is specifically tailored to security issues in cryptographic protocols. In our work, however, we are much more interested in the expressiveness of quantum distributed computations and the behavioural properties of distributed primitives. In a way, we take the inverse approach, assuming that computations are well-defined and investigating what programming concepts are at work, instead of the other way around.

While the fact that formal verification tools for distributed quantum computation are currently under development may suggest that a mature distributed paradigm already exists, this is actually not the case. Distributed protocols are still very much conceived on intuition, and there is no good notion or formalisation of the programming concepts that are at work. We therefore require adequate formal tools with which to explore and evolve the distributed quantum computing paradigm. In this paper we explore the structure and applicability of the Distributed Measurement Calculus (DMC) (D'Hondt, 2005; Danos et al., 2005), an assembly language for distributed measurement-based quantum computations. We describe its syntax and semantics, both operational and denotational, and state several properties that are crucial to the practical usability of our language, such as equivalence of our semantics, as well as compositionality and context-freeness of DMC programs. We show how to put these properties to use by constructing a composite program that implements distributed controlled operations, and demonstrate that the semantics of this program does not change under the various composition operations. Finally, we elaborate on an implementation for the DMC language which we developed under the form of a virtual machine, a platform-independent programming environment that abstracts away details of the underlying hardware or operating system. This virtual framework is a crucial step in providing DMC as an experimentation platform for distributed quantum computations, as reasoning within the formal model by hand proves quite cumbersome for even small computations. At the basis of our work lies the measurement calculus (Danos et al., 2007), a low-level quantum programming language for measurement-based quantum computations from which we explore the distributed paradigm. Next to the obvious advantage of starting from a proper formal framework, measurement-based models are well-suited as a starting point for distributed quantum computations because they are inherently distributed. What is important here is that the well-known physical framework of quantum computation is ported to an equivalent computer science framework, opening up the field towards investigations from this branch of science as well.

The structure of this paper is as follows. In the next section we discuss the syntax of our language, while Sec. 3 covers the semantics of DMC and lists its properties. Sec. 4 applies the previous to a concrete example, a composite protocol implementing distributed controlled gates. We discuss the implementation of the quantum virtual machine for DMC in Sec. 5, and conclude in Sec. 6. Some basic knowledge of quantum computation is assumed; for the reader not familiar with the domain we refer to the excellent (Nielsen and Chuang, 2000). Our approach in this article is to explain the notions of the model by example, rather than providing a series of formal definitions which are hard to interpret

and for which space it too limited. However, we stress that the model is a rigorous one, and, while this paper is a stand-alone document, refer the interested reader to the appendix for the full formal semantics of the DMC language and (D'Hondt, 2005; Danos *et al.*, 2005) for complete definitions.

## 2 Syntax

The language we propose is, broadly speaking, an assembly language for distributed measurement-based quantum computations. It is an assembly language in that it provides syntax only for the most basic operations while at the same time being universal. It is measurement-based as we build our language on top of the measurement calculus (MC) (Danos et al., 2007), an assembly language for measurement-based quantum computations. The latter depart from the usual circuit-based approach to quantum computing, where unitary transformations are the driving force of computations. While measurement operations were long seen as a necessary but disruptive part of quantum computing, in algorithms such as teleportation they act as an essential part of the computation. This gave rise to models where the computation is steered by pre-established generic entanglement combined with measurements, such as the one-way quantum computer (Raussendorf et al., 2003). Because measurements are inherently probabilistic, correction operations are required that are conditionally applied depending on previous measurement outcomes, thus rendering the computation deterministic. All of this is nicely captured in the measurement calculus. For this reason, as well as the inherent distributive aspect of measurement-based models, MC is an ideal basis from which to develop our language.

We describe our language model in three layers. The base layer consists of MC patterns, which describe local quantum computations. These are combined with distribution primitives into the notion of agents in the middle layer. Finally, agents and their shared entanglement resources are bundled into networks in the top layer.

**Patterns.** In MC a pattern is defined by a sequence of commands together with sets of qubits for working, input and output memory. As an example consider the following pattern, which, as we explain below, implements the Hadamard operation,

$$\mathcal{H}(1,2) := (\{1,2\},\{1\},\{2\},X_2^{s_1}M_1^X E_{12}). \tag{1}$$

All qubits are uniquely identified using numbers. The first argument denotes that the pattern has a computation space of two qubits, referenced by 1 and 2. The next two arguments specify the pattern's inputs and outputs. Working qubits that are not input qubits are initialised to  $|+\rangle = |0\rangle + |1\rangle$ . The last argument is the pattern's command sequence, a list of operations taken from a basic set and executed from right to left, analogous to matrix applications. Subscripts denote qubit arguments of the operation, while corrections are conditionally executed depending on their superscript. Concretely, one runs the pattern  $\mathcal H$  by preparing the first qubit in some input state  $|\psi\rangle$  and the second

in state  $|+\rangle$ , then entangling both qubits with the controlled-Z operation to obtain  $CZ_{12}(|\psi\rangle\otimes|+\rangle)$  (syntax:  $E_{12}$ ). Next, the first qubit is measured in the  $\{|+\rangle,|-\rangle\}$  basis (syntax:  $M_1^X$ ), where  $|-\rangle=|0\rangle-|1\rangle$ . Finally, an X correction is applied on the output qubit if the measurement outcome was  $s_1$  (syntax:  $X_2^{s_1}$ ). Here  $s_1$  defines a signal-simply 0 or 1-coming from the X-basis measurement on qubit 1 in  $M_1^X$ . If the signal is 0, the operation is not performed. Measurements are considered to be destructive, which is why qubit 1 does not appear in the output set.

A general pattern is denoted  $\mathcal{P}(V, I, O, \mathcal{A})$ , with computation space V, inputs I and outputs O (together called the pattern type), and command sequence  $\mathcal{A}$  that consists of entanglements  $E_{ij}$ , measurements  ${}^t[M_i^{\alpha}]^s$ , or corrections  $X_i^s$  or  $Z_i^s$ , where  $i, j \in V$ ,  $\alpha \in [0, 2\pi]$  and  $s, t \in \mathbb{Z}_2$ . Allowed measurements are those in the XY-plane of the Bloch sphere, and are specified by the angle  $\alpha$  ( $M_1^X$  is actually syntactic sugar for  $M_1^0$ ). Measurement angles may also be conditioned by signals, written  ${}^t[M_i^{\alpha}]^s$ , with  $(-1)^s\alpha + t\pi$  being the actual measurement angle. The four basic instructions together with signal conditioning suffice to make the model universal (Raussendorf et al., 2003; Danos et al., 2007).

Patterns can be combined into larger ones to create arbitrary quantum computations. The *sequential composition* of patterns  $\mathcal{P}_1 = (V_1, I_1, O_1, \mathcal{A}_1)$  and  $\mathcal{P}_2 = (V_2, I_2, O_2, \mathcal{A}_2)$ , with  $O_1 = I_2$ , is defined as

$$\mathcal{P}_2 \mathcal{P}_1 := (V_1 \cup V_2, I_1, O_2, \mathcal{A}_2 \mathcal{A}_1) , \qquad (2)$$

while the parallel composition of the same patterns is defined as

$$\mathcal{P}_1 \otimes \mathcal{P}_2 := (V_1 \uplus V_2, I_1 \uplus I_2, O_1 \uplus O_2, \mathcal{A}_1 \mathcal{A}_2) . \tag{3}$$

Note that one can always rename qubits for parallel composition to become possible, while sequential composition also needs  $I_2$  and  $O_1$  to have the same cardinality.

As an example, here is the pattern to create a 3-fold GHZ-state  $|000\rangle + |111\rangle$ ,

$$GHZ_{123} = (\{1, 2, \hat{2}, 3, \hat{3}\}, \cdot, \{1, 2, 3\}, \mathcal{H}(\hat{3}, 3)E_{2\hat{3}}\mathcal{H}(\hat{2}, 2)E_{1\hat{2}}),$$
(4)

where  $\hat{2}$  and  $\hat{3}$  are working qubits.

MC is equipped with a powerful standardisation theorem which provides a procedure for bringing any pattern, such as the one above, into EMC-form, i.e. with entanglements first and corrections last. This is important from an experimental implementation perspective, and also corresponds nicely with the typical structure of a distributed quantum protocol where one starts out with a shared entanglement resource. In fact within MC we can already express distributed computations such as teleportation,

$$(\{1,2,3\},\{1\},\{3\},X_3^{s_2}Z_3^{s_1}M_2^XM_1^XE_{12}E_{23}).$$
 (5)

While the pattern describes the intended computation, we find no notion of separate parties participating in the teleportation, and neither of the required communication between them. Purely by convention, and because everybody knows the protocol, we can say that qubits 1 and 2 belong to Alice and qubit

 $<sup>^1\</sup>mathrm{A}$  controlled-Z operation on two qubits applies the Z operation to the second qubit provided the first is set to 1.

3 to Bob. It is still hard to see if and what both parties have to communicate, e.g. Bob needs  $s_1$  and  $s_2$  but has no access to qubits 1 and 2, thus Alice needs to perform the measurements and communicate the results to Bob. We want this information to be explicit and obvious in the language, which in turn makes it easier to describe and reason about distributed quantum programs. To see that this is important, try to identify what is implemented with the following distributed protocol,

$$(\{0,\bar{0},1,\bar{1},2,\bar{2}\},\cdot,\{0,1,2\},Z_2^{s_{\bar{2}}}Z_1^{s_{\bar{1}}}Z_0^{s_{\bar{0}}}M_{\bar{0}\bar{1}\bar{2}}^{GHZ}E_{0\bar{0}}E_{1\bar{1}}E_{2\bar{2}}). \tag{6}$$

Here  $M^{GHZ}$  is a pattern for GHZ-measurement. Note that pattern (4) implements the unitary transformation between the diagonal basis and the GHZ-basis  $\{|i\rangle+|\overline{i}\rangle\}_{i=0}^{2^n-1}$ . Hence a GHZ-measurement is executed by applying the inverse pattern followed by a measurement in the diagonal basis. We will get back to this pattern in Sec. 4.

Agents. An agent embodies a single computation node running in isolation in a distributed algorithm, i.e. a processor or a party such as Alice or Bob. Each agent has a local command sequence, which operates on a set of resources contained within its environment. Agent instructions are either measurement patterns or communication primitives. For example, the teleportation pattern (5) really needs the following agent definitions for Alice and Bob, respectively,

$$\mathbf{A}: \{1, 2\}. \mathbf{c}! s_2. \mathbf{c}! s_1. M_2^X M_1^X E_{12} \text{ and } \mathbf{B}: \{3\}. X_3^{s_2} Z_3^{s_1}. \mathbf{c}? s_2. \mathbf{c}? s_1$$
. (7)

As we can see an agent definition consists of a type in curly braces, which specifies what qubits an agent owns, and an instruction sequence. Next to local quantum operations, written in MC language, we specify the exchange of the signals  $s_1$  and  $s_2$  between Alice and Bob via the communication primitives c? and c!. What is missing here as compared to pattern (5) is the prior shared entanglement  $E_{23}$ , as indeed it is impossible to factorise this part of the protocol into agent definitions. There is no other option than to specify shared entanglement separately in the network definition for the full protocol, as we shall see below. Of course the qubits are still local to the agents – as reflected in the type – only their description is not.

Formally, an agent is defined by an expression  $\mathbf{A}(i,o):Q.\mathcal{E}$ , where the type Q is a set of qubit references and  $\mathcal{E}$  is a finite instruction sequence composed of pattern command sequences  $\mathcal{A}$ , classical message reception  $\mathbf{c}$ ?s and sending  $\mathbf{c}$ !s, where s is a signal, and quantum reception  $\mathbf{q}\mathbf{c}$ ?q and sending  $\mathbf{q}\mathbf{c}$ !q, where q is a qubit reference. Corresponding communication actions are synchronised, meaning that an agent executing a receive will pause its program until the required agent has sent a message on the same channel and vice versa. Classical inputs i and outputs o are used in protocols such as superdense coding, and are not mentioned when there are none. Also, working qubits required by patterns are initialised to  $|+\rangle$  as before and are not specified in the type Q.

**Networks.** A network of agents consists of several agents that execute their command sequence concurrently. Typically, quantum resources are shared between agents – indeed most distributed quantum protocols benefit from some type of shared entanglement being present prior to the start of the protocol. We

have no way of splitting the representation of these states over agent definitions, and instead have to specify agents separately in a network's definition. This is why a network specification is more than just a collection of agent definitions, a feature specific to quantum concurrent frameworks and absent in classical concurrency. To keep track of how a shared resource is distributed each agent's type Q specifies which qubits of a shared resource are local to that agent.

At this point we can finally write down the full teleportation protocol, including all relevant distributed aspects.

$$TP := \mathbf{A} : \{1, 2\}. \mathsf{c!} s_2. \mathsf{c!} s_1. M_2^X M_1^X E_{12}$$

$$\mid \mathbf{B} : \{3\}. X_3^{s_2} Z_3^{s_1}. \mathsf{c?} s_2. \mathsf{c?} s_1$$

$$\parallel E_{23}$$

$$(8)$$

Analogous to process algebraic notation we use a vertical bar | to separate concurrently executing agents. Shared network resources are specified after the double bar ||. For teleportation this is the state  $E_{23}$ , which was produced and handed out to Alice and Bob prior to the execution of the protocol. If so desired the establishment of shared resources can itself be seen as a distributed protocol involving a server agent polled for entanglement services. We now likewise lift the pattern given in (6) to a distributed setting, obtaining the following

$$ES := \mathbf{L} : \{\bar{0}, \bar{1}, \bar{2}\}. \mathbf{c}_{0}! s_{\bar{0}}. \mathbf{c}_{1}! s_{\bar{1}}. \mathbf{c}_{2}! s_{\bar{2}}. M_{\bar{0}\bar{1}\bar{2}}^{GHZ} .$$

$$| \mathbf{A}_{0} : \{0\}. Z_{0}^{x_{\bar{0}}}. \mathbf{c}_{0}? x_{\bar{0}}$$

$$| \mathbf{A}_{1} : \{1\}. Z_{1}^{x_{\bar{1}}}. \mathbf{c}_{1}? x_{\bar{1}}$$

$$| \mathbf{A}_{2} : \{2\}. Z_{2}^{x_{\bar{2}}}. \mathbf{c}_{2}? x_{\bar{2}}$$

$$|| E_{0\bar{0}} E_{1\bar{1}} E_{2\bar{2}} ,$$

$$(9)$$

where  $M_{\overline{0}\overline{1}\overline{2}}^{GHZ}$  is a pattern executing a GHZ-measurement. This is the the entanglement swapping protocol for three agents (Zukowski *et al.*, 1993; Bose *et al.*, 1998), which produces a GHZ-state shared between  $\mathbf{A}_0$ ,  $\mathbf{A}_1$  and  $\mathbf{A}_2$ . Because the resulting GHZ state is in the diagonal basis we have Z rather than X corrections in the network specification.

Arbitrary networks are denoted  $\mathcal{N} = |_i \mathbf{A}_i(\mathbf{i}_i, \mathbf{o}_i) : Q_i.\mathcal{E}_i \parallel \sigma$ . For simplicity, we assume that networks of agents satisfy a number of definiteness conditions which ensure that the computation is well-defined. For example, an agent may only operate on qubits he owns, and every communication event needs to have a corresponding dual event in the network. Of course these issues are important but we are glossing over them here as we assume checking these occurs in a pre-compilation step rather than at runtime.

Our extensions to the measurement calculus have made distributed notions explicit. We can now see directly from a protocol specification what agents have to communicate to whom as well as which assumptions are made about non-local entanglement resources. Also, because each agent's instructions are expressed separately, we do not impose any particular execution order for local quantum operations. This makes it clear which parts of the protocols can and cannot run concurrently.

## 3 Semantics

In the previous section we established the syntax of DMC. The next step is to establish the meaning of a program written down in this language, i.e. its semantics. Formal semantics is a means of assigning objects to chunks of code so that one can reason with these objects rather than with the code itself. In this section we describe an operational as well as a denotational semantics for DMC programs and, with this in hand, state some important properties of our framework.

**Definition.** A network's operational semantics reflects how it affects the state of a distributed system on which the network is run. That is, we associate with each agent a local state (quantum and classical), and specify how the network specification updates these. Since we always have to take the shared entanglement resources into account the operational semantics does not decompose into state transformers for each of the agents separately. On top of this a crucial ingredient of the semantics specifies how quantum resources move throughout the system. Concurrency is not really an issue for the semantics because, due to the linearity of quantum mechanics, the order in which local operations are applied is unimportant. Hence we pick an order arbitrarily and derive the semantics for this case.

Concretely, the operational semantics of a network is found by collapsing individual small-step semantical rules into one transition system. The full smallstep semantics of DMC can be found in the appendix; it consists of a number of quantum mechanical rules - corresponding to measurement pattern commands - and a number of communication rules. These affect the global quantum state (EMC commands), the local resources (quantum communication), and the local memory (M commands and classical communication). Since measurement is probabilistic, so are the small-step rules. The formal semantics of measurement pattern commands is well established (Danos et al., 2007; D'Hondt, 2005), and just needs to be lifted to the DMC setting. We elaborate here the semantics of typically distributed concepts introduced in the previous section. The formal semantics of MC relies on the component  $\Gamma$ , the outcome map, which contains a number of bindings from signal names to measurement outcomes. In DMC the outcome map is lifted to a local agent memory recording also the values of classical messages. Concretely, a classical communication event between two agents has the following semantics, given that  $\Gamma_b$ , the local memory of Bob, evaluates the name y to the value v,

$$\mathbf{A}: (Q_a, \Gamma_a).\mathcal{E}_a.\mathsf{c}?x \mid \mathbf{B}: (Q_b, \Gamma_b).\mathcal{E}_b.\mathsf{c}!y$$

$$\Longrightarrow \mathbf{A}: (Q_a, \Gamma_a[x \to v]).\mathcal{E}_a \mid \mathbf{B}: (Q_b, \Gamma_b).\mathcal{E}_b \quad . \tag{10}$$

Here  $\Gamma_a[x \downarrow v]$  means that we assign a new binding of x to v in the local memory of Alice. These rules are typically much harder to read and write down formally than to apply concretely. What we are saying here is that if a classical communication takes place, the value is looked up by the sending agent, and received and bound to a new name by the receiving agent, after which the computation ( $\mathcal{E}_a$  and  $\mathcal{E}_b$ ) continues. Sending and receiving qubits over a quantum channel changes the types Q of the agents involved. An agent sending a qubit can no longer perform any operations on it, therefore the corresponding

qubit reference is removed from Q, while it is added to the receiving agent. Formally, the rule for the exchange of a qubit with reference i is given by

$$\mathbf{A}: (Q_a, \Gamma_a).\mathcal{E}_a.\operatorname{qc}?i \mid \mathbf{B}: (Q_b, \Gamma_b).\mathcal{E}_b.\operatorname{qc}!i$$

$$\Longrightarrow \mathbf{A}: (Q_a \cup \{i\}, \Gamma_a).\mathcal{E}_a \mid \mathbf{B}: (Q_b \setminus \{i\}, \Gamma_b).\mathcal{E}_b . \tag{11}$$

It goes without saying that the way in which the actual communication of a qubit is implemented is more intricate than this simple rule. However, by providing quantum communication as primitive syntax we are precisely choosing not to get into these implementation matters. Both rules mentioned here do not affect the global quantum state, which is why it is not mentioned. This is not the case for the small-step semantics for pattern commands, which moreover may be probabilistic. The full small-step semantics can be found in the appendix or in (D'Hondt, 2005; Danos  $et\ al.$ , 2005).

We obtain the operational semantics by defining computation paths in the usual way and gathering all small steps in a computation path in one big step from initial to final state in that path. As such the operational semantics of a quantum distributed network is essentially a probabilistic transition system (PTS). However, since resource allocation is crucial, we have to augment this PTS with information on how qubits move throughout the network. This is formalised as a type signature. We denote the operational semantics of a network  $\mathcal{N}$  by  $[\![\mathcal{N}]\!]_{op}$ . For example, the operational semantics of TP is given by the deterministic transition system

$$[TP]_{op}: (\{1\}, \cdot) \to (\cdot, \{3\}) \cdot \rho_1 \Longrightarrow \rho_3 ,$$
 (12)

where  $\rho$  is the density matrix specifying the input quantum state to be teleported and subscripts indicate qubit systems. The operational semantics of the ES network given in (9) is also deterministic.

$$[ES]_{op}: (\cdot, \cdot, \cdot, \cdot) \to (\cdot, \{0\}, \{1\}, \{2\}) \cdot \mathbf{0} \Longrightarrow GHZ_{012}^D, \tag{13}$$

where the superscript D denotes that the resulting state is in the diagonal basis. Note that we only write real quantum inputs in the type while not mentioning entanglement resources. The good thing about operational semantics is that in principle it can be derived automatically by induction on the small-step rules.

Denotational semantics is a second means of assigning a formal meaning to a chunk of code, this time by way of mathematical objects. If we look at the skeleton of a distributed protocol, i.e. the equivalent non-distributed pattern, we find a multi-local probabilistic quantum operation, which is mathematically represented by (a special type of) completely positive map (CPM). Since protocols with different distribution of resources are observationally different, we have to pair this CPM with a function mapping input to output resources, formally represented by a type signature. We denote this type of semantics by  $[\![\mathcal{N}]\!]_{de}$  for a network  $\mathcal{N}$ . For example, the denotational semantics of TP is given by the map

$$[TP]_{de}: (\{1\}, \cdot) \to (\cdot, \{3\}) \cdot \mathcal{I} ,$$
 (14)

i.e. TP implements the identity map from qubit 1 to qubit 3. Here we see the importance of the type signature in specifying the semantics. The subtle difference between both types of semantics becomes more apparent once one starts

investigating more complex protocols involving mixed states and probabilistic semantics.

As we see from the examples both types of semantics are quite similar, and indeed we have the following result, which we state without proof.

**Proposition 3.1** There is a precise correspondence between the operational and the denotational semantics of networks of agents, that is to say

$$\forall \mathcal{N}_1, \mathcal{N}_2 : \mathcal{N}_1 \equiv_{op} \mathcal{N}_2 \iff \mathcal{N}_1 \equiv_{de} \mathcal{N}_2 . \tag{15}$$

Two networks are semantically equivalent if they have the same semantics.<sup>2</sup> While the equivalence between semantics may seem obvious (certainly from the example), it is crucial to prove this statement formally, as it allows one to switch between semantics where appropriate without giving the issue further thought. We note that in general this equivalence is not guaranteed.

Building larger protocols. With syntax and semantics in hand we can now consider constructing larger networks from smaller components. We consider parallel as well as sequential composition of networks, denoted by  $\otimes$  and  $\circ$  respectively. These operations are defined in the obvious way so we will not spell out concrete definitions here. Suffice to say that one needs to pay attention to agent as well as qubit names to ensure that networks are connected as desired. More importantly we need to make sure that these forms of composition are consistent with the semantics, as is indeed the case.

**Proposition 3.2** The semantics of networks is compositional, i.e.

$$[\mathcal{N}_2 \circ \mathcal{N}_1] = [\mathcal{N}_2] \circ [\mathcal{N}_1]. \tag{16}$$

$$[\mathcal{N}_1 \otimes \mathcal{N}_2] = [\mathcal{N}_2] \otimes [\mathcal{N}_1]$$

$$\tag{17}$$

Here we have a first application of Prop. 3.1, as first we do not need to specify which type of semantics we mean in the statement of Prop. 3.2, and second we can choose the most convenient semantics for the proof, which is the denotational semantics. While the compositions on the left hand side are at the level of agent and network definitions, the composition of network semantics on the right hand side is a functional one at the level of type signatures, PTS's and CPMs. For the full proof we refer to (D'Hondt, 2005); an example is given in Sec. 4.

A second important notion when constructing larger protocols, closely related to compositionality, is that of *contexts*. Informally speaking a context is a program with a "hole" in which a network specification can be inserted, typically denoted  $C[\cdot]$ . On some occasions, one finds that programs that are considered equivalent in isolation no longer behave in the same way when placed within the context of a larger program. This is particularly the case in concurrent systems. A historical example is the *Brock-Ackerman anomaly* (Brock and Ackerman,

<sup>&</sup>lt;sup>2</sup>For example, one can prove that qubit communication between two agents is semantically equivalent to the teleportation network.

1981), by which it was realised that simple input-output behaviour is not enough to reason about equivalence of concurrent systems. In our framework contexts  $C[\mathcal{N}]$  are given by arbitrary network compositions C in which our network  $\mathcal{N}$  is placed. What is different with ordinary compositionality is that the quantum resources of the network  $\mathcal{N}$  are no longer considered to be independent of those of its context. Indeed in Prop. 3.2 quantum resources are considered to be provided independently, such that composite networks operate on disentangled inputs. While this is sensible when each of the networks operate in isolation, it is less so when a network is only one factor in a complex compositional structure. Hence we first need to show that our semantics is independent of so-called entanglement contexts. This is a consequence of the following proposition, which we state here without proof.

**Proposition 3.3** Suppose  $\mathcal{L}: \rho_A \longrightarrow \rho_B = \sum_k L_k \rho_A L_k^{\dagger}$  is a completely positive map. Then for all quantum states  $\rho_{AC}$  applying  $\mathcal{L}$  to the A-part of  $\rho_{AC}$  results in

$$\rho_{BC} = \sum_{k} (L_k \otimes I_C) \rho_{AC} (L_k^{\dagger} \otimes I_C). \tag{18}$$

The proof, though easy, is not trivial. Using this proposition and compositionality we have the following important result.

**Corollary 3.4** Equivalence holds in arbitrary contexts, that is, if  $\mathcal{N}_1 \equiv \mathcal{N}_2$ , then for all network contexts  $C[\cdot]$  we have that  $C[\mathcal{N}_1] \equiv C[\mathcal{N}_2]$ .

## 4 Applications

We now show how all of the formal ingredients can be put to use in a concrete example. The distributed primitive we will investigate is that of a distributed remotely controlled gate (Yimsiriwattana and Lomonaco, 2005). Near-future quantum computers are expected to have only a limited number of qubits per machine. Even in quantum simulating environments the current qubit limit is only about 36 (Raedt et al., 2006). Hence one can imagine that quantum computations need to proceed much like cluster computations today, with resources spread over different processors in order to make them feasible. One common situation would be where the central processor needs to execute a controlled operation with the target qubits spread over a group of agents. Following the ideas in (Yimsiriwattana and Lomonaco, 2005), once we have a GHZ resource we can execute a distributed controlled gate. Calling the central processor L (for Leader), and assuming there are two subordinate processors A and B, the trick is to establish a shared control qubit between target agents, which is achieved through the share control (SC) protocol as follows,

$$SC := \mathbf{L} : \{c, 0\}. c_{1,2}! s_0. M_0^X E_{0c} .$$

$$| \mathbf{A}_1 : \{1\}. X_1^{x_0}. c_1? x_0$$

$$| \mathbf{A}_2 : \{2\}. X_2^{x_0}. c_2? x_0$$

$$| GHZ_{012}^{D}$$
(19)

Here qubit c is the input control qubit which is in the state  $\alpha|0\rangle + \beta|1\rangle$ . That the protocol indeed establishes its goal may be seen from its semantics, which

can be derived unambiguously to be

$$[SC]_{op}: (\{c\}, \cdot, \cdot) \to (\{c\}, \{1\}, \{2\})$$

$$(\alpha|0\rangle + \beta|1\rangle)_c \Longrightarrow (\alpha|000\rangle + \beta|111\rangle)_{c12}.$$

$$(20)$$

Once this type of shared entanglement is in place each target agent just has to execute a local controlled unitary gate with as control its qubit in the shared entanglement resource. Note that because we have context-independence of the semantics, target qubits may be entangled over different agents. However, in order for the distribution approach to be possible at all, the controlled unitary CU must have  $U = U_1 \otimes U_2$  where,  $U_1$  and  $U_2$  operate on a number of target qubits smaller or equal than the maximum available qubits for each agent.

We see that the SC protocol requires GHZ entanglement. It is probably realistic to assume that in a quantum network each agent can ask for Bell-state entanglement with the central server. It is not so realistic to assume that groups of agents can demand direct GHZ entanglement whenever they need it; rather, we expect GHZ entanglement to be produced via the entanglement swapping protocol. Note that, since we need GHZ entanglement between  $\mathbf{L}$ ,  $\mathbf{A}_1$  and  $\mathbf{A}_2$ , we need to compose  $\mathbf{L}$  with  $\mathbf{A}_0$  in the ES protocol as presented earlier in (9), that is

$$\mathbf{L} := \mathbf{A}_0 \circ \mathbf{L} : \{0, \bar{0}, \bar{1}, \bar{2}\}.Z_0^{x_{\bar{0}}}.c_1!s_{\bar{1}}.c_2!s_{\bar{2}}.M_{\bar{0}\bar{1}\bar{2}}^{GHZ} . \tag{21}$$

What we are actually doing here is composing the SC protocol with no shared resource with the ES protocol to establish the resource, and our semantics ensures that this is something we can do unambiguously. Indeed, we have

$$\begin{aligned}
&[SC \circ ES] = [SC] \circ [ES] \\
&= (\{c\}, \cdot, \cdot) \to (\{c\}, \{1\}, \{2\}) \\
&\cdot (\alpha |0\rangle + \beta |1\rangle)_c \Longrightarrow (\alpha |000\rangle + \beta |111\rangle)_{c12},
\end{aligned} (22)$$

with [SC] and [ES] given in (20) and (13) respectively.

More agents. The networks that we defined for implementing the distributed remote gate protocol can be generalised to n agents. This requires generalised procedures for GHZ-measurement, entanglement swapping and establishing a shared control qubit. First, an n-fold GHZ state is produced by generalising the 3-GHZ-pattern given in (4). That is, through the pattern with no input qubits, output qubits  $\{1, 2, \ldots, n\}$  and an event sequence interleaving E and H operations, as follows,

$$GHZ_{1...n} = \mathcal{H}(\hat{n}, n)E_{(n-1)\hat{n}} \dots E_{3\hat{4}}\mathcal{H}(\hat{3}, 3)E_{2\hat{3}}\mathcal{H}(\hat{2}, 2)E_{1\hat{2}}, \qquad (23)$$

where the hatted qubits are again working qubits. Again, a GHZ-measurement is executed by applying the inverse pattern followed by a diagonal-basis measurement. This leads to the pattern

$$M_{1...n}^{GHZ} = M_{\hat{1}}^{0} M_{\hat{2}}^{0} \dots M_{\hat{n}}^{0} E_{1\hat{2}} \mathcal{H}(2, \hat{2}) E_{2\hat{3}} \mathcal{H}(3, \hat{3}) E_{3\hat{4}} \dots E_{(n-1)\hat{n}} \mathcal{H}(n, \hat{n}) , \quad (24)$$

with input qubits the qubits  $\{1, 2, ..., n\}$  to be measured and with no outcome qubits. Using this sub-pattern we can establish GHZ-entanglement between the

leader **L** and agents  $\mathbf{A}_1$  through to  $\mathbf{A}_n$  by the generalised entanglement swapping protocol (Zukowski *et al.*, 1993; Bose *et al.*, 1998), which has the following network specification,

$$ES := \mathbf{L} : \{0, \bar{0}, \dots, \bar{n}\}. Z_0^{x_{\bar{0}}}. (\mathbf{c}_{\underline{i}}! s_{\bar{i}})_{i=1}^n. M_{\bar{0}...\bar{n}}^{GHZ}$$

$$|_{i=1}^n \mathbf{A}_i : \{i\}. Z_i^{x_{\bar{i}}}. \mathbf{c}_{\underline{i}}? x_{\bar{i}}$$

$$|| \otimes_{i=0}^n E_{i\bar{i}}.$$
(25)

This is just the generalisation of network (9) where we have merged agents **L** and  $\mathbf{A}_0$  for the purpose of establishing a shared control qubit as before. The signal  $s = s_{\overline{0}} \dots s_{\overline{n}}$  corresponds to a projection on the GHZ-state  $|s\rangle + |\overline{s}\rangle$ . This network has the following semantics,

$$[\![ES]\!]_{op}: (\cdot, \dots, \cdot) \to (\{0\}, \{1\}, \dots, \{n\}) \cdot \mathbf{0} \Longrightarrow GHZ_{0\dots n}^D$$
 (26)

After establishing GHZ-entanglement between agents a shared control qubit is obtained through the following protocol.

$$SC := \mathbf{L} : \{c, 0\}. \mathbf{c_i} ! s_0. M_0^X E_{0c} .$$

$$|_{i=1}^n \mathbf{A}_i : \{i\}. X_i^{x_0}. \mathbf{c_i} ? x_0$$

$$|| GH Z_0^{D_n}$$
(27)

The semantics of this network is given by

$$[SC]_{op}: (\{c\}, \cdot, \dots, \cdot) \to (\{c\}, \{1\}, \dots, \{n\})$$

$$(\alpha|0\rangle + \beta|1\rangle)_c \Longrightarrow (\alpha|0\rangle^{\otimes(n+1)} + \beta|1\rangle^{\otimes(n+1)})_{c1...n}.$$
(28)

Control qubit c can now indirectly control unitary operations at the sites of all agents by having agent  $\mathbf{A}_i$  execute a local pattern for a controlled unitary gate where the control is its qubit i and the targets are locally available qubits. Again, for this approach to be possible the controlled unitary CU must have  $U = U_1 \otimes \ldots \otimes U_n$  where each of the sub-unitary  $U_i$  is executable by agent  $\mathbf{A}_i$ .

### 5 Virtualisation

In the above we introduced a formal language for distributed quantum computation. Providing a number of tools for constructing higher-level programs, it should be seen as a first step in a bottom-up construction of a distributed quantum programming paradigm. However, the previous section clearly shows how cumbersome it is to describe and evaluate computations purely within the formal model. Rather, these developments are only really useful when one thinks of the language DMC in terms of a quantum virtual machine (QVM). A virtual machine is a platform-independent programming environment that abstracts away details of the underlying hardware or operating system. In our setting it is a low-level language abstraction layer which executes DMC programs independent of the actual implementation of quantum operations, which could be executed by any of several existing quantum simulators or even by a physical quantum computer. As a mediator between a set of low-level basic quantum gates and the construction of more complex quantum programs, a QVM forms a crucial

layer in a tiered quantum computation architecture (Svore *et al.*, 2006). It is especially valuable if one is concerned with developing higher-level distributed quantum computation languages through experimentation and abstraction, for which a user-friendly and flexible programming environment is invaluable.

Translating the formal model into a virtual machine has the obvious benefits of automating program execution and composition. Furthermore, low-level inspection during a pre-compilation step can automatically determine a number of issues such as well-definedness of code. While the formal language we have discussed in the above provides the backbone for our virtual machine, its actual implementation is far from trivial. Issues such as well-definedness of programs, naming of variables within local computations as well as in larger program contexts and efficiency all come into play in a more concrete sense that is absent in the more abstract formal model. In what follows below we give details on our QVM implementation. We first discuss the QVM for sequential, MC-based computations, then extending the platform further towards a distributed version thereof. We use the adjectives formal and virtual to differentiate between similar objects in the formal model and the virtual one whenever the context is unclear.

### 5.1 The quantum virtual machine

The first step in the development of an execution environment for the measurement calculus is translating formal measurement patterns into structured data for the computer to work with. We represent this data with *symbolic* or *s-expressions* (McCarthy, 1960), as popularised by the programming language Lisp. At the same time we chose to adhere insofar as possible to the notation used in the formal model. Before discussing the full virtual syntax below, we give the example of the Hadamard pattern from Eq.(1), which is expressed virtually as

$$\mathcal{H} := ((?i\ ?o)\ (?i)\ (?o)\ ((E\ ?o\ ?i)\ (M\ ?o\ 0)\ (X\ ?o\ (s\ ?i))))\ . \ (29)$$

Similarly to the formal setting we see a pattern expressed as a list of three qubit sets (V, I, O) and a command sequence. Question marks in qubit names indicate that they are variables, subject to renaming and instantiation to concrete qubit references. Important to note is that in the command sequence operations are executed from left to right, in the formal model notation this was the reverse. This change is due to both implementation reasons and computer science tradition.

Our current execution environment is split in two separate layers: an execution and composition layer. The execution layer takes a command sequence as an input and is in charge of performing the quantum operations it specifies. This command sequence is obtained after the data representation of a measurement pattern is assembled into the low-level language that the execution layer can understand – essentially by replacing qubit variables with concrete qubit references. Hence in terms of abstractness this layer lies below the one in which ordinary patterns are defined. The composition layer, on the other hand, adds an abstraction layer on top of pattern definition in terms of the pattern composition structures from MC. It compiles the composition of patterns into a single pattern data representation, by merging arbitrary compositions of multiple patterns into a single new pattern. In the rest of this section we discuss

implementation details of each of these components, focusing on respectively the execution layer and its pattern assembler, as well as the composition layer and its compiler.

#### 5.1.1 Execution layer

One arrives at the execution layer whenever a pattern has to run, either directly as a result of its definition being called or indirectly after a composition of patterns was compiled to an expression ready for execution. The execution layer consists of an interpreter which directly executes each operation in its input language. This interpreter is the core of our quantum virtual machine (QVM), so called to stress the low-level and machine-generated nature of its input language. It is essentially an automatic version of the operational semantics described earlier in Sec. 3 and specified in full in the appendix. The input language syntax for the QVM is obtained by turning MC notation into an s-expression form, specified as follows in a BNF notation:

```
({ <instruction>})
                ::=
<sequence>
                     <correction> | <measurement> | <entanglement>
<instruction>
<correction>
                     ( X <quref> [ <signal> ] ) |
                     ( Y <quref> [ <signal> ] )
                     ( M <quref> <angle> [ <signal> ] [ <signal> ])
<measurement>
                ::=
<entanglement>
                     ( E <quref> <quref>)
<signal>
                     0 | 1 | <input-name> |
                     ( s <quref>) | ( + <signal> { signal } )
```

where the symbols in boldface indicate syntax (round brackets and literals), while square and curly brackets are meta-syntax symbols which denote option (zero or one occurrence) and repetition (occur any finite number of times) respectively. The only real difference with the formal syntax is the notation ( $s \leq quref >$ ) for signals, referring explicitly to measurement outcomes rather than relying on the naming convention  $s_n$  to denote the outcome on qubit n. The execution layer interpreter takes any command sequence written in this language and executes it according to the semantics of MC. Note that in this way we already obtain a low-level quantum programming language. For example we can write and execute the Hadamard gate with the following command sequence:

$$((E \ 0 \ 1) \ (M \ 0 \ 0) \ (X \ 1 \ (s \ 0)))$$
. (31)

Transforming a pattern definition Eq. (29) to its assembled version Eq. (31) is relatively simple. Concretely, the pattern assembler replaces each qubit variable name in a pattern's command sequence by concrete qubit references. For every distinct variable name used we choose a unique integer and replace every occurrence of the variable by it. At this stage only the input qubit set is used, in particular for initialisation of the (non-input) auxiliary qubits to the default  $|+\rangle$ -state. As we shall see below both input and output set qubit sets are used when combining patterns in the composition layer.

Practically, the interpreter takes besides a command sequence also an *envi*ronment, which contains the current state of execution. After an initialisation step the first operation in the sequence is evaluated, after which the interpreter recursively calls itself with an environment modified by this operation. Just as in the formal model, the environment for the interpreter consists of a quantum as well as a classical part, respectively denoted Q and  $\Gamma$  below. They are defined as follows, where T stands for tangle.

$$\Gamma ::= (o, i) \tag{32}$$

$$Q ::= T \mid T \otimes Q \tag{33}$$

$$T ::= (\{q_1, q_2, \dots q_n\}, |\psi_{q_1, q_2, \dots q_n}\rangle), n \in \mathbb{N}$$
(34)

Differences with the formal model arise from practical considerations. The classical state  $\Gamma$  is split into the outcome map o and the input map i to simplify lookup of qubit and input names respectively. Effectively, evaluating an expression of the form (s n), where n is a number, becomes a lookup in o. In the formal model the quantum state (called  $\rho$  in the semantics) is presented as a single qubit state by creating a tensor product of all computation qubits during initialisation. Given the computational explosion of operations on larger tensors, this is highly impractical in a classical simulation environment. For this reason we introduce the concept of a tangle, a set of qubits which is known to be disentangled from the full quantum state and thus allowing a more compact representation. Hence in the QVM the full qubit state Q is composed of a number of individual tangle objects T, which are, evidently, updated during execution. Concretely, during initialisation a new tangle  $T_i:(\{q_i\},|\psi\rangle)$  is created for each qubit, where  $q_i$  a single qubit reference, and auxiliary qubits have  $|\psi\rangle = |+\rangle$ . After interaction due to an entanglement operation  $E_{q_iq_j}$  qubit references  $q_i$  and  $q_j$  are put into the same tangle  $T_{ij...}:(\{q_i,q_j,\ldots\},|\psi_{q_iq_j...}\rangle)$ , while the original tangles  $T_i$  and  $T_j$  are destroyed. Here the dots represent other qubits which have interacted with qubits i or j earlier in the computation. After initialisation each command in the command sequence is executed closely following the semantics of the formal model, albeit updated to the slight variations in auxiliary structures as exhibited above. In our proof-of-concept implementation, we have implemented the semantics directly using numerical linear algebra. This implementation was based on an Lisp-based simulation environment developed earlier in (Desmet et al., 2006).

#### 5.1.2 Composition layer

The language understood by our execution layer is essentially MC without composition operators. Writing larger quantum algorithms like this by hand quickly becomes a tedious and error-prone process. The measurement calculus introduces measurement patterns and pattern composition which simplify this task. Pattern definition enables local reasoning while composition allows the creation of more complex programs out of smaller parts. In this way more general patterns can be defined by using parameters, and algorithms can be programmed by composing them out of smaller building blocks. All these abstractions are taken care of by the composition layer which we explain now. Its compiler effectively machine-generates command sequences in the language we see above.

Patterns are command sequences where all qubit names are categorised as either input, output or working qubits. Defining patterns is a straightforward process: the command sequence syntax is the same as the execution layer language,

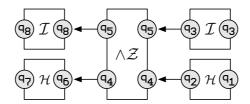


Figure 1: Graphical representation of  $\wedge \mathcal{X}$ 's pattern composition.

but instead of concrete qubit references variable names may be used. An example is shown in Eq. (29). On top of this a pattern definition can be parametrised, which is necessary for patterns such as the  $\mathcal{J}(\alpha)$ -pattern (Danos *et al.*, 2007).

By using qubit variables we facilitate the rewriting of concrete qubit names, which in turn allows patterns to be composed by concatenating their command sequences while sharing the correct qubit names. Composing two patterns into a larger one is defined in Eqs. (2) and (3). These definitions rely on the programmer to prepare the composition by renaming qubit names and, in case of sequential composition, tensoring identity patterns where needed to make sure the condition  $V_1 \cap V_2 = O_1 = I_2$  is met. In a setting where we wish the programmer to be able to compose arbitrary patterns in intricate configurations it is crucial to automatise this renaming process. Our pattern compiler does precisely that. It allows the programmer to declare his intentions in non-trivial composition cases, whereas a standard rule is applied in absence thereof. We explain the general composition case first, as we define the standard rule in terms thereof later on.

As an example let us consider the controlled-not pattern  $\wedge \mathcal{X}$ , which is defined in terms of the  $\mathcal{H}$  and  $\wedge \mathcal{Z} = (\{1,2\},\{1,2\},\{1,2\},E_{12})$  patterns, as follows.

$$\wedge \mathcal{X} := (\mathcal{I}(1) \otimes \mathcal{H}(3,4)) \circ \wedge \mathcal{Z}(1,3) \circ (\mathcal{I}(1) \otimes \mathcal{H}(2,3)) \tag{35}$$

where  $\mathcal{I}$  is the identity pattern, used as filler to match the number of in and output qubits. This composition pattern simply follows from the analogous matrix identity. The same intent can also be expressed without using concrete qubit names as follows

$$\wedge \mathcal{X} := \{ (\mathcal{I}(q_8) \otimes \mathcal{H}(q_6, q_7)) \circ \wedge \mathcal{Z}(q_5, q_4) \circ (\mathcal{I}(q_3) \otimes \mathcal{H}(q_1, q_2)) \\
\text{with } q_2 = q_4, q_3 = q_5, q_4 = q_6, q_5 = q_1) \},$$
(36)

which is just the textual version of Figure 1 (in fact, as we shall see below, our framework also provides a graphical tool for specifying constraints, much as in the figure). By simply matching variable names, we can derive Eq. (35) automatically from Eq. (36). In fact we can simplify Eq. (36) further by getting rid of identity patterns and only expressing non-trivial matching qubit names, which is specified in our (left-to-right) syntax as follows,

$$\{(\mathcal{H}(q_1, q_2)), \land \mathcal{Z}(q_5, q_4), \mathcal{H}(q_6, q_7)\}, \{(q_2, q_4), (q_4, q_6)\}\}.$$
 (37)

Note that the syntax uses tuples for constraints, and we have reverted to our left-to-right evaluation order as everywhere in the virtual model.

In our approach the programmer has to list patterns in order and give pairs of qubits to specify how patterns are to be composed. Note that if no pairs are specified this corresponds qubits not being linked up, or in other words to a parallel composition rather than a sequential one. One can then compile pattern composition by instantiating the involved measurement patterns with fresh variable names and subsequently matching the correct variable pairs. Now that we have this generalised notation for pattern composition in place, we can derive an automatic renaming procedure corresponding to the standard rules for composing patterns (e.g. linking up the first pattern's output qubits one by one with the second pattern's input qubits for sequential composition). This allows us to automate the tedious part of the qubit renaming process, while at the same time reducing the chance for human error on larger pattern compositions. However, at the same time more arbitrary composition structures may still be defined through the qubit pair syntax provided to the programmer. Nevertheless, one should view the latter as a more advanced use of the QVM, while in most cases qubit bindings are generated as per the standard case.

Once a set of qubit pairs has been specified for a certain pattern composition – either automatically or by the user himself – the automated renaming process works by generating a set of bindings that map the qubit variable names in every pattern in the composition to new ones, such that the chosen names are equal if they appear in the same pair. This is essentially the same process as one is assumed to execute manually when composing patterns in the formal model. To ensure that the construction of new bindings occurs in a well-defined and finite manner, elementary compositions in a composite structure are processed in topological order. To be precise, an arbitrary pattern composition expression such as the one in Eq. (37), is viewed as a graph, with patterns as nodes and qubit variable pairs as edges. Since we allow only pattern combinations that form directed acyclic graphs, there is a unique topological ordering on the list of nodes, and this is the order in which we evaluate composition bindings. The full set of bindings B is constructed iteratively through the following rules on the pairs of every elementary pattern composition in topological order.

$$\frac{B(q) = B(q') = \varnothing, q_c \text{ fresh}}{(q, q') \to B[q_c/q][q_c/q']}$$
(38)

$$\frac{B(q) = q_c}{(q, q') \to B[q_c/q']} \tag{39}$$

$$\frac{B(q') = q_c}{(q, q') \to B[q_c/q]} \tag{40}$$

In other words, when both names in the pair do not appear in the binding list, rule (38) will trigger. A fresh qubit variable name  $q_c$  is chosen and added as binding for both variable names in the pair. Rules (39) & (40) ensure that if a binding already exists for one of the variable names in the pair, the other will use the same binding. The topological sort will ensure that at all times only one of the three rules will execute.

Once the binding set B is constructed qubit variables in each of the composing patterns are substituted with their bound value,  $P'_i = \{B(q) | \forall q \in P_i\}$ . After this renaming process all patterns in the composition can finally be merged pairwise by joining qubit sets in the right way and appending command sequences.

Since our composition mechanism encompasses both definitions in the formal model our rules for joining qubits sets are also slightly more complicated (as they are more general). Concretely, we have the following definition, where patterns are assumed to have passed the renaming process already.

**Definition 5.1** The composition of patterns  $\mathcal{P}_1 = (V_1, I_1, O_1, \mathcal{A}_1)$  and  $\mathcal{P}_2 = (V_2, I_2, O_2, \mathcal{A}_2)$  is defined as the pattern  $\mathcal{P} = (V_1 \cup V_2, I, O, \mathcal{A}_1 \mathcal{A}_2)$  where

$$I = I_1 \cup (I_2 \backslash O_1) \tag{41}$$

$$O = (O_1 \setminus I_2) \cup O_2 . \tag{42}$$

Indeed, qubit variables from the old input and output sets that were matched become auxiliary qubits and hence are no longer represented in the final input and output set.

We stress again that except in the most complex cases, the programmer does not have to explicitly denote qubit bindings when creating pattern compositions. Indeed, in most situations these can be created automatically or one can use implicit notation for qubit bindings as in the formal model. For example, we automatically derive regular sequential composition as specified in Eq. (2) when both patterns have distinct and unique variable names while the number of output and input qubits are the same. The automatic renaming process ensures that the condition  $I_2 = O_1$  holds by automatically creating a composition expression of  $P_1$  and  $P_2$  while identifying each qubit in  $O_1$  with the corresponding one in  $I_2$  (i.e. in sequential order). Likewise, parallel composition as specified in Eq. (3) can be trivially derived, simply by not specifying any qubit pairs. Without being renamed, the qubit names of each pattern in the parallel composition remain unique. Because of this, Eqs. (41)-(42) essentially merge the corresponding qubit spaces in the right way.

With these shortcuts for the original composition operations in hand, we may now return to our example pattern  $\wedge \mathcal{X}$ . Indeed, while one way of specifying this composition is given in Eq. (37), by relying on the shortcut for sequential composition we can also express this composition without qubit names as:

$$\wedge \mathcal{X} := (\mathcal{I} \otimes \mathcal{H}) \circ \wedge \mathcal{Z} \circ (\mathcal{I} \otimes \mathcal{H}) . \tag{43}$$

The compiler will instantiate each pattern with fresh variable names, which, since they are unique, will precisely lead to the desired compositions as specified in the expressions above.

#### 5.2 The distributed quantum virtual machine

Now that our virtual machine for sequential MC computations is in place, we move on to its extension into a distributed version for the DMC language elaborated earlier in this article. Incorporating distribution into our framework comes down to extending each of the abstraction layers. Our philosophy has been to view network definitions as essentially specialised forms of pattern composition. That is, they are dealt with in a generalisation of the composition layer from Sec. 5.1.2, and compiled towards a set of distributed pattern definitions, i.e. patterns which also allow communication commands. This set of distributed patterns, together with information on how the network is set up in terms of

shared resources and agent channels agents, is then assembled to an executable form and sent to a distributed extension of the execution layer from Sec. 5.1. We discuss both the execution layer and the composition or *network layer* of the distributed quantum virtual machine (DQVM in short) in Secs. 5.2.1 and 5.2.2 below, first giving a quick overview here.

In the formal model agents are defined through a pattern with distribution extensions, while the shared resources may be viewed as a regular non-distributed pattern. Before discussing the full virtual syntax below, we give the example of the leader agent  $\mathbf{L}$  in the entanglement swapping protocol in Eq. (9). Its data representation in our QVM is as follows:

$$\begin{split} \mathbf{L} := & \left( \left( ?\mathbf{q}_{0}, ?\mathbf{q}_{1}, ?\mathbf{q}_{2} \right) \left( ?\mathbf{ch}_{0}, ?\mathbf{ch}_{1}, ?\mathbf{ch}_{2} \right) \\ & \left( M^{GHZ}_{?q_{0}, ?q_{1}, ?q_{2}} \left( \texttt{send} \ ?\mathbf{ch}_{0} \left( \texttt{s} \ ?\mathbf{q}_{0} \right) \right) \left( \texttt{send} \ ?\mathbf{ch}_{1} \left( \texttt{s} \ ?\mathbf{q}_{1} \right) \right) \left( \texttt{send} \ ?\mathbf{ch}_{2} \left( \texttt{s} \ ?\mathbf{q}_{2} \right) \right) \right) \right) . \end{split}$$

As before, an agent pattern consists of a qubit sort, specifying which qubits it owns, and a command sequence in which pattern commands as well as communication operations (here send instead of !) can be used. The main difference we see is that agents now also have a *channel sort*, in line with the resource sensitivities of distributed computations and with our naming mechanisms for variables, be it qubits or channels.

A network definition is essentially a grouping of multiple agent patterns (the single bar |) together with a shared resource pattern (the double bar |). In our setting we have to augment this with a so-called *network configuration*, a list of qubit and channel variable pairs to organise composition of the network components in the right way. That is, qubit pairs specify how the shared resource pattern needs to compose with the different agents patterns, as in Sec. 5.1.2, while channel pairs indicate how communication between agents is organised. Concretely, a network definition is compiled into a list of multiple command sequences that are to be executed concurrently. The network configuration will enable the automatic renaming process to match the names of various qubit and channel names in these command sequences. We explain this process in more detail in Sec. 5.2.2 below.

#### 5.2.1 Execution layer

The language understood by the execution layer of our distributed quantum virtual machine (DQVM in short) is the lowest-level language in our framework. It is arrived at after a list of distributed patterns, compiled from a network definition at a higher abstraction layer, is in turn assembled to a list of command sequences devoid of variable names, much as in Sec. 5.1. Each of these sequences is a data representation of an agent's command sequence in a distributed network, and the interpreter for DMC needs to run these sequences concurrently. Note that the shared resource pattern (behind the double bars ||) is dealt with during a compilation step executed prior to arriving at the execution layer, so that qubits are initialised in the right way once we are ready for actual execution. While the agent abstraction is not explicitly present in the execution layer, it does provide the required distributed functionality, namely channel communication operations and concurrent execution of multiple sequences.

We extend the input language syntax for the QVM (specified in Eq. (30)) to support DMC computations as follows:

This is again a mirror of the formal syntax, bar the use of the labels send and recv instead of! and? to indicate message sending and reception respectively – this to avoid confusion with variable names. Also communication commands now carry an extra variable specifying the channel to be used for communication, an improvement over the formal model where the use of channels was somewhat implicit. Finally, we see no syntax for the entanglement resource, which is instead viewed as part of the initialisation procedure carried out in the network layer of our virtual machine. That is, during compilation from the network layer we first run the pattern to construct the shared resources through the QVM. The quantum state resulting from this execution is then passed on to the initial environment of the execution layer.

The main functionality of the DQVM's interpreter over the QVM's is to do with communication operations and the scheduling of multiple command sequences. Practically, this means that the current state of execution retained in the interpreter's environment not only contains information on the quantum and classical states of the network (Q and  $\Gamma$  respectively) but also on issues to do with channel usage and scheduling. We call the latter the *network state* Nand add it to the interpreter's environment. The network state  $N = \{C, P\}$ holds the channel map C and the network program P. The latter is a simple collection of each agent's command sequence at each point in the computation, and is used by he interpreter to keep track of each command sequence as it is dynamically scheduled for execution. Each command sequence is executed to the point where it is empty or a send or receive operation blocks. A next command sequence will then take turn via a round robin selection and be executed in the same fashion. This round robin system ensures that sequences with blocking operations become available for execution at later time, successfully executing the send or receive operation in question if a matching receive or send has been performed by a previous sequence's turn. This process halts when all command sequences are empty. Deadlocks or similar errors will cause the interpreter to get stuck in a loop. We note that the formal semantics of a network is the same for any chosen schedule and for this reason we may choose a schedule at will without having to worry about the network's behaviour deviating from the intended one (Danos et al., 2005).

The second component of the network state is the channel map C. It simply maps channel names to values (sent along the channel) and is used to implement the communication primitives in a straightforward way:

$$\frac{C(ch) = \varnothing \qquad \Gamma(s) = v}{\Gamma, C, ((\text{send } ch \ s) \ \mathcal{E})} \qquad \to \qquad \Gamma, C[v/ch], (\mathcal{E})$$

$$\tag{46}$$

$$\frac{C(ch) = v}{\Gamma, C, ((\texttt{recv}\ ch\ name)\ \mathcal{E}) \quad \rightarrow \quad \Gamma[v/name], C[\varnothing/ch], (\mathcal{E})} \ . \tag{47}$$

Here ch is a channel name,  $\Gamma(s)$  denotes the value of the signal s with respect to the given outcome map and  $\mathcal{E}$  is the rest of the execution sequence. In other words, the send operation associates the value to be sent to the appropriate channel name in the channel map C, while the corresponding receive operation removes the value from C and stores it in the classical state. In case where a channel already contains a value during a send operation along a particular channel, the send operation blocks, and likewise for a receive operation on an empty channel:

$$\frac{C(ch) = v}{C, ((\text{send } ch \ v') \ \mathcal{E})} \rightarrow C, ((\text{send } ch \ v') \ \mathcal{E})$$

$$(48)$$

$$\frac{C(ch) = \varnothing}{C, ((\texttt{recv}\ ch\ name)\ \mathcal{E})} \rightarrow C, ((\texttt{recv}\ ch\ name)\ \mathcal{E})$$
 (49)

Note that these rules constitute a concretisation of the semantical rule for classical communication in the formal model (see Eq. 10). Here we have chosen a concrete semantics using a form of *blocking*, in the sense that execution of a particular command sequence may halt until certain requirements are met, embodied in the equations above.

Next to the network state the interpreter's environment also keeps track of the quantum and the classical state of the network. For practical reasons we reuse the original QVM's computation state as described in Sec. 5.1.1, i.e. we have one global quantum state Q and one global classical state  $\Gamma$  for the entire network. This simplifies the structure of the DQVM, allowing it to focus on the network-specific features while the QVM is used for the execution of regular MC operations. Indeed, the quantum state Q is passed through unchanged to the QVM which executes regular MC (i.e. quantum) operations. The execution schedule described above imposes a synchronised access to this global state  $(Q,\Gamma)$ , while our naming procedure ensures that names are unique so that there are no clashes between classical variable names of different agents. Note that a more sophisticated implementation would adhere to the DMC's formal semantics where each agent has its own local state. For local quantum states this is to some extent captured by our use of tangles for disentangled quantum states (cfr. Eq. (32)). For classical states this would require a separate environment for each agent such that a received value is stored in the local classical state of the receiving agent.

#### 5.2.2 Network layer

The network layer is the place where the programmer defines distributed programs, such as for example the network for entanglement sharing specified formally in Eq. (9). A network consists of a list of agent definitions, a shared resource pattern, and a network configuration which specifies how resources are distributed and connected. Concretely we need to specify a list of patterns, namely the extended patterns of all involved agents and a single regular pattern for the shared resources. For example, to define the ES network from Eq. (9) we

need the leader agent pattern given in Eq. 44 together with the agent pattern  $\mathbf{A}$  and the resource pattern  $\mathcal{R}$ :

$$\mathbf{A} := ((?q) \ (?ch) \ ((recv ?ch \ (s \ ?v)) \ (X \ ?q \ (s \ ?v)))) \tag{50}$$

$$\mathcal{R} := (V \quad I \quad O \quad ((E ?a ?b) (E ?c ?d) (E ?e ?f))), \qquad (51)$$

where V = I = 0 = (?a ?b ?c ?d ?e ?f). The full network is specified by the expression  $ES := (\mathcal{R}, \{\mathbf{L}, \mathbf{A}, \mathbf{A}, \mathbf{A}\}, NC)$ , where NC is the network configuration. As in Sec. 5.1.2, each of these patterns is instantiated with distinct qubit and channel variable names. The network configuration specified drives the composition of patterns in a network, which steers an automatic renaming process much as in Sec. 5.1.2. This network configuration contains two things. First, a list of qubit variable pairs linking output qubit variables of the shared resource pattern to input qubit variables of some agent pattern. Second, besides qubit variables there are now also channel variables to be matched between agent patterns, due to the communication extension to their command sequence. However, matching of channel names can be performed by the same renaming process defined earlier.

After the automated renaming process has finished, agent patterns are assembled and collected in a list, forming the multiple command sequences list that is used as input to the DQVM's interpreter. The pattern for shared resources does not appear in the command sequences list for the DQVM. Instead, it is assembled and executed by the QVM as an the initialisation step, and passed on as initial quantum state Q in the environment to the DQVM when it starts executing the agent programs of the network.

#### 5.3 A graphical user interface

In the above we gave an overview on the design and architecture of our virtual machine for the DMC language. While the subject of this article has to do with the inner structure of our framework, the goal of the latter is nevertheless to provide a user-friendly programming environment. For this reason we have developed a graphical user interface (GUI) on top of the virtual machine to facilitate experimentation. This GUI tool currently supports only regular (non-distributed) patterns, allowing definition of patterns from scratch as well as in terms of compositions of existing patterns. Figure 5.3 shows the GUI being used to define a pattern, the  $W_3$  entanglement pattern, as a composition of known patterns; at the bottom we find the compiled version of that same pattern. As seen previously in this section our framework allows for a general, more explicit form of composition, such that explicit set of qubit name pairs subsumes the implicit method of matching qubits by manual rewriting. While expressing these pairs in writing is still a feat for large programs, a graphical notation similar to Figure 1 is much more expressive. Our GUI uses a similar notation where black boxes denote in- and outputs and the user may connect these in arbitrary ways to concretise the desired connections. This technique allows a quantum programmer to express complex quantum algorithms more easily, certainly for composition structures with non-trivial connections between component patterns.

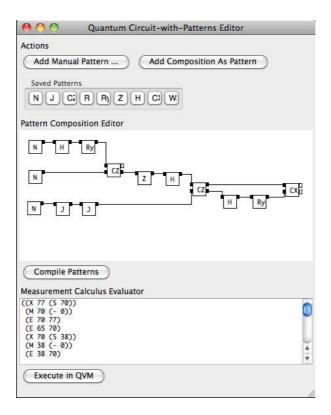


Figure 2: Graphical User Interface of the QVM's design tool, showing the composition and compilation of the 3-qubit W entanglement state preparation pattern.

## 6 Conclusion

In the above we describe an assembly language for distributed measurement-based quantum (or DMC) computations in all its aspects. While the first half of this article deals with the formal model, the second half elaborates on a virtual framework developed in close relationship with the formal model, i.e. a programming environment for the DMC language.

DMC programs satisfy several formal properties crucial to the practical usability of the language, such as compositionality and context-freeness. We showed how to put these properties to use by formally implementing a composite program to control operations in a distributed setting, and demonstrated that the semantics does not change under the various composition operations. DMC was developed with expressiveness in mind, a crucial property when establishing the first layer in a distributed quantum programming paradigm. Indeed only through experimentation combined with abstraction can one hope to move towards a higher level in the language hierarchy. Our first experiments on paper already prove that DMC is indeed capable of expressing more complicated programs and that the formal features are necessary and sufficient in determining their functionality. This is very different in flavour from earlier formal frameworks in this area (Gay and Nagarajan, 2004; Jorrand and Lalire, 2005),

which are much more concerned with issues such as verification, and focus on providing tools such as type systems to facilitate investigations of this nature.

While the first half or this article shows that it is possible in principle to write DMC programs within the formal model, the limits of its usability are acutely felt even for the relatively small example of distributed controlled operations. Indeed, as an experimentation platform our formal framework is only really useful in terms of a quantum virtual machine (QVM), a programming environment for the DMC language automating execution and composition of DMC programs. In the second half of this work we discuss a first implementation of such a virtual machine, a proof-of-concept execution environment effectively virtualising the formal DMC model. The benefits of having a QVM are many: as a first layer of a tiered quantum computation architecture (Svore et al., 2006), a platform for automated verification and model-checking, and maybe most importantly, a basis from which to develop higher-order distributed quantum computation languages through experimentation and abstraction. Our QVM is built up in terms of several abstraction layers, most importantly a platform-independent execution layer to deal with the low-level semantics of basic patterns, and a composition layer to create and compose larger programs. The composition layer comes with an associated compiler that translates any compositional structure to one single pattern definition. Any pattern definition, be it specified directly by the user or produced by the compiler, is assembled into an expression that the execution layer understands, essentially a command sequence where all variable names have been replaced by concrete references in the intended way. While the execution layer is defined independently of any actual implementation, we used a Lisp-based simulation environment in our experiments to evaluate executionlayer expressions. The QVM is extended with distributed structures at all layers into a distributed quantum virtual machine (DQVM) which allows specification and execution of arbitrary distributed networks. Finally, a graphical user interface (GUI) is added to facilitate usability of the framework. We note that due to the fact that some aspects of the formal model necessitate further concretisation in a virtual model, there are some differences in syntax and semantics between the two. For example the implicit naming conventions for variables, channels, and in the composition of programs, required a concrete design in the virtual

The virtual machine developed in this article, while covering almost all aspects of the formal model, is a first implementation and as such there are several avenues for improvement. We list these here moving from higher-abstraction layers to lower ones. First, the distributed layer requires several extensions to be fully compatible with the formal model, most importantly by giving agents a more prominent role and allowing network composition. In order to develop the GUI into a fully-featured development tool, it also needs to be lifted to a distributed setting: adding agents and networks to the graphical notation is top priority. Ultimately the goal is to make the graphical notation into a self-contained language, allowing the programmer to specify arbitrary programs without needing to create ex-nihilo patterns, which involves writing command sequence code by hand. Second, much work can be done at the level of optimisation of the execution layer, which now relies on conventional rather than optimal data structures for its implementation. We are in the process of developing tailor-built data structures, so enhancing the performance of our framework by relying on domain-specific measurement calculus optimisations. For example,

we are currently looking into the stabiliser calculus as a means to efficiently execute MC operations where possible, switching to a different representation when operations unsupported by the stabiliser calculus are performed. Optimisations at the level of command sequences are also possible, such as detecting and directly initialising cluster states rather than constructing them incrementally by executing the required entanglements. On the other hand, classically simulating large entangled states means dealing with an exponential blow-up of computational time and space. For this reason the so-called EMC form (Danos et al... 2007) (which puts entanglements first) is not always the preferred one in a simulation setting, since we need to minimise the size of entangled states during the length of the computation. Finding the right balance between direct cluster state generation and exploiting classical resources at their fullest is an exercise which is currently underway. Finally, we are heavily looking into parallel computing techniques to improve the simulation of quantum operations, at the moment carried out by straightforward linear algebraic techniques. Concretely, we are investigating the compilation of command sequences into a dataflow network (Gordon et al., 2006). In such a network quantum states are represented by a long stream of amplitudes, which has the double benefit of exposing the inherent parallelism while at the same time relaxing the need to fit entire vectors inside the same computer memory. This line of research has already lead to a first implementation of a parallelised simulation environment in (Verhaegen, 2009).

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## 8 Appendix

Here we give the full operational semantics of the distributed measurement calculus. We are concerned here with the small-step rules, indicating how atomic expressions in the language are evaluated. The big-step operational semantics of a program is found by grouping the pertaining small-step rules into one transition for the whole program. We use the standard notation of sequents and rules. A sequent  $\Gamma \vdash E \downarrow$  is read as "given environment  $\Gamma$ , the expression E evaluates to the value v". In case that the environment itself changes during evaluation of an expression E we write  $\Gamma, E \longrightarrow \Gamma'$ . We write  $\Gamma(x)$  for the value of x in  $\Gamma$  and  $\Gamma[v/x]$  for the environment  $\Gamma$  with the added binding of x to v. Sequents can be combined into rules  $\frac{S}{S'}$  which are just a different notation for  $S \Rightarrow S'$ .

There are four groups of rules, dealing with classical values (signals and angles), measurement patterns and distributed measurement patterns respectively. Each group builds on top of the previous one. The first group of rules is to do with the evaluation of signals.

$$\Gamma \vdash 0 \downarrow 0$$
 and  $\Gamma \vdash 1 \downarrow 1$  (52)

$$\frac{\Gamma \vdash s_i \perp \Gamma(i)}{\Gamma \vdash s_i \perp \Gamma(i)} \tag{53}$$

$$\frac{\overline{\Gamma \vdash s_i \downarrow \Gamma(i)}}{\Gamma[v/s_i] \vdash s_i \downarrow v} \quad \text{and} \quad \frac{\Gamma[v/s_i] \vdash s_j \downarrow \Gamma(j)}{\overline{\Gamma[v/s_i] \vdash s_j \downarrow \Gamma(j)}} \quad \text{if } i \neq j \tag{54}$$

$$\frac{\Gamma \vdash s \downarrow v \quad \Gamma \vdash t \downarrow u}{\Gamma \vdash s + t \downarrow v \oplus u} \tag{55}$$

$$\frac{\Gamma \vdash s \downarrow v \quad \Gamma \vdash t \downarrow u}{\Gamma \vdash s + t \downarrow v \oplus u} \tag{55}$$

Here  $\Gamma$  is the outcome map or classical state, and  $\oplus$  denotes addition in  $\mathbb{Z}_2$ . Angles, which can have signal dependencies percolated through via dependent measurements, are also purely classical. Values of signals are looked up in  $\Gamma$  via the ruleset for signals in order to determine the actual value of a measurement angle. This procedure is summarised in the following rules.

$$\frac{\Gamma \vdash \alpha \perp \alpha}{\Gamma} \tag{56}$$

$$\frac{\Gamma \vdash s \downarrow v \quad \Gamma \vdash t \downarrow u}{\Gamma \vdash {}^{t}[\alpha]^{s} \downarrow (-1)^{v}.\alpha + u.\pi}$$

$$(57)$$

$$\Gamma \vdash [\alpha]^s \downarrow {}^0[\alpha]^s$$
 and  $\Gamma \vdash {}^t[\alpha] \downarrow {}^t[\alpha]^0$  (58)

Having defined how signals and angles are evaluated, we can now move on to the operational semantics of the basic commands. These commands operate on an environment consisting of a quantum state  $\rho$  as well as a classical state  $\Gamma$ . We have presented these rules here for density matrices; in pure state derivations we often use state transitions for brevity. Specifically, for a pure state we have  $\rho =$  $|\psi\rangle\langle\psi|$ , which is mapped to  $L|\psi\rangle\langle\psi|L^{\dagger}$ , with L any of the entanglement, Pauli or projection operators below. A pure state transition can then be alternatively specified as mapping  $|\psi\rangle$  to  $L|\psi\rangle$ .

$$\frac{1}{\rho, \Gamma, E_{ij} \longrightarrow \wedge Z_{ij} \rho \wedge Z_{ij}, \Gamma}$$

$$\tag{59}$$

$$\frac{\Gamma \vdash s \downarrow v}{\rho, \Gamma, X_i^s \longrightarrow X_i^v \rho X_i^v, \Gamma} \tag{60}$$

$$\frac{\Gamma \vdash s \downarrow v}{\rho, \Gamma, Z_i^s \longrightarrow Z_i^v \rho Z_i^v, \Gamma} \tag{61}$$

$$\frac{\Gamma \vdash {}^{t}[\alpha]^{s} \downarrow \beta}{\rho, \Gamma, {}^{t}[M_{i}^{\alpha}]^{s} \longrightarrow_{\lambda_{0}} \langle +_{\beta}|_{i} \rho \mid +_{\beta} \rangle_{i}, \Gamma[0/i]}, \lambda_{0} = \frac{\operatorname{tr}(|+_{\beta}\rangle\langle +_{\beta}|_{i} \rho)}{\operatorname{tr}\rho}$$
(62)

$$\frac{\Gamma \vdash {}^{t}[\alpha]^{s} \downarrow \beta}{\rho, \Gamma, {}^{t}[M_{i}^{\alpha}]^{s} \longrightarrow_{\lambda_{1}} \langle -\beta|_{i} \rho \mid -\beta \rangle_{i}, \Gamma[1/i]} \quad , \lambda_{1} = \frac{\operatorname{tr}(|-\beta\rangle\langle -\beta|_{i} \rho)}{\operatorname{tr}\rho}$$
(63)

$$\frac{\rho, \Gamma, C_1 \longrightarrow_{\lambda} \rho', \Gamma'}{\rho, \Gamma, C_2 C_1 \longrightarrow_{\lambda} \rho', \Gamma', C_2} \tag{64}$$

The first three commands are purely quantum and straightforward. The measurement command is the only command that affects the quantum state as well as the classical state. First, the measurement angle has to be evaluated, which in turn requires evaluating the X- and Z-signals by the previous sets of rules. Measurement commands are also the only nondeterministic commands, as the measured qubit is projected onto either  $|+_{\alpha}\rangle$  or  $|-_{\alpha}\rangle$  with transition probabilities as stated. Usually, the convention is to renormalise the state after measurement, but we do not adhere to it here, as in this way the probability of reaching a given state can be read off its norm, and moreover the overall treatment is simpler. The last rule is for a composition of commands.

Finally, we need a set of rules to deal with distributed extensions to patterns. Essentially these transitions describe how agents  $\mathbf{A}(i,o):Q.\mathcal{E}$  and networks  $\mathcal{N}=|_{i}\mathbf{A}_{i}(\mathbf{i}_{i},\mathbf{o}_{i}):Q_{i}.\mathcal{E}_{i}\parallel\sigma$  evolve over different time steps. We adopt a shorthand notation for agents, leaving out classical inputs and output, which do not change with small-step reductions.

$$\mathbf{a}_{i} = \mathbf{A}_{i} : Q_{i}.\mathcal{E}_{i}$$

$$\mathbf{a}_{i}.E = \mathbf{A}_{i} : Q_{i}.[\mathcal{E}_{i}.E]$$

$$\mathbf{a}^{-q} = \mathbf{A} : Q \setminus \{q\}.\mathcal{E}$$

$$\mathbf{a}^{+q} = \mathbf{A} : Q \uplus \{q\}.\mathcal{E}[q/x] ,$$
(65)

where E is some event, and  $\mathcal{E}_i$  and  $\mathcal{E}'_i$  are event sequences. A configuration is given by the system state  $\sigma$  together with a set of agent programs, and their states, specifically

$$\sigma, |_{i}\Gamma_{i}, \mathbf{a}_{i} = \sigma, \Gamma_{1}, \mathbf{a}_{1} \mid \Gamma_{2}, \mathbf{a}_{2} \mid \dots \mid \Gamma_{m}, \mathbf{a}_{m}.$$

$$(66)$$

The small-step rules for configuration transitions, denoted  $\Longrightarrow$ , are specified below; we provide some explanations afterwards. When the system state is not changed in an evaluation step, we stress this by preceding a rule by  $\sigma \vdash$ .

$$\frac{\sigma, \mathcal{P}(V, I, O, \mathcal{A}), \Gamma \longrightarrow_{\lambda} \sigma', \Gamma'}{\sigma, \Gamma, \mathbf{A} : I \uplus R.[\mathcal{E}.\mathcal{P}] \Longrightarrow_{\lambda} \sigma', \Gamma', \mathbf{A} : O \uplus R.\mathcal{E}}$$
(67)

$$\frac{\Gamma_2(y) = v}{\sigma \vdash (\Gamma_1, \mathbf{a}_1, \mathbf{c}?x \mid \Gamma_2, \mathbf{a}_2, \mathbf{c}!y \Longrightarrow \Gamma_1[x \mapsto v], \mathbf{a}_1 \mid \Gamma_2, \mathbf{a}_2)}$$
(68)

$$\frac{}{\sigma \vdash (\Gamma_1, \mathbf{a}_1.\mathsf{qc}?x \mid \Gamma_2, \mathbf{a}_2.\mathsf{qc}!q \Longrightarrow \Gamma_1, \mathbf{a}_1^{+q} \mid \Gamma_2, \mathbf{a}_2^{-q})} \tag{69}$$

$$\frac{L \Longrightarrow_{\lambda} M}{L \mid N \Longrightarrow_{\lambda} M \mid N} \tag{70}$$

Implicit in these rules is a sequential composition rule, which ensures that all events in an agent's event sequence are executed one after the other. The first rule is for local operations; we have written the full pattern instead of only its command sequence here to make pattern input and output explicit. Because a pattern's big-step semantics is given by a probabilistic transition system described by  $\longrightarrow$ , we pick up a probability  $\lambda$  here. Furthermore, an agent changes its sort depending on pattern's output O. The next rule is for

classical rendezvous and is straightforward. For quantum rendezvous, we need to substitute q for x in the event sequence of the receiving agent, and furthermore adapt qubit sorts. The last rule is a metarule, which is required to express that any of the other rules may fire in the context of a larger system. L and R stand for any of the possible left-, respectively right-hand sides of any of the previous rules, while L' is an arbitrary configuration. Note that we might need to rearrange terms in the parallel composition of agents in order to be able to apply the context rule. This can always be done since the order of agents in a configuration is arbitrary. In derivations of network execution, we often do not explicitly write reductions as specified by (70), but rather specify in which order the other rules fire for the network at hand. It is precisely in this last rule that introduces nondeterminism at the network level, that is, several agent transitions may be possible within the context of a network at the same time.

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