# Fredholm equations for non-symmetric kernels, with applications to iterated integral operators. 

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#### Abstract

We give the Jordan form and the Singular Value Decomposition for an integral operator $\mathcal{N}$ with a non-symmetric kernel $N(y, z)$. This is used to give solutions of Fredholm equations for non-symmetric kernels, and to determine the behaviour of $\mathcal{N}^{n}$ and $\left(\mathcal{N} \mathcal{N}^{*}\right)^{n}$ for large $n$.


## 1 Introduction and summary

Suppose that $\Omega \subset R^{p}$ and that $\mu$ is a $\sigma$-finite measure on $\Omega$. Consider a $s_{1} \times s_{2}$ complex matrix function $N(y, z)$ on $\Omega \times \Omega$. Generally we shall assume that $N \in L_{2}(\mu \times \mu)$ and is non-trivial, that is,

$$
0<\|\mathcal{N}\|_{2}^{2}=\iint\|N(y, z)\|^{2} d \mu(y) d \mu(z)<\infty
$$

The integral operator asociated with $(N, \mu)$ is $\mathcal{N}$ defined by

$$
\begin{equation*}
\mathcal{N} q(y)=\int N(y, z) q(z) d \mu(z), p(z)^{*} \mathcal{N}=\int p(y)^{*} N(y, z) d \mu(y) \tag{1.1}
\end{equation*}
$$

where $p: \Omega \rightarrow C^{s_{1}}$ and $q: \Omega \rightarrow C^{s_{2}}$ are any functions for which these integrals exist, for example $p, q \in L_{2}(\mu)$, that is $\int|p|^{2} d \mu<\infty$, and similarly for $q$. (All integrals are over $\Omega$. * denotes the transpose of the complex conjugate.)

Section 2 reviews Fredholm theory for Hermitian kernels, that is, when $N(y, z)^{*}=N(z, y)$, so that $s_{1}=s_{2}$. For this case, $\mathcal{N}^{n}=O\left(r_{1}^{n}\right)$ for large $n$ where $r_{1}$ is the magnitude of the largest eigenvalue.

Section 3 extends this to non-Hermitian kernels for the case of diagonal Jordan form. Again $\mathcal{N}^{n}=O\left(r_{1}^{n}\right)$ for large $n$ for $r_{1}$ as before.

Section 4 deals with the case of non-diagonal Jordan form. In this case $\mathcal{N}^{n}=O\left(r_{1}^{n} n^{M-1}\right)$ for large $n$ for $r_{1}$ as before where $M$ is the largest multiplicity of those eigenvalues with modulus $r_{1}$.

Section 5 gives the Singular Value Decomposition (SVD) for a non-symmetric kernel $N(y, z)$. In this case one has results such as $\left(\mathcal{N} \mathcal{N}^{*}\right)^{n}=O\left(\theta_{1}^{2 n}\right)$ for large $n$ where $\theta_{1}$ is the largest singular value..

## 2 Hermitian kernels

## Matrix theory

First consider a Hermitian matrix $N^{*}=N \in C^{s \times s}$. Its eigenvalues $\nu_{1}, \cdots, \nu_{s}$ are the roots of $\operatorname{det}(N-\nu I)=0$. They are real. Corresponding to $\nu_{j}$ is an eigenvector $p_{j}$ satisfying $N p_{j}=\nu_{j} p_{j}$. These are orthonormal: $p_{j}^{*} p_{k}=\delta_{j k}$ where $\delta_{j j}=1$ and $\delta_{j k}=0$ for $j \neq k$. Set $P=\left(p_{1}, \cdots, p_{s}\right)$. If $N$ and its eigenvalues are real, then $P$ can be taken to be real. The spectral decomposition of $N$ in terms of its eigenvalues and eigenvectors is

$$
\begin{align*}
N & =P \Lambda P^{*}=\sum_{j=1}^{s} \nu_{j} p_{j} p_{j}^{*} \text { where } \Lambda=\operatorname{diag}\left(\nu_{1}, \cdots, \nu_{s}\right),  \tag{2.1}\\
\sum_{j=1}^{s} \nu_{j} p_{j} p_{j}^{*} & =P P^{*}=I_{s}=P^{*} P=\left(p_{j}^{*} p_{k}\right) .
\end{align*}
$$

So for $\alpha \in C$

$$
\begin{equation*}
N^{\alpha}=P \Lambda^{\alpha} P^{*}=\sum_{j=1}^{s} \nu_{j}^{\alpha} p_{j} p_{j}^{*} \tag{2.2}
\end{equation*}
$$

provided that if $\operatorname{det}(N)=0$, then $\alpha$ has non-negative real part. So for $\operatorname{det}(N) \neq 0$,

$$
N f=g \Rightarrow f=N^{-1} g=\sum_{j=1}^{s} \nu_{j}^{-1} p_{j}\left(p_{j}^{*} g\right) .
$$

Similarly for $\nu$ not an eigenvalue and $f, g \in C^{s}$,

$$
(\nu I-N) f=g \Rightarrow f=(\nu I-N)^{-1} g=\sum_{j=1}^{s}\left(\nu-\nu_{j}\right)^{-1} p_{j}\left(p_{j}^{*} g\right)
$$

For large $n$, if

$$
\begin{equation*}
r_{1}=\left|\nu_{1}\right|=\cdots=\left|\nu_{M}\right|>r_{0}=\max _{j=M+1}^{S}\left|\nu_{j}\right| \tag{2.3}
\end{equation*}
$$

then

$$
\begin{equation*}
N^{n}=r_{1}^{n} C_{n}+O\left(r_{0}^{n}\right) \text { where } C_{n}=\sum_{j=1}^{M}\left[\operatorname{sign}\left(\nu_{j}\right)\right]^{n} p_{j} p_{j}^{*}=O(1) \tag{2.4}
\end{equation*}
$$

assuming that $s$ does not depend on $n$.

## Function theory

Now consider a function $N(y, z): \Omega^{2} \rightarrow C^{s \times s}$ and a $\sigma$-finite measure $\mu$ on $\Omega \subset R^{p}$. Its integral operator with respect to $\mu$ is $\mathcal{N}$ defined by (1.1). Suppose that the kernel $N$ is Hermitian, that is,

$$
N(y, z)^{*}=N(z, y) .
$$

Then the analogues of the matrix results above are as follows. Suppose that $\|\mathcal{N}\|_{2}^{2}>0$ and

$$
\sum_{j=1}^{s} \int \mid N_{j j}(x, x) d \mu(x)<\infty
$$

The spectral decomposition of $N$ in terms of its eigenvalues and vector eigenfunctions $\left\{p_{j}(y)\right\}: \Omega \rightarrow$ $C^{s}$, is

$$
\begin{equation*}
N(y, z)=P(y) \Lambda P(z)^{*}=\sum_{j=1}^{\infty} \nu_{j} p_{j}(y) p_{j}(z)^{*} \tag{2.5}
\end{equation*}
$$

$$
\text { where } \Lambda=\operatorname{diag}\left(\nu_{1}, \nu_{2}, \cdots\right), P(y)=\left(p_{1}(y), p_{2}(y), \cdots\right) \text {. }
$$

The eigenfunctions are orthonormal with respect to $\mu$ :

$$
\int p_{j}^{*} p_{k} d \mu=\delta_{j k}
$$

In Fredholm theory the convention is to call $\left\{\lambda_{j}=\nu_{j}^{-1}\right\}$ the eigenvalues, rather than $\left\{\nu_{j}\right\}$.
$\mathcal{I}_{y z}=P(y) P(z)^{*}$ is generally divergent but can be thought of as a generalized Dirac function: $\int \mathcal{I}_{y z} f(z) d \mu(z)=f(y) . P(z)^{*} P(y)=\left(p_{j}(z)^{*} p_{k}(y)\right)$ satisfies $\int P(z)^{*} P(z) d \mu(z)=I_{\infty}$.

For $n=1,2, \cdots, N_{n}(y, z)=\mathcal{N}^{n-1} N(y, z)$ satisfies

$$
N_{n}(y, z)=P(y) \Lambda^{n} P(z)^{*}=\sum_{j=1}^{s} \nu_{j}^{n} p_{j}(y) p_{j}(z)^{*} .
$$

For large $n$, if (2.3) holds then

$$
N_{n}(y, z)=r_{1}^{n} C_{n}(y, z)+O\left(r_{0}^{n}\right) \text { where } C_{n}(y, z)=\sum_{j=1}^{M}\left[\operatorname{sign}\left(\nu_{j}\right)\right]^{n} p_{j}(y) p_{j}(z)^{*}=O(1) .
$$

Conditions for (2.5) to hold pointwise and uniformly are given in Withers (1974, 1975, 1978.) It is known as Mercer's Theorem.

## The resolvent

Given functions $f, g: \Omega \rightarrow C^{s}$, the Fredholm integral equation of the second kind,

$$
\begin{equation*}
p(y)-\lambda \mathcal{N} p(y)=f(y) \tag{2.6}
\end{equation*}
$$

can be solved for $\lambda$ not an eigenvalue using

$$
(I-\lambda \mathcal{N})^{-1}=I+\lambda \mathcal{N}_{\lambda}, \text { that is, } \mathcal{N}_{\lambda}=(I-\lambda \mathcal{N})^{-1} \mathcal{N}
$$

where

$$
\mathcal{N}_{\lambda} f(y)=\int N_{\lambda}(y, z) f(z) d \mu(z), g(z) \mathcal{N}_{\lambda}=\int g(y) N_{\lambda}(y, z) d \mu(y)
$$

and the resolvent of $\mathcal{N}$,

$$
N_{\lambda}(y, z)=(I-\lambda \mathcal{N})^{-1} N(y, z): \mathcal{C} \times \Omega^{2} \rightarrow \mathcal{C}^{s \times s}
$$

with operator $\mathcal{N}_{\lambda}$ is the unique solution of

$$
(I-\lambda \mathcal{N}) \mathcal{N}_{\lambda}=\mathcal{N}=\mathcal{N}_{\lambda}(I-\lambda \mathcal{N}),
$$

that is,

$$
\lambda \mathcal{N} N_{\lambda}(y, z)=N(y, z)-N_{\lambda}(y, z)=\lambda N_{\lambda}(y, z) \mathcal{N} .
$$

If this can be solved analytically or numerically, then one has a solution of (2.6) without the need to compute the eigenvalues and eigenfunctions of $\mathcal{N}$.

The resolvent satisfies

$$
\begin{equation*}
N_{\lambda}(y, z)=\sum_{j=1}^{\infty} p_{j}(y) p_{j}(z)^{*} /\left(\lambda_{j}-\lambda\right) . \tag{2.7}
\end{equation*}
$$

Conditions for this to hold are given by Corollary 3 of Withers (1975). The Fredholm equation of the second kind, (2.6), has solution

$$
p(y)=f(y)+\lambda \sum_{j=1}^{\infty} p_{j}(y) \int p_{j}^{*} f d \mu /\left(\lambda_{j}-\lambda\right) .
$$

The resolvent exists except for $\lambda=\lambda_{j}$, an eigenvalue. The eigenvalues of $\mathcal{N}$ are the zeros of its Fredholm determinant

$$
\begin{equation*}
D(\lambda)=\Pi_{j=1}^{\infty}\left(1-\lambda / \lambda_{j}\right)=\exp \left\{-\int_{0}^{\lambda} d \lambda \int \operatorname{trace} N_{\lambda}(x, x) d \mu(x)\right\} \tag{2.8}
\end{equation*}
$$

The Fredholm integral equation of the first kind

$$
\lambda \mathcal{N} p(x)=p(x)
$$

has a solution provided that $\lambda$ is an eigenvalue. For $\nu$ an eigenvalue, its general solution $p(x)$ is a linear combination of the eigenfunctions $\left\{p_{j}(x)\right\}$ corresponding to $\lambda_{j}=\lambda$.

Example 2.1 Suppose that $Y, Z \in R$ and $\binom{Y}{Z} \sim \mathcal{N}_{2}(0, V)$ where $V=\left(\begin{array}{ll}I & r \\ r & I\end{array}\right)$. So $V$ is the correlation matrix for $\binom{Y}{Z}$. For $j \in N_{+}, x \in R$ set $p_{j}(x)=H_{j}(x) / j!^{1 / 2}$ where $H_{j}(x)$ is the standard univariate Hermite polynomial. Then $\int p_{j} p_{k} \phi_{I}=\delta_{j k}$ and

$$
\sum_{j=0}^{\infty} r^{j} p_{j}(y) p_{j}(z)=\phi_{C}(y, z) / \phi_{I}(y) \phi_{I}(z)
$$

This is Mehler's expansion for the standard bivariate normal distribution. Pearson gave an integrated version and Kibble extended it to an expansion for $\phi_{V}(x) / \phi_{I}(x)$ for $x \in R^{k}$ and $V$ a correlation matrix. See for example (45.52) p127 and p321-2 of Kotz, Balakrishnan and Johnson (2000).

## 3 Functions of two variables with diagonal Jordan form

## Diagonal Jordan form for matrices

Consider $N \in C^{s \times s}$ with eigenvalues $\nu_{1}, \cdots, \nu_{s}$, the roots of $\operatorname{det}(N-\nu I)=0 . N$ is said to have diagonal Jordan form (DJF) if

$$
N=P \Lambda Q^{*} \text { where } P Q^{*}=I \text { and } \Lambda=\operatorname{diag}\left(\nu_{1}, \cdots, \nu_{s}\right) .
$$

So

$$
\begin{aligned}
N & =\sum_{j=1}^{s} \nu_{j} p_{j} q_{j}^{*}, q_{j}^{*} p_{k}=\delta_{j k} \text { where } P=\left(p_{1}, \cdots, p_{s}\right), Q=\left(q_{1}, \cdots, q_{s}\right), \\
\sum_{j=1}^{s} p_{j} q_{j}^{*} & =P Q^{*}=I_{s}=Q^{*} P=\left(q_{j}^{*} p_{k}\right) .
\end{aligned}
$$

If $N, \Lambda$ are real then $P, Q$ can be taken as real. Also

$$
\begin{equation*}
N P=P \Lambda, \quad N p_{j}=\nu_{j} p_{j}, \quad N^{*} Q=Q \bar{\Lambda}, \quad N^{*} q_{j}=\bar{\nu}_{j} q_{j}, \tag{3.1}
\end{equation*}
$$

and for any complex $\alpha$,

$$
N^{\alpha}=P \Lambda^{\alpha} Q^{*}=\sum_{j=1}^{s} \nu_{j}^{\alpha} p_{j} q_{j}^{*}
$$

provided that if $\operatorname{det}(N)=0$, then $\alpha$ has non-negative real part. Suppose that

$$
\begin{equation*}
\nu_{j}=r_{j} e^{i \theta_{j}} \text { and } r_{1}=\cdots=r_{R}>r_{0}=\max _{j=R+1}^{s} r_{j} . \tag{3.2}
\end{equation*}
$$

Then

$$
\begin{equation*}
N^{n}=r_{1}^{n} C_{n}+O\left(r_{0}^{n}\right) \text { where } C_{n}=\sum_{j=1}^{R} e^{i n \theta_{j}} p_{j} q_{j}^{*}=O(1) \tag{3.3}
\end{equation*}
$$

Taking $\alpha=-1$ gives the inverse of $N$ when this exists.

## Diagonal Jordan form for functions

Now consider a function $N(y, z): \Omega^{2} \rightarrow C^{s \times s}$. When $N$ has diagonal Jordan form (for example when its eigenvalues are all different), then the Fredholm equations of the first kind,

$$
\lambda \mathcal{N} p(y)=p(y), \bar{\lambda} \mathcal{N}^{*} q(z)=q(z)
$$

or equivalently for $\nu=\lambda^{-1}$,

$$
\mathcal{N} p(y)=\nu p(y), \mathcal{N}^{*} q(z)=\bar{\nu} q(z),
$$

also have only a countable number of solutions, say $\left\{\lambda_{j}=\nu_{j}^{-1}, p_{j}(y), q_{j}(z), j \geq 1\right\}$ up to arbitrary constant multipliers for $\left\{p_{j}(y), j \geq 1\right\}$, satisfying the bi-orthogonal conditions

$$
\int q_{j}^{*} p_{k} d \mu=\delta_{j k} .
$$

These are called the eigenvalues and right and left eigenfunctions of $(N, \mu)$ or $\mathcal{N}$. Also

$$
\begin{equation*}
N(y, z)=\sum_{j=1}^{\infty} \nu_{j} p_{j}(y) q_{j}(z)^{*}=P(y) \Lambda Q(z)^{*} \tag{3.4}
\end{equation*}
$$

$$
\text { where } \Lambda=\operatorname{diag}\left(\nu_{1}, \nu_{2}, \cdots\right), P(y)=\left(p_{1}(y), p_{2}(y), \cdots\right), Q(z)=\left(q_{1}(y), q_{2}(y), \cdots\right) \text {. }
$$

with convergence in $L_{2}(\mu \times \mu)$, or pointwise and uniform under stronger conditions. If $N$ is a real function and $\Lambda$ is real, then $P, Q$ can be taken as real functions

For $n \geq 1$,

$$
\begin{equation*}
N_{n}(y, z)=\mathcal{N}^{n-1} N(y, z) \tag{3.5}
\end{equation*}
$$

satisfies

$$
\begin{align*}
& N_{n}(y, z)=\sum_{j=1}^{\infty} \nu_{j}^{n} p_{j}(y) q_{j}(z)^{*}=P(y) \Lambda^{n} Q(z)^{*}  \tag{3.6}\\
& \mathcal{N}^{n} p(y)=\sum_{j=1}^{\infty} \nu_{j}^{n} p_{j}(y) \int q_{j}^{*} p d \mu
\end{align*}
$$

If (3.2) holds then

$$
\begin{align*}
& N_{n}(y, z)=r_{1}^{n} C_{n}+O\left(r_{0}^{n}\right) \text { where } C_{n}=\sum_{j=1}^{R} e^{i n \theta_{j}} p_{j}(y) q_{j}(z)^{*}=O(1)  \tag{3.7}\\
& \mathcal{N}^{n} p(y)=r_{1}^{n} c_{n}+O\left(r_{0}^{n}\right) \text { where } c_{n}=\sum_{j=1}^{R} e^{i n \theta_{j}} p_{j}(y) \int q_{j}^{*} p d \mu=O(1)
\end{align*}
$$

The resolvent satisfies the equations of Section 2 except that (2.7) is replaced by

$$
\begin{equation*}
N_{\lambda}(y, z)=\sum_{j=1}^{\infty} p_{j}(y) q_{j}(z)^{*} /\left(\lambda_{j}-\lambda\right) \tag{3.8}
\end{equation*}
$$

The Fredholm determinant is again given by (2.8). If only a finite number of eigenvalues are nonzero, the kernel $N(y, z)$ is said to be degenerate. (For example this holds if $\mu$ puts weight only at $n$ points.) If $R=1$, that is,

$$
\begin{equation*}
\left|\lambda_{1}\right|<\left|\lambda_{j}\right| \text { for } j>1 \tag{3.9}
\end{equation*}
$$

then as $n \rightarrow \infty$,

$$
\mathcal{N}^{n+1} f(y) / \mathcal{N}^{n} f(y) \rightarrow \lambda_{1}^{-1}, f(y) \mathcal{N}^{n+1} / f(y) \mathcal{N}^{n} \rightarrow \lambda_{1}^{-1}
$$

This is one way to obtain the first eigenvalue $\lambda_{1}$ arbitrarily closely. Another is to use

$$
\begin{align*}
& \lambda_{1}^{-1}=\sup \left\{\int g \mathcal{N} h d \mu: \int g h d \mu=1\right\} \text { if } \lambda_{1}>0  \tag{3.10}\\
& \lambda_{1}^{-1}=\inf \left\{\int g \mathcal{N} h d \mu: \int g h d \mu=1\right\} \text { if } \lambda_{1}<0 \tag{3.11}
\end{align*}
$$

The maximising/minimising functions are the first eigenfunctions $g=g_{1}, h=h_{1}$. These are unique up to a constant multiplier if (3.9) holds. If $\lambda_{1}$ is known, one can use

$$
\left(\lambda_{1} \mathcal{N}\right)^{n} f(y) \rightarrow p_{1}(y) \int q_{1}^{*} f d \mu, f(y)^{*}\left(\lambda_{1} \mathcal{N}\right)^{n} \rightarrow q_{1}(y) \int f^{*} p_{1} d \mu,
$$

for any function $f: \Omega \rightarrow \mathcal{C}^{s}$, to approximate $p_{1}(y), q_{1}(y)$. One may now repeat the procedure on the operator $\mathcal{N}_{1}$ corresponding to

$$
N_{1}(y, z)=N(y, z)-\nu_{1} p_{1}(y) q_{1}(z)^{*}
$$

to approximate $\lambda_{2}, p_{2}(y), q_{2}(z)$ assuming that the next eigenvalue in magnitude, $\lambda_{2}$, has multiplicity 1.

For further details see Withers $(1974,1975,1978)$ and references.

## 4 Fredholm theory for non-diagonal Jordan form

## Non-diagonal Jordan form for matrices

For $N \neq N^{*}$ a matrix in $C^{s \times s}$, its general Jordan form is

$$
\begin{align*}
N & =P J P^{-1} \text { where } J=\operatorname{diag}\left(J_{1}, \cdots, J_{r}\right), J_{j}=J_{m_{j}}\left(\lambda_{j}\right), \sum_{j=1}^{r} m_{j}=s, \\
J_{m}(\lambda) & =\lambda I_{m}+U_{m}=\left(\begin{array}{lllll}
\lambda & 1 & 0 & \cdots & 0 \\
0 & \lambda & 1 & \cdots & 0 \\
0 & 0 & 0 & \cdots & \lambda
\end{array}\right), \tag{4.1}
\end{align*}
$$

for some matrix $P$, and $U_{m}$ is the $m \times m$ matrix with 1 s on the superdiagonal and 0 s elsewhere:

$$
\left(U_{m}\right)_{j k}=\delta_{j, k-1} .
$$

(See [1] for example. If $N$ and its eigenvalues are real, then $P$ can be taken as real.) So for $n \geq 1$,

$$
N^{n}=P J^{n} P^{-1} \text { where } J^{n}=\operatorname{diag}\left(J_{1}^{n}, \cdots, J_{r}^{n}\right)
$$

By the Binomial Theorem,

$$
J_{m}(\lambda)^{n}=\sum_{a=0}^{n}\binom{n}{a} \lambda^{n-a} U_{m}^{a} \text { and }\left(U_{m}^{a}\right)_{j k}=\delta_{j, k-a}
$$

So $U_{m}^{m}=0$. For example

$$
J_{2}(\lambda)^{n}=\lambda^{n} I_{2}+n \lambda^{n-1} U_{2}=\left(\begin{array}{cc}
\lambda^{n} & n \lambda^{n-1} \\
0 & \lambda^{n}
\end{array}\right)
$$

So $N^{n}$ can be expanded in block matrix form

$$
\left(N^{n}\right)_{j k}=\sum_{c=1}^{r} P_{j c} J_{c}^{n} P^{c k}
$$

where we partition $P$ and its inverse as

$$
P=\left(P_{j k}: j, k=1, \cdots, r\right), P^{-1}=\left(P^{j k}: j, k=1, \cdots, r\right)
$$

with elements $P_{j k}$ and $P^{j k}$ matrices in $C^{m_{j} \times m_{k}}$.
Alternatively setting

$$
Q^{*}=P^{-1},\left(P_{1}, \cdots, P_{r}\right)=P,\left(Q_{1}, \cdots, Q_{r}\right)=Q
$$

with $P_{j}, Q_{j} \in C^{s \times m_{j}}$, we have

$$
\begin{align*}
N^{n} & =P J^{n} Q^{*}=\sum_{j=1}^{r} P_{j} J_{j}^{n} Q_{j}^{*},  \tag{4.2}\\
\sum_{j=1}^{r} P_{j} Q_{j}^{*} & =P Q^{*}=I_{s}=Q^{*} P=\left(Q_{j}^{*} P_{k}\right)
\end{align*}
$$

so that

$$
Q_{j}^{*} P_{j}=I_{m_{j}}, Q_{j}^{*} P_{k}=0 \in C^{m_{j} \times m_{k}} \text { if } j \neq k .
$$

$P$ can be obtained as follows. Let $p_{j k}$ be the $k$ th column of $P_{j}$ for $k=1, \cdots, m_{j}$. Then

$$
\begin{equation*}
N P=P J \Rightarrow N P_{j}=P_{j} J_{j} \Rightarrow N p_{j k}=\lambda_{j} p_{j k}+p_{j, k-1} \text { where } p_{j 0}=0 . \tag{4.3}
\end{equation*}
$$

So one first obtains $p_{j 1}$, the right eigenvector of $N$, then $p_{j 2}, \cdots, p_{j m_{j}}$. This is called the Jordan chain. $Q$ can either be obtained by inverting $P$ or using

$$
\begin{equation*}
N^{*} Q=Q J^{*} \Rightarrow N^{*} Q_{j}=Q_{j} J_{j}^{*} \Rightarrow N^{*} q_{j k}=\bar{\lambda}_{j} q_{j k}+q_{j, k+1} \text { where } q_{j, m_{j}+1}=0 . \tag{4.4}
\end{equation*}
$$

So one first computes $q_{j, m_{j}}$, the right eigenvector of $N^{*}$ then $q_{j, m_{j}-1}, \cdots, q_{j 1}$. For large $n$ and $\lambda \neq 0$,

$$
J_{n}(\lambda)^{n}=\binom{n}{m-1} \lambda^{n-m+1}\left[U_{m}^{m-1}+O(1)\right]
$$

and $U_{m}^{m-1}$ is a matrix of 0 's except for a 1 in its upper right corner. So if (3.2) holds and

$$
m_{1}=\cdots=m_{M}>\max _{j=M+1}^{R} m_{j}
$$

then

$$
\begin{aligned}
\left(N^{n}\right)_{j k} & =\binom{n}{M-1} r_{1}^{n-M+1}\left[D_{n}+O\left(n^{-1}\right)\right] \\
\text { where } D_{n} & =\sum_{c=1}^{M} P_{j c} e^{i(n-M+1) \theta_{c}} U_{M}^{M-1} P^{c k}=\sum_{c=1}^{M} P_{j c} e^{i(n-M+1) \theta_{c}} U_{M}^{M-1} P^{c k}=O(1) .
\end{aligned}
$$

See Withers and Nadarajah (2008) for more details.
Non-diagonal Jordan form for functions
Now consider $N: \Omega^{2} \rightarrow C^{s \times s}$. Suppose that $\mu$ is a $\sigma$-finite measure on $\Omega$ and that $N$ is not Hermitian, that is $N(y, z)^{*} \neq N(z, y)$. Its Jordan form is

$$
\begin{equation*}
N(y, z)=P(y) J P(z)^{-1}=\text { where } J=\operatorname{diag}\left(J_{1}, J_{2}, \cdots\right), J_{j}=J_{m_{j}}\left(\lambda_{j}\right) \tag{4.5}
\end{equation*}
$$

for $P(y): \Omega \rightarrow C^{s \times \infty}$ and $J_{m}(\lambda)$ of (4.1) above. So partitioning

$$
P(y)=\left(P_{j k}(y): j, k=1,2, \cdots\right), P(z)^{-1}=\left(P^{j k}(z): j, k=1,2, \cdots\right),
$$

with elements $P_{j k}(y)$ and $P^{j k}(z)$ matrix functions in $\Omega \rightarrow C^{m_{j} \times m_{k}}$, we can partition the $n$th iterated kernel, $N_{n}(y, z)=\mathcal{N}^{n-1} N(y, z)$ as

$$
\left[N_{n}(y, z)\right]_{j k}=\sum_{c=1}^{\infty} P_{j c}(y) J_{c}^{n} P^{c k}(z)
$$

Alternatively setting

$$
Q(z)^{*}=P(z)^{-1},\left(P_{1}(y), P_{2}(y), \cdots\right)=P(y),\left(Q_{1}(z), Q_{2}(z), \cdots\right)=Q(z)
$$

with $P_{j}(y), Q_{j}(z): \Omega \rightarrow C^{s \times m_{j}}$, we have

$$
\begin{equation*}
N_{n}(y, z)=P(y) J^{n} Q(z)^{*}=\sum_{j=1}^{\infty} P_{j}(y) J_{j}^{n} Q_{j}(z)^{*} \tag{4.6}
\end{equation*}
$$

$P(y)$ can be obtained as follows. Let $p_{j k}(y)$ be the $k$ th column of $P_{j}(y)$ for $k=1, \cdots, m_{j}$. Then

$$
\begin{equation*}
\mathcal{N} P(y)=P(y) J \Rightarrow \mathcal{N} P_{j}(y)=P_{j}(y) J_{j} \Rightarrow \mathcal{N} p_{j k}(y)=\lambda_{j} p_{j k}(y)+p_{j, k-1}(y) \tag{4.7}
\end{equation*}
$$

where $p_{j 0}(y)=0$. So one first obtains $p_{j 1}(y)$, the right eigenfunction of $N$, then $p_{j 2}(y), \cdots, p_{j m_{j}}(y)$. $Q(z)$ can either be obtained by inverting $P(z)$ or using

$$
\begin{equation*}
\mathcal{N}^{*} Q(z)=Q(z) J^{*} \Rightarrow \mathcal{N}^{*} Q_{j}(z)=Q_{j}(z) J_{j}^{*} \Rightarrow \mathcal{N}^{*} q_{j k}(z)=\bar{\lambda}_{j} q_{j k}(z)+q_{j, k+1}(z) \tag{4.8}
\end{equation*}
$$

where $q_{j, m_{j}+1}(z)=0$. So one first computes $q_{j, m_{j}}(z)$, the right eigenfunction of $N^{*}$ then $q_{j, m_{j}-1}(z), \cdots, q_{j 1}(z)$.
So if (3.2) holds and

$$
m_{1}=\cdots=m_{M}>\max _{j=M+1}^{R} m_{j}
$$

then

$$
\begin{aligned}
\quad\left(N_{n}(y, z)\right)_{j k} & =\binom{n}{M-1} r_{1}^{n-M+1}\left(D_{j k n}(y, z)+O\left(n^{-1}\right)\right) \\
\text { where } D_{j k n}(y, z) & =\sum_{c=1}^{M} e^{i(n-M+1) \theta_{c}} P_{j c}(y) U_{M}^{M-1} P^{c k}(z)=O(1)
\end{aligned}
$$

has $(a, b)$ element $\sum_{c=1}^{M} e^{i(n-M+1) \theta_{c}}\left[P_{j c}(y)\right]_{a 1}\left[P^{c k}(z)\right]_{M b}$.

## Example 4.1

## 5 The SVD for functions of two variables

## The SVD for matrices

Suppose that $N \in C^{s_{1} \times s_{2}}$. That is, $N$ is a $s_{1} \times s_{2}$ complex matrix. Denote its complex conjugate transpose by $N^{*}$. Its SVD is

$$
\begin{align*}
N & =P D Q^{*}=\sum_{j=1}^{r} \theta_{j} p_{j} q_{j}^{*} \text { where } P P^{*}=I, Q Q^{*}=I, r=\min \left(s_{1}, s_{2}\right),  \tag{5.1}\\
P & =\left(p_{1}, \cdots, p_{s_{1}}\right) \in C^{s_{1} \times s_{1}}, Q=\left(q_{1}, \cdots, q_{s_{2}}\right) \in C^{s_{2} \times s_{2}}, \theta_{1} \geq \cdots \geq \theta_{r}>0
\end{align*}
$$

and for $s_{1}=s_{2}, s_{1}>s_{2}, s_{1}<s_{2}$

$$
D=\Lambda,\binom{\Lambda}{0},(\Lambda, 0) \text { respectively where } \Lambda=\operatorname{diag}\left(\theta_{1}, \cdots, \theta_{r}\right)
$$

If $N$ is real, then so are $P$ and $Q$.
So for $s_{1}>s_{2}$,

$$
D D^{*}=\left(\begin{array}{cc}
\Lambda^{2} & 0 \\
0 & 0
\end{array}\right), D^{*} D=\Lambda^{2}
$$

and for $s_{1}<s_{2}$,

$$
D D^{*}=\Lambda^{2}, D^{*} D=\left(\begin{array}{cc}
\Lambda^{2} & 0 \\
0 & 0
\end{array}\right)
$$

Compare this with (3.1). Also for $1 \leq j \leq r$,

$$
N q_{j}=\theta_{j} p_{j}, \quad N^{*} p_{j}=\theta_{j} q_{j}
$$

for $r<j \leq s_{1}, N q_{j}=0$, and for $r<j \leq s_{2}, N^{*} p_{j}=0$. Also since

$$
N N^{*} P=P D D^{*}, N^{*} N Q=Q D^{*} D,
$$

the $p_{j}$ is a right eigenvector of $N N^{*}$ with eigenvalue $\theta_{j}^{2}$ (or 0 if $r<j \leq s_{1}$ ) and the $q_{j}$ is a right eigenvector of $N^{*} N$ with eigenvalue $\theta_{j}^{2}$ (or 0 if $r<j \leq s_{2}$ ). So (or by Section 2),

$$
\begin{align*}
\left(N N^{*}\right)^{n} & =\sum_{j=1}^{r} \theta_{j}^{2 n} p_{j} p_{j}^{*},\left(N^{*} N\right)^{n}=\sum_{j=1}^{r} \theta_{j}^{2 n} q_{j} q_{j}^{*} \text { for } n \geq 1, \\
\left(N N^{*}\right)^{n} N & =\sum_{j=1}^{r} \theta_{j}^{2 n+1} p_{j} q_{j}^{*},\left(N^{*} N\right)^{n} N^{*}=\sum_{j=1}^{r} \theta_{j}^{2 n+1} q_{j} p_{j}^{*} \text { for } n \geq 0 . \tag{5.2}
\end{align*}
$$

These do not depend on the vectors $\left\{p_{j}, q_{j}, j \geq r\right\}$.
So if $\theta_{1}=\cdots=\theta_{M}>\theta_{M+1}$, then we have the approximations for $n \geq 0$,

$$
\begin{align*}
\left(N N^{*}\right)^{n} & =\theta_{1}^{2 n} \sum_{j=1}^{M} p_{j} p_{j}^{*}+O\left(\theta_{M+1}^{2 n}\right),\left(N^{*} N\right)^{n}=\theta_{1}^{2 n} \sum_{j=1}^{M} q_{j} q_{j}^{*}+O\left(\theta_{M+1}^{2 n}\right) \text { for } n \geq 1, \\
\left(N N^{*}\right)^{n} N & =\theta_{1}^{2 n+1} \sum_{j=1}^{r} p_{j} q_{j}^{*}+O\left(\theta_{M+1}^{2 n+1}\right),\left(N^{*} N\right)^{n} N^{*}=\theta_{1}^{2 n+1} \sum_{j=1}^{r} q_{j} p_{j}^{*}+O\left(\theta_{M+1}^{2 n+1}\right), \tag{5.3}
\end{align*}
$$

But

$$
I_{s_{1}}=\left(N N^{*}\right)^{0}=\sum_{j=1}^{s_{1}} p_{j} p_{j}^{*}, I_{s_{1}}=\left(N^{*} N\right)^{0}=\sum_{j=1}^{s_{2}} q_{j} q_{j}^{*}
$$

If $s_{1}=s_{2}$ and $N$ is non-singular, its inverse is

$$
N^{-1}=Q \Lambda^{-1} P^{*}
$$

However unlike Jordan form, the SVD does not give a nice form for powers of $N$.
Now suppose $\Omega \subset R^{p}$ and that $\mu$ is a $\sigma$-finite measure on $\Omega$. Consider a function $N(y, z): \Omega^{2} \rightarrow$ $C^{s_{1} \times s_{2}}$.

The equations

$$
\mathcal{N} q(y)=\theta p(y), \mathcal{N}^{*} p(z)=\theta q(z)
$$

have a countable number of solutions, say $\left\{\theta_{j}, p_{j}(y), q_{j}(z), j \geq 1\right\}$ satisfying

$$
\int p_{j}^{*} p_{k} d \mu=\int q_{j}^{*} q_{k} d \mu=\delta_{j k} .
$$

The singular values $\left\{\theta_{j}\right\}$ may be taken as real, non-negative and non-increasing. (For convenience we have included $\theta_{j}=0$.) $\left\{p_{j}(y)\right\}$ and $\left\{q_{j}(y)\right\}$ are the right eigenfunctions of $\mathcal{N} \mathcal{N}^{*}$ and $\mathcal{N}^{*} \mathcal{N}$ respectively, with eigenvalues $\left\{\theta_{j}^{2}\right\}$. Also in $L_{2}(\mu \times \mu)$

$$
\begin{equation*}
N(y, z)=\sum_{j=1}^{\infty} \theta_{j} p_{j}(y) q_{j}(z)^{*} . \tag{5.4}
\end{equation*}
$$

If $N$ is real, then so are $\left\{p_{j}, q_{j}\right\}$. By (5.4), for $n \geq 0$,

$$
\begin{aligned}
\left(\mathcal{N} \mathcal{N}^{*}\right)^{n} N(y, z) & =\sum_{j=1}^{\infty} \theta_{j}^{2 n+1} p_{j}(y) q_{j}(z)^{*}, \\
\left(\mathcal{N}^{*} \mathcal{N}\right)^{n} N(y, z)^{*} & =\sum_{j=1}^{\infty} \theta_{j}^{2 n+1} q_{j}(z) p_{j}(y)^{*}, \\
\mathcal{N}^{*}\left(\mathcal{N} \mathcal{N}^{*}\right)^{n} N(y, z) & =\sum_{j=1}^{\infty} \theta_{j}^{2 n+2} q_{j}(y) q_{j}(z)^{*}, \\
\mathcal{N}\left(\mathcal{N}^{*} \mathcal{N}\right)^{n} N(y, z)^{*} & =\sum_{j=1}^{\infty} \theta_{j}^{2 n+2} p_{j}(z) p_{j}(y)^{*}, \\
\left(\mathcal{N} \mathcal{N}^{*}\right)^{n} p(y) & =\sum_{j=1}^{\infty} \theta_{j}^{2 n} p_{j}(y) \int p_{j}^{*} p d \mu, \\
\left(\mathcal{N}^{*} \mathcal{N}\right)^{n} q(z) & =\sum_{j=1}^{\infty} \theta_{j}^{2 n} q_{j}(z) \int q_{j}^{*} q d \mu, \\
\left(\mathcal{N} \mathcal{N}^{*}\right)^{n} \mathcal{N} q(y) & =\sum_{j=1}^{\infty} \theta_{j}^{2 n+1} p_{j}(y) \int q_{j}^{*} q d \mu, \\
\left(\mathcal{N}^{*} \mathcal{N}\right)^{n} \mathcal{N}^{*} p(y) & =\sum_{j=1}^{\infty} \theta_{j}^{2 n+1} q_{j}(y) \int p_{j}^{*} p d \mu, \\
\left.\int \operatorname{trace} \mathcal{N}^{*}\left(\mathcal{N} \mathcal{N}^{*}\right)^{n} N(y, z)\right|_{z=y} d \mu(y) & =\left.\int \operatorname{trace} \mathcal{N}\left(\mathcal{N}^{*} \mathcal{N}\right)^{n} N(y, z)\right|_{z=y} d \mu(y)=\sum_{j=1}^{\infty} \theta_{j}^{2 n+2} .
\end{aligned}
$$

So if $\theta_{1}=\cdots=\theta_{M}>\theta_{M+1}$, then we have approximations such as

$$
\left(\mathcal{N N}^{*}\right)^{n} p(y)=\theta_{1}^{2 n} \sum_{j=1}^{\infty} p_{j}(y) \int p_{j}^{*} p d \mu+O\left(\theta_{M+1}^{2 n}\right)
$$

if $p_{j}(y) \int p_{j}^{*} p d \mu=O(1)$ for $j>M$. So for iterations of $\mathcal{N}^{*} \mathcal{N}$ or $\mathcal{N} \mathcal{N}^{*}$ the most important parameter is the largest singular value.

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