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# **Accepted Manuscript**

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# asasim: Adaptive Sampling for Electromagnetic Simulations

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## Abstract

For simulations of electromagnetic resonance spectra, where the locations of spectral features are unknown, and for wide-band simulation. in general, a substantial number of wavelengths must be simulated for acceptable resolution, increasing computation time. This problem is exacerbated for spectra containing narrow-band features, as a high spectral resolution is required to even detect them. To address this challenge, a heuristic algorithm is presented for electromagnetic simulations, which adaptively refines the local resolution of spectral features druing a simulation. The method supports parallel processing and plugs in the existing simulation systems, such as rigorous coupled-wave analysis (RCWA). It can routinely reduce the computational load by two orders of magnitude.

Keywords: lorentzian; adaptive; resolution, rcwa; PACS 40; PACS 41

## PROGRAM SUMMARY

Program Title: asasim

Program Files doi: http://dv.doi.org/10.17632/d6gty7kr2x.1

Licensing provisions: CC By ' 0

Programming language: MATLAL

Nature of problem: Similations are challenging when information is needed both on a long scale (broad intermine) and on a short scale (high local resolution), such as wide-band electromagnetic spectral containing narrow-band features. When resolution is insufficient, narrow-band features may be downright absent from the spectrum, if neighboring points are simulated on with a side of a narrow peak. When local resolution is sufficient, it will necessarily be excessive in flat regions, wastefully increasing computation time.

Solution method: Tl  $\circ$  presented method enables adaptive resolution, which ensures that

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all peaks of a given minimum width are always detected and maximally resolved, while feature-less regions remain minimally resolved. An optimum point spanne is derived for lorentzian peaks (descriptive of optical resonances) and is applied to optimize computation time.

#### 1. Introduction

The ability to accurately simulate light-matter interactions in nanostructures has enabled breakthroughs in areas as diverse as optical bittensets [1], pigment-free coloration [2], and solar cells [3]. Rigorous coupled-wave an arge's (RCWA) is a popular semi-analytical method for electromagnetic simulations or grantly described by Moharam and Gaylord in 1981 [4]. However, the method is computationally demanding, and this can be a limitation for high-resolution, wide-band simulations. This can be particularly problematic for optimization, methods, such as particle swarm optimization [5] or genetic algorithms [6], where an extensive number of simulations in a many-dimensional parameter space should are performed. Such endeavors would benefit from increased simulation efficiency:

The challenges of multiscale simulations is we similarly been encountered in other fields, where more intricate schemes have been demonstrated, e.g., for elastodynamic shock propagation[7] or particle-particle interactions[8].



Figure 1: Illustration of detection problem. At insufficient resolution, the peak at 547 nm is not registered at all, as the neighboring points may randomly fall on either side of it. The much wider peak at 553 nm is certain, to be detected at this resolution, but it is still poorly resolved.

Simulation time can be reduced by simply reducing the spectral resolution. However, when simulation resolution becomes too low, narrow peaks may not even be detected, as illu trated in figure 1. Depending on whether a data point happens to fall on the narrow peak, it may or may not register as a bump, but there will be a risk that the peak is completely absent from the simulation,  $w_{1,c_1}$  can be quite misleading.

As also illustrated in figure 1, peaks of sufficient width will "urer, be detected, as at least one point will always fall on the peak. However, "eson, "ion may still be too low to properly resolve its lineshape, which is commonly a limulation goal. An obvious solution is to increase the spectral resolution of the simulation, proportionally increasing simulation time, but this would result in an unnecessarily high resolution in the flat parts of the spectrum. Figure 2A in ustrat is this central issue, i.e., spectrally flat regions are over-emphasized, whereas regions with features are under-emphasized. Ideally, an initial simulation should unly be fine-grained enough that the presence of a peak would always be detected, and peak regions should then be further resolved to the desired resolution.



Figure 2: The same peak simulated with . 79 points using a) static resolution, and b) adaptive resolution. The static method yield, a high information density off-peak, whereas the adaptive method emphasizes the peak regin.

Here, a heuristic MATLA? Figorithm is demonstrated for achieving adaptive resolution in electromatic ic simulations. This effectively reverses the information emphasis to lay on the spectral features rather than the background, as illustrated in figure 2b. Furth received, optimal parameters are derived, and the speed of the method is evaluated.

### 2. Examples and installation

At its cole, the **sasim** algorithm simply replaces the per-wavelength **for**-loop of a typical ensure nagnetic simulation system. Thus, instead of statically looping over all vavelengths in an interval with identical spacing, as is the common approach, **asasim** engluate, which regions to refine during runtime, thereby only refining regions where features are present. Therefore, in principle, any script that calls a limitation function from a for-loop can be integrated with asasim. The system performs best for "needle-in-a-haystack"-type simulations, where narrow-band patters are found in a broad interval. As this is common for electromagnetic resonance spectra, these are the focus of this paper, although the system will likely also be applicable to many other topics of simulation. This section explains he v at asim is integrated with existing simulation systems in MATLAB 2016b, running on a MacBook Pro (2.4 GHz Intel Core i5, 8 GB 1600 MHz DDR3 RAM).

#### 2.1. A simple test



Figure 3: Output of example script, whi  $\therefore$  init as an electromagnetic simulation at a random position a) within a wide interval, with b) na row line width. The resolution varies from 4.6 pm on the peak, to 4.6 nm in the flat parts of the spectrum. Achieving the same resolution with static spacing would require  $370 \times$  as many point

As actual electromagnetic simul. "ons can be complex to set up and time-consuming to run, a test-function was written that imitates an actual simulation function, but returns values from an araly. A evaluation of lorentzians. The script asasim\_Example\_1\_Imit\_\_or.m exemplifies how the asasim method may be integrated with an arbitrary simulation function. A lorentzian of narrow line width

(0.5 nm) is simulated, at a random location within the broad interval 300–1100 nm, on a slightly sloped  $\sim$  kground, given by a broad lorentzian. The time to run is  $\sim$ 1 second, and f gure 3 shows the results.

#### 2.2. Photonic crystal lab sensor at varying angles of incidence

GD-Calc is a NATLAB-package for RCWA, which can be downloaded from http://kjinn.uation.com/GD-Calc.html and installed as per the instructions given on the workhorse of the GD-Calc package is the function gdc.m, which takes three inputs: a Grating-struct with the relevant optical grating parameters and





Figure 4: Example simulation of PCS sensor at varying angles C incidence. Each simulation has a variable resolution between 2 pm and 2342 pm depending on the local region.

geometric layout, an IncField-struct with in formation about the incoming field, and an order-matrix, specifying the diffraction order, to be used in the calculations. The function gdcWrapper has been written to much a sasim with GD-Calc. It defines these necessary three inputs and calls gdc.m o get a reflectivity value R at the given input wavelength w. An additional Fulle. + put struct contains all diffraction efficiencies. asasim uses this reflectivit where R to evaluate whether points neighboring w should be further resolved. The asas m system was recently used with RCWA to substantiate experimental observations of waveguide core swelling [9] in a photonic crystal slab (PCS) sensor. Here, the simulation parameters describe a linear grating of period 368 nm, duty cycle 50% and grating height 100 nm, illuminated at an angle  $\theta$ . The model incorporates refractive index dispersion data for all three materials constituting the sensor, na. 1/a (adding layer of Efiron PC409AP (Luvantix, Korea), a nano-structured ore lay. of HI01XP (micro resist technologies, Germany) and water as superstrate. In the example given here, simulations are performed in a broad wavelength interval of 450–850 nm, with angles of incidence in steps of 0.02 rad between 0 and 0.1 r d. The full example code is given in asasim\_Example\_2\_GDC.m, and the result is shown in figure 4. It should be noted that asasim is incapable of displaying a progrest bar during runtime due to parallel processing working asynchronously, and be au e the wavelengths to simulate are being refined adaptively, and are thus not known a priori. Instead, the total number of points calculated is displayed for each r und to indicate progress. On the computer used here, the total run time  $f_{\mathcal{A}}$  all  $\upsilon$  angles is  $\sim 10$  min.



Figure 5: Simulation performed using the mnpbem-package. a, "Ising ".e default static spacing places 200 points, many of them off-peak. b) With adaptive resolution, the peaks are better resolved using only 165 points.

To demonstrate the versatility of asasim it was also integrated with the excellent class mnpbem[10], available at http://cpc.cs.q.b.ac.uk/summaries/AEKJ\_v3\_0.html, which is used for calculating scattering crc suctions of plasmonic nanoparticles near surfaces. An example script, based the demo. stration demospecstat1.m of that toolbox, is given in asasim\_Example\_3\_MNP.1. <sup>14</sup>erc, the per-wavelength for-loop is quite simply replaced by a call to asasir using mnpWrapper, which is functionally three lines long. As mentioned in section 2.4, the wrapper must take two inputs and return two outputs, and it serves as an example of how simply a wrapper function can be written. The resulting graph is shown in figure 5, and although there is perhaps limited gain from using asasin for sir lulating such wide peaks in a limited interval, the example illustrates how additional simulation systems can be interfaced using a wrapper function. The cod h is a typical runtime of  $\sim 3$  seconds.

#### 2.4. Installation

The contents of as \_\_im.zip should be decompressed to a folder on the MATLAB search path, such as MA" LAB/Toolboxes/asasim. With the MATLAB-folder as the working directory, scrip., should always contain the command addpath(genpath(pwd)) in order to add fil files in all subfolders to the search path. This line is included in all example files.

The main system for adaptive resolution is now installed, and is ready to be interfaced wi h an e listing simulation system. In general, conversion of a simulation script (listing 1) into a version utilizing adaptive resolution (listing 2), is a three-step operatio1:

```
Listing 1: Original code

1 dx = 1; % desired x-axis resolution

2 x = 400:dx:700;

3 for ix = 1:numel(x)

4 y(ix) = simulateR(x(ix),input1,input2,input?);

5 end

6 plot(x,y);
```

Listing 2: asasim-implementation, illustrating how a wrappe -function may be defined.

1	$\min HalfWidth = 0.1;$
2	dy = 0.01; % desired y-axis resolution
3	dx = optimalLorentzianSpacing(minHalf Vic' .n, y);
4	xInitial = 400: dx: 700;
5	wrapFunction = $@(x,S)$ simulateR(x,S.in. S.in2,S.in3);
6	SimInput = struct('in1', input1, 'in2', input2, 'in3', input3);
7	$[x, y] = asasim(xInitial, dy, wrap\Gammaunct, m, SimInput);$
8	$plot(\mathbf{x}, \mathbf{y});$

- 1. Define wrapper. Copy the content of the per-wavelength for-loop into a wrapper function.
- 2. Determine initial point spacing. For peaks described as lorentzians ranging from 0 to 100% intensity, the point spacing can be calculated automatically using the optimalLoren.gianSpacing function.
- 3. Replace the per-wavelength for-loop by a call to asasim.

Each of these steps will be elaborate l next in the context of concrete application examples.

The purpose of the wrop r function is to bridge the simulation system with asasim, and must take two inputs, namely the simulation coordinate x (e.g., wavelength) and a struct of additional parameters. The wrapper function passes these parameters on to the simulation function(s), using the specific syntax of that simulation system. In general, the content of a per-wavelength for-loop can often just be excised and placed does a wrapper function, and the for-loop itself is then replaced by a call to ass sim. This was first exemplified in the simple case above, where only a single hype direct al function simulateR was called from within the for-loop. Subsequent examples demonstrated how more involved cases could be handled.

For simulations where the peak-shape resembles a lorentzian with a maximum intensity of 100.% .ne optimum initial point spacing dx can be calculated by calling optimal orenvianSpacing, where minHalfWidth is the minimum lorentzian halfwidth to be defected, and yRes is the desired y-axis resolution. Using this point

spacing, a vector of points to be calculated initially is then defined as "Initial = x1:dx:x2, where x1 and x2 represent the extremes of the interval to simulate. Alternatively, if the spectrum does not contain lorentzian peaks a 10.7% maximum, arbitrary values for dy and xInitial can be defined manually, using a sustom point spacing dx. The following section explains how and why an ordinal point spacing is found.

### 3. Background



Figure 6: Working princip's 'whind the algorithm. An optimal point spacing is calculated using equation 5, and an initial roug's simulation is performed in order to detect all relevant features. Then, each peak 's to ther resolved by adaptive subdivisioning until the desired resolution is achieved. At this point, a symmetry-check is performed in order to ensure that two points with similar values are not just proceed symmetrically around a peak.

A simplified in trajeon of the working principle is presented in figure 6, with the goal of producing  $\varepsilon$  ( $\alpha$ , R)-spectrum using  $\Delta R_{\rm max}$  (goal resolution) and  $\gamma_{\rm min}$  (minimum peak h. lf-widt i) as input parameters. First, an optimum initial point spacing  $\Delta \lambda_{\rm max}$  is entoneitically calculated, such that any peak of a given minimum half-width  $\gamma_{\rm n-n}$  is certain to cause a perturbation exceeding the threshold  $\Delta R_{\rm max}$  within

the interval, flagging the region for further refinement. Then, a rough simulation is performed, with the purpose of detecting all spectral features of index st.

After each round of simulations, the difference  $\Delta R$  between  $\operatorname{aig'}$  boring points is evaluated, to identify regions that exceed  $\Delta R_{\max}$ . These will  $\operatorname{big'}$  function refined by subdivisioning. New points inherit the  $\Delta R$ -value of its parer in that round, such that subdivisioning does not continue indefinitely.

When a point no longer exceeds  $\Delta R_{\text{max}}$ , it is subdivided one last time as a symmetry-check. It is entirely possible for two points to be placed symmetrically around a peak, in which case their difference  $\Delta R$  could be zero, without the upper part of the peak having been resolved. The symmology-check safe-guards against this. If the difference still does not exceed  $\Delta R_{\text{max}}$ , the region is considered fully resolved. Thus, once a peak is detected, i.e., at lead one point satisfies the criterion  $\Delta R > \Delta R_{\text{max}}$ , the entire peak always becomes fully resolved.

For the sake of argument, consider a lorentz. n at an arbitrary location, i.e.,  $\lambda_0 = 0$ , normalized so that  $R(\lambda_0) = 100\%$ :

$$R(\lambda) = \frac{1}{1 + \frac{\lambda^2}{\gamma^2}} \tag{1}$$

$$\lambda = \sqrt{\gamma} \left(\frac{1}{R} - 1\right) \tag{2}$$

Because of the subdivisioning-schen. employed, whenever a peak is detected, it is certain to come out fully resolved. Detection in this context entails that the perturbation from a peak causes two neighboring points, spaced apart by  $\Delta\lambda$  on the first axis, to have a sufficient difference on the second axis,  $\Delta R > \Delta R_{\text{max}}$ . The narrowest peak of half-width  $\gamma$  that is certain to be detected is then a peak that is so narrow, that its perturbation of ally just causes  $\Delta R$  between any two neighboring points to exceed  $\Delta R_{\text{me}}$  even when the peak is placed right between those two neighboring points, such as points a and b in figure 7. In this case,  $\Delta R$  is zero between them, so for acception, the difference to the *next* neighbor (point c) must instead satisfy  $\Delta R > \Delta R_{\text{max}}$ . If the distance from the peak center to the first symmetrically placed neighbor (point b) is  $\lambda_1 = \Delta\lambda/2$ , then the distance to its next neighbor (point  $\gamma$ ) r ust be  $3\lambda_1$ . Thus,

$$\Delta R = R_1 - R_2 \tag{3}$$

$$= \frac{1}{1 + \frac{\lambda_1^2}{\gamma^2}} - \frac{1}{1 + \frac{(3\lambda_1)^2}{\gamma^2}}$$
(4)

This does not have a simple solution for the optimal point spacing  $\Delta \lambda = 2\lambda_1$ , but it can be isolated as



Figure 7: Illustration of points placed in the first three rounds of simulations. a-d) An initial rough-simulation is performed in order to dete + all relevant features. e-g) A symmetry-check is performed in order to ensure that two points with similar values are not just placed symmetrically around the peak (like points a and b). h-j) Peak refinement continues until a desired resolution is achieved. Note that here, the spect. In beyond  $\lambda > 531.3$  nm is therefore not refined after round two.

#### 4. Discussion

The optimum ratio between the parameters  $\Delta\lambda_{\rm max}$  and  $\Delta R_{\rm max}$  for a given  $\gamma_{\rm min}$ is given by equation 5, and shown in figure 8. Parameter-sets below the curve are sub-optimal in the sender that they cause more simulation points to be calculated than necessary, wastefully increasing computation time. Sets above the curve will only serendipitously re-olve peaks of a given half-width  $\gamma_{\rm min}$ . As an example, if a simulation is to be performed with a 5% resolution on the y-axis, and the expected minimum half-width  $\gamma_{\rm min}$  is 1 nm, the optimum spacing on the x-axis is 8.2 nm. At a resolution of 0.5%, the spacing can be 26.6 nm, while still detecting the perturbation from the narrow peak.



Figure 8: Optimum ratio of simulation par ame ers.

Decreasing  $\Delta R_{\text{max}}$  causes more points to be calculated on the *R*-axis, but this also increases  $\Delta \lambda_{\text{max}}$ , reducing the number of point, to be calculated initially on the  $\lambda$ -axis. In the simplest possible model, consider a prectrum only containing a single lorentzian peak with a half-width of  $\gamma$ . The prime of points simulated on the fully resolved peak is  $N_R = 2/\Delta R$ , and the number of points to be simulated statically across the spectrum is  $N_{\lambda} = (\lambda_{\text{max}} - \lambda_{\min})/2\lambda$ . The total number of points to be simulated adaptively is

$$N = \frac{2}{\Delta R_{\max}} - \frac{\lambda_{\max} - \lambda_{\min}}{\Delta \lambda_{\max}}$$
(6)

Using this equation for estimating the number of simulation points, the two methods were compared for speed as shown in figure 9. It is clear that the adaptive method is generally a couple of orders of a agnitude faster than the static method. The figure also illustrates the computational optimum (curve minimum) for the  $\Delta R$  and thus  $\Delta \lambda$  parameters, whilm depend on the minimum necessary peak half-width  $\gamma_{\min}$ . Time-optimal parameters could be determined by combination of equations 5 and 6 and solving  $\frac{dN}{d\Delta R_{\max}} = 0$  but this becomes rather unwieldy. As figure 9 indicates, the total number of simulate 1 points does not vary steeply for similar values of  $\Delta R$ , and so the choice of resolution is perhaps more a question of preference.

For comparison, in order to achieve the same *R*-axis resolution with static sampling as with add, ive sampling, the static first-axis point spacing must equal the smallest distance between two points separated by  $\Delta R_{\text{max}}$  on the second axis. For example, to resolve a peak of half-width  $\gamma = 0.5$  nm at  $\Delta R = 1\%$  resolution on the steepest part in  $\gamma$ . 800 nm interval, the adaptive resolution varies between 0.0046 nm and 4.6 nm, depending on the local spectral features. To achieve a static resolution of 0.0045 nm, more than 170,000 points would be required. With adaptive



resolution, the same is achieved with 468 points, making the simulation  $3.1 \times \text{faster}$ .

Figure 9: Number of points simulated as function f *R*-axis resolution.

Apart from reducing computation time, a main advantage of **asasim** is the decreased amount of *a priori* information needed appetrally flat regions are computed very fast, the precise spectral positions of features need not be known beforehand in order to simulate the narrow region of malevance. Furthermore, whereas one might have to iteratively adjust first-axis resolution in order to achieve the desired second-axis resolution, with this method, mesolution is decided for the second-axis directly.

While the amount of necessary  $a p_{1}$  or information is reduced, a rough estimate of the smallest realistic half-width is still required. The consequences of choosing a poor value for this parameter wis discussed in section 4. For very dense spectra, e.g., containing many closely spaced poaks such as interference patterns, the same amount of points may end up being simulated as using static spacing. Furthermore, when the background is strongly sloped, e.g. the resonance peak of interest resembles a bump on a larger and much broader poaks, the background also becomes highly resolved. This is partially the call of the figure 5.

Because the x-axis \_ solution varies across all spectra, direct comparison between spectra or presentat on / f data as an image will require interpolation. This is quite simply achieved using up built-in MATLAB-function interp1, as exemplified below.

#### 5. Summary

In summary, the presented method allows high-speed, high-resolution simulation of narrow-bond potential features in a broad range, with no *a priori* information about the location of prectral features. In one example, the number of points necessary to simulate was reduced from >170,000 to 468, with an accompanying reduction in

```
Listing 3: Interpolation
1
           dx = 0.01; % desired x-axis resolutior
2
           interpolatedX = min([Sim.xs]):dx:max([.'n.xs]);
3
           interpolatedY = zeros(numel(Sim),nume)
              interpolatedX));
4
           for iSim = 1:numel(Sim)
             interpolatedY(iSim,:) = interp*(Sim'iSim).xs,
5
                Sim(iSim).ys, interpolatedX 'lin ar');
6
           end
7
           imagesc(interpolatedX,1:nume?(Sin),interpolatedY);
```

computation time from  $\sim 20$  hours to  $\sim 3$  minutes. The solution is tailored to simulations of electromagnetic spectra containing loren  $\sim i$  in features, but can trivially be adapted to other simulation functions and other lineshapes. The solution easily plugs into existing simulation systems, and n. erfacing to three different simulation systems was demonstrated by examples.

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