NOTES ON LIOUVILLE TYPE THEOREMS FOR THE STATIONARY COMPRESSIBLE NAVIER-STOKES EQUATIONS

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ABSTRACT. In this paper, we investigate the three dimensional stationary compressible Navier-Stokes equations, and obtain Liouville type theorems if a smooth solution (ρ , **u**) satisfies some suitable conditions. In particular, our results improve and generalize the corresponding result of Li and Yu (2014) [8].

Keywords: Liouville type theorem; Compressible Navier-Stokes equations; Lorentz space *AMS Subject Classification (2000):* 35Q30, 76N10.

1. INTRODUCTION

This paper is concerned with the following stationary barotropic compressible Navier-Stokes equations on \mathbb{R}^3

(1.1)
$$\begin{cases} \operatorname{div}(\rho \mathbf{u}) = 0, \\ \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) - \nu \Delta \mathbf{u} - (\lambda + \nu) \nabla \operatorname{div} \mathbf{u} + \nabla P = 0, \end{cases}$$

where the vector **u** denotes the flow velocity field and the scalar function ρ represents the density of the fluid. The pressure P is given by the γ -law:

$$P(\rho) = a\rho^{\gamma},$$

where a > 0 and $\gamma > 1$ are physical constants. The shear viscosity ν and bulk viscosity λ are both constants and satisfy

$$\nu > 0 \ , \ \lambda + \frac{2}{3}\nu > 0.$$

For more physical explanations about (1.1), see [5, 10].

In the past decades, the time-dependent compressible Navier-Stokes equations have been studied by many authors. See [3, 14] for global existence and blow-up criteria of solutions, and [4, 12] for regularity criterion of weak solutions. In general, Liouville type theorems appear naturally when studying the regularity of the time-dependent Navier-Stokes equations. However, Liouville type theorems have not been solved yet, as far as we know, even in the stationary case.

Here we consider the Liouville type theorems for the stationary compressible Navier-Stokes equations (1.1). In [2], Chae showed the Liouville type theorems for the compressible Navier-Stokes equations on \mathbb{R}^n . In particular, he stated that if the smooth solution (ρ, \mathbf{u}) satisfies

(1.2)
$$\|\rho\|_{L^{\infty}(\mathbb{R}^3)} + \|\nabla \mathbf{u}\|_{L^2(\mathbb{R}^3)} + \|\mathbf{u}\|_{L^{\frac{3}{2}}(\mathbb{R}^3)} < \infty,$$

then $u \equiv 0$ and $\rho = constant$. Later, Li and Yu in [8] pointed out that the condition

$$\|
ho\|_{L^{\infty}(\mathbb{R}^3)} + \|
abla \mathbf{u}\|_{L^2(\mathbb{R}^3)} < \infty$$

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is very natural because most physical flows have bounded density and finite enstrophy, and improved the result of Chae, assuming that

(1.3)
$$\|\rho\|_{L^{\infty}(\mathbb{R}^3)} + \|\nabla \mathbf{u}\|_{L^2(\mathbb{R}^3)} + \|\mathbf{u}\|_{L^{\frac{9}{2}}(\mathbb{R}^3)} < \infty.$$

The interested readers can also refer [9], which proved a different result. We note that the condition (1.3) is weaker than (1.2) in the sense that **u** can decay more slowly at infinity. Thus, a natural and interesting problem is how to further weaken the integrability condition $\mathbf{u} \in L^{\frac{9}{2}}(\mathbb{R}^3)$. The purpose of this paper is to give a positive answer. We state the first result of this paper as follows:

Theorem 1.1. Suppose that (ρ, \mathbf{u}) is a smooth solution to (1.1) with $\rho \in L^{\infty}(\mathbb{R}^3)$, $\nabla \mathbf{u} \in L^2(\mathbb{R}^3)$ and $\mathbf{u} \in L^{p,q}(\mathbb{R}^3)$ for $3 , <math>3 \le q \le \infty$ or p = q = 3. Then $\mathbf{u} \equiv 0$ and $\rho = \text{constant on } \mathbb{R}^3$.

Now one can ask what happens for the values $p \ge \frac{9}{2}$? We apply some ideas of the papers [6, 13] dealing with the Liouville type theorems for stationary incompressible Navier-Stokes equations. Let us define

$$M_{p,q}(R) := R^{\frac{2}{3} - \frac{3}{p}} \|\mathbf{u}\|_{L^{p,q}(R \le |x| \le 2R)}.$$

Our second result is

Theorem 1.2. Suppose that (ρ, \mathbf{u}) is a smooth solution to (1.1) with $\rho \in L^{\infty}(\mathbb{R}^3)$ and $\nabla \mathbf{u} \in L^2(\mathbb{R}^3)$. For $p \geq \frac{9}{2}$, $3 \leq q \leq \infty$, assume that

$$\liminf_{R\to\infty} M_{p,q}(R) < \infty,$$

then

(1.4)
$$D(u) := \int_{\mathbb{R}^3} |\nabla \mathbf{u}|^2 \, dx \le C_0 \liminf_{R \to \infty} M^3_{p,q}(R).$$

If moreover assume

(1.5)
$$\liminf_{R \to \infty} M^3_{p,q}(R) \le \delta D(\mathbf{u})$$

for some $0 < \delta < 1/C_0$, then $\mathbf{u} \equiv 0$ and $\rho = constant$ on \mathbb{R}^3 .

Indeed, we establish the Liouville type theorems in the setting of Lorentz spaces, which can be regarded as a natural generalisation of Lesbesgue spaces. Roughy speaking, we prove that a smooth solution (ρ, \mathbf{u}) is trivial if $\rho \in L^{\infty}(\mathbb{R}^3)$, $\nabla \mathbf{u} \in L^2(\mathbb{R}^3)$ and $\mathbf{u} \in L^{p,q}$ with $p \geq 3$. Compared with the case in [6, 13], we investigate the Liouville type theorems for the stationary compressible Navier-Stokes equations (1.1). Note that in (1.1), we need the condition $\nabla \mathbf{u} \in L^2(\mathbb{R}^3)$ in order to deal with the pressure term. On the other hand, we more carefully discuss $\mathbf{u} \in L^{p,q}$ for the two cases: $3 \leq p < \frac{9}{2}$ and $p \geq \frac{9}{2}$ in contrast to the result given in [6, 13].

Remark 1.1. (i). It should be noted that if we consider the integrability condition $\mathbf{u} \in L^{p,q}(\mathbb{R}^3)$ with $3 , <math>3 \le q \le \infty$ or p = q = 3 in Theorem 1.1, then we do not need any additional assumption which is similar to (1.5).

(ii). For the value $p = \frac{9}{2}$ and $q = \infty$ in Theorem 1.2, our result can be regarded as a relaxation to the corresponding result of Li and Yu [8].

The organization of this paper is as follows: In Section 2, we collect some elementary facts. Section 3 is devoted to obtaining a prior estimate, which is the key of our proof. Finally, the proofs of Theorem 1.1 and 1.2 are given in Section 4.

2. Preliminaries

We recall the definition of Lorentz space. Given $1 \le p < \infty, 1 \le q \le \infty$, we say that a measurable function $f \in L^{p,q}(\mathbb{R}^3)$ if $||f||_{L^{p,q}(\mathbb{R}^3)} < \infty$, where

$$||f||_{L^{p,q}(\mathbb{R}^3)} := \begin{cases} \left(\int_0^\infty t^{q-1} |\{x \in \mathbb{R}^3 : |f(x)| > t\}|^{\frac{q}{p}} dt \right)^{\frac{1}{q}}, & \text{if } q < +\infty \\ \sup_{t>0} t |\{x \in \mathbb{R}^3 : |f(x)| > t\}|^{\frac{1}{p}}, & \text{if } q = +\infty. \end{cases}$$

The space satisfies the continuous embeddings

$$L^{p}(\mathbb{R}^{3}) = L^{p,p}(\mathbb{R}^{3}) \hookrightarrow L^{p,q}(\mathbb{R}^{3}) \hookrightarrow L^{p,\infty}(\mathbb{R}^{3}) , \ p \le q < \infty.$$

It should be stressed that $\|\cdot\|_{L^{p,q}}$ is a quasi-norm, namely, $\|\cdot\|_{L^{p,q}}$ do not satisfy the usual triangle inequality. Instead, we have

$$||f + g||_{L^{p,q}} \le C(p,q)(||f||_{L^{p,q}} + ||g||_{L^{p,q}}),$$

with $C(p,q) = 2^{1/p} \max(1, 2^{(1-q)/q})$. See [7] for details.

The following inequalities in Lorentz spaces are useful.

Lemma 2.1 (Hölder inequality, [11]). Let $f \in L^{p_1,q_1}(\mathbb{R}^3)$ and $g \in L^{p_2,q_2}(\mathbb{R}^3)$ with $1 \le p_1, p_2 \le \infty$, $1 \le q_1, q_2 \le \infty$. Then $fg \in L^{p,q}(\mathbb{R}^3)$ with $\frac{1}{p} = \frac{1}{p_1} + \frac{1}{p_2}, \frac{1}{q} \le \frac{1}{q_1} + \frac{1}{q_2}$ and

$$\|fg\|_{L^{p,q}(\mathbb{R}^3)} \le C \|f\|_{L^{p_1,q_1}(\mathbb{R}^3)} \|g\|_{L^{p_2,q_2}(\mathbb{R}^3)}$$

for a constant C > 0.

Lemma 2.2 (Calderón-Zygmund inequality, [1]). Let Ω be a bounded domain in \mathbb{R}^n , $1 , <math>1 < q \le \infty$ and $f \in L^{p,q}(\Omega)$. Then

$$\|\nabla^2 (-\Delta)^{-1} f\|_{L^{p,q}(\Omega)} \le C \|f\|_{L^{p,q}(\Omega)},$$

where the constant C > 0 is independent of Ω .

3. A priori estimate

In this section, we derive a local estimate of $\nabla \mathbf{u}$ by means of \mathbf{u} :

Proposition 3.1. Let (ρ, \mathbf{u}) be a smooth solution to (1.1) with $\rho \in L^{\infty}(\mathbb{R}^3)$ and $\nabla u \in L^2(\mathbb{R}^3)$. If p > 3, $3 \le q \le \infty$ or p = q = 3, then we have

(3.1)
$$\int_{|x| \le R} |\nabla \mathbf{u}|^2 \, dx + \int_{|x| \le R} |\operatorname{div} \mathbf{u}|^2 \, dx \le C_0 D_1,$$

where $D_1 := R^{1-\frac{6}{p}} \|\mathbf{u}\|_{L^{p,q}(R \le |x| \le 2R)}^2 + R^{2-\frac{9}{p}} \|\mathbf{u}\|_{L^{p,q}(R \le |x| \le 2R)}^3 + R^{\frac{1}{2}-\frac{3}{p}} \|\mathbf{u}\|_{L^{p,q}(R \le |x| \le 2R)}$ and the constant C_0 is independent of R > 0.

Proof. Let $\varphi \in C_0^{\infty}(\mathbb{R}^3)$ be a radial cut-off function satisfying

$$\varphi(|x|) = \begin{cases} 1, & \text{if } |x| < 1, \\ 0, & \text{if } |x| > 2, \end{cases}$$

and $0 \le \varphi(|x|) \le 1$ for $1 \le |x| \le 2$. For each given R > 0, we define $\varphi_R(x) := \varphi(\frac{|x|}{R})$ satisfying $\|\nabla^k \varphi_R\|_{L^{\infty}} \le CR^{-k}$

for k = 0, 1, 2 with some positive constant C independent of $x \in \mathbb{R}^3$.

Taking the inner product of $(1.1)_2$ with $\mathbf{u}\varphi_R^2$ and integrating by parts over \mathbb{R}^3 , it follows

$$\begin{split} \nu \int_{\mathbb{R}^3} \nabla \mathbf{u} : \nabla (\mathbf{u}\varphi_R^2) \, dx + (\lambda + \nu) \int_{\mathbb{R}^3} \operatorname{div} \mathbf{u} \operatorname{div}(\mathbf{u}\varphi_R^2) \, dx + \int_{\mathbb{R}^3} \operatorname{div}(\rho \mathbf{u} \otimes \mathbf{u}) \cdot \mathbf{u}\varphi_R^2 \, dx \\ + \int_{\mathbb{R}^3} \nabla P \cdot \mathbf{u}\varphi_R^2 \, dx = 0. \end{split}$$

Applying the fact $\operatorname{div}(\rho \mathbf{u}) = 0$, we deduce that

(3.2)

$$\begin{aligned}
\nu \int_{\mathbb{R}^{3}} |\nabla \mathbf{u}|^{2} \varphi_{R}^{2} dx + (\lambda + \nu) \int_{\mathbb{R}^{3}} |\operatorname{div} \mathbf{u}|^{2} \varphi_{R}^{2} dx \\
&= -2\nu \int_{\mathbb{R}^{3}} \varphi_{R} \nabla \mathbf{u} : \mathbf{u} \otimes \nabla \varphi_{R} dx - 2(\lambda + \nu) \int_{\mathbb{R}^{3}} \varphi_{R} \operatorname{div} \mathbf{u} \cdot \nabla \varphi_{R} dx \\
&- \int_{\mathbb{R}^{3}} \rho \mathbf{u} \cdot \nabla \mathbf{u} \cdot \mathbf{u} \varphi_{R}^{2} dx - \int_{\mathbb{R}^{3}} \nabla P \cdot \mathbf{u} \varphi_{R}^{2} dx \\
&= \sum_{i=1}^{4} I_{i}.
\end{aligned}$$

In the following, we estimate I_i term by term. we assume p > 3, $3 \le q \le \infty$ or p = q = 3. For I_1 , we get

$$I_{1} = \nu \int_{\mathbb{R}^{3}} \mathbf{u} \cdot \operatorname{div}(\mathbf{u} \otimes \nabla(\varphi_{R}^{2})) dx$$
$$= \nu \int_{\mathbb{R}^{3}} \mathbf{u} \cdot (\nabla \mathbf{u} \cdot \nabla(\varphi_{R}^{2}) + \mathbf{u}\Delta(\varphi_{R}^{2})) dx$$
$$= \nu \int_{\mathbb{R}^{3}} |\mathbf{u}|^{2} (\varphi_{R} \Delta \varphi_{R} + |\nabla \varphi_{R}|^{2}) dx.$$

Using Lemma 2.1 implies that

(3.3)

$$|I_{1}| \leq C(\nu) \int_{R \leq |x| \leq 2R} |\mathbf{u}|^{2} (|\varphi_{R} \Delta \varphi_{R}| + |\nabla \varphi_{R}|^{2}) dx$$

$$\leq CR^{-2} |||\mathbf{u}|^{2}||_{L^{\frac{p}{2}, \frac{q}{2}}(R \leq |x| \leq 2R)} ||1||_{L^{\frac{p}{p-2}, \frac{q}{q-2}}(R \leq |x| \leq 2R)}$$

$$\leq CR^{1-\frac{6}{p}} ||\mathbf{u}||_{L^{p,q}(R \leq |x| \leq 2R)}^{2}.$$

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For I_2 , an application of Young inequality yields

$$|I_{2}| \leq C(\lambda + \nu) \int_{R \leq |x| \leq 2R} |\varphi_{R} \operatorname{div} \mathbf{u}| |\mathbf{u}| |\nabla \varphi_{R}| dx$$

$$\leq C(\lambda + \nu) R^{-1} ||\varphi_{R} \operatorname{div} \mathbf{u}||_{L^{\frac{p}{p-1}, \frac{q}{q-1}}(R \leq |x| \leq 2R)} ||\mathbf{u}||_{L^{p,q}(R \leq |x| \leq 2R)}$$

$$(3.4) \leq C(\lambda + \nu) R^{-1} ||\varphi_{R} \operatorname{div} \mathbf{u}||_{L^{2}(R \leq |x| \leq 2R)} ||\mathbf{u}||_{L^{p,q}(R \leq |x| \leq 2R)} ||\mathbf{1}||_{L^{\frac{2p}{p-2}, \frac{2q}{q-2}}(R \leq |x| \leq 2R)}$$

$$\leq C(\lambda + \nu) R^{\frac{1}{2} - \frac{3}{p}} ||\varphi_{R} \operatorname{div} \mathbf{u}||_{L^{2}(R \leq |x| \leq 2R)} ||\mathbf{u}||_{L^{p,q}(R \leq |x| \leq 2R)}$$

$$\leq \frac{(\lambda + \nu)}{2} ||\varphi_{R} \operatorname{div} \mathbf{u}||_{L^{2}(\mathbb{R}^{3})}^{2} + CR^{1 - \frac{6}{p}} ||\mathbf{u}||_{L^{p,q}(R \leq |x| \leq 2R)}^{2}.$$

For I_3 , we obtain from $\operatorname{div}(\rho \mathbf{u}) = 0$ that

$$I_{3} = -\int_{\mathbb{R}^{3}} \rho \mathbf{u} \cdot \nabla \mathbf{u} \cdot \mathbf{u} \varphi_{R}^{2} dx = -\frac{1}{2} \int_{\mathbb{R}^{3}} \rho \mathbf{u} \cdot \nabla |\mathbf{u}|^{2} \varphi_{R}^{2} dx = \frac{1}{2} \int_{\mathbb{R}^{3}} |\mathbf{u}|^{2} \operatorname{div}(\rho \mathbf{u} \varphi_{R}^{2}) dx$$
$$= \frac{1}{2} \int_{\mathbb{R}^{3}} |\mathbf{u}|^{2} (\varphi_{R}^{2} \operatorname{div}(\rho \mathbf{u}) + 2\varphi_{R} \nabla \varphi_{R} \cdot \rho \mathbf{u}) dx = \int_{\mathbb{R}^{3}} |\mathbf{u}|^{2} \varphi_{R} \nabla \varphi_{R} \cdot \rho \mathbf{u} dx,$$

thus

(3.5)
$$|I_{3}| \leq CR^{-1} \|\rho\|_{L^{\infty}} \||\mathbf{u}|^{3}\|_{L^{\frac{p}{3},\frac{q}{3}}(R \leq |x| \leq 2R)} \|1\|_{L^{\frac{p}{p-3},\frac{q}{q-3}}(R \leq |x| \leq 2R)} \leq CR^{2-\frac{9}{p}} \|\mathbf{u}\|_{L^{p,q}(R \leq |x| \leq 2R)}^{3}.$$

To obtain the estimate for the pressure term, we recall the following lemma.

Lemma 3.1 (see [8]). Let $P \in L^{\infty}(\mathbb{R}^3)$, $p_1 \in L^{r_1}(\mathbb{R}^3)$, $p_2 \in L^{r_2}(\mathbb{R}^3)$ with $1 \leq r_1, r_2 < \infty$. Suppose that $P - p_1 - p_2$ is weakly harmonic, that is

$$\Delta(P - p_1 - p_2) = 0$$

in the sense of distribution, then there exists a constant c such that

$$P - p_1 - p_2 = c$$
 a.e. $x \in \mathbb{R}^3$

If furthermore $P(x) \ge 0$ a.e., then we also have $c \ge 0$.

With Lemma 3.1 in hand, we give the estimate of I_4 . Taking the divergence on both sides of $(1.1)_2$, we have

$$\Delta(P - p_1 - p_2) = 0,$$

where $p_1 := (-\Delta)^{-1} \partial_i \partial_j (\rho u_i u_j)$ and $p_2 := (\lambda + 2\nu) \operatorname{div} \mathbf{u}$.

Using the assumption $\nabla \mathbf{u} \in L^2(\mathbb{R}^3)$ and the Sobolev embedding $\dot{H}^1(\mathbb{R}^3) \hookrightarrow L^6(\mathbb{R}^3)$, it follows

$$p_1 \in L^3(\mathbb{R}^3) , \ p_2 \in L^2(\mathbb{R}^3)$$

Due to Lemma 3.1, there exists a constant $c \ge 0$ such that

$$a\rho^{\gamma} = P = c + p_1 + p_2.$$

Considering the function

$$P_1 := \rho^{\gamma - 1} - \left(\frac{c}{a}\right)^{\frac{\gamma - 1}{\gamma}} = \left(\frac{c + p_1 + p_2}{a}\right)^{\frac{\gamma - 1}{\gamma}} - \left(\frac{c}{a}\right)^{\frac{\gamma - 1}{\gamma}},$$

we have

$$\nabla P = \nabla(a\rho^{\gamma}) = \frac{a\gamma}{\gamma - 1}\rho\nabla(\rho^{\gamma - 1}) = \frac{a\gamma}{\gamma - 1}\rho\nabla P_{1}$$

and

(3.6)
$$|P_1\rho| \le C(a, \|\rho\|_{L^{\infty}})(|p_1| + |p_2|).$$

With respect to (3.6), more detailed arguments see [8].

Hence, making use of integration by parts, it implies

$$\begin{split} I_4 &= -\int_{\mathbb{R}^3} \frac{a\gamma}{\gamma - 1} \rho \nabla P_1 \cdot \mathbf{u} \varphi_R^2 \, dx \\ &= \frac{a\gamma}{\gamma - 1} \int_{\mathbb{R}^3} P_1 \operatorname{div}(\rho \mathbf{u} \varphi_R^2) \, dx \\ &= \frac{a\gamma}{\gamma - 1} \int_{\mathbb{R}^3} P_1 \operatorname{div}(\rho \mathbf{u}) \varphi_R^2 \, dx + \frac{a\gamma}{\gamma - 1} \int_{\mathbb{R}^3} P_1 \rho \mathbf{u} \cdot \nabla(\varphi_R^2) \, dx \\ &= \frac{2a\gamma}{\gamma - 1} \int_{\mathbb{R}^3} P_1 \rho \mathbf{u} \cdot (\varphi_R \nabla \varphi_R) \, dx, \end{split}$$

and then

$$|I_4| \le C(a, \gamma, \|\rho\|_{L^{\infty}}) R^{-1} \int_{R \le |x| \le 2R} (|p_1| + |p_2|) |\mathbf{u}| \, dx.$$

By Lemmas 2.1 and 2.2, we have

$$(3.7) \qquad C(a,\gamma,\|\rho\|_{L^{\infty}})R^{-1}\int_{R\leq|x|\leq 2R}|p_{1}||\mathbf{u}|\,dx$$
$$\leq CR^{-1}\|p_{1}\|_{L^{\frac{p}{2},\frac{q}{2}}(R\leq|x|\leq 2R)}\|\mathbf{u}\|_{L^{p,q}(R\leq|x|\leq 2R)}\|1\|_{L^{\frac{p}{p-3},\frac{q}{q-3}}(R\leq|x|\leq 2R)}$$
$$\leq CR^{-1}\|\rho\|_{L^{\infty}}\||\mathbf{u}|^{2}\|_{L^{\frac{p}{2},\frac{q}{2}}(R\leq|x|\leq 2R)}\|\mathbf{u}\|_{L^{p,q}(R\leq|x|\leq 2R)}\|1\|_{L^{\frac{p}{p-3},\frac{q}{q-3}}(R\leq|x|\leq 2R)}$$
$$\leq CR^{2-\frac{9}{p}}\|\mathbf{u}\|_{L^{p,q}(R\leq|x|\leq 2R)}^{3}$$

and

$$(3.8) \qquad C(a,\gamma,\|\rho\|_{L^{\infty}})R^{-1}\int_{R\leq|x|\leq 2R}|p_{2}||\mathbf{u}|\,dx$$
$$\leq C(a,\gamma,\|\rho\|_{L^{\infty}})R^{-1}\|p_{2}\|_{L^{\frac{p}{p-1},\frac{q}{q-1}}(R\leq|x|\leq 2R)}\|\mathbf{u}\|_{L^{p,q}(R\leq|x|\leq 2R)}$$
$$\leq CR^{-1}\|\nabla\mathbf{u}\|_{L^{2}(R\leq|x|\leq 2R)}\|\mathbf{u}\|_{L^{p,q}(R\leq|x|\leq 2R)}\|1\|_{L^{\frac{2p}{p-2},\frac{2q}{q-2}}(R\leq|x|\leq 2R)}$$
$$\leq CR^{\frac{1}{2}-\frac{3}{p}}\|\nabla\mathbf{u}\|_{L^{2}(\mathbb{R}^{3})}\|\mathbf{u}\|_{L^{p,q}(R\leq|x|\leq 2R)}$$

Therefore,

(3.9)
$$|I_4| \le CR^{2-\frac{9}{p}} \|\mathbf{u}\|_{L^{p,q}(R \le |x| \le 2R)}^3 + CR^{\frac{1}{2}-\frac{3}{p}} \|\mathbf{u}\|_{L^{p,q}(R \le |x| \le 2R)}$$

Substituting the estimates of I_1 , I_2 , I_3 and I_4 into (3.2) leads to

(3.10)
$$\nu \int_{|x| \le R} |\nabla \mathbf{u}|^2 dx + \frac{(\lambda + \nu)}{2} \int_{|x| \le R} |\operatorname{div} \mathbf{u}|^2 dx \le C \left(R^{1 - \frac{6}{p}} \|\mathbf{u}\|_{L^{p,q}(R \le |x| \le 2R)}^2 + R^{2 - \frac{9}{p}} \|\mathbf{u}\|_{L^{p,q}(R \le |x| \le 2R)}^3 + R^{\frac{1}{2} - \frac{3}{p}} \|\mathbf{u}\|_{L^{p,q}(R \le |x| \le 2R)} \right)$$

for all R > 0 with the constant C independent of R. Thanks to $\lambda + \frac{2}{3}\nu \ge 0, \nu > 0$, we see that $\lambda + \nu \ge 0$. Thus, Proposition 3.1 is proved.

4. The Proofs of Theorems

Now, we are ready to complete the proofs of Theorem 1.1 and 1.2.

Proof of Theorem 1.1. Assume that $\rho \in L^{\infty}(\mathbb{R}^3)$, $\nabla \mathbf{u} \in L^2(\mathbb{R}^3)$ and $u \in L^{p,q}(\mathbb{R}^3)$ with $3 , <math>3 \le q \le \infty$ or p = q = 3. Passing $R \to +\infty$ in (3.1), we get

$$\lim_{R \to +\infty} D_1 = 0$$

Thus, it gives

$$\lim_{R \to +\infty} \left(\int_{|x| \le R} |\nabla \mathbf{u}|^2 \, dx + \int_{|x| \le R} |\operatorname{div} \mathbf{u}|^2 \, dx \right) = 0$$

By virtue of the Lebesgue dominated convergence theorem, it leads to

(4.1)
$$\int_{\mathbb{R}^3} |\nabla \mathbf{u}|^2 \, dx + \int_{\mathbb{R}^3} |\operatorname{div} \mathbf{u}|^2 \, dx = 0$$

Hence, **u** is a constant vector on \mathbb{R}^3 , which follows from (4.1). Since $\mathbf{u} \in L^{p,q}(\mathbb{R}^3)$ with 3 or <math>p = q = 3, we conclude that $\mathbf{u} \equiv 0$. On the other hand, by means of $(1.1)_2$, we know that $\nabla(a\rho^{\gamma}) = 0$, which implies that ρ =constant on \mathbb{R}^3 . The proof of Theorem 1.1 is ended.

Proof of Theorem 1.2. Due to $M_{p,q}(R) := R^{\frac{2}{3}-\frac{3}{p}} \|\mathbf{u}\|_{L^{p,q}(R \leq |x| \leq 2R)}$, we see that

$$D_1 = R^{-\frac{1}{3}} M_{p,q}^2(R) + M_{p,q}^3(R) + R^{-\frac{1}{6}} M_{p,q}(R)$$

Thus, using the assumption and passing $R \to \infty$ in (3.1), it gets

$$\lim_{R \to +\infty} \left(\int_{|x| \le R} |\nabla \mathbf{u}|^2 \, dx + \int_{|x| \le R} |\operatorname{div} \mathbf{u}|^2 \, dx \right) \le C_0 \liminf_{R \to \infty} M^3_{p,q}(R).$$

The dominated convergence theorem implies

$$\int_{\mathbb{R}^3} |\nabla \mathbf{u}|^2 \, dx + \int_{\mathbb{R}^3} |\operatorname{div} \mathbf{u}|^2 \, dx \le C_0 \liminf_{R \to \infty} M^3_{p,q}(R).$$

Consequently, we prove (1.4). Noting the condition (1.5), we have

(4.2)
$$D(u) := \int_{\mathbb{R}^3} |\nabla \mathbf{u}|^2 \, dx \le C_0 \liminf_{R \to \infty} M^3_{p,q}(R) \le C_0 \delta D(\mathbf{u}).$$

Since $0 < C_0 \delta < 1$, we conclude that $D(\mathbf{u}) = 0$ and $\mathbf{u} \equiv 0$ on \mathbb{R}^3 . Again using $(1.1)_2$, we obtain ρ =constant on \mathbb{R}^3 . This ends the proof of Theorem 1.2.

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References

- [1] J. Bergh, J. Löfström, Interpolation Spaces: An Introduction. Berlin-New York: Springer-Verlag, 1976.
- [2] D. Chae, Remarks on the liouville type results for the compressible navier-stokes equations in ℝ^N, Nonlinearity, 25 (2012) 1345–1349.
- [3] R. Danchin, Global existence in critical spaces for compressible navier-stokes equations. *Inventiones Mathematicae*, 141 (2000) 579-614.
- [4] B. Desjardins, Regularity of Weak Solutions of the Compressible isentropic Navier-Stokes Equations, Communications in Partial Differential Equations, 22 (1997) 977-1008.
- [5] E. Feireisl, Dynamics of Viscous Compressible Fluids, (Oxford Lecture Series in Mathematics and its Applications vol 26) (Oxford: Clarendon) 2004.
- [6] H. Kozono, Y. Terasawa, Y. Wakasugi, A remark on Liouville-type theorems for the stationary Navier-Stokes equations in three space dimensions, *Journal of Functional Analysis*, 272 (2017) 804–818.
- [7] G. G. Lorentz, Some new functional spaces, Ann. Math., 1 (1950) 37–55.
- [8] D. Li, X. Yu, On some Liouville type theorems for the compressible Navier-Stokes equations, *Discrete Contin Dyn Syst.*, 34 (2014) 4719-4733.
- [9] Z. Li, P. Niu, Liouville type Theorems for the 3-D stationary Hall-MHD equations, Z. Angew. Math. Mech., (2020), https://doi.org/10.1002/zamm.201900200.
- [10] P. -L. Lions, Mathematical Topics in Fluid Mechanics, (Oxford Lecture Series in Mathematics and its Applications vol 2) (Oxford: Clarendon) 1998.
- [11] R. O'Neil, Convolution Operators and $L^{p,q}$ Spaces, Duke Math., **30** (1963) 129–142.
- [12] Y. Qin, L. Huang, S. Deng, Z. Ma, X. Su, X Yang, Interior regularity of the compressible Navier-Stokes equations with degenerate viscosity coefficient and vacuum. *Discrete and Continuous Dynamical Systems*, *Series-S*, 2 (2009) 163-192.
- [13] G. A. Seregin, W. Wang, Sufficient conditions on Liouville type theorems for the 3D steady Navier-Stokes equations, arXiv:1805.02227.
- [14] Z. Xin, Blow-up of smooth solutions to the compressible Navier-Stokes equation with compact density, Commun. Pure Appl. Math., 51 (1998) 229-40.