# Large-signal stability conditions for semi-quasi-Z-source inverters: switched and averaged models 

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#### Abstract

The recently introduced semi-quasi-Z-source inverter can be interpreted as a DC-DC converter whose inputoutput voltage gain may take any value in the interval $(-\infty, 1)$ depending on the applied duty cycle. In order to generate a sinusoidal voltage waveform at the output of this converter, a time-varying duty cycle needs to be applied. Application of a time-varying duty cycle that produces large-signal behavior requires careful consideration of stability issues. This paper provides stability results for both the large-signal averaged and the switched models of the semi-quasi-Z-source inverter operating in continuous conduction mode. We show that if the load is linear and purely resistive then the boundedness and ultimate boundedness of the state trajectories is guaranteed provided some reasonable operation conditions are ensured. These conditions amount to keeping the duty cycle away from the extreme values 0 or 1 (averaged and switched models), and limiting the maximum PWM switching period (switched model). The results obtained can be used to give theoretical justification to the inverter operation strategy recently proposed by Cao et al. in [1].


## I. Introduction

The semi-quasi-Z-source inverter introduced in [1] is a single-phase single-stage low-cost (only two active components) transformerless inverter whose input and output terminals share the same ground. This semi-quasi-Z-source inverter is especially suited for renewable-energy distributedgeneration photovoltaic applications and its name derives from the Z-source inverter [2], [3] because it also contains an LC network, the distinguishing feature of the Z-source inverter. However, the shoot-through state responsible for the boost capability of the Z-source inverter is not applicable to the semi-quasi-Z-source inverter and hence the principle of operation of the latter inverter is quite different.

The semi-quasi-Z-source inverter contains two active components (such as IGBTs or MOSFETs) which conduct in a complementary manner during normal operation (continuous conduction mode, CCM). The steady-state input-output voltage gain of this inverter can theoretically be any value in the interval $(-\infty, 1)$, provided the circuit is operated at a constant duty cycle with a sufficiently high switching frequency. To obtain a sinusoidal voltage waveform at the output of the inverter, [1] asserts that if the frequency of the desired output sine wave is small enough, then the steadystate input-output gain equation would be approximately valid at every time instant (after a possible initial transient)

[^0]and hence the required time-varying duty cycle can be deduced from this gain equation. This inverter operation strategy was tested on a 40W prototype connected to a linear purely resistive load. Operating the inverter in this manner necessarily produces large-signal behavior, meaning that the linearized averaged model ([4], [5], [6]) is not an accurate model of the evolution of the circuit variables. Also, note that the inverter operation strategy as tested on the prototype is an open-loop strategy since no feedback from the circuit variables is involved in the computation of the applied duty cycle.

Although the inverter operation strategy proposed in [1] was experimentally tested and showed acceptable results, the fact that such a strategy is open-loop and time-varying raises the following issues. First, extinction of the initial transient, i.e., convergence of the state trajectory to the steady-state (time-varying) one, should be established for every possible initial condition and every possible value of the circuit capacitors, inductors and load. Second, operation of the converter under closed-loop control, which is a necessity in order to compensate for input voltage and load variations, should be ensured not to cause unstable behavior.

In this work, we address the aforementioned issues by showing that if the load is linear and purely resistive then neither the averaged nor the switched models' state trajectories can become unbounded irrespective of the (time-varying) duty cycle applied (even if it involves feedback), provided the duty cycle does not reach the extreme values 0 or 1 and, for the switched model, the maximum pulse-width modulator (PWM) switching period is suitably limited. Based on these results, we show that convergence to the periodic steady-state behavior that corresponds to a periodic duty cycle (such as that proposed in [1]) is ensured when the load is linear and purely resistive.

For both the averaged and the switched models, our stability results are derived by exploiting the natural energy function of the circuit, even though the quadratic term of its derivative along trajectories is only negative semidefinite. In the case of the averaged model, we show how to modify the natural energy function in order to analytically construct a suitable quadratic Lyapunov function. The stability results for the averaged model provide a starting point for more complicated stability analyses, since it is known that if the switching frequency is sufficiently high, then the averaged model trajectories will be close to those of the switched model [7]. In the case of the switched model, we remark that neither of the switching modes (neither of the subsystems, employing switched systems terminology [8]) is asymptot-
ically stable nor detectable from the output voltage, and hence stability is dependent on the limitations imposed on switching. In addition, since neither mode is asymptotically stable, then a dwell-time condition is not sufficient in order to ensure asymptotic stability.

The remainder of the paper is organized as follows. In Section II. we provide a brief description of the semi-quasi-Z-source inverter, derive the switched and the averaged models, and explain the duty cycle computation method of [1]. In Section III, we derive the stability results for the averaged model and in Section IV those for the switched model. Conclusions and future research directions are outlined in Section V Most proofs are given in the Appendix.

Notation. For a matrix $M, \boldsymbol{\rho}(M)$ denotes its spectral radius and $M^{\prime}$ its transpose. The $i$-th column of the identity matrix is denoted $e_{i}$ and $\|\cdot\|$ denotes the (induced) 2-norm.

## II. SEMI-QUASI-Z-SOURCE INVERTER

Figure 1] shows the semi-quasi-Z-source inverter connected to a resistive load. The two active components are transistors $S_{1}$ and $S_{2}$. In CCM, $S_{1}$ and $S_{2}$ operate in a complementary manner so that either $S_{1}$ is on and $S_{2}$ is off (Mode I) or $S_{1}$ is off and $S_{2}$ is on (Mode II). By discontinuous conduction mode (DCM) we refer to the case when both $S_{1}$ and $S_{2}$ are on because either the antiparallel diode of $S_{2}$ becomes forward biased during Mode I, or that of $S_{1}$ during Mode II. The duty cycle of $S_{1}$ is denoted $d$ and hence the duty cycle of $S_{2}$ is $1-d$ if operating in CCM.


Fig. 1: Semi-quasi-Z-source inverter.

## A. Switched and averaged models

Defining the state vector $z$ as

$$
z:=\left[\begin{array}{llll}
i_{L_{1}} & i_{L_{2}} & v_{C_{1}} & v_{C_{2}} \tag{1}
\end{array}\right]^{\prime}
$$

with the positive convention for each variable as shown in Figure 1, the state equations for the switched model in CCM can readily be written as

$$
\begin{array}{rlrl}
\dot{z}(t) & =\left[\begin{array}{lll}
A_{\mathrm{I}} z(t)+b_{\mathrm{I}}
\end{array}\right] r(t)+\left[A_{\mathrm{II}} z(t)+b_{\mathrm{II}}\right][1-r(t)], \\
A_{i} & =P^{-1} A_{i}^{q} / 2, & b_{i}=P^{-1} b_{i}^{q} / 2, & i \in\{\mathrm{I}, \mathrm{II}\}, \\
A_{\mathrm{I}}^{q} & =\left[\begin{array}{cccc}
0 & 0 & 0 & 0 \\
0 & 0 & 1 & 1 \\
0 & -1 & 0 & 0 \\
0 & -1 & 0 & -1 / R
\end{array}\right], & A_{\mathrm{II}}^{q}=\left[\begin{array}{cccc}
0 & 0 & -1 & 0 \\
0 & 0 & 0 & 1 \\
1 & 0 & 0 & 0 \\
0 & -1 & 0 & -1 / R
\end{array}\right], \\
b_{\mathrm{I}}^{q} & =\left[\begin{array}{llll}
V_{\mathrm{in}} & 0 & 0 & 0
\end{array}\right]^{\prime}, & b_{\mathrm{II}}^{q}=\left[\begin{array}{ccc}
0 & -V_{\mathrm{in}} & 0
\end{array}\right]^{\prime}, \\
P & :=\frac{1}{2} \operatorname{diag}\left(L_{1}, L_{2}, C_{1}, C_{2}\right), \tag{6}
\end{array}
$$

with $r(t)=1$ for Mode I and $r(t)=0$ for Mode II. These equations correspond to the case of a purely resistive load of resistance $R$. Neither $A_{\mathrm{I}}=P^{-1} A_{\mathrm{I}}^{q} / 2$ nor $A_{\mathrm{II}}=$ $P^{-1} A_{\mathrm{II}}^{q} / 2$ are Hurwitz since they contain eigenvalues with zero real part. Using the switched model $\sqrt{26}-(6)$, the standard averaged model is derived

$$
\begin{equation*}
\langle z \dot{(t)}\rangle=\left[A_{\mathrm{I}}\langle z(t)\rangle+b_{\mathrm{I}}\right] d(t)+\left[A_{\mathrm{II}}\langle z(t)\rangle+b_{\mathrm{II}}\right][1-d(t)] \tag{7}
\end{equation*}
$$

where $d(t)$ is interpreted as the instantaneous duty cycle of Mode I and is allowed to take any value between 0 and 1. From this averaged model, the steady-state equations for a constant duty cycle $d(t) \equiv D$ are given by

$$
\begin{equation*}
0=\left[A_{\mathrm{I}} D+A_{\mathrm{II}}(1-D)\right] \overline{\langle z\rangle}+b_{\mathrm{I}} D+b_{\mathrm{II}}(1-D) \tag{8}
\end{equation*}
$$

Solving (8) yields

$$
\begin{aligned}
\overline{\langle z\rangle} & =\left[\begin{array}{llrr}
\overline{\left\langle i_{L_{1}}\right\rangle} & \overline{\left\langle i_{L_{2}}\right\rangle} & \overline{\left\langle v_{C_{1}}\right\rangle} & \overline{\left\langle v_{C_{2}}\right\rangle}
\end{array}\right]^{\prime}(9 \\
& =\left[\begin{array}{llll}
\frac{\overline{\left\langle i_{L_{2}}\right\rangle} D}{1-D} & -\frac{\overline{\left\langle v_{C_{2}}\right\rangle}}{R} & V_{\mathrm{in}} \frac{D}{1-D} & V_{\mathrm{in}} \frac{1-2 D}{1-D}
\end{array}\right]^{\prime},
\end{aligned}
$$

$$
\begin{equation*}
\frac{\overline{\left\langle v_{o}\right\rangle}}{V_{\mathrm{in}}}=\frac{\overline{\left\langle v_{C_{2}}\right\rangle}}{V_{\mathrm{in}}}=\frac{1-2 D}{1-D} \tag{10}
\end{equation*}
$$

Equation (10) gives the steady-state input-output voltage gain for a constant duty cycle $D$. Since the averaged models corresponding to $d(t) \equiv D=0$ or $d(t) \equiv D=1$ are not stable, then $D$ cannot reach the extreme values 0 or 1 . Therefore, $D \in(0,1)$ and $\overline{\left\langle v_{o}\right\rangle} / V_{\text {in }} \in(-\infty, 1)$. To obtain $\overline{\left\langle v_{o}\right\rangle} \in\left(-V_{\mathrm{in}}, V_{\mathrm{in}}\right)$ then $D \in\left(0, \frac{2}{3}\right)$.

## B. Inverter operation

We next briefly explain the inverter operation strategy proposed in [1], which is based on the steady-state gain equation (10]. In [1], the authors state that if a time-varying duty cycle is applied with a sufficiently slow time variation, then the steady-state gain equation would be approximately valid at every time instant (after the effect of the possible initial transients becomes negligible). Consequently, if the desired sinusoidal output voltage is given by

$$
\begin{equation*}
\left\langle v_{o}(t)\right\rangle=V_{o} \sin (\omega t) \tag{11}
\end{equation*}
$$

and if $\omega$ is small enough, then application of the time-varying duty cycle $d(t)$ that satisfies

$$
\begin{equation*}
\frac{\left\langle v_{o}(t)\right\rangle}{V_{\mathrm{in}}}=\frac{V_{o} \sin (\omega t)}{V_{\mathrm{in}}}=\frac{1-2 d(t)}{1-d(t)} \tag{12}
\end{equation*}
$$

would approximately yield the desired sinusoidal waveform at the output. Solving (12) for $d(t)$ yields

$$
\begin{equation*}
d(t)=\frac{1-M \sin (\omega t)}{2-M \sin (\omega t)} \quad \text { where } M=\frac{V_{o}}{V_{\mathrm{in}}} \tag{13}
\end{equation*}
$$

The circuit component values should be selected so that the harmonic distortion introduced by this steady-state-equationbased operation strategy is acceptable for the required output frequency $\omega$. This operation strategy is tested on a 40 W prototype [1], showing good results when connected to a purely resistive linear load. In the next sections, we derive results that give theoretical justification to this inverter operation strategy for all possible values of $L_{1}, L_{2}, C_{1}, C_{2}$, and $R$.

## III. AVERAGED MODEL STABILITY

In this section, we will show that if the load is linear and purely resistive, then the trajectories of the averaged model will remain bounded irrespective of the duty cycle evolution, provided the duty cycle does not reach the extreme values 0 or 1 . We then employ these results to show that if an open-loop duty cycle evolution, e.g. (13), is applied to the inverter, then the difference between any two trajectories having different initial conditions will converge to zero exponentially, and uniformly with respect to different duty cycle evolutions.

Consider the averaged model (7) and its steady-state solutions 10. Let $\overline{\langle z\rangle}_{0.5}$ denote the steady-state solution corresponding to $D=0.5$ :

$$
\overline{\langle z\rangle}_{0.5}=\left[\begin{array}{llll}
0 & 0 & V_{\mathrm{in}} & 0 \tag{14}
\end{array}\right]^{\prime}
$$

and define

$$
\begin{align*}
& x(t):=\langle z(t)\rangle-\overline{\langle z\rangle}_{0.5}  \tag{15}\\
& \mu(t):=d(t)-0.5 \tag{16}
\end{align*}
$$

The averaged model in the variable $x$ can be written as

$$
\begin{equation*}
\dot{x}(t)=A(\mu(t)) x(t)+B \mu(t) \tag{17}
\end{equation*}
$$

where

$$
\begin{align*}
& A(\mu(t))=P^{-1}\left[A_{0}^{q}+E_{0}^{q} \mu(t)\right] / 2,  \tag{18}\\
& B=P^{-1}\left[\begin{array}{llll}
2 V_{\text {in }} & 2 V_{\text {in }} & 0 & 0
\end{array}\right]^{\prime} / 2,  \tag{19}\\
& A_{0}^{q}=\left[\begin{array}{cccc}
0 & 0 & -0.5 & 0 \\
0 & 0 & 0.5 & 1 \\
0.5 & -0.5 & 0 & 0 \\
0 & -1 & 0 & -1 / R
\end{array}\right], \quad E_{0}^{q}=\left[\begin{array}{cccc}
0 & 0 & 1 & 0 \\
0 & 0 & 1 & 0 \\
-1 & -1 & 0 & 0 \\
0 & 0 & 0 & 0
\end{array}\right], \tag{20}
\end{align*}
$$

Since $d(t) \in[0,1]$, then according to 16$) \mu(t) \in[-0.5,0.5]$. Note that $A(-0.5)=A_{\mathrm{II}}$ and $A(0.5)=A_{\mathrm{I}}$, with $A_{\mathrm{I}}$ and $A_{\mathrm{II}}$ as in (3)-(4). For any given $0<\varepsilon \leq 0.5$, we will consider duty cycle evolutions $d(t) \in[\varepsilon, 1-\varepsilon]$, so that $\mu(t) \in \mathbb{U}$ with

$$
\begin{equation*}
\mathbb{U}:=[-\bar{\mu}, \bar{\mu}] \quad \text { and } \quad \bar{\mu}:=0.5-\varepsilon \tag{21}
\end{equation*}
$$

If $V_{\mathrm{in}}=0$ (passivation of the input voltage source), then the averaged model 17) reduces to

$$
\begin{equation*}
\dot{x}(t)=A(\mu(t)) x(t) \tag{22}
\end{equation*}
$$

We will also consider a perturbed version of system (22, given by

$$
\begin{equation*}
\dot{x}(t)=A(\mu(t)) x(t)+B u(t) \tag{23}
\end{equation*}
$$

with $B$ as in 19) and $u(t)$ an arbitrary perturbation input.
Our main result for the averaged model is the following.
Theorem 1: Let $0<\varepsilon \leq 0.5$ and consider the set $\mathbb{U}$ in (21). Let $H$ be a symmetric matrix defined as in 24) and whose nonzero entries satisfy 25.

$$
\begin{array}{r}
H:=\left[\begin{array}{cccc}
0 & 0 & h_{13} & h_{14} \\
0 & 0 & h_{23} & h_{24} \\
h_{13} & h_{23} & 0 & 0 \\
h_{14} & h_{24} & 0 & 0
\end{array}\right] \\
h_{13}<0, \quad h_{14}=-\frac{C_{2}\left(h_{13}-h_{23}\right)}{2 C_{1}} \tag{25a}
\end{array}
$$

$$
\begin{align*}
& h_{23}<\frac{L_{2}}{L_{1}} \frac{1-\varepsilon}{\varepsilon} h_{13}<0  \tag{25b}\\
& h_{24}>\frac{C_{2}}{C_{1}} \frac{\left(h_{23}-h_{13}\right)^{2} \bar{\mu}^{2}+h_{13} h_{23}}{-h_{13} 4 \varepsilon}>0 \tag{25c}
\end{align*}
$$

For every $\xi>0$, consider the quadratic function

$$
\begin{equation*}
V_{\xi}(x)=x^{\prime} \tilde{P}_{\xi} x, \quad \text { with } \quad \tilde{P}_{\xi}=P+\xi H \tag{26}
\end{equation*}
$$

and $P$ as in (6). Then,
a) for every $\xi>0$ sufficiently small and all $\mu \in \mathbb{U}$, we have $\tilde{P}_{\xi}>0$ and

$$
\begin{equation*}
A(\mu)^{\prime} \tilde{P}_{\xi}+\tilde{P}_{\xi} A(\mu) \leq-\alpha(\xi) \mathrm{I}<0 \tag{27}
\end{equation*}
$$

Consequently, the origin of system 22 is globally uniformly exponentially stable, where uniformity is with respect to the duty cycle $\mu$.
b) The trajectories of the perturbed system (23) satisfy

$$
\begin{equation*}
\|x(t)\| \leq K\|x(0)\| e^{-\lambda t}+G \sup _{0 \leq \tau \leq t}\|u(\tau)\| \tag{28}
\end{equation*}
$$

for some positive constants $K, \lambda$ and $G$, whenever $\mu(t) \in \mathbb{U}$ for all $t \geq 0$. Consequently, the system (23) is input-to-state stable (ISS) with respect to the input $u$, with exponential decay rate and linear input-to-state asymptotic gain, uniformly with respect to the duty cycle $\mu$.
c) The trajectories of the averaged model (17) satisfy

$$
\begin{equation*}
\|x(t)\| \leq K\|x(0)\| e^{-\lambda t}+G \bar{\mu} \tag{29}
\end{equation*}
$$

with the same positive constants $K, \lambda$ and $G$ as in item b) above, and provided $\mu(t) \in \mathbb{U}$ for all $t \geq 0$.
Theorem 1 establishes the global uniform boundedness and global uniform ultimate boundedness of the state of the averaged model (17) or (7), irrespectively of the duty cycle evolution but provided that $\mu(t) \in \mathbb{U}$ for all $t \geq 0$. We highlight that the global nature of this result is only theoretical because it is based on the assumption that the inverter operates in CCM. As we previously mentioned, in order for the inverter to operate in CCM, then the instantaneous $C_{1}$ voltage cannot become lower than $-V_{\text {in }}$.

The importance of Theorem 1 lies on the fact that it establishes the boundedness of the trajectories for every possible evolution of the duty cycle (provided it does not reach the extreme values) and all positive values of $L_{1}, L_{2}$, $C_{1}, C_{2}$ and $R$. Theorem 1 can also be used to show that the difference between any two trajectories starting from different initial conditions but corresponding to the same duty cycle evolution will ultimately converge to zero.

Corollary 1: Let $0<\varepsilon \leq 0.5$ and $\mu^{*}: \mathbb{R} \rightarrow \mathbb{U}$ with $\mathbb{U}$ as in 21. Let $\phi\left(t, x_{o}\right)$ denote the solution to 17) that satisfies $\dot{\phi}\left(t, x_{o}\right)=A\left(\mu^{*}(t)\right) \phi\left(t, x_{o}\right)+B \mu^{*}(t)$ and $\phi\left(0, x_{o}\right)=x_{o}$. Then, there exist positive constants $K$ and $\lambda$ such that for all $t \geq 0$ and all $x_{o}, y_{o}$,

$$
\begin{equation*}
\left\|\phi\left(t, x_{o}\right)-\phi\left(t, y_{o}\right)\right\| \leq K e^{-\lambda t}\left\|x_{o}-y_{o}\right\| \tag{30}
\end{equation*}
$$

Proof: Let $\delta(t)=\phi\left(t, x_{o}\right)-\phi\left(t, y_{o}\right)$. Then, $\dot{\delta}(t)=$ $A\left(\mu^{*}(t)\right) \delta(t)$ and apply Theorem 1/a.

Corollary 1 becomes especially useful in the case of a periodic duty cycle, such as that proposed in [1] and shown in 13). When $d(t)$ (and hence $\mu(t)$ ) is periodic and there exists a periodic steady-state trajectory $x^{*}(t)$ corresponding to such a periodic duty cycle evolution, then Corollary 1 establishes that the state trajectory will converge to the steady-state periodic trajectory when started from any initial condition (provided the inverter operates in CCM at all times). This latter result establishes the rationale for the inverter operation strategy proposed in [1], since it shows that the start-up transient will become extinct and hence the system will (eventually) operate along the periodic steady-state trajectory (which generates a sinusoidal waveform at the output if its frequency is low enough).

## IV. SWITCHED MODEL STABILITY

In this section, we provide results analogous to those in Theorem 1 and Corollary 1 but for the switched model of the semi-quasi-Z-source inverter.

Consider again the switched model (2)-6. We define

$$
\begin{equation*}
x(t):=z(t)-\overline{\langle z\rangle}_{0.5}, \quad \mu(t):=r(t)-0.5 \tag{31}
\end{equation*}
$$

with $\overline{\langle z\rangle}_{0.5}$ as in 14 . In the new variables, the switched model of the inverter coincides with (17)-20), the only difference being that $\mu(t)$ can only take the values 0.5 (Mode I) or -0.5 (Mode II). We consider that the switching function $\mu(t)$ is the output of a PWM. We thus consider the following class of switching signals.

Definition 1: A signal $\mu: \mathbb{R} \rightarrow\{-0.5,0.5\}$ is said to be of class $\mathbb{P W M}(T, \epsilon)$ with $0<2 \epsilon \leq T$, if it is rightcontinuous and for every integer $k$ there exists a time instant $t_{k}$ satisfying $k T+\epsilon \leq t_{k} \leq(k+1) T-\epsilon$ and

$$
\mu(t)= \begin{cases}0.5 & \text { if } k T \leq t<t_{k},  \tag{32}\\ -0.5 & \text { if } t_{k} \leq t<(k+1) T .\end{cases}
$$

The quantity $T$ will be called the period of $\mu$ (even though $\mu$ need not be periodic) and $\epsilon$ the dwell-time.

Our main results for the switched model are given below as Theorem 2 and Corollary 2 These results require the following lemma.

Lemma 1: Consider the matrix $P$ as in (6), and $A_{\mathrm{I}}, A_{\mathrm{II}}$ as in (3)-(4). Let $t_{\mathrm{I}}, t_{\mathrm{II}}$ be positive and for every $\epsilon \geq 0$ define

$$
\begin{equation*}
M_{\epsilon}:=e^{A_{\mathrm{I}} \epsilon} e^{A_{\mathrm{II}} t_{\mathrm{II}}} e^{A_{\mathrm{I}} t_{\mathrm{I}}} \tag{33}
\end{equation*}
$$

a) If $t_{\text {II }}<\pi \sqrt{L_{1} C_{1}}$, then $\boldsymbol{\rho}\left(M_{0}\right)<1$ and

$$
\begin{equation*}
M_{\epsilon}^{\prime} P M_{\epsilon}-P<0, \quad \text { for all } \epsilon>0 \tag{34}
\end{equation*}
$$

b) If $t_{\text {II }}=k \pi \sqrt{L_{1} C_{1}}$ for some positive integer $k$, then $\rho\left(M_{\epsilon}\right)=1$ for all $\epsilon \geq 0$.
Theorem 2: Consider $A(\cdot)$ and $B$ as in $18-20$ with $P$ as in (6) and let $\epsilon$ and $T$ satisfy $0<2 \epsilon \leq T<\pi \sqrt{L_{1} C_{1}}$. Then, for all $\mu \in \mathbb{P W M}(T, \epsilon)$, we have that
a) there exist positive constants $K$ and $\lambda$ so that for every initial condition $x(0)$ and all $t \geq 0$, the trajectories of (22) satisfy

$$
\begin{equation*}
\|x(t)\| \leq K\|x(0)\| e^{-\lambda t} \tag{35}
\end{equation*}
$$

b) there exist positive constants $K, \lambda$ and $G$ such that for every initial condition $x(0)$ and all $t \geq 0$, the trajectories of (23) satisfy

$$
\begin{equation*}
\|x(t)\| \leq K\|x(0)\| e^{-\lambda t}+G \sup _{0 \leq \tau \leq t}\|u(\tau)\| \tag{36}
\end{equation*}
$$

c) the trajectories of the switched model (17) satisfy

$$
\begin{equation*}
\|x(t)\| \leq K\|x(0)\| e^{-\lambda t}+\frac{G}{2} \tag{37}
\end{equation*}
$$

with the same $K, \lambda$ and $G$ as for item $b$ above.
Corollary 2: Consider $A(\cdot)$ and $B$ as in $18-20$ with $P$ as in (6). Let $\epsilon$ and $T$ satisfy $0<2 \epsilon \leq T<\pi \sqrt{L_{1} C_{1}}$ and consider $\mu^{*} \in \mathbb{P W M}(T, \epsilon)$. Let $\phi\left(t, x_{o}\right)$ denote the solution to 17) that satisfies $\dot{\phi}\left(t, x_{o}\right)=A\left(\mu^{*}(t)\right) \phi\left(t, x_{o}\right)+B \mu^{*}(t)$ and $\phi\left(0, x_{o}\right)=x_{o}$. Then, there exist positive constants $K$ and $\lambda$ such that for all $t \geq 0$ and all $x_{o}, y_{o}$,

$$
\begin{equation*}
\left\|\phi\left(t, x_{o}\right)-\phi\left(t, y_{o}\right)\right\| \leq K e^{-\lambda t}\left\|x_{o}-y_{o}\right\| \tag{38}
\end{equation*}
$$

Proof: Let $\delta(t)=\phi\left(t, x_{o}\right)-\phi\left(t, y_{o}\right)$. Then, $\dot{\delta}(t)=$ $A\left(\mu^{*}(t)\right) \delta(t)$ and apply Theorem 2 ab.
Theorem (2/a) establishes the global exponential stability of the switched system (22), uniformly over switching signals of PWM type, provided the PWM period $T$ is suitably limited. Since neither $A(-0.5)=A_{\text {II }}$ nor $A(0.5)=A_{\mathrm{I}}$ are Hurwitz, then the exponential stability of the switched system 22, is not possible under arbitrary switching, nor by imposing only a minimum dwell-time on each mode [9], [10]. By Lemma 1 1 D, it follows that imposing a maximum dwell-time for Mode II is essential in order to ensure the exponential stability.

The switched system 22 has a natural quadratic Lyapunov function given by the circuit's stored energy: $V(x)=$ $x^{\prime} P x$ with $P$ as in (6). The derivative of this energy function along the trajectories of 22 has the form $-x^{\prime} C^{\prime} C x$, where

$$
\begin{aligned}
A_{\mathrm{I}}^{\prime} P+P A_{\mathrm{I}} & =A_{\mathrm{II}}^{\prime} P+P A_{\mathrm{II}}=-C^{\prime} C \\
C & =\left[\begin{array}{llll}
0 & 0 & 0 & R^{-1 / 2}
\end{array}\right],
\end{aligned}
$$

and is hence only negative semidefinite. Even though several existing results extend LaSalle's Invariance Principle to switched systems [11], [12], [13], none of these can be employed to ascertain the uniform asymptotic or exponential stability of 22 by means of such natural Lyapunov function. The main difficulty is that neither $\left(A_{\mathrm{I}}, C\right)$ nor $\left(A_{\mathrm{II}}, C\right)$ are detectable, as can be easily checked. Theorem 8 of [11] relaxes this observability assumption. When applied to system 22) employing the natural quadratic Lyapunov function $V(x)$, Theorem 8 of [11] establishes that for a sufficiently large type of switching signals (identified as $\mathcal{S}_{p \text {-dwell }}$ in [11]), the trajectories of (22) will converge to the smallest subspace $\mathcal{M}$ that is invariant under both $A_{\mathrm{I}}$ and $A_{\mathrm{II}}$, and contains both the unobservable subspaces of $\left(A_{\mathrm{I}}, C\right)$ and $\left(A_{\mathrm{II}}, C\right)$. The problem in the current case is that $\mathcal{M}=\mathbb{R}^{4}$, which coincides with the whole state space. In view of the aforementioned difficulties, one main feature of our proof
of Theorem 2 is that all stability results are established by exploiting the natural Lyapunov function $V(x)=x^{\prime} P x$.

Analogously to Theorem 11 in Section III, Theorem 2b) establishes that the perturbed switched system (23) is ISS with respect to the perturbation input $u$, having linear input-to-state gain and exponential decay, uniformly over switching signals in $\mathbb{P W W M}(T, \epsilon)$. Theorem 2 $2 \mathbb{E}$ ) establishes that all trajectories of the inverter circuit remain bounded, and ultimately bounded uniformly over the switching signals considered. Finally, Corollary 2 shows that the difference between trajectories corresponding to different initial conditions but the same switching signal will exponentially decrease to zero, and that the exponential convergence rate can be uniform over all the switching signals considered. As previously mentioned in the case of the averaged model, this result gives theoretical justification to the inverter operation strategy proposed in [1].

## V. Conclusions

We have provided large-signal stability results for both the standard averaged model and the switched model of the semi-quasi-Z-source inverter operating in continuous conduction mode. Specifically, we have established that the state trajectories are bounded and ultimately bounded with exponential decay rate, uniformly over all duty cycle evolutions provided reasonable operating conditions are ensured. Our stability results hold for every possible value of the circuit inductors, capacitors and linear resistive load, and can be used to show that all trajectories corresponding to the same duty cycle evolution but different initial conditions will converge to the same steady-state trajectory. This result gives theoretical justification to the inverter operation strategy proposed in [1], in the sense that under reasonable restrictions on the PWM signal, the state trajectory will converge exponentially to the steady state one from any initial condition. Our stability results are derived by exploiting the natural energy function of the circuit, although the quadratic term in its derivative along the system trajectories is only negative semidefinite.

Future work will be focused on establishing similar stability results for more general, including nonlinear, load types, and on obtaining less conservative transient and ultimate bounds.

## Appendix

## Proof of Theorem 1

We require the following lemma, whose proof is straightforward and hence omitted.

Lemma 2: Let $\bar{A}$ be Hurwitz and suppose that there exists $P=P^{\prime}>0$ such that:

$$
\begin{equation*}
\bar{A}^{\prime} P+P \bar{A}=-Q \leq 0 . \tag{39}
\end{equation*}
$$

Then, there exists $\tilde{P}=\tilde{P}^{\prime}>0$ such that $\bar{A}^{\prime} \tilde{P}+\tilde{P} \bar{A}<0$ if and only if there exists $H=H^{\prime}$ satisfying:

$$
\begin{equation*}
X(\bar{A}):=S_{Q}^{\prime}\left(\bar{A}^{\prime} H+H \bar{A}\right) S_{Q}<0 \tag{40}
\end{equation*}
$$

where the columns of $S_{Q}$ constitute an orthonormal basis for the null space of $Q$, i.e. $S_{Q} \in \mathbb{R}^{n \times r}, S_{Q}^{\prime} S_{Q}=\mathrm{I}, Q S_{Q}=0$,
and $r$ is the greatest such integer. If (40) holds, then every $\tilde{P}$ of the form $\tilde{P}=P+\xi H$ with $\xi>0$ sufficiently small will satisfy $\bar{A}^{\prime} \tilde{P}+\tilde{P} \bar{A}<0$.
a) First, note that $A(\mu)$ is Hurwitz for all $\mu \in \mathbb{U}$. The proof is based on showing that $X(A(\mu))<0$ for all $\mu \in \mathbb{U}$ for the given matrix $H$. Note that for $\bar{A}=A(\mu)$, 39) is satisfied with $P$ as in 6 for all $\mu \in \mathbb{U}$. We have $\operatorname{ker} Q=\operatorname{Im} S_{Q}$, $S_{Q}=\left[e_{1}, e_{2}, e_{3}\right]$. Let $X_{i}$ denote the principal $i$-th order determinant of $X$, so that $X<0$ if and only if $X_{i}<0$ for odd $i$ and $X_{i}>0$ for even $i$. Direct computation gives

$$
X(A(\mu))_{1}=\frac{h_{13}(2 \mu+1)}{C_{1}}<0
$$

for all $\mu \in \mathbb{U}$, where the inequality follows from 25a. Also,

$$
\begin{aligned}
X(A(\mu))_{2} & =-\frac{b}{C_{1}}-\frac{2 X(A(\mu))_{1} h_{24}}{C_{2}} \\
b & =\frac{\left(h_{13}-h_{23}\right)^{2} \mu^{2}+h_{13} h_{23}}{C_{1}}
\end{aligned}
$$

and $X(A(\mu))_{2}>0$ for all $\mu \in \mathbb{U}$ follows from 25c). Finally,

$$
\begin{aligned}
X(A(\mu))_{3} & =\operatorname{det}(X(A(\mu)))=-c X(A(\mu))_{2} \\
c & =\frac{L_{2} h_{13}(2 \mu+1)+L_{1} h_{23}(2 \mu-1)}{L_{1} L_{2}} .
\end{aligned}
$$

We have $c>0$ for all $\mu \in \mathbb{U}$ from (25b), and hence $X(A(\mu))_{3}>0$ for all $\mu \in \mathbb{U}$.

By application of Lemma 2, we know that there exist $\xi_{1}, \xi_{2}>0$ such that $A(\bar{\mu})^{\prime} \tilde{P}_{\xi}+\tilde{P}_{\xi} A(\bar{\mu})<0$ for all $0<$ $\xi<\xi_{1}$ and $A(-\bar{\mu})^{\prime} \tilde{P}_{\xi}+\tilde{P}_{\xi} A(-\bar{\mu})<0$ for all $0<\xi<\xi_{2}$. Since $A(\mu)$ is a convex combination of $A(-\bar{\mu})$ and $A(\bar{\mu})$, it follows that $A(\mu)^{\prime} \tilde{P}_{\xi}+\tilde{P}_{\xi} A(\mu)<0$ for all $\mu \in \mathbb{U}$ and all $0<\xi<\min \left\{\xi_{1}, \xi_{2}\right\}$.
b) Let $\xi>0$ be such that $\tilde{P}_{\xi}>0$ and 27 holds for all $\mu \in \mathbb{U}$. Along trajectories of (23) we have that for all $\mu \in \mathbb{U}$,

$$
\dot{V}_{\xi}(x, \mu, u) \leq-\alpha(\xi)\|x\|^{2}+2\left\|\tilde{P}_{\xi} B\right\|\|x\|\|u\|
$$

and the proof follows from standard ISS arguments [14].
c) Just apply Theorem $1 \mid$ bith $u(t)=\mu(t)$.

## Proof of Lemma 1

We have

$$
\begin{equation*}
A_{\mathrm{I}}^{\prime} P+P A_{\mathrm{I}}=A_{\mathrm{II}}^{\prime} P+P A_{\mathrm{II}}=-Q \leq 0 \tag{41}
\end{equation*}
$$

with

$$
\begin{equation*}
Q=\operatorname{diag}\left(0,0,0, \frac{1}{R}\right) \tag{42}
\end{equation*}
$$

Consequently, for every $\epsilon \geq 0$ we can prove that

$$
\begin{equation*}
M_{\epsilon}^{\prime} P M_{\epsilon}-P \leq 0, \quad\left(M_{\epsilon}^{2}\right)^{\prime} P M_{\epsilon}^{2}-P \leq 0 \tag{43}
\end{equation*}
$$

and

$$
\begin{align*}
& x^{\prime}\left(M_{\epsilon}^{\prime} P M_{\epsilon}-P\right) x<0 \Longrightarrow \\
& x^{\prime}\left(\left(M_{\epsilon}^{2}\right)^{\prime} P M_{\epsilon}^{2}-P\right) x<0 \tag{44}
\end{align*}
$$

Let $x_{o}$ satisfy

$$
\begin{equation*}
x_{o}^{\prime}\left(M_{\epsilon}^{\prime} P M_{\epsilon}-P\right) x_{o}=0 . \tag{45}
\end{equation*}
$$

The vector $x_{o}$ satisfies 45 if and only if

$$
\begin{align*}
& e^{A_{\mathrm{I}} t} x_{o} \in \mathcal{S}_{\mathrm{I}}, \text { for all } t \in\left[0, t_{\mathrm{I}}\right),  \tag{46}\\
& e^{A_{\mathrm{II}} t} e^{A_{1} t_{\mathrm{I}}} x_{o} \in \mathcal{S}_{\mathrm{II}}, \quad \text { for all } t \in\left[0, t_{\mathrm{II}}\right) \quad \text { and }  \tag{47}\\
& e^{A_{\mathrm{I}} t} e^{A_{\mathrm{I}} t_{\mathrm{II}}} e^{A_{1} t_{\mathrm{I}}} x_{o} \in \mathcal{S}_{\mathrm{I}}, \quad \text { for all } t \in[0, \epsilon) . \tag{48}
\end{align*}
$$

where $\mathcal{S}_{\mathrm{I}}=\operatorname{span}\left\{e_{1}\right\}$ and $\mathcal{S}_{\text {II }}=\operatorname{span}\left\{e_{1}, e_{3}\right\}$ are the largest subspaces invariant under $A_{\mathrm{I}}$ or $A_{\mathrm{II}}$, respectively, that are contained in ker $Q$. From (3)-(4) and (46), it follows that $x_{o}=e^{A_{1} t} x_{o}=\alpha e_{1}$ for some $\alpha \in \mathbb{R}$ and for all $t \in\left[0, t_{\mathrm{I}}\right]$. From (3)- 4 ) and (47), then $e^{A_{\mathrm{II}} t_{\mathrm{II}}} e^{A_{1} t_{\mathrm{I}}} x_{o}=$ $\alpha\left[\cos \left(\omega t_{\mathrm{II}}\right) e_{1}+\sqrt{L_{1} / C_{1}} \sin \left(\omega t_{\mathrm{II}}\right) e_{3}\right]$, with $\omega=1 / \sqrt{L_{1} C_{1}}$. Since $t_{\text {II }}<\pi \sqrt{L_{1} C_{1}}$, then $\sin \left(\omega t_{\text {II }}\right) \neq 0$. Hence, if $\epsilon>0$, from (48) we must have $x_{1}:=M_{0} x_{o}=e^{A_{I I} t_{I I}} e^{A_{1} t_{\mathrm{I}}} x_{o} \in \mathcal{S}_{\mathrm{I}}$, which implies that $x_{o}=0$. Therefore, if $\epsilon>0$, 45) implies that $x_{o}=0$ and 34 is established.

For $\epsilon=0$, suppose that $x_{o} \neq 0$ satisfies 45). By the previous argument, $x_{1}=M_{0} x_{o} \notin \mathcal{S}_{\text {I }}$ and

$$
\begin{aligned}
x_{1}^{\prime}\left(M_{0}^{\prime} P M_{0}-P\right) x_{1} & =x_{o}^{\prime}\left(\left(M_{0}^{2}\right)^{\prime} P M_{0}^{2}-M_{0} P M_{0}\right) x_{o} \\
& =x_{o}^{\prime}\left(\left(M_{0}^{2}\right)^{\prime} P M_{0}^{2}-P\right) x_{o}<0
\end{aligned}
$$

where the latter inequality follows because, repeating the previous argument, $x_{1}^{\prime}\left(M_{0}^{\prime} P M_{0}-P\right) x_{1}=0$ would imply that $x_{1} \in \mathcal{S}_{1}$, a contradiction. Recalling (44), then we conclude that $\boldsymbol{\rho}\left(M_{0}^{2}\right)<1$ and hence $\boldsymbol{\rho}\left(M_{0}\right)<1$.
(b) By (43), then $\rho\left(M_{\epsilon}\right) \leq 1$ for all $\epsilon \geq 0$. Take $x_{o}=e_{1}$. Then, $e^{A_{1} t_{\mathrm{I}}} x_{o}=x_{o}$ and $e^{A_{\mathrm{II}} t_{\mathrm{II}}} x_{o}=\cos \left(\omega t_{\mathrm{II}}\right) e_{1}+$ $\sqrt{L_{1} / C_{1}} \sin \left(\omega t_{\mathrm{II}}\right) e_{3}=(-1)^{k} e_{1}=(-1)^{k} x_{o}$. Therefore, $M_{\epsilon} x_{o}=(-1)^{k} x_{o}$, showing that $\lambda=(-1)^{k}$ is an eigenvalue of $M_{\epsilon}$ with $|\lambda|=1$.

## Proof of Theorem 2

Consider $V(x)=x^{\prime} P x$. For all $x$, we have

$$
\begin{align*}
\kappa_{1}\|x\|^{2} & \leq V(x)  \tag{49}\\
\dot{V}(x, \mu) & =\kappa_{2}\|x\|^{2}  \tag{50}\\
\dot{V}(x) & =-x^{\prime} Q x \leq 0
\end{align*}
$$

with $\kappa_{1}, \kappa_{2}>0$ and $Q$ as in (42) being positive semidefinite. Consider time instants of the form $t_{k}=k T+\epsilon / 2$ for all nonnegative integers $k$.
a) We have

$$
\begin{align*}
x\left(t_{k+1}\right) & =M\left(t_{\mathrm{I}, k}, t_{\mathrm{II}, k}\right) x\left(t_{k}\right)  \tag{51}\\
M\left(t_{\mathrm{I}, k}, t_{\mathrm{II}, k}\right) & :=e^{A_{\mathrm{I}} \epsilon / 2} e^{A_{\mathrm{II}} t_{\mathrm{II}, k}} e^{A_{\mathrm{I}} t_{\mathrm{I}, k}} \tag{52}
\end{align*}
$$

with $t_{\mathrm{II}, k} \in[\epsilon, T-\epsilon]$ and $t_{I, k}=T-t_{\mathrm{II}, k}-\epsilon / 2$ because $\mu \in \mathbb{P W M}(T, \epsilon)$. Since $T<\pi \sqrt{L_{1} C_{1}}$, using Lemma 1 it follows that

$$
\begin{equation*}
R\left(t_{\mathrm{I}, k}, t_{\mathrm{II}, k}\right):=M\left(t_{\mathrm{I}, k}, t_{\mathrm{II}, k}\right)^{\prime} P M\left(t_{\mathrm{I}, k}, t_{\mathrm{II}, k}\right)-P<0 \tag{53}
\end{equation*}
$$

Define

$$
\begin{equation*}
\kappa_{3}:=\sup _{\substack{t_{\mathrm{II}} \in[\epsilon, T-\epsilon] \\ t_{\mathrm{I}}=T-t_{\mathrm{II}}-\epsilon / 2}} \lambda_{\min }\left[-R\left(t_{\mathrm{I}}, t_{\mathrm{II}}\right)\right] \tag{54}
\end{equation*}
$$

By 53, and since $M(\cdot, \cdot)$ is continuous on its arguments and the supremum above is taken over a compact set, then
$\kappa_{3}>0$. We thus have

$$
\begin{align*}
& V\left(x\left(t_{k+1}\right)\right)-V\left(x\left(t_{k}\right)\right)= \\
& x\left(t_{k}\right)^{\prime}\left[M\left(t_{\mathrm{I}, k}, t_{\mathrm{II}, k}\right)^{\prime} P M\left(t_{\mathrm{I}, k}, t_{\mathrm{II}, k}\right)-P\right] x\left(t_{k}\right) \\
& \leq-\kappa_{3}\left\|x\left(t_{k}\right)\right\|^{2} \leq-\frac{\kappa_{3}}{\kappa_{2}} V\left(x\left(t_{k}\right)\right), \quad \text { for all } k \geq 0 \tag{55}
\end{align*}
$$

Consequently, $V\left(x\left(t_{k}\right)\right) \leq\left(1-\frac{\kappa_{3}}{\kappa_{2}}\right)^{k} V\left(x\left(t_{0}\right)\right)$. Taking 50p and 55) into account, we can obtain the following bound

$$
\begin{align*}
V(x(t)) & \leq e^{r(\epsilon / 2+T)} e^{-r t} V(x(0)), \quad \text { for all } t \geq 0  \tag{56}\\
r & :=-\frac{\log \left(1-\kappa_{3} / \kappa_{2}\right)}{T}>0 \tag{57}
\end{align*}
$$

whence, for all $t \geq 0$,

$$
\begin{equation*}
\|x(t)\| \leq \sqrt{\frac{\kappa_{2}}{\kappa_{1}} e^{r(\epsilon / 2+T)}} e^{-r t / 2}\|x(0)\| \tag{58}
\end{equation*}
$$

Therefore, we have established (35) with $K=$ $\sqrt{\frac{\kappa_{2}}{\kappa_{1}}} e^{r(\epsilon / 2+T)}$ and $\lambda=r / 2$.
b) Let $\Phi\left(t ; t_{\mathrm{I}}, t_{\mathrm{II}}, \epsilon\right)$ be defined as follows

$$
\Phi\left(t ; t_{\mathrm{I}}, t_{\mathrm{II}}, \epsilon\right):= \begin{cases}e^{A_{\mathrm{I}} t} & \text { if } 0 \leq t<t_{\mathrm{I}},  \tag{59}\\ e^{A_{\mathrm{II}}\left(t-t_{1}\right)} e^{A_{1} t_{\mathrm{I}}} & \text { if } t_{\mathrm{t}} \leq t<t_{\mathrm{I}}+t_{\mathrm{II}}, \\ e^{A_{\mathrm{I}} t} e^{A_{\mathrm{II}} t_{\mathrm{II}} e^{A_{\mathrm{I}} t_{\mathrm{I}}}} & \text { if } t_{\mathrm{l}}+t_{\mathrm{II}} \leq t \leq t_{\mathrm{I}}+t_{\mathrm{II}}+\epsilon .\end{cases}
$$

For all $t_{k} \leq t \leq t_{k+1}$, we have

$$
\begin{align*}
x(t) & =\Phi\left(t-t_{k} ; t_{\mathrm{I}, k}, t_{\mathrm{II}, k}, \epsilon / 2\right) x\left(t_{k}\right)+\Psi_{k}(t), \text { with } \\
\Psi_{k}(t) & :=\int_{t_{k}}^{t} \Phi\left(t-\tau ; t_{\mathrm{I}, k}, t_{\mathrm{II}, k}, \epsilon / 2\right) B u(\tau) d \tau \tag{61}
\end{align*}
$$

where $t_{\mathrm{II}, k} \in[\epsilon, T-\epsilon]$ and $t_{\mathrm{I}, k}=T-t_{\mathrm{II}, k}-\epsilon / 2$ for all $k$. For all $t_{k} \leq t \leq t_{k+1}$, we can bound the norm of $\Psi_{k}(t)$ as follows

$$
\begin{align*}
\left\|\Psi_{k}(t)\right\| & \leq \bar{\Phi}\|B\| \bar{u}_{k},  \tag{62}\\
\bar{\Phi} & :=\sup _{\substack{t \in[0, T] \\
t_{\mathrm{II}} \in[\epsilon, T-\epsilon] \\
t_{1}=T-t_{\mathrm{II}}-\epsilon / 2}}\left\|\Phi\left(t ; t_{\mathrm{I}}, t_{\mathrm{II}}, \epsilon / 2\right)\right\|,  \tag{63}\\
\bar{u}_{k} & :=\int_{t_{k}}^{t_{k+1}}\|u(t)\| d t . \tag{64}
\end{align*}
$$

Using (60)-64), we can bound the increment of $V$ as

$$
\begin{align*}
V\left(x\left(t_{k+1}\right)\right) & -V\left(x\left(t_{k}\right)\right) \\
& \leq-\kappa_{3}\left\|x\left(t_{k}\right)\right\|^{2}+c_{1} \bar{u}_{k}\left\|x\left(t_{k}\right)\right\|+c_{2} \bar{u}_{k}^{2} \tag{65}
\end{align*}
$$

with $\kappa_{3}$ as in (52-(54) and $c_{1}, c_{2}>0$. It follows that for every $0<\theta<1$,

$$
\begin{align*}
V\left(x\left(t_{k+1}\right)\right)-V\left(x\left(t_{k}\right)\right) & \leq-\kappa_{3}(1-\theta)\left\|x\left(t_{k}\right)\right\|^{2} \\
& \leq-\frac{\kappa_{3}(1-\theta)}{\kappa_{2}} V\left(x\left(t_{k}\right)\right) . \tag{66}
\end{align*}
$$

whenever

$$
\left\|x\left(t_{k}\right)\right\| \geq \frac{c_{1}+\sqrt{c_{1}^{2}+4 \kappa_{3} \theta c_{2}}}{2 \kappa_{3} \theta} \bar{u}_{k}=: L_{1} \bar{u}_{k}
$$

Following standard Lyapunov arguments [14], we can reach the following inequality

$$
\begin{equation*}
\left\|x\left(t_{k}\right)\right\| \leq c_{3} \alpha^{k}\left\|x\left(t_{0}\right)\right\|+L_{2} \sup _{0 \leq \tilde{k} \leq k-1} \bar{u}_{\tilde{k}} \tag{67}
\end{equation*}
$$

where $0<\alpha<1$. From (60), it follows that

$$
\begin{equation*}
\|x(t)\| \leq \bar{\Phi}\left\|x\left(t_{k}\right)\right\|+\bar{\Phi}\|B\| \int_{t_{k}}^{t}\|u(\tau)\| d \tau \tag{68}
\end{equation*}
$$

for all $t_{k} \leq t \leq t_{k+1}$ for all $k \geq 0$. Combining 67)-68),

$$
\begin{align*}
\|x(t)\| \leq & \bar{\Phi} c_{3} \alpha^{k}\left\|x\left(t_{0}\right)\right\| \\
& +\bar{\Phi} L_{2} \sup _{0 \leq \tilde{k} \leq k-1} \bar{u}_{\tilde{k}}+\bar{\Phi}\|B\| \int_{t_{k}}^{t}\|u(\tau)\| d \tau \tag{69}
\end{align*}
$$

for all $t_{k} \leq t \leq t_{k+1}$. Define

$$
\begin{equation*}
\left\|u_{[a, b]}\right\|:=\sup _{\tau \in[a, b]}\|u(\tau)\| \tag{70}
\end{equation*}
$$

Then,

$$
\begin{equation*}
\bar{u}_{k} \leq T\left\|u_{\left[t_{k}, t_{k+1}\right]}\right\| \leq T\left\|u_{\left[t_{0}, t_{k+1}\right]}\right\| \tag{71}
\end{equation*}
$$

Using (71) in 69) we arrive at

$$
\begin{equation*}
\|x(t)\| \leq \bar{\Phi} c_{3} \alpha^{k}\left\|x\left(t_{0}\right)\right\|+c_{5}\left\|u_{\left[t_{0}, t\right]}\right\| \tag{72}
\end{equation*}
$$

for all $t_{k} \leq t \leq t_{k+1}$, for all $k \geq 0$. Therefore,

$$
\begin{equation*}
\|x(t)\| \leq c_{6} e^{-\lambda t}\left\|x\left(t_{0}\right)\right\|+c_{5}\left\|u_{\left[t_{0}, t\right]}\right\| \tag{73}
\end{equation*}
$$

with $\lambda>0$. Since $t_{0}=\epsilon / 2$, it follows that

$$
\begin{equation*}
\left\|x\left(t_{0}\right)\right\| \leq\left\|e^{A_{1} \epsilon / 2}\right\|\|x(0)\|+\bar{\Phi}\|B\|\left\|u_{[0, \epsilon / 2]}\right\| \epsilon / 2 \tag{74}
\end{equation*}
$$

Combining (73)-(74), the result follows.
c) Just apply Theorem 2bp with $u(t)=\mu(t)$.

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