# A $W^{*}$-CORRESPONDENCE APPROACH TO MULTI-DIMENSIONAL LINEAR DISSIPATIVE SYSTEMS 

J.A. BALL AND S. TER HORST


#### Abstract

Recent work of the operator algebraists P. Muhly and B. Solel, primarily motivated by the theory of operator algebras and mathematical physics, delineates a general abstract framework where system theory ideas appear in disguised form. These system-theory ingredients include: system matrix for an input/state/output linear system, $Z$-transform from a "time domain" to a "frequency domain", and $Z$-transform of the output signal given by an observation function applied to the initial condition plus a transfer function applied to the $Z$-transform of the input signal. Here we set down the definitions and main results for the general Muhly-Solel formalism and illustrate them for two specific types of multi-dimensional linear systems: (1) dissipative FornasiniMarchesini state-space representations with transfer function equal to a holomorphic operator-valued function on the unit ball in $\mathbb{C}^{d}$, and (2) noncommutative dissipative Fornasini-Marchsini linear systems with evolution along a free semigroup and with transfer function defined on the noncommutative ball of strict row contractions on a Hilbert space.


## 1. Introduction

In the last decade operator algebraists have worked on a sophisticated theory, based on the notions of $C^{*}$ - and $W^{*}$-correspondences, that encompasses many operator theoretical results, for instance, from characteristic model and interpolation theory [12, 11, both in time-varying and multivariable settings, as well as for a range of more exotic examples. In the recent paper [14] a transfer function realization result was obtained for the generalized Schur class (introduced in [13]) in this $W^{*}$-correspondence setting; see also 4]. It is the purpose of this note to enrich the developed theory of [14] with a system theory interpretation, and show how the $W^{*}$ correspondences can be specified to include some examples from multi-dimensional system theory.

In Section 2 we delineate the two examples that we will obtain as special cases of the general Muhly-Solel framework: (1) commutative Fornasini-Marchesini systems with evolution along the positive-orthant integer lattice $\mathbb{Z}_{+}^{d}$, and (2) noncommutative Fornasini-Marchesini systems with evolution along the free semigroup generated by $d$ letters. In Sections 3 and 4 we lay out the somewhat formidable array of definitions needed to formulate the general Muhly-Solel setup. The definition and basic properties of the abstract input/state/output linear systems arising in the general Muhly-Solel framework are given in Section 5. Section 6 is dedicated to the abstract point-evaluation and the associated abstract $Z$-transform, while in Section 7 we arrive at the frequency-domain version of the system and basic properties of the associated observation function and transfer function. An other type of linear systems which fits the Muhly-Solel framework (1-D time-varying systems) and some other types which do not are mentioned in the final Section 8. Throughout
the exposition, new ideas and concepts are illustrated for the two basic examples as they are introduced.

The notation is mostly standard, but we mention here a few conventions. For any index set $\Omega$ we write $\ell^{2}(\Omega)$ for Hilbert space of complex-valued square summable functions over $\Omega$ :

$$
\left\{f: \Omega \rightarrow \mathbb{C}: \sum_{\omega \in \Omega}|f(\omega)|^{2}<\infty\right\}
$$

while $\ell^{\infty}(\Omega)$ denotes the space of bounded complex-valued functions over $\Omega$. For a Hilbert space $\mathcal{H}$ we use $\ell_{\mathcal{H}}^{2}(\Omega)$ and $\ell_{\mathcal{H}}^{\infty}(\Omega)$ as short-hands for $\ell^{2}(\Omega) \otimes \mathcal{H}$ and $\ell^{\infty}(\Omega) \otimes \mathcal{H}$, respectively.

## 2. Two examples of multidimensional Linear systems and their TRANSFER FUNCTIONS

In this section we review dissipative systems of Fornasini-Marchesini type; both commutative and noncommutative. In a somewhat more general form the noncommutative case was studied in [5, 6] with connections to robust control discussed in [7. See also [3] for an overview of these types of systems and their applications.
2.1. Commutative Fornasini-Marchesini systems. Consider a system evolving over $\mathbb{Z}_{+}^{d}$, the set of $d$-tuples of nonnegative integers, given by the system equations:

$$
\left\{\begin{align*}
x(n)= & A_{1} x\left(n-e_{1}\right)+\cdots+A_{d} x\left(n-e_{d}\right)  \tag{2.1}\\
& \quad+B_{1} u\left(n-e_{1}\right)+\cdots+B_{d} u\left(n-e_{d}\right) \\
y(n)= & C x(n)+D u(n)
\end{align*}\right.
$$

where $e_{1}, \ldots, e_{d}$ are the standard basis vectors in $\mathbb{R}^{d}$ and we set $x(n)=0$ for all $n \in \mathbb{Z}^{d}$ that are not in $\mathbb{Z}_{+}^{d}$, and with a system matrix of the form

$$
\left[\begin{array}{cc}
A_{1} & B_{1}  \tag{2.2}\\
\vdots & \vdots \\
A_{d} & B_{d} \\
C & D
\end{array}\right]:\left[\begin{array}{l}
\mathcal{X} \\
\mathcal{U}
\end{array}\right] \rightarrow\left[\begin{array}{c}
\mathcal{X}^{d} \\
\mathcal{Y}
\end{array}\right]
$$

Here $A_{j}, B_{j}, C$ and $D$ are bounded linear operators and $\mathcal{X}, \mathcal{U}$ and $\mathcal{Y}$ are Hilbert spaces. We further assume the system matrix (2.2) to be contractive. In this case we speak of a dissipative system of commutative Fornasini-Marchasini type (see [9]). Note that the output sequence $\{y(n)\}_{n \in \mathbb{Z}_{+}^{d}}$ and state sequence $\{x(n)\}_{n \in \mathbb{Z}_{+}^{d}}$ are completely determined by the input sequences $\{u(n)\}_{n \in \mathbb{Z}_{+}^{d}}$ and the initial state $x(0)$. Moreover, the input, output and state sequences satisfy the energy balance inequalities: For each $N \in \mathbb{Z}_{+}$we have

$$
\begin{aligned}
\sum_{n \in \mathbb{Z}_{+}^{d} ;|n| \leq N} \omega(n)\|y(n)\|^{2}+\sum_{n \in \mathbb{Z}_{+}^{d} ;|n|=N} \omega(n)\|x(n)\|^{2} \leq \\
\leq \sum_{n \in \mathbb{Z}_{+}^{d} ;|n| \leq N} \omega(n)\|u(n)\|^{2}+\|x(0)\|^{2}
\end{aligned}
$$

where $\omega(n)=\frac{|n|!}{n!}$ with $|n|=n_{1}+\cdots+n_{d}$ and $n!=n_{1}!\cdots n_{d}$ ! for each $n=$ $\left(n_{1}, \ldots, n_{d}\right) \in \mathbb{Z}_{+}^{d}$. For a Hilbert space $\mathcal{V}$ we write $\ell_{\mathcal{V}, \omega}^{2}\left(\mathbb{Z}_{+}^{d}\right)$ for the weighted
$\ell^{2}$-space of sequences $v=\{v(n)\}_{n \in \mathbb{Z}_{+}^{d}}$ with $v(n) \in \mathcal{U}$ and

$$
\|v\|^{2}=\sum_{n \in \mathbb{Z}_{+}^{d}} \omega(n)\|v(n)\|^{2}<\infty .
$$

It then follows from the energy balance inequalities that if the input sequence $u$ is in $\ell_{\mathcal{U}, \omega}^{2}\left(\mathbb{Z}_{+}^{d}\right)$, the output sequence $y$ is in $\ell_{\mathcal{Y}, \omega}^{2}\left(\mathbb{Z}_{+}^{d}\right)$ and $\|y\|^{2} \leq\|u\|^{2}+\|x(0)\|^{2}$, while the state sequence $x$ is bounded; more precisely $\sup _{N \in \mathbb{Z}_{+}} \sum_{|n|=N} \omega(n)\|x(n)\|^{2}<\infty$. The state sequence is in $\ell_{\mathcal{X}, \omega}^{2}\left(\mathbb{Z}_{+}^{d}\right)$ if an additional stability assumption on state operators $A_{1}, \ldots, A_{d}$ is satisfied.

We write $\mathbb{B}^{d}$ for the open unit ball in $\mathbb{C}^{d}$ :

$$
\begin{equation*}
\mathbb{B}^{d}:=\left\{z=\left(z_{1}, \ldots, z_{d}\right) \in \mathbb{C}^{d}: \sum_{j=1}^{d}\left|z_{j}\right|^{2}<1\right\} \tag{2.3}
\end{equation*}
$$

and for a $z=\left(z_{1}, \ldots, z_{d}\right)$ in $\mathbb{C}^{d}$ we use the notation $z^{n}$ for $z_{1}^{n_{1}} \cdots z_{d}^{n_{d}}$ in case $n=\left(n_{1}, \ldots, n_{d}\right) \in \mathbb{Z}_{+}^{d}$. Given a bounded sequence $v=\{v(n)\}_{n \in \mathbb{Z}_{+}^{d}}$, with $v(n)$ in a Hilbert space $\mathcal{V}$, we define a $\mathcal{V}$-valued function $\widehat{v}$ on $\mathbb{B}^{d}$ via the $Z$-transform:

$$
\begin{equation*}
\widehat{v}(z)=\sum_{n \in \mathbb{Z}_{+}^{d}} z^{n} v(n) \tag{2.4}
\end{equation*}
$$

Provided that the input sequence is in $\ell_{\mathcal{U}, \omega}^{2}\left(\mathbb{Z}_{+}^{d}\right)$, the $Z$-transforms of the input, output and state sequences are related through the frequency domain equations:

$$
\left\{\begin{array}{l}
\widehat{x}(z)=x(0)+\left(\sum_{k=1}^{d} z_{k} A_{k}\right) \widehat{x}(z)+\left(\sum_{k=1}^{d} z_{k} B_{k}\right) \widehat{u}(z) \\
\widehat{y}(z)=C \widehat{x}(z)+D \widehat{u}(z)
\end{array}\right.
$$

from which one finds that the input/inital state to output map is given by

$$
\begin{equation*}
\widehat{y}(z)=F(z) \widehat{u}(z)+W(z) x(0) \tag{2.5}
\end{equation*}
$$

where $F$ and $W$ are the transfer function and observation function given by

$$
\begin{align*}
F(z)= & D+C\left(I-z_{1} A_{1}-\cdots-z_{d} A_{d}\right)^{-1} \times \\
& \times\left(z_{1} B_{1}+\cdots+z_{d} B_{d}\right)  \tag{2.6}\\
W(z)= & C\left(I-z_{1} A_{1}-\cdots-z_{d} A_{d}\right)^{-1}
\end{align*}
$$

2.2. Noncommutative Fornasini-Marchesini systems. We write $\mathcal{F}_{d}$ for the free semigroup generated by the $d$-letter alphabet $\{1, \ldots, d\}$; elements of $\mathcal{F}_{d}$ are words of the form $\alpha=i_{N} \cdots i_{2} i_{1}$ with $i_{j} \in\{1, \ldots, d\}$ for $j=1, \ldots, N$. We also include the empty word $\emptyset$ that serves as the unit element with respect to the semigroup action: concatenation of words. With $|\alpha|$ we indicate the length of the word $\alpha$.

Now consider a system that evolves along the free semigroup $\mathcal{F}_{d}$ given by the equations

$$
\left\{\begin{align*}
x(k \cdot \alpha) & =A_{k} x(\alpha)+B_{k} u(\alpha), k=1, \ldots, d,  \tag{2.7}\\
y(\alpha) & =C x(\alpha)+D u(\alpha)
\end{align*}\right.
$$

where we have a system matrix of the form (2.2) that we assume to be contractive. For notational simplicity we assume $\mathcal{U}=\mathcal{Y}=\mathbb{C}$. Again, the output sequence $\{y(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$ and state sequence $\{x(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$ are completely determined by the input
sequence $\{u(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$ and an initial condition: $x(\emptyset)=x_{\emptyset}$, and in this case we have energy balance inequalities of the form

$$
\begin{aligned}
& \sum_{\alpha \in \mathcal{F}_{d} ;|\alpha| \leq N}\|y(\alpha)\|^{2}+ \sum_{\alpha \in \mathcal{F}_{d} ;|\alpha|=N}\|x(\alpha)\|^{2} \leq \\
& \leq \sum_{\alpha \in \mathcal{F}_{d} ;|\alpha| \leq N}\|u(\alpha)\|^{2}+\|x(\emptyset)\|^{2} \\
&
\end{aligned}
$$

for each $N \in \mathbb{Z}_{+}$. Hence an input sequence $u$ from $\ell^{2}\left(\mathcal{F}_{d}\right)$ produces an output sequence $y$ in $\ell^{2}\left(\mathcal{F}_{d}\right)$ so that $\|y\|^{2} \leq\|u\|^{2}+\|x(\emptyset)\|^{2}$, and the state sequence $x$ is bounded; more precisely $\sup _{N \in \mathbb{Z}_{+}} \sum_{|\alpha|=N}\|x(\alpha)\|^{2}<\infty$. In case the state operators $A_{1}, \ldots, A_{d}$ satisfy an additional stability the state sequence $x$ is in $\ell_{\mathcal{X}}^{2}\left(\mathcal{F}_{d}\right)$.

For this example we shall use a frequency domain consisting of operator-tuples rather than tuples of complex numbers. We therefore assume that we are given an additional Hilbert space $\mathcal{K}$ and a $d$-tuple $T=\left(T_{1}, \ldots, T_{d}\right)$ of operators on $\mathcal{K}$. Such a $d$-tuple $T$ is said to be a strict row contraction if the operator matrix $\left[\begin{array}{lll}T_{1} & \cdots & T_{d}\end{array}\right]$ defines a strict contraction from $\mathcal{K}^{d}$ into $\mathcal{K}$. For a strict row contraction $T$ we define the noncommutative functional calculus

$$
\begin{equation*}
T^{\alpha}=T_{i_{N}} \cdots T_{i_{1}} \text { if } \alpha=i_{N} \cdots i_{1}, \quad T^{\emptyset}=I_{\mathcal{K}} \tag{2.8}
\end{equation*}
$$

Then for $\mathcal{W}=\mathbb{C}, \mathcal{X}$ and $w=\{w(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$ a bounded $\mathcal{W}$-valued sequence we define a function $\widehat{w}$ on the set of strict row contractions via the noncommutative $Z$ transform: For each strict row contraction $T$ we set

$$
\begin{equation*}
\widehat{w}(T)=\sum_{\alpha \in \mathcal{F}_{d}} w(\alpha) \otimes T^{\alpha} \in \mathcal{L}(\mathcal{K}, \mathcal{W} \otimes \mathcal{K}) . \tag{2.9}
\end{equation*}
$$

Then the system (2.7) is described in the frequency domain by the equations

$$
\left\{\begin{align*}
& \widehat{x}(T)=x(\emptyset)+\left(\sum_{k=1}^{d} A_{k} \otimes T_{k}\right) \widehat{x}(T)+  \tag{2.10}\\
&+\left(\sum_{k=1}^{d} B_{k} \otimes T_{k}\right) \widehat{u}(T) \\
& \widehat{y}(T)=\left(C \otimes I_{\mathcal{K}}\right) \widehat{x}(T)+\left(D \otimes I_{\mathcal{K}}\right) \widehat{u}(T)
\end{align*}\right.
$$

and the input/intial state-to-output map is given by

$$
\begin{equation*}
\widehat{y}(T)=F(T) \widehat{u}(T)+W(T) x(\emptyset) \tag{2.11}
\end{equation*}
$$

where we define the transfer function $F$ and observation function $W$ on the set of strict row contractions $T$ by

$$
\begin{align*}
F(T)= & D+C\left(I-A_{1} \otimes T_{1}-\cdots-A_{d} \otimes T_{d}\right)^{-1} \times \\
& \times\left(B_{1} \otimes T_{1}+\cdots+B_{d} \otimes T_{d}\right) \\
= & D+\sum_{k=1}^{d} \sum_{\alpha \in \mathcal{F}_{d}} C A^{\alpha} B_{k} T^{\alpha \cdot k}  \tag{2.12}\\
W(T)= & C\left(I-A_{1} \otimes T_{1}-\cdots-A_{d} \otimes T_{d}\right)^{-1} \\
= & \sum_{\alpha \in \mathcal{F}_{d}}\left(C A^{\alpha}\right) \otimes T^{\alpha}
\end{align*}
$$

It is possible to symmetrize the noncommutative Fornasini-Marchesini system equations (2.7) to recover the commutative Fornasini-Marchesini system equations
(2.1). To describe this in detail, we introduce the abelianization map

$$
\begin{equation*}
\mathbf{a}(\alpha)=\left(n_{1}, \ldots, n_{d}\right) \quad\left(\alpha \in \mathcal{F}_{d}\right) \tag{2.13}
\end{equation*}
$$

where, for $j=1, \ldots, d$, the integer $n_{j}$ is equal to the number of times the letter $j$ appears in the word $\alpha$. Let $\left\{u(\alpha), x(\alpha), y(\alpha): \alpha \in \mathcal{F}_{d}\right\}$ be a system trajectory satisfying (2.7), for each $n \in \mathbb{Z}_{+}^{d}$ define $\bar{u}(n)=\sum_{\alpha: \mathbf{a}(\alpha)=n} u(\alpha), \bar{x}(n)=\sum_{\alpha: \mathbf{a}(\alpha)=n} x(\alpha)$, $\bar{y}(n)=\sum_{\alpha: \mathbf{a}(\alpha)=n} y(\alpha)$. Then one can check that the more detailed system equations (2.7) imply that the aggregate quantities $\left\{\bar{u}(n), \bar{x}(n), \bar{y}(n): n \in \mathbb{Z}_{+}^{d}\right\}$ satisfy the system equations (2.1). Moreover, in general if the $d$-tuple $T=\left(T_{1}, \ldots, T_{d}\right)$ of operators on $\mathcal{K}$ is commutative, then $T^{\alpha}=T^{\alpha^{\prime}}=: T^{n}$ whenever $\mathbf{a}(\alpha)=\mathbf{a}\left(\alpha^{\prime}\right)=$ $n \in \mathbb{Z}_{+}^{d}$ and the value of the $Z$-transform $\widehat{w}(T)$ in (2.9) can be expressed as $\widehat{w}(T)=\sum_{n \in \mathbb{Z}_{+}^{d}} \bar{w}(n) T^{n}$. This applies in particular to the case where we take $\mathcal{K}=\mathbb{C}$ to be one-dimensional; then $T=\left(T_{1}, \ldots, T_{d}\right)$ becomes a $d$-tuple of scalars $z=\left(z_{1}, \ldots, z_{d}\right) \in \mathbb{B}^{d}$, and the noncommutative $Z$-transform (2.9) collapses to the commutative $Z$-transform (2.4) (with $\bar{w}$ in place of $v$ ).

## 3. $C^{*}$-CORRESPONDENCE PRELIMINARIES

Let $\mathcal{A}$ and $\mathcal{B}$ be $C^{*}$-algebras and $E$ a linear space. We say that $E$ is an $(\mathcal{A}, \mathcal{B})$ correspondence when $E$ is a bi-module with respect to a right $\mathcal{B}$-action and a left $\mathcal{A}$-action, and $E$ is endowed with a $\mathcal{B}$-valued inner product $\langle,\rangle_{E}$ satisfying the following axioms: For any $\lambda, \mu \in \mathbb{C}, \xi, \eta, \zeta \in E, a \in \mathcal{A}$ and $b \in \mathcal{B}$, it is the case that
(1) $\langle\lambda \xi+\mu \zeta, \eta\rangle_{E}=\lambda\langle\xi, \eta\rangle_{E}+\mu\langle\zeta, \eta\rangle_{E} ;$
(2) $\langle\xi \cdot b, \eta\rangle_{E}=\langle\xi, \eta\rangle_{E} b$;
(3) $\langle a \cdot \xi, \eta\rangle_{E}=\left\langle\xi, a^{*} \cdot \eta\right\rangle_{E}$;
(4) $\langle\xi, \eta\rangle_{E}^{*}=\langle\eta, \xi\rangle_{E}$;
(5) $\langle\xi, \xi\rangle_{E} \geq 0$ (in $\mathcal{B}$ );
(6) $\langle\xi, \xi\rangle_{E}=0$ implies that $\xi=0$,
and that $E$ is a Banach space with respect to the norm

$$
\|\xi\|_{E}=\left\|\langle\xi, \xi\rangle_{E}\right\|_{\mathcal{B}}^{\frac{1}{2}} \quad(\xi \in E)
$$

where $\left\|\|_{\mathcal{B}}\right.$ denotes the norm of $\mathcal{B}$. We also impose that

$$
(\lambda \xi) \cdot b=\xi \cdot(\lambda b) \quad \text { and } \quad(\lambda a) \cdot \xi=a \cdot(\lambda \xi)
$$

for any $\lambda \in \mathbb{C}, a \in \mathcal{A}, b \in \mathcal{B}, \xi \in E$. In practice we usually write $\langle$,$\rangle and \|\|$ for the inner product and norm on $E$, and in case $\mathcal{A}=\mathcal{B}$ we say that $E$ is an $\mathcal{A}$-correspondence.

Example 3.1. Part A: For the commutative Fornasini-Marchensini system interpretation, we take $\mathcal{A}=\mathcal{L}(\mathcal{U})$ and $E=\operatorname{col}_{j=1}^{d}[\mathcal{L}(\mathcal{U})]=\mathcal{L}\left(\mathcal{U}, \mathcal{U}^{d}\right)$, where $\mathcal{U}$ is a given Hilbert space. Then $E$ can be seen as an $\mathcal{A}$-correspondence with left and right $\mathcal{A}$-action and inner product given by

$$
\begin{gathered}
a \cdot \operatorname{col}_{j=1}^{d}\left[T_{j}\right] \cdot a^{\prime}=\operatorname{col}_{j=1}^{d}\left[a T_{j} a^{\prime}\right] \\
\left\langle\operatorname{col}_{j=1}^{d}\left[T_{j}\right], \operatorname{col}_{j=1}^{d}\left[R_{j}\right]\right\rangle=\sum_{j=1}^{d} R_{j}^{*} T_{j} .
\end{gathered}
$$

Part B: For the noncommutative Fornasini-Marchensini system interpretation we take the same correspondence, but then with $\mathcal{U}=\mathbb{C}$.

Given two $(\mathcal{A}, \mathcal{B})$-correspondences $E$ and $F$, the set of bounded linear operators from $E$ to $F$ is denoted by $\mathcal{L}(E, F)$. It may happen that a $T \in \mathcal{L}(E, F)$ is not adjointable, i.e., it is not necessarily the case that

$$
\langle T \xi, \gamma\rangle=\left\langle\xi, T^{*} \gamma\right\rangle \quad(\xi \in E, \gamma \in F)
$$

for some $T^{*} \in \mathcal{L}(F, E)$. We use $\mathcal{L}^{a}(E, F)$ to indicate the set of operators in $\mathcal{L}(E, F)$ which are adjointable. As usual we have the abbreviations $\mathcal{L}(E)$ and $\mathcal{L}^{a}(E)$ in case $F=E$.

The third inner-product axiom implies that the left $\mathcal{A}$-action can be identified with a $*$-representation (meaning a nondegenerate $*$-homomorphism) $\varphi$ of $\mathcal{A}$ into the $C^{*}$-algebra $\mathcal{L}^{a}(E)$. In case this $*$-representation $\varphi$ is specified we will occasionally write $\varphi(a) \xi$ instead of $a \cdot \xi$.

Furthermore, an operator $T \in \mathcal{L}(E, F)$ is said to be a right $\mathcal{B}$-module map if

$$
T(\xi \cdot b)=T(\xi) \cdot b \quad(\xi \in E, b \in \mathcal{B})
$$

and a left $\mathcal{A}$-module map whenever

$$
T(a \cdot \xi)=a \cdot T(\xi) \quad(\xi \in E, a \in \mathcal{A})
$$

We will have a need for various constructions which create new correspondences out of given correspondences.

Given two $(\mathcal{A}, \mathcal{B})$-correspondences $E$ and $F$, we define the direct-sum correspondence $E \oplus F$ to be the direct-sum vector space $E \oplus F$ together with the left $\mathcal{A}$-action and right $\mathcal{B}$-action and the direct-sum $\mathcal{B}$-valued inner-product defined by setting for each $\xi, \xi^{\prime} \in E, \gamma, \gamma^{\prime} \in F, a \in \mathcal{A}$ and $b \in \mathcal{B}$ :

$$
\begin{gathered}
a \cdot(\xi \oplus \gamma)=(a \cdot \xi) \oplus(a \cdot \gamma) \\
(\xi \oplus \gamma) \cdot b=(\xi \cdot b) \oplus(\gamma \cdot b), \\
\left\langle\xi \oplus \gamma, \xi^{\prime} \oplus \gamma^{\prime}\right\rangle_{E \oplus F}=\left\langle\xi, \xi^{\prime}\right\rangle_{E}+\left\langle\gamma, \gamma^{\prime}\right\rangle_{F} .
\end{gathered}
$$

Bounded linear operators between direct-sum correspondences admit operator matrix decompositions in precisely the same way as in the Hilbert space case $(\mathcal{B}=\mathbb{C})$, while adjointability and the left and right module map property of such an operator corresponds to the operators in the decomposition being adjointable, or left or right module maps, respectively.

Now suppose we are given three $C^{*}$-algebras $\mathcal{A}, \mathcal{B}$ and $\mathcal{C}$, an $(\mathcal{A}, \mathcal{B})$-correspondence $E$ and a ( $\mathcal{B}, \mathcal{C}$ )-correspondence $F$. Then we define the tensor-product correspondence $E \otimes F$ to be the completion of the linear span of all tensors $\xi \otimes \gamma$ (with $\xi \in E$ and $\gamma \in F)$ subject to the identification

$$
\begin{equation*}
(\xi \cdot b) \otimes \gamma=\xi \otimes(b \cdot \gamma) \quad(\xi \in E, \gamma \in F, b \in \mathcal{B}) \tag{3.14}
\end{equation*}
$$

with left $\mathcal{A}$-action, right $\mathcal{C}$-action and the $\mathcal{C}$-valued inner-product defined, initially just for pure tensors, by setting for each $\xi, \xi^{\prime} \in E, \gamma, \gamma^{\prime} \in F, a \in \mathcal{A}$ and $c \in \mathcal{C}$ :

$$
\begin{gathered}
a \cdot(\xi \otimes \gamma)=(a \cdot \xi) \otimes \gamma, \quad(\xi \otimes \gamma) \cdot c=\xi \otimes(\gamma \cdot c), \\
\left\langle\xi \otimes \gamma, \xi^{\prime} \otimes \gamma^{\prime}\right\rangle_{E \otimes F}=\left\langle\left\langle\xi, \xi^{\prime}\right\rangle_{E} \cdot \gamma, \gamma^{\prime}\right\rangle_{F},
\end{gathered}
$$

and then extending by linearity and continuity to $E \otimes F$. In case the left action on $F$ is given by the $*$-representation $\varphi$ we occasionally emphasize this by writing $E \otimes_{\varphi} F$ for $E \otimes F$. It is more complicated to characterize the bounded linear operators between tensor-product correspondences. One way to construct such operators is as follows. Let $E$ and $E^{\prime}$ be $(\mathcal{A}, \mathcal{B})$-correspondences and $F$ and $F^{\prime}$ $(\mathcal{B}, \mathcal{C})$-correspondences. Furthermore, let $X \in \mathcal{L}\left(E, E^{\prime}\right)$ be a right module map and
$Y \in \mathcal{L}\left(F, F^{\prime}\right)$ a left module map. We then write $X \otimes Y$ for the tensor product operator in $\mathcal{L}\left(E \otimes F, E^{\prime} \otimes F^{\prime}\right)$ determined by

$$
\begin{equation*}
X \otimes Y(\xi \otimes \gamma)=(X \xi) \otimes(Y \gamma) \quad(\xi \otimes \gamma \in E \otimes F) \tag{3.15}
\end{equation*}
$$

If, in addition, $X$ is also a left module map, then $X \otimes Y$ is a left module map, while $Y$ also being a right module map guarantees that $X \otimes Y$ is a right module map. Moreover, $X \otimes Y$ is adjointable when $X$ and $Y$ are both adjointable, and in that case $(X \otimes Y)^{*}=X^{*} \otimes Y^{*}$. Notice that the left action on $E \otimes F$ can now be written as $a \mapsto \varphi(a) \otimes I_{F} \in \mathcal{L}^{a}(E \otimes F)$, where $I_{F}$ denotes the identity operator on $F$ and $\varphi$ indicates the left $\mathcal{A}$-action on $E$.

## 4. Correspondence-representation pairs and their duals

In the remainder of the paper we consider the situation where $\mathcal{A}=\mathcal{B}$, i.e, $E$ is an $\mathcal{A}$-correspondence. We also restrict our attention to the case where $\mathcal{A}$ is a von Neumann algebra and $E$ an $\mathcal{A}-W^{*}$-correspondence. This means that $E$ is an $\mathcal{A}$-correspondence which is also self-dual in the sense that any right $\mathcal{A}$-module map $\rho: E \rightarrow \mathcal{A}$ is given by taking the inner-product against some element $e_{\rho}$ of $E$ :

$$
\begin{equation*}
\rho(e)=\left\langle e, e_{\rho}\right\rangle_{E} \in \mathcal{A} . \tag{4.16}
\end{equation*}
$$

Suppose that we are also given an auxiliary Hilbert space $\mathcal{E}$ and a *-representation $\sigma: \mathcal{A} \rightarrow \mathcal{L}(\mathcal{E})$; as this will be the setting for much of the analysis to follow, we refer to such a pair $(E, \sigma)$ as a Correspondence-Representation pair, or $C R$-pair for short. We further assume that $\sigma$ is faithful (injective) and normal ( $\sigma$-weakly continuous). The Hilbert space $\mathcal{E}$ becomes an $(\mathcal{A}, \mathbb{C})$-correspondence with left $\mathcal{A}$-action given by $\sigma$ :

$$
a \cdot y=\sigma(a) y \text { for all } a \in \mathcal{A} \text { and } y \in \mathcal{E}
$$

Thus we can form the tensor-product $(\mathcal{A}, \mathbb{C})$-correspondence $E \otimes_{\sigma} \mathcal{E}$. We then write $E^{\sigma}$ for the set of bounded linear operators $\mu: \mathcal{E} \rightarrow E \otimes_{\sigma} \mathcal{E}$ which are also left $\mathcal{A}$-module maps:

$$
\begin{equation*}
E^{\sigma}=\left\{\mu: \mathcal{E} \rightarrow E \otimes_{\sigma} \mathcal{E}: \mu \sigma(a)=\left(\varphi(a) \otimes I_{\mathcal{E}}\right) \mu\right\} . \tag{4.17}
\end{equation*}
$$

It turns out that $E^{\sigma}$ is itself a $W^{*}$-correspondence, not over $\mathcal{A}$ but over the $W^{*}$ algebra

$$
\sigma(\mathcal{A})^{\prime}=\{b \in \mathcal{L}(\mathcal{E}): b \sigma(a)=\sigma(a) b \text { for all } a \in \mathcal{A}\} \subset \mathcal{L}(\mathcal{E})
$$

(the commutant of the image $\sigma(\mathcal{A})$ of the $*$-representation $\sigma$ in $\mathcal{L}(\mathcal{E})$ ) with left and right $\sigma(\mathcal{A})^{\prime}$-action and $\sigma(\mathcal{A})^{\prime}$-valued inner-product $\langle\cdot, \cdot\rangle_{E^{\sigma}}$ given by

$$
\mu \cdot b=\mu b, \quad b \cdot \mu=\left(I_{E} \otimes b\right) \mu, \quad\langle\mu, \nu\rangle_{E^{\sigma}}=\nu^{*} \mu
$$

Let $\iota: \sigma(\mathcal{A})^{\prime} \rightarrow \mathcal{L}(\mathcal{E})$ be the identity $*$-representation: $\iota(b)=b$ for each $b \in \sigma(\mathcal{A})^{\prime}$. Then $\left(E^{\sigma}, \iota\right)$ is a CR-pair that is said to be dual to $(E, \sigma)$; see Section 3 in 11 for further details. In particular, $\mathcal{E}$ can be seen both as an $(\mathcal{A}, \mathbb{C})$-correspondence (with left action given by $\sigma$ ) and a $\left(\sigma(\mathcal{A})^{\prime}, \mathbb{C}\right)$-correspondence (with left action given by $\iota$ ).

Example 4.1 (Continuation of Example 3.1).
Part A: We take $\mathcal{E}=\mathcal{U}$ and for the $*$-representation $\sigma$ the identity map: $\sigma(T)=T$ for each $T \in \mathcal{L}(\mathcal{U})$. Then $E \otimes \mathcal{E}=\mathcal{U}^{d}, \sigma(\mathcal{A})^{\prime}=\mathbb{C} I_{\mathcal{U}}:=\left\{\mu I_{\mathcal{U}}: \mu \in \mathbb{C}\right\}$ and

$$
E^{\sigma}=\left\{\operatorname{col}_{j=1}^{d}\left[\mu_{j} I_{\mathcal{U}}\right]: \mu_{j} \in \mathbb{C}\right\} \subset \mathcal{L}\left(\mathcal{U}, \mathcal{U}^{d}\right)
$$

The spaces $\sigma(\mathcal{A})^{\prime}$ and $E^{\sigma}$ shall be identified with $\mathbb{C}$ and $\mathbb{C}^{d}$, respectively.
Part B: Let $\mathcal{K}$ be some Hilbert space. We then take $\mathcal{E}=\mathcal{K}$ and for $\sigma: \mathbb{C} \rightarrow \mathcal{L}(\mathcal{K})$ the $*$-representation $\sigma(\lambda)=\lambda I_{\mathcal{K}}$ for each $\lambda \in \mathbb{C}$. Then $E \otimes \mathcal{E}=\mathcal{K}^{d}, \sigma(\mathcal{A})^{\prime}=\mathcal{L}(\mathcal{K})$ and

$$
E^{\sigma}=\left\{\operatorname{col}_{j=1}^{d}\left[T_{j}\right]: T_{j} \in \mathcal{L}(\mathcal{K})\right\}=\mathcal{L}\left(\mathcal{K}, \mathcal{K}^{d}\right)
$$

In these two examples the duality mentioned above takes the following form. In Part A we start with $E=\mathcal{L}\left(\mathcal{U}, \mathcal{U}^{d}\right), \mathcal{A}=\mathcal{L}(\mathcal{U})$ and $\sigma$ the identity map on $\mathcal{L}(\mathcal{U})$. We then find the dual CR-pair $\left(E^{\sigma}, \iota\right)$ with $E^{\sigma}=\mathbb{C}^{d}, \sigma(\mathcal{A})^{\prime}=\mathbb{C}$ (after the identification) and $\iota(\lambda)=\lambda I_{\mathcal{U}}$. This is precisely the CR-pair we use in Part 2 (if we take $\mathcal{K}=\mathcal{U})$. It then follows from Part 2 that forming the CR-pair dual to ( $\left.E^{\sigma}, \iota\right)$ we return to $(E, \sigma)$. In this sense the two examples are also each others dual.

Since $E$ is an $\mathcal{A}$-correspondence, we can inductively define the self-tensor-product $\mathcal{A}$-correspondences $E^{\otimes n}$ via $E^{\otimes 1}=1$ and $E^{\otimes n+1}=E \otimes\left(E^{\otimes n}\right)$ for $n=1,2, \ldots$. Formally we set $E^{\otimes 0}=\mathcal{A}$. Using the notation in (3.15), the left $\mathcal{A}$-action on $E^{\otimes n}$ can be written as $\varphi(a) \otimes I_{E^{\otimes n-1}}$ if we use $\varphi: \mathcal{A} \rightarrow \mathcal{L}^{a}(E)$ to indicate the left $\mathcal{A}$ action of $E$. We may then also form the $(\mathcal{A}, \mathbb{C})$-correspondences $E^{\otimes n} \otimes_{\sigma} \mathcal{E}$. Besides the left $\mathcal{A}$-action, these Hilbert spaces also admit a left $\sigma(\mathcal{A})^{\prime}$-action via the $*$ representation $b \mapsto I_{E^{\otimes n}} \otimes b$ for $b \in \sigma(\mathcal{A})^{\prime}$. Furthermore, for a $\mu \in E^{\sigma}$ we can form the tensor product operator $I_{E^{\otimes n}} \otimes \mu$ that maps $E^{\otimes n} \otimes \mathcal{E}$ into $E^{\otimes n+1} \otimes \mathcal{E}$. In a similar way we define the $\sigma(\mathcal{A})^{\prime}$-correspondences $\left(E^{\sigma}\right)^{\otimes n}$ and $\left(\sigma(\mathcal{A})^{\prime}, \mathbb{C}\right)$-correspondences $\left(E^{\sigma}\right)^{\otimes n} \otimes_{\iota} \mathcal{E}$ for $n \in \mathbb{Z}_{+}$.

One of the consequences of the duality between $(E, \sigma)$ and $\left(E^{\sigma}, \iota\right)$ is the following result from [11.
Proposition 4.2. For each $n \in \mathbb{Z}_{+}$the map $\Phi_{n}:\left(E^{\sigma}\right)^{\otimes n} \otimes_{\iota} \mathcal{E} \rightarrow E^{\otimes n} \otimes_{\sigma} \mathcal{E}$ induced by the identity

$$
\Phi_{n}\left(\mu_{n} \otimes \cdots \otimes \mu_{1} \otimes e\right)=\left(I_{E \otimes n-1} \otimes \mu_{n}\right) \cdots\left(I_{E} \otimes \mu_{2}\right) \mu_{1} e
$$

is a unitary left $\sigma(\mathcal{A})^{\prime}$-module map.
Example 4.3 (Continuation of Example 4.1).
Part A: For $n \in \mathbb{Z}_{+}$set $d_{n}=\#\left\{\alpha \in \mathcal{F}_{d},|\alpha|=n\right\}=d^{n}$. Then

$$
E^{\otimes n}=\mathcal{L}\left(\mathcal{U}, \mathcal{U}^{d_{n}}\right) \quad \text { and } \quad E^{\otimes n} \otimes_{\sigma} \mathcal{E}=\mathcal{U}^{d_{n}}
$$

while

$$
\left(E^{\sigma}\right)^{\otimes n}=\mathcal{L}\left(\mathbb{C}, \mathbb{C}^{d_{n}}\right)=\mathbb{C}^{d_{n}} \quad \text { and } \quad\left(E^{\sigma}\right)^{\otimes n} \otimes_{\iota} \mathcal{E}=\mathcal{U}^{d_{n}}
$$

In particular, $E^{\otimes n} \otimes_{\sigma} \mathcal{E}$ and $\left(E^{\sigma}\right)^{\otimes n} \otimes_{\iota} \mathcal{E}$ are the same Hilbert space, and moreover, the left $\sigma(\mathcal{A})^{\prime}=\mathbb{C}$-actions on both spaces is just scalar multiplication. The left $\mathcal{A}=\mathcal{L}(\mathcal{U})$-action on $\mathcal{U}^{d_{n}}$ is given by

$$
\begin{equation*}
T\left(\oplus_{j=1}^{d_{n}} e_{j}\right)=\oplus_{j=1}^{d_{n}} T e_{j} \tag{4.18}
\end{equation*}
$$

for $T \in \mathcal{L}(\mathcal{U})$ and $e_{j} \in \mathcal{U}$. Moreover, the map $\Phi_{n}$ defined in Proposition 4.2 is the identity map on $\mathcal{U}^{d_{n}}$.
Part B: Due to the duality between the two examples, as observed in Example 4.1, we find:

$$
\begin{gathered}
E^{\otimes n}=\mathcal{L}\left(\mathbb{C}, \mathbb{C}^{d_{n}}\right)=\mathbb{C}^{d_{n}} \quad \text { and } \quad E^{\otimes n} \otimes_{\sigma} \mathcal{E}=\mathcal{K}^{d_{n}} \\
\left(E^{\sigma}\right)^{\otimes n}=\mathcal{L}\left(\mathcal{K}, \mathcal{K}^{d_{n}}\right) \quad \text { and } \quad\left(E^{\sigma}\right)^{\otimes n} \otimes_{\iota} \mathcal{E}=\mathcal{K}^{d_{n}}
\end{gathered}
$$

with left $\sigma(\mathcal{A})^{\prime}=\mathcal{L}(\mathcal{K})$-action on $\mathcal{K}^{d_{n}}$ given by (4.18), with $T \in \mathcal{L}(\mathcal{K}), e_{j} \in \mathcal{K}$. Moreover, $\Phi_{n}$ is again the identity map on $\mathcal{K}^{d_{n}}$.

## 5. Linear dissipative discrete-time systems

Let $(E, \sigma)$ be a CR-pair as in Section 4. A (linear) system associated with the pair $(E, \sigma)$ is a set $\{A, B, C, D, \pi\}$ consisting of four left $\sigma(\mathcal{A})^{\prime}$-module maps:

$$
\left[\begin{array}{ll}
A & B  \tag{5.19}\\
C & D
\end{array}\right]:\left[\begin{array}{l}
\mathcal{X} \\
\mathcal{E}
\end{array}\right] \rightarrow\left[\begin{array}{c}
E^{\sigma} \otimes \mathcal{X} \\
\mathcal{E}
\end{array}\right]
$$

and a $*$-representation $\pi: \sigma(\mathcal{A})^{\prime} \rightarrow \mathcal{L}(\mathcal{X})$ that provides a left $\sigma(\mathcal{A})^{\prime}$-action on the Hilbert space $\mathcal{X}$, called the state space, making $\mathcal{X}$ into a $\left(\sigma(\mathcal{A})^{\prime}, \mathbb{C}\right)$-correspondence. In this paper we are primarily interested in the case that the system matrix (5.19) is contractive, in which case we speak of a dissipative system. For simplicity we only consider here the case that the input space and output space are the same, i.e., both equal to $\mathcal{E}$; various techniques exist to extend to the case with different input and output spaces, cf., [14, 4].

Given a system $\{A, B, C, D, \pi\}$, for each $n \in \mathbb{Z}_{+}$we define the operators:

$$
\begin{gathered}
A_{n}=I_{\left(E^{\sigma}\right)^{\otimes n}} \otimes A, \quad B_{n}=\left(I_{\left(E^{\sigma}\right)^{\otimes n}} \otimes B\right) \Phi_{n}^{*} \\
C_{n}=\Phi_{n}\left(I_{\left(E^{\sigma}\right)^{\otimes n}} \otimes C\right), \quad D_{n}=\Phi_{n}\left(I_{\left(E^{\sigma}\right)^{\otimes n}} \otimes D\right) \Phi_{n}^{*}
\end{gathered}
$$

where $\Phi_{n}$ is as defined in Proposition 4.2, and associate with $\{A, B, C, D, \pi\}$ the following linear system of equations:

$$
\Sigma=\left\{\begin{align*}
x(n+1) & =A_{n} x(n)+B_{n} u(n),  \tag{5.20}\\
y(n) & =C_{n} x(n)+D_{n} u(n),
\end{align*} \quad\left(n \in \mathbb{Z}_{+}\right)\right.
$$

where $u(n), y(n) \in E^{\otimes n} \otimes \mathcal{E}$, and $x(n) \in\left(E^{\sigma}\right)^{\otimes n} \otimes \mathcal{X}$. Note that an input sequence $\{u(n)\}_{n \in \mathbb{Z}_{+}}$and initial state $x(0)$ uniquely determine the output sequence $\{y(n)\}_{n \in \mathbb{Z}_{+}}$and state sequence $\{x(n)\}_{n \in \mathbb{Z}_{+}}$.

Example 5.1 (Continuation of Example 4.3).
Part A: Let $\{A, B, C, D, \pi\}$ be a dissipative system associated with the CR-pair of Example 4.1, Part A. Then $\pi: \mathbb{C} \rightarrow \mathcal{L}(\mathcal{X})$ must be equal to $\pi(\lambda)=\lambda I_{\mathcal{X}}$, and thus $E^{\pi} \otimes \mathcal{X}=\mathcal{X}^{d}$. The system then has the form

$$
\left[\begin{array}{cc}
A & B  \tag{5.21}\\
C & D
\end{array}\right]=\left[\begin{array}{c|c}
A_{1} & B_{1} \\
\vdots & \vdots \\
A_{d} & B_{d} \\
\hline C & D
\end{array}\right]:\left[\begin{array}{l}
\mathcal{X} \\
\mathcal{U}
\end{array}\right] \rightarrow\left[\begin{array}{c}
\mathcal{X}^{d} \\
\mathcal{U}
\end{array}\right]
$$

and, since $\sigma(\mathcal{A})^{\prime}=\mathbb{C}$, the $\sigma(\mathcal{A})^{\prime}$-module map properties are satisfied trivially. The input and output sequences $\mathbf{u}=\{\widetilde{u}(n)\}_{n \in \mathbb{Z}_{+}}$and $\mathbf{y}=\{\widetilde{y}(n)\}_{n \in \mathbb{Z}_{+}}$have elements $\widetilde{u}(n)$ and $\widetilde{y}(n)$ in $\mathcal{U}^{k_{n}}$, while the elements of the state sequence $\mathbf{x}=\{\widetilde{x}(n)\}_{n \in \mathbb{Z}_{+}}$are in $\mathcal{X}^{k_{n}}$. The system equations take the form

$$
\Sigma:=\left\{\begin{align*}
\widetilde{x}(n+1) & =\tilde{A}_{n} \widetilde{x}(n)+\tilde{B}_{n} \widetilde{u}(n),  \tag{5.22}\\
\widetilde{y}(n) & =\tilde{C}_{n} \widetilde{x}(n)+\tilde{D}_{n} \widetilde{u}(n),
\end{align*} \quad\left(n \in \mathbb{Z}_{+}\right)\right.
$$

with

$$
\begin{aligned}
& \tilde{A}_{n}=\operatorname{diag}_{i=1}^{d_{n}}\left(\left[\begin{array}{c}
A_{1} \\
\vdots \\
A_{d}
\end{array}\right]\right), \quad \tilde{B}_{n}=\operatorname{diag}_{i=1}^{d_{n}}\left(\left[\begin{array}{c}
B_{1} \\
\vdots \\
B_{d}
\end{array}\right]\right), \\
& \tilde{C}_{n}=\operatorname{diag}_{i=1}^{d_{n}}(C), \quad \tilde{D}_{n}=\operatorname{diag}_{i=1}^{d_{n}}(D),
\end{aligned}
$$

where diag forms a block diagonal operator, i.e., $\tilde{A}_{n} \in \mathcal{L}\left(\mathcal{X}^{d_{n}}, \mathcal{X}^{d_{n+1}}\right), \tilde{B}_{n} \in$ $\mathcal{L}\left(\mathcal{U}^{d_{n}}, \mathcal{X}^{d_{n+1}}\right)$, etc. After making the identification of an element $\widetilde{u}(n)$ in $\mathcal{U}^{d_{n}}$ with a set of elements of $\mathcal{U}$ of the form $\left\{u(\alpha): \alpha \in \mathcal{F}_{d},|\alpha|=n\right\}$ via the bijective map

$$
\begin{gather*}
\nu_{n}:\left\{\alpha \in \mathcal{F}_{d}:|\alpha|=n\right\} \rightarrow\left\{1, \ldots, d_{n}\right\},  \tag{5.23}\\
\nu_{n}(\alpha)=1+\sum_{j=1}^{n}\left(i_{j}-1\right) d^{j-1} \quad \text { if } \quad \alpha=i_{n} \cdots i_{1}
\end{gather*}
$$

and similarly for $\mathcal{X}^{d_{n}}$, and after unpacking the system equations (5.22), we arrive at the noncommutative system

$$
\left\{\begin{align*}
x(k \cdot \alpha) & =A_{k} x(\alpha)+B_{k} u(\alpha), k=1, \ldots, d  \tag{5.24}\\
y(\alpha) & =C x(\alpha)+D u(\alpha)
\end{align*}\right.
$$

that evolves over $\mathcal{F}_{d}$. We obtain the commutative system (2.1) after applying the abelianization procedure.
Part B: In this case the system matrix of a dissipative system $\{A, B, C, D, \pi\}$ associated with the CR-pair of Example 4.1, Part B, is as in (5.21), but with $\mathcal{U}=\mathcal{K}$, and $\pi$ is a $*$-representation from $\mathcal{L}(\mathcal{K})$ into $\mathcal{L}(\mathcal{X})$. Moreover, the operators in $\{A, B, C, D, \pi\}$ satisfy the following intertwining relations:

$$
\begin{array}{cl}
\pi(T) A_{j}=A_{j} \pi(T), & \pi(T) B_{j}=B_{j} T \\
T C=C \pi(T), & T D=D T \tag{5.25}
\end{array}
$$

for any $T \in \mathcal{L}(\mathcal{K})$ and $j=1, \ldots, d$. In particular, $D$ is a scalar multiple of $I_{\mathcal{K}}$. Assume that $\pi$ is a scalar multiple of the identity representation; i.e., there exists a Hilbert space $\overline{\mathcal{X}}$ so that $\mathcal{X}=\overline{\mathcal{X}} \otimes \mathcal{K}$ and $\pi(T)=I_{\overline{\mathcal{X}}} \otimes T$ for each $T \in \mathcal{L}(\mathcal{K})$. In fact, if $\mathcal{K}$ is finite dimensional, then any representation $\pi$ has this form (see [2]). Then (5.25) simplifies to

$$
\begin{equation*}
A_{j}=\bar{A}_{j} \otimes I_{\mathcal{K}}, \quad B_{j}=\bar{B}_{j} \otimes I_{\mathcal{K}}, \quad C=\bar{C} \otimes I_{\mathcal{K}} \tag{5.26}
\end{equation*}
$$

with $\bar{A}_{j} \in \mathcal{L}(\overline{\mathcal{X}}), \bar{B}_{j} \in \mathcal{L}(\mathbb{C}, \overline{\mathcal{X}})$ and $\bar{C} \in \mathcal{L}(\overline{\mathcal{X}}, \mathbb{C})$ for $j=1, \ldots, d$. Due to this special form of the system operators, we see that, if the input signal $\{u(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$, identified with the actual input signal $\{\widetilde{u}(n)\}_{n \in \mathbb{Z}_{+}}$as in Part A, and initial state $x(0)$ are restricted to have the form $\left\{\bar{u}(\alpha) k_{0}\right\}_{\alpha \in \mathcal{F}_{d}}$ and $\bar{x}(0) \otimes k_{0}$, respectively, where $\bar{u}(\alpha)$ is scalar-valued, $\bar{x}(0) \in \overline{\mathcal{X}}$ and $k_{0}$ is a fixed vector in $\mathcal{K}$, then the output sequence $\{y(\alpha)\}_{\alpha \in \mathcal{F}_{d}}=\left\{\bar{y}(\alpha) k_{0}\right\}_{\alpha \in \mathcal{F}_{d}}$ will have the same form as the input sequence, i.e., $\bar{y}(\alpha) \in \mathbb{C}$, while the state sequence is of the form $\{x(\alpha)\}_{\alpha \in \mathcal{F}_{d}}=\left\{\bar{x}(\alpha) \otimes k_{0}\right\}_{\alpha \in \mathcal{F}_{d}}$ with $\bar{x}(\alpha) \in \overline{\mathcal{X}}$. From now on, we shall restrict the inputs and initial states to vectors of this form for a fixed vector $k_{0} \in \mathcal{K}$. In that case, the transformation to the noncommutative system equations occurs in the same way as in Part A, and we arrive at (2.7) with $A_{j}, B_{j}$ and $C$ replaced with $\bar{A}_{j}, \bar{B}_{j}$ and $\bar{C}$.

We now define $\mathcal{F}^{2}(E)=\oplus_{n=0}^{\infty} E^{\otimes n}$, which is an $\mathcal{A}$-correspondence with left $\mathcal{A}$ action given by the $*$-representation $\varphi_{\infty}: \mathcal{A} \rightarrow \mathcal{L}^{a}\left(\mathcal{F}^{2}(E)\right)$ that is defined via

$$
\varphi_{\infty}(a)=\operatorname{diag}_{n \in \mathbb{Z}_{+}}\left(\varphi(a) \otimes I_{E^{\otimes n-1}}\right)
$$

where $\varphi$ is used to indicate the $*$-representation that provides the left $\mathcal{A}$-action on $E$. Then we can form the tensor product space $\mathcal{F}^{2}(E, \sigma)=\mathcal{F}^{2}(E) \otimes_{\sigma} \mathcal{E}$, which in turn is an $(\mathcal{A}, \mathbb{C})$-correspondence. We refer to $\mathcal{F}^{2}(E)$ and $\mathcal{F}^{2}(E, \sigma)$ as the Fock spaces associated with the correspondence $E$ and the CR-pair $(E, \sigma)$, respectively.

Besides the left $\mathcal{A}$-action, the Hilbert space $\mathcal{F}^{2}(E, \sigma)$ is also equipped with a left $\sigma(\mathcal{A})^{\prime}$-action that is given by the $*$-representation $b \mapsto I_{\mathcal{F}^{2}(E)} \otimes b \in \mathcal{L}\left(\mathcal{F}^{2}(E, \sigma)\right)$ for $b \in \sigma(\mathcal{A})^{\prime}$.

In a similar way as above, we define the Fock spaces $\mathcal{F}^{2}\left(E^{\sigma}\right)=\oplus_{n=0}^{\infty}\left(E^{\sigma}\right)^{\otimes n}$ and $\mathcal{F}^{2}\left(E^{\sigma}, \pi\right)=\mathcal{F}^{2}\left(E^{\sigma}\right) \otimes_{\pi} \mathcal{X}$. Moreover, with $\mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)$ we denote the space of bounded sequences $\{x(n)\}_{n \in \mathbb{Z}_{+}}$with $x(n) \in\left(E^{\sigma}\right)^{\otimes n} \otimes_{\pi} \mathcal{X}$.

Example 5.2. (Continuation of Example 5.1)
Part A: As observed in Example 4.3, $E^{\otimes n} \otimes_{\sigma} \mathcal{E}=\mathcal{U}^{d_{n}}$, which, via the map $\nu_{n}$ in (5.23), can be identified with $\oplus_{\alpha \in \mathcal{F}_{d},|\alpha|=n} \mathcal{U}$. It thus follows that $\mathcal{F}^{2}(E, \sigma)$ can be identified with $\ell_{\mathcal{U}}^{2}\left(\mathcal{F}_{d}\right)$, and similarly $\mathcal{F}^{2}\left(E^{\sigma}, \pi\right)$ with $\ell_{\mathcal{X}}^{2}\left(\mathcal{F}_{d}\right)$. For both spaces the left $\mathbb{C}$-action is just scalar multiplication. Via the same identification, $\mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)$ can be seen as the subspace of $\ell_{\mathcal{X}}^{\infty}\left(\mathcal{F}_{d}\right)$ consisting of elements $\{x(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$ from $\ell_{\mathcal{X}}^{\infty}\left(\mathcal{F}_{d}\right)$ such that $\sup _{N \in \mathbb{Z}_{+}} \sum_{|\alpha|=N}\|x(\alpha)\|^{2}<\infty$.
Part B: Similarly as in part A, $\mathcal{F}^{2}(E, \sigma)$ and $\mathcal{F}^{2}\left(E^{\sigma}, \pi\right)$ can be identified with $\ell_{\mathcal{K}}^{2}\left(\mathcal{F}^{d}\right)$ and $\ell_{\mathcal{X}}^{2}\left(\mathcal{F}_{d}\right)$, respectively, and $\mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)$ with a subspace of $\ell_{\mathcal{X}}^{\infty}\left(\mathcal{F}_{d}\right)$. However, in this case we have a left $\mathcal{L}(\mathcal{K})$-action that is given by entry-wise multiplication with $T \in \mathcal{L}(\mathcal{K})$ (resp. $\pi(T)$ ) on the left.
From here on we shall identify the spaces $\mathcal{F}^{2}(E, \sigma), \mathcal{F}^{2}\left(E^{\sigma}, \pi\right)$ and $\mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)$, in both parts A and B , with the appropriate $\ell^{2}$ and $\ell^{\infty}$ spaces over $\mathcal{F}_{d}$.

With the state operator $A$ and two nonnegative integers $n, m \in \mathbb{Z}_{+}$we associate the following generalized power:

$$
\left(A_{n}\right)^{m}=\left(I_{\left(E^{\sigma}\right)^{\otimes m-1}} \otimes A_{n}\right) \cdots\left(I_{E^{\sigma}} \otimes A_{n}\right) A_{n}
$$

which defines an operator in $\mathcal{L}\left(\left(E^{\sigma}\right)^{\otimes n} \otimes \mathcal{E},\left(E^{\sigma}\right)^{\otimes n+m} \otimes \mathcal{E}\right)$. Here $A_{n}^{0}$ is to be understood as $I_{\left(E^{\sigma}\right)^{\otimes n} \otimes \mathcal{E}}$. Note that $\left(A_{n}\right)^{m}=I_{\left(E^{\sigma}\right)^{\otimes n}} \otimes A_{1}^{m}=\left(A^{m}\right)_{n}$. Thus we can unambiguously write $A_{n}^{m}$ instead of $\left(A_{n}\right)^{m}$ or $\left(A^{m}\right)_{n}$.

Theorem 5.3. Let $y=\{y(n)\}_{n \in \mathbb{Z}_{+}}$and $x=\{x(n)\}_{n \in \mathbb{Z}_{+}}$be the output and state sequences generated by the dissipative system (5.20) with input sequence $u=\{u(n)\}_{n \in \mathbb{Z}_{+}}$and initial state $x(0)$. Assume that $u \in \mathcal{F}^{2}(E, \sigma)$. Then $y \in$ $\mathcal{F}^{2}(E, \sigma), x \in \mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)$ and

$$
y=T_{\Sigma} u+W_{\Sigma} x(0)
$$

with $T_{\Sigma} \in \mathcal{L}\left(\mathcal{F}^{2}(E, \sigma)\right)$ and $W_{\Sigma} \in \mathcal{L}\left(\mathcal{X}, \mathcal{F}^{2}(E, \sigma)\right)$ the left $\sigma(\mathcal{A})^{\prime}$-module maps given by

$$
T_{\Sigma}=\left[\begin{array}{cccc}
D_{0} & 0 & 0 & \cdots \\
C_{1} A_{1}^{0} B_{0} & D_{1} & 0 & \cdots \\
C_{2} A_{1}^{1} B_{0} & C_{2} A_{2}^{0} B_{1} & D_{2} & \cdots \\
\vdots & \vdots & \vdots & \ddots
\end{array}\right] \quad W_{\Sigma}=\left[\begin{array}{c}
C_{0} A_{0}^{0} \\
C_{1} A_{0}^{1} \\
C_{2} A_{0}^{2} \\
\vdots
\end{array}\right] .
$$

Moreover, the operator matrix [ $\left.\begin{array}{cc}T_{\Sigma} & W_{\Sigma}\end{array}\right]$ is contractive, and thus $\|y\|^{2} \leq\|u\|^{2}+$ $\|x(0)\|^{2}$. Finally, if, in addition, $\|A\|<1$, then $x \in \mathcal{F}^{2}\left(E^{\sigma}, \pi\right)$.

The fact that $T_{\Sigma}$ is a $\sigma(\mathcal{A})^{\prime}$-module map means that it commutes with the operators $I_{\mathcal{F}^{2}(E)} \otimes b$ for each $b \in \sigma(\mathcal{A})^{\prime}$. In addition, it is not hard to show (see
the proof of Theorem 3.5 in [14]) that $T_{\Sigma}$ also commutes with $I_{\mathcal{F}^{2}(E)} \otimes \mu$ for each $\mu \in E^{\sigma}$. In order to make sense of $I_{\mathcal{F}^{2}(E)} \otimes \mu$ as an operator on $\mathcal{F}^{2}(E, \sigma)$, because $\mu$ maps $\mathcal{E}$ into $E \otimes \mathcal{E}$, we have to make the identification $\mathcal{F}^{2}(E) \otimes E=\mathcal{F}^{2}(E)$. These commutative relations are one way of defining the analytic Toeplitz algebra $\mathcal{T}_{+}(E, \sigma)$ associated with the CR-pair $(E, \sigma)$ of [10]. There are various other ways to characterize $\mathcal{T}_{+}(E, \sigma)$; originally in [10] the space was generated by a collection of creation operators, while in [8] a Toeplitz structure characterization is given. The definition used here stems from 11.

## 6. Point-evaluation

Given a dissipative system $\{A, B, C, D, \pi\}$ as in (5.19), there is a notion of point evaluation giving rise to a $Z$-transform enabling the passage from the time domain to the frequency domain for this abstract setting. In the $W^{*}$-correspondence setting studied here, points are pairs $(\eta, b)$ with $\eta$ an element of the generalized disk

$$
\mathbb{D}\left(E^{\sigma *}\right):=\left\{\eta: \eta^{*} \in E^{\sigma},\|\eta\|<1\right\} .
$$

and $b \in \sigma(\mathcal{A})^{\prime}$. With a sequence $f=\{f(n)\}_{n \in \mathbb{Z}_{+}} \in \mathcal{F}^{2}(E, \sigma)$ (so $f(n) \in E^{\otimes n} \otimes \mathcal{E}$ ), we associate a function $f^{\wedge}$ on $\mathbb{D}\left(E^{\sigma *}\right) \times \sigma(\mathcal{A})^{\prime}$ with values in $\mathcal{E}$ defined by

$$
f^{\wedge}(\eta, b)=\sum_{n=0}^{\infty} \eta^{n}\left(I_{E \otimes n} \otimes b\right) f_{n} \quad\left((\eta, b) \in \mathbb{D}\left(E^{\sigma *}\right) \times \sigma(\mathcal{A})^{\prime}\right)
$$

Here we use $\eta^{n}$ to indicate the $n^{\text {th }}$ generalized power of $\eta$ :

$$
\eta^{n}=\eta\left(I_{E} \otimes \eta\right) \cdots\left(I_{E^{\otimes n-1}} \otimes \eta\right): E^{\otimes n} \otimes \mathcal{E} \rightarrow \mathcal{E}
$$

where for $n=0$ we set $\eta^{0}=I_{\mathcal{E}}$. Note that the generalized powers $\eta^{n}$ are well defined because $\eta$ is a left $\mathcal{A}$-module map; $\eta$ inherits this property from the fact that $\eta^{*} \in E^{\sigma}$ is a left $\mathcal{A}$-module map. The sum in the point evaluation is well defined because $\left\|\eta^{n}\right\| \leq\|\eta\|^{n}$ for each $n$.

With a element $x=\{x(n)\}_{n \in \mathbb{Z}_{+}} \in \mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)$ we also associate a function $x^{\wedge}$ on $\mathbb{D}\left(E^{\sigma *}\right) \times \sigma(\mathcal{A})^{\prime}$, but this point evaluation is essentially different from that in $\mathcal{F}^{2}(E, \sigma)$. Given $\mu \in E^{\sigma}$, we define the operator $L_{\mu}: \sigma(\mathcal{A})^{\prime} \rightarrow E^{\sigma}$ by

$$
L_{\mu} b=\mu \otimes b=\mu b \quad\left(b \in \sigma(\mathcal{A})^{\prime}\right)
$$

Then $L_{\mu}$ is a right $\sigma(\mathcal{A})^{\prime}$-module map, and thus, for $n=1,2,3, \ldots$, we can define the $n^{\text {th }}$ generalized power:

$$
L_{\mu}^{n}=\left(L_{\mu} \otimes I_{\left(E^{\sigma}\right)^{\otimes n-1}}\right) \cdots\left(L_{\mu} \otimes I_{E^{\sigma}}\right) L_{\mu}
$$

which is an operator from $\sigma(\mathcal{A})^{\prime}$ into $\left(E^{\sigma}\right)^{\otimes n}$; we set $L_{\mu}^{0}$ equal to the identity map on $\sigma(\mathcal{A})^{\prime}$. We further define $L_{\mu, \pi}=L_{\mu} \otimes I_{\mathcal{X}}$ and $L_{\mu, \pi}^{n}=L_{\mu}^{n} \otimes I_{\mathcal{X}}$. Then one can compute that $\left\|L_{\mu, \pi}^{n}\right\| \leq\|\mu\|^{n}$ for each $n \in \mathbb{Z}_{+}$. In particular, for each $\eta \in \mathbb{D}\left(E^{\sigma *}\right)$ we have the operators $L_{\eta *, \pi}$ and $L_{\eta *, \pi}^{n}$. Then, with a state sequence $x=\{x(n)\}_{n \in \mathbb{Z}_{+}} \in \mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)$ we associate a function $x^{\wedge}$ on $\mathbb{D}\left(E^{\sigma *}\right) \times \sigma(\mathcal{A})^{\prime}$ with values in $\mathcal{X}$ that is defined by

$$
x^{\wedge}(\eta, b)=\sum_{n=0}^{\infty}\left(L_{\eta^{*}, \pi}^{n}\right)^{*}\left(b \otimes I_{\left(E^{\sigma}\right)^{\otimes n-1} \otimes \mathcal{X}}\right) x(n)
$$

for each $(\eta, b) \in \mathbb{D}\left(E^{\sigma *}\right) \times \sigma(\mathcal{A})^{\prime}$. In a different context, and restricted to the case $x=\{x(n)\}_{n \in \mathbb{Z}_{+}} \in \mathcal{F}^{2}\left(E^{\sigma}, \pi\right)$, this form of point evaluation was first studied in [15]; see also [8].

Example 6.1 (Continuation of Example 5.2).
Part A: In this case $\mathbb{D}\left(E^{\sigma *}\right)$ can be identified with the unit ball $\mathbb{B}^{d}$ defined in (2.3) (formally elements of $\mathbb{D}\left(E^{\sigma *}\right)$ are row vectors) and $\sigma(\mathcal{A})^{\prime}=\mathbb{C}$. Let $\eta=$ $\left(z_{1}, \ldots, z_{d}\right) \in \mathbb{B}^{d}$. Then the generalized power $\eta^{n}$ corresponds to the vector in $\mathbb{C}^{d_{n}}$ whose $j^{t h}$ entry is equal to

$$
z^{\alpha}=z_{i_{n}} \cdots z_{i_{1}} \text { if } \alpha:=\nu_{n}(j)=i_{n} \cdots i_{1}
$$

Since $z_{1}, \ldots, z_{d}$ are scalars, and thus commute, the entries in $\eta^{n}$ corresponding to $\alpha_{1}$ and $\alpha_{2}$ (via the map $\nu_{n}$ ) that are mapped to the same element in $\mathbb{Z}_{+}^{d}$ via the abelianization map $\mathbf{a}$ in (2.13) are the same. Then for $u=\{u(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$ in $\mathcal{F}^{2}(E, \sigma)=\ell_{\mathcal{U}}^{2}\left(\mathcal{F}_{d}\right)$, the point evaluation in a point $(z, b) \in \mathbb{B}^{d} \times \mathbb{C}$ is given by

$$
\widehat{u}(z, b)=\sum_{\alpha \in \mathcal{F}_{d}} z^{\alpha} b u(\alpha)=b \sum_{n \in \mathbb{Z}_{+}^{d}} z^{n} \bar{u}(n)
$$

where $\bar{u}(n)=\sum_{\mathbf{a}(\alpha)=n} u(\alpha)$ for each $n \in \mathbb{Z}_{+}^{d}$. In particular, the part of the point evaluation that comes from $\sigma(\mathcal{A})^{\prime}$ plays no real role, and will be left out in the sequel, i.e., we only consider $b=1 \in \mathbb{C}$ and shall write $\widehat{u}(z)$ instead of $\widehat{u}(z, 1)$. Using that $\mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)$ can be identified with a subspace of $\ell_{\mathcal{X}}^{\infty}\left(\mathcal{F}_{d}\right)$, it follows in a similar way that the point-evaluation for an $x=\{x(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$ in $\mathcal{F}^{2}\left(E^{\sigma}, \pi\right)$ is defined by:

$$
\widehat{x}(z)=\sum_{n \in \mathbb{Z}_{+}^{d}} z^{n} \bar{x}(n) \quad\left(z \in \mathbb{B}^{d}\right)
$$

where $\bar{x}(n)=\sum_{\mathbf{a}(\alpha)=n} x(\alpha)$ for each $n \in \mathbb{Z}_{+}^{d}$.
Part B: In the noncommutative setting $\mathbb{D}\left(E^{\sigma *}\right)$ corresponds to the set of strict row contractions $T=\left(T_{1}, \ldots, T_{d}\right)$ with $T_{j} \in \mathcal{L}(\mathcal{K})$ and $\sigma(\mathcal{A})^{\prime}=\mathcal{L}(\mathcal{K})$. Similar to the case for Part A, the $n^{t h}$ generalized power of $T$ is defined via a noncommutative functional calculus: the $j^{\text {th }}$ entry of $T^{n}$ is given by $T^{\alpha}$ in (2.8) in case $\alpha:=\nu_{n}(j)$. Since the elements $T_{j}$ of $T$ do not (necessarily) commute pairwise, we obtain a noncommutative point evaluation for an element $u=\{u(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$ in $\mathcal{F}^{2}(E, \sigma)=$ $\ell_{\mathcal{K}}^{2}\left(\mathcal{F}_{d}\right):$

$$
\hat{u}(T, b)=\sum_{\alpha \in \mathcal{F}_{d}} T^{\alpha} b u(\alpha)
$$

where $T$ is a strict row contraction and $b \in \mathcal{L}(\mathcal{K})$. The point evaluation of an element $x=\{x(\alpha)\}_{\alpha \in \mathcal{F}_{d}}$ in $\mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)=\ell_{\mathcal{X}}^{2}\left(\mathcal{F}_{d}\right)$ goes through the $*$-representation $\pi$ :

$$
\hat{x}(T, b)=\sum_{\alpha \in \mathcal{F}_{d}} \pi\left(T^{\alpha} b\right) x(\alpha)
$$

with again $T$ a strict row contraction and $b \in \mathcal{L}(\mathcal{K})$. Now, under the restrictions of Example 5.1 we have input sequences of the form $u=\left\{\bar{u}(\alpha) k_{0}\right\}_{\alpha \in \mathcal{F}_{d}}$, with $\bar{u}(\alpha) \in \mathbb{C}$, and state sequences of the form $x=\left\{\bar{x}(\alpha) \otimes k_{0}\right\}_{\alpha \in \mathcal{F}_{d}}$, with $\bar{x}(\alpha) \in \overline{\mathcal{X}}$, where $k_{0}$ is a fixed vector in $\mathcal{K}$. In that case $\hat{u}(T, b)=\sum_{\alpha \in \mathcal{F}_{d}} \bar{u}(\alpha) T^{\alpha} b k_{0}$ and $\hat{x}(T, b)=\sum_{\alpha \in \mathcal{F}_{d}} \bar{x}(\alpha) \otimes\left(T^{\alpha} b k_{0}\right)$, and replacing $k_{0}$ with $b k_{0}$ we see that we may restrict the point evaluation to the case that $b=I_{\mathcal{K}}$.

## 7. Frequency-domain equations

With the point evaluations as defined in the previous section we have the following result.
Theorem 7.1. Let $\{A, B, C, D, \pi\}$ be a dissipative system. Given an input sequence $u=\{u(n)\}_{n \in \mathbb{Z}_{+}} \in \mathcal{F}^{2}(E, \sigma)$ and an initial state $x(0)$, let $y=\{y(n)\}_{n \in \mathbb{Z}_{+}} \in$ $\mathcal{F}^{2}(E, \sigma)$ and $x=\{x(n)\}_{n \in \mathbb{Z}_{+}} \in \mathcal{F}^{\infty}\left(E^{\sigma}, \pi\right)$ be the output and state sequences given by the system of equations (5.20). Then the $Z$-transforms $u^{\wedge}, x^{\wedge}$ and $y^{\wedge}$ of $u, x$ and $y$ satisfy the frequency domain equations:

$$
\begin{align*}
& x^{\wedge}(\eta, b)=x(0)+L_{\eta^{*}, \pi}^{*} A x^{\wedge}(\eta, b)+L_{\eta^{*}, \pi}^{*} B u^{\wedge}(\eta, b), \\
& y^{\wedge}(\eta, b)=C x^{\wedge}(\eta, b)+D u^{\wedge}(\eta, b), \tag{7.27}
\end{align*}
$$

and we have

$$
\begin{equation*}
y^{\wedge}(\eta, b)=F_{\Sigma}(\eta) u^{\wedge}(\eta, b)+W_{\Sigma}(\eta) x(0) \tag{7.28}
\end{equation*}
$$

where the transfer function $F_{\Sigma}$ and observation function $W_{\Sigma}$ are given by

$$
\begin{gather*}
F_{\Sigma}(\eta)=D+C\left(I-L_{\eta^{*}, \pi}^{*} A\right)^{-1} L_{\eta^{*}, \pi}^{*} B  \tag{7.29}\\
W_{\Sigma}(\eta)=C\left(I-L_{\eta^{*}, \pi}^{*} A\right)^{-1}
\end{gather*}
$$

Example 7.2 (Continuation of Example 6.1).
Part A: For this example the system matrix has the form (5.21). An element $\eta$ of $\mathbb{D}\left(E^{\sigma *}\right)$ is identified with a point $z=\left(z_{1}, \ldots, z_{d}\right) \in \mathbb{B}^{d}$, and in that case the operator $L_{\eta^{*}}^{*} A \in \mathcal{L}(\mathcal{X})$ works out to be $z_{1} A_{1}+\cdots+z_{d} A_{d}$. Similarly $L_{\eta^{*}}^{*} B=z_{1} B_{1}+\cdots z_{d} B_{d}$ and (7.29) gives us

$$
\begin{aligned}
F_{\Sigma}(z)= & D+C\left(I-z_{1} A_{1}-\cdots-z_{d} A_{d}\right)^{-1} \times \\
& \times\left(z_{1} B_{1}+\cdots+z_{d} B_{d}\right) \\
W_{\Sigma}(z)= & C\left(I-z_{1} A_{1}-\cdots-\cdots z_{d} A_{d}\right)^{-1}
\end{aligned}
$$

for the transfer function and the observation function, in agreement with (2.6), and (7.28) reduces to (2.5).

Part B: For this example, $\sigma(\mathcal{A})^{\prime}=\mathcal{L}(\mathcal{K})$ is nontrivial, $\pi: \sigma(\mathcal{A})^{\prime}=\mathcal{L}(\mathcal{K}) \rightarrow \mathcal{L}(\mathcal{X})$ is a nontrivial representation, and $\mathbb{D}\left(E^{\sigma *}\right)$ is identified with the set of strict row contractions $T=\left[\begin{array}{lll}T_{1} & \cdots & T_{d}\end{array}\right] \in \mathcal{L}\left(\mathcal{K}^{d}, \mathcal{K}\right)$. Nevertheless $E^{\sigma} \otimes_{\pi} \mathcal{X}$ is again identified with $\mathcal{X}^{d}$ and the system matrix has the form (5.21). Following the assumptions made in Example 5.1, i.e., $\mathcal{X}=\overline{\mathcal{X}} \otimes \mathcal{K}$ and the system operators are of the form (5.26), if $\eta \in \mathbb{D}\left(E^{\sigma *}\right)$ is identified with the strict row contraction $T=\left[\begin{array}{lll}T_{1} & \cdots & T_{d}\end{array}\right]$, then it works out that the operators $L_{\eta^{*}}^{*} A \in \mathcal{L}(\mathcal{X})$ and $L_{\eta^{*}}^{*} B \in \mathcal{L}(\mathcal{E}, \mathcal{X})$ can be identified with

$$
\begin{aligned}
& L_{\eta^{*}}^{*} A=\bar{A}_{1} \otimes T_{1}+\cdots+\bar{A}_{d} \otimes T_{d} \\
& L_{\eta^{*}}^{*} B=\bar{B}_{1} \otimes T_{1}+\cdots+\bar{B}_{d} \otimes T_{d}
\end{aligned}
$$

whence the transfer function and observation function assume the form

$$
\begin{gathered}
F_{\Sigma}(T)=D+\bar{C} \otimes I_{\mathcal{K}}\left(I-\bar{A}_{1} \otimes T_{1}-\cdots-\bar{A}_{d} \otimes T_{d}\right)^{-1} \times \\
\times\left(\bar{B}_{1} \otimes T_{1}+\cdots \bar{B}_{d} \otimes T_{d}\right), \\
W_{\Sigma}(T)=\bar{C} \otimes I_{\mathcal{K}}\left(I-\bar{A}_{1} \otimes T_{1}-\cdots-\bar{A}_{d} \otimes T_{d}\right)^{-1}
\end{gathered}
$$

in agreement with (2.12), and, moreover, (7.28) collapses to (2.11).

## 8. OTHER EXAMPLES AND OTHER TYPES OF GENERALIZED SYSTEMS

Our main purpose is to illustrate the general Muhly-Solel framework of tensor algebras, generalized Schur class, and Hardy algebras arising from a CR pair for two specific cases, namely: (1) commutative Fornasini-Marchesini input/state/output systems and their associated transfer functions as holomorphic functions on the unit ball, and (2) noncommutative Fornasini-Marchesini systems with associated transfer functions defined on the noncommutative unit ball. We mention that the Muhly-Solel formalism has at least one other special case which is a familiar setting in system theory, specifically, time-varying systems with the Muhly-Solel Schur-class function identifiable as the input-output map for a discrete-time linear time-varying system. Details on this can be found in 4].

Still other types of systems and associated transfer functions have been introduced and studied recently. Specifically there is a system theory over hyperholomorphic functions (with quaternions playing the role of the complex numbers) and a system theory for stochastic systems where the system parameters are themselves random variables; in these instances the familiar algebra centering around pointwise multiplication of functions is replaced by something more exotic: CauchyKovalevskaya product of hyperholomorphic functions or Wick product of stochastic variables. An introductory survey to these matters can be found in [1].

## References

[1] D. Alpay and D. Lavanony, Rational functions associated with the white noise space and related topics, Potential Anal. 29 (2008), 195-220.
[2] W. Arveson, An Invitation to $C^{*}$-Algebras, Graduate Texts in Mathematics, No. 39, Springer-Verlag, New York-Heidelberg, 1976.
[3] J.A. Ball, Dissipative noncommutative multidimensional linear systems and robust control theory, in: Proceedings of the International Workshop on Multidimensional ( $n D$ ) Systems, (Aveiro, 2007)pp. 123-129, www.ieeexplore.org.
[4] J.A. Ball, A. Biswas, Q. Fang, and S. ter Horst, Multivariable generalizations of the Schur class: positive kernel characterization and transfer function realization, in: Recent Advances in Operator Theory and Applications, pp. 17-79, OT 187, Birkhäuser, Basel, 2008.
[5] J.A. Ball, G. Groenewald and T. Malakorn, Structured noncommutative multidimensional linear systems, SIAM Journal of Control and Optimization 44 (2005), 1474-1528.
[6] J.A. Ball, G. Groenewald and T. Malakorn, Conservative structured noncommutative linear systems, in: The State Space Method: Generalizations and Applications, pp. 179-223, OT 161 Birkhäuser, Basel, 2006.
[7] J.A. Ball, G. Groenewald and T. Malakorn, Bounded real lemma for structured noncommutative multidimensional linear systems and robust control, Multidimens. Syst. Signal Process. 17 (2006), 119-150.
[8] J.A. Ball and S. ter Horst, Multivariable operator-valued Nevanlinna-Pick interpolation: a survey, Proceedings of IWOTA 2007, Potchestroom, South Africa, Birkhäuser volume to appear.
[9] E. Fornasini and G. Marchesini, Doubly-indexed dynamical systems: state-space models and structural properties, Mathematical Systems Theory 12 (1978), 59-72.
[10] P.S. Muhly and B. Solel, Tensor algebras over $C^{*}$-correspondences: representations, dilations and $C^{*}$-envelopes, J. Funct. Anal. 158 (1998), 389-457.
[11] P.S. Muhly and B. Solel, Hardy algebras, $W^{*}$-correspondences and interpolation theory, Math. Annalen 330 (2004), 353-415.
[12] P.S. Muhly and B. Solel, Canonical models for representations of Hardy algebras, Integral Equations Operator Theory 53 (2005), 411-452.
[13] P.S. Muhly and B. Solel, Hardy algebras associated with $W^{*}$-correspondences (point evaluation and Schur class functions), in: Operator Theory, Systems Theory and Scattering Theory: Multidimensional Generalizations, pp. 221-241, OT 157, Birkhäuser-Verlag, Basel, 2005.
[14] P.S. Muhly and B. Solel, Schur class operator functions and automorphisms of Hardy algebras, preprint.
[15] P.S. Muhly and B. Solel, The Poisson kernel for Hardy algebras, Complex Anal. Oper. Theory, available online.

Department of Mathematics, Virginia Tech, Blacksburg VA 24061, USA
E-mail address: ball@math.vt.edu
Department of Mathematics, Virginia Tech, Blacksburg VA 24061, USA
E-mail address: terhorst@math.vt.edu

