

The Method of Double Averaging: An Approach for Modeling Power-Factor-Correction Switching Converters

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Abstract—This paper describes the modeling of power-factor-correction converters under average-current-mode control, which are widely used in switch-mode power supply applications. The objective is to identify stability boundaries in terms of major circuit parameters for facilitating design of such converters. The approach employs a double averaging procedure, which first applies the usual averaging over the switching period and subsequently applies generalized averaging over the mains period. The resulting model, after two averaging steps and application of a harmonic balance procedure, is nonlinear and capable of describing the low-frequency nonlinear dynamics of the system. The parameter ranges within which stable operation is guaranteed can be accurately and easily found using this model. Experimental measurements are provided for verification of the analytical results.

Index Terms—Averaged model, closed-form stability condition, double averaging, power factor correction (PFC), switching converters.

I. INTRODUCTION

POWER-FACTOR-CORRECTION (PFC) boost stages are widely used in off-line switching power supplies. The basic configuration of a PFC boost stage consists of 1) an average-current-mode (ACM) control loop which “forces” the input current to follow the waveform of the input voltage; and 2) a slow output voltage feedback loop which controls the amplitude of the “shaped” input current according to the power demanded by the load. The aim of the control is to achieve a unity power factor at the input side and to maintain power balance while keeping the output voltage at a nearly constant level [1], [2]. Previously, several attempts have been made to model the dynamics of power converters based on dynamic phasors [3]–[5], which can be simplified in some specific applications to averaging and linearization [6]–[10]. However, because the actual process is

nonlinear, stability analyzes based on linearized models are very restricted and often provide stability information (e.g., parameter ranges for stable operation) which are not necessarily consistent with reality. In particular, low-frequency period-doubling phenomena in PFC boost preregulators have been reported by Orabi *et al.* [11]–[13]. Also, fast-scale instability has also been reported by Iu *et al.* [15]. In this paper we make use of a hierarchical modeling approach, a basic form of which was proposed in [6], and apply generalized averaging together with a harmonic balance procedure to identify the low-frequency nonlinear dynamics of the PFC boost converter. This method results in a model that is capable of revealing the phenomena found in [11]–[13]. Essentially we apply averaging twice in succession, over the switching period and the mains period, to obtain a double-averaged model in the form of a system of simple continuous-time differential equations which can be readily analyzed and numerically simulated. Stability information generated from this model is accurate for the purpose of locating the ranges of parameters within which the PFC boost converter operates in its normally expected regime. This technique was also used in Gordillo *et al.* [14] for the study of the dynamics of a power factor preregulator under a special control. The work of Gordillo *et al.* [14], however, focuses on varying the output voltage. More comprehensive stability information that is presented in terms of major circuit parameters is unavailable. Furthermore, as is common to switching converters, operation may assume either continuous conduction mode (CCM) or discontinuous conduction mode (DCM). While much of the previous work has focused on the CCM case [11]–[15], very few has addressed the case of DCM [16]. Moreover, the problem of bifurcation of a supposedly CCM operation into DCM due to variation of some parameter(s) is a complex issue, which requires further investigation. In this paper we restrict ourselves to CCM operation, as it is the usual operating condition for most of medium to high power applications.

This paper is organized as follows. In Section II, we will explain the usual operation of the boost converter as employed in typical PFC applications, with emphasis on the salient operational conditions in the presence of a slowly varying input voltage. Our main proposal is described in Section III, where we explain the mathematical technique of *doubling averaging*, and apply it to produce a model for the PFC boost converter under average current-mode control. The derivation of the steady-state solution will be given in Section IV. Application of the double-averaged model for analysis of the stability of PFC converters

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for $n\pi \leq \omega_m t < (n+1)\pi$, $n = 0, 1, 2, \dots$. It should be noted that $p(t)$ is connected with v_o via the feedback circuit. Therefore, (8) is a continuous-time differential equation in v_o .

Thus, the system under study can be described by the set of equations consisting of (4) and (8). To analyze this system, we will make use of the fact that it describes a system of *forced oscillation*, where the variables are “driven” periodically, containing frequency components up to twice the mains frequency.

B. Second Averaging

We now consider further application of averaging to the system of equations under study, which consists of (4) and (8). Specifically, we take *moving average* over the mains period [19]. For any variable $a(t)$, we may assume that it takes the following form:

$$a(t) = A_0 + A_1 \sin(\omega_m t - \theta_1(t)) + A_2 \sin(2\omega_m t + \theta_2(t)). \quad (9)$$

We define each component of the moving average of $a(t)$, denoted by $[a(t)]_k$ or simply $[a]_k$ for $k = 0, 1, 2$, by

$$[a(t)]_k = [a]_k = \frac{\omega_m}{2\pi} \int_{t-\frac{2\pi}{\omega_m}}^t a(\tau) \exp(-jk\omega_m\tau) d\tau. \quad (10)$$

For brevity and without confusion, we simply write $[a]_k$ as a_k . In other words, a_0 is the dc component of the moving average of $a(t)$, a_1 is the component of the moving average of $a(t)$ at ω_m , and a_2 is the component of the moving average of $a(t)$ at $2\omega_m$. Thus, we may write

$$a(t) = a_0(t) + a_1(t)e^{j\omega_m t} + [a_1(t)e^{j\omega_m t}]^* + a_2(t)e^{2j\omega_m t} + [a_2(t)e^{2j\omega_m t}]^* \quad (11)$$

where superscript $*$ denotes complex conjugation. Upon expanding, we get

$$a_0(t) = a_0 = A_0 \quad (12)$$

$$a_1(t) = a_{1r} + ja_{1i} = -\frac{A_1}{2} \sin \theta_1 - j\frac{A_1}{2} \cos \theta_1 \quad (13)$$

$$a_2(t) = a_{2r} + ja_{2i} = -\frac{A_2}{2} \sin \theta_2 - j\frac{A_2}{2} \cos \theta_2. \quad (14)$$

It is readily shown that the moving average of the time-derivative of each component is related to the time-derivative of the component by

$$\begin{aligned} \left[\frac{da(t)}{dt} \right]_k &= \frac{\omega_m}{2\pi} \int_{t-\frac{2\pi}{\omega_m}}^t \frac{da(\tau)}{d\tau} \exp(-jk\omega_m\tau) d\tau \\ &= \frac{da_k(t)}{dt} + jk\omega_m a_k(t) \end{aligned} \quad (15)$$

where $[]_k$ denotes consistently the k th component of the moving average, with $k = 0$ corresponding to the dc component, $k = 1$

the ω_m component, and $k = 2$ the $2\omega_m$ component. Furthermore, the following relations are readily obtained:

$$[a(t) \cos 2\omega_m t]_k = \frac{1}{2}(a_{k-2} + a_{k+2}) \quad (16)$$

$$[a^2(t)]_0 = a_0^2 + 2|a_1|^2 + 2|a_2|^2 \quad (17)$$

$$[a^2(t)]_1 = 2a_0 a_1 + 2(a_{1i} a_{2i} + a_{1r} a_{2r} + j(a_{1r} a_{2i} - a_{1i} a_{2r})) \quad (18)$$

$$[a^2(t)]_2 = (a_1^*)^2 + 2a_0 a_2. \quad (19)$$

Remarks: It should be noted that the essential procedure of taking the moving average differs from the basic double-averaging method of [6] where a constant steady-state point is assumed. Specifically, the method of [6] results in a model that reduces all variables varying at the mains frequency to fixed steady-state points. Thus, the subsequent small-signal analysis only reflects stability of the mains cycle but falls short of identifying oscillations at other subharmonics which have been found experimentally. As will be shown shortly, a moving-averaged model, coupled with the application of a harmonic balance procedure at any specific harmonic of interest, permits a quick identification of forced oscillation which may occur at half, quarter, etc. of the mains frequency. A similar application was also reported in Gordillo *et al.* [14], where bifurcation resulting from varying the output voltage was studied. However, practical stability information that is presented in terms of major circuit parameters is still unavailable.

C. Derivation of the Double Averaged Model for PFC Boost Converter

In this subsection we apply the foregoing averaging to (4) and (8). We expect to get three equations from each of (4) and (8), corresponding to the dc, ω_m and $2\omega_m$ components. To avoid possible confusion due to mixup in subscripts, we define

$$\begin{cases} x = v_o(t) \\ y = p(t). \end{cases} \quad (20)$$

In the following, the use of subscripts will be consistent with the description of the second averaging in Section III-B, i.e., subscript 0 for dc components, subscript 1 for mains frequency components, and subscript 2 for twice mains frequency components. Now, application of second averaging to (8) gives

$$\begin{aligned} \frac{C}{2} \frac{d}{dt} (x_0^2 + 2x_{1r}^2 + 2x_{1i}^2 + 2x_{2r}^2 + 2x_{2i}^2) \\ + \frac{(x_0^2 + 2x_{1r}^2 + 2x_{1i}^2 + 2x_{2r}^2 + 2x_{2i}^2)}{R} \\ = y_0 - y_{2r} \end{aligned} \quad (21)$$

$$\begin{aligned} \frac{C}{2} \frac{d}{dt} (x_0 x_1 + x_{1r} x_{2r} + x_{1i} x_{2i} + j(x_{1r} x_{2i} - x_{1i} x_{2r})) \\ + \left(\frac{j\omega_m C}{2} + \frac{1}{R} \right) (x_0 x_1 + x_{1r} x_{2r} + x_{1i} x_{2i} \\ + j(x_{1r} x_{2i} - x_{1i} x_{2r})) \\ = \frac{y_1}{2} - \frac{y_{-1}}{4} \end{aligned} \quad (22)$$

$$\begin{aligned} \frac{C}{2} \frac{d}{dt} (2x_0 x_2 + (x_1^*)^2) + \left(j\omega_m C + \frac{1}{R} \right) (2x_0 x_2 + (x_1^*)^2) \\ = y_2 - \frac{y_0}{2}. \end{aligned} \quad (23)$$

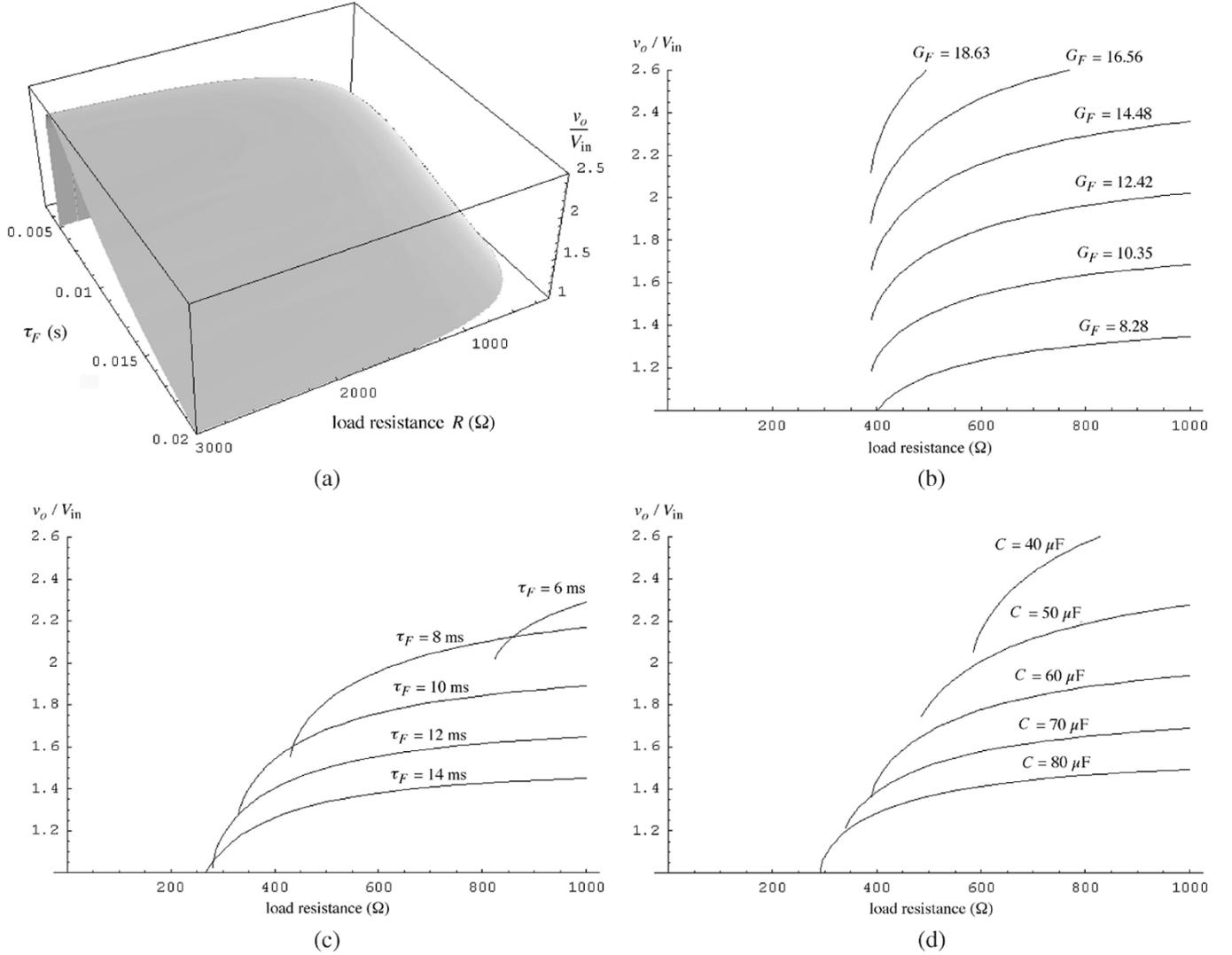


Fig. 3. Stability boundaries with: (a) $G_F = 12.8$ A, $C = 60$ μF , and $\omega_m = 2\pi(60)$ rad/s; (b) $\tau_F = 8.6$ ms, $C = 60$ μF , and $\omega_m = 2\pi(60)$ rad/s; (c) $G_F = 12.8$ A, $C = 60$ μF , and $\omega_m = 2\pi(60)$ rad/s; (d) $\tau_F = 8.6$ ms, $G_F = 12.8$ A, and $\omega_m = 2\pi(60)$ rad/s. Region above the boundary surface or curve corresponds to stable operating region. At the boundary, period-doubling occurs to disrupt normal operation. Region under the boundary surface or curve corresponds to “unstable operation” in the sense that normal operation cannot be maintained.

Note that y_{-1} is equal to y_1^* . Moreover, (4) is decomposed into

$$\tau_F \frac{dy_0}{dt} + y_0 = -G_F(x_0 - V_{\text{ref}}) \quad (24)$$

$$\tau_F \frac{dy_1}{dt} + (j\omega_m \tau_F + 1)y_1 = -G_F x_1 \quad (25)$$

$$\tau_F \frac{dy_2}{dt} + (j2\omega_m \tau_F + 1)y_2 = -G_F x_2. \quad (26)$$

Thus, (21)–(26) are the double averaged equations for the PFC boost converter under study.

IV. STEADY-STATE SOLUTION

To find the steady-state solution, we put all time-derivatives to zero, giving

$$x_1 = y_1 = 0 \quad (27)$$

$$x_{2r} = \left[\frac{V_{\text{ref}} - x_0}{2\Delta} \right] \left[1 + 2x_0 \frac{1 + 4\omega_m^2 \tau_F^2}{G_F R} \right] \quad (28)$$

$$x_{2i} = \left[\frac{\omega_m (V_{\text{ref}} - x_0)}{\Delta} \right] \left[-\tau_F + x_0 \frac{1 + 4\omega_m^2 \tau_F^2}{G_F} \right] \quad (29)$$

$$y_{2r} = \left[\frac{V_{\text{ref}} - x_0}{\Delta} \right] \left[G_F + 2x_0 \frac{1 - 2\omega_m^2 C R \tau_F}{R} \right] \quad (30)$$

$$y_{2i} = \left[\frac{V_{\text{ref}} - x_0}{\Delta} \right] \left[\omega_m C + \frac{2\omega_m \tau_F}{R} \right] x_0 \quad (31)$$

where

$$\Delta = 1 + \frac{4(1 - 2\omega_m^2 C R \tau_F) x_0}{G_F R} + \frac{4(1 + \omega_m^2 C^2 R^2)(1 + 4\omega_m^2 \tau_F^2) x_0^2}{G_F^2 R^2}. \quad (32)$$

Putting the above equations in (21), and using (24), we get

$$\begin{aligned} & 8(1 + \omega_m^2 C^2 R^2)(1 + 4\omega_m^2 \tau_F^2) x_0^4 \\ & + 8G_F R [2 - 2\omega_m^2 C R \tau_F + 4\omega_m^2 \tau_F^2 \\ & + \omega_m^2 C^2 R^2 (1 + 4\omega_m^2 \tau_F^2)] x_0^3 \\ & + G_F R [-8V_{\text{ref}} (1 + \omega_m^2 C^2 R^2)(1 + 4\omega_m^2 \tau_F^2) \\ & + G_F R (9 - 12\omega_m^2 C R \tau_F + 4\omega_m^2 \tau_F^2)] x_0^2 \\ & + G_F^2 R^2 [G_F R + 2V_{\text{ref}} (-4 + 6\omega_m^2 C R \tau_F - 4\omega_m^2 \tau_F^2)] x_0 \\ & + G_F^2 R^2 V_{\text{ref}} [-G_F R + V_{\text{ref}} (1 + 4\omega_m^2 \tau_F^2)] = 0 \end{aligned} \quad (33)$$

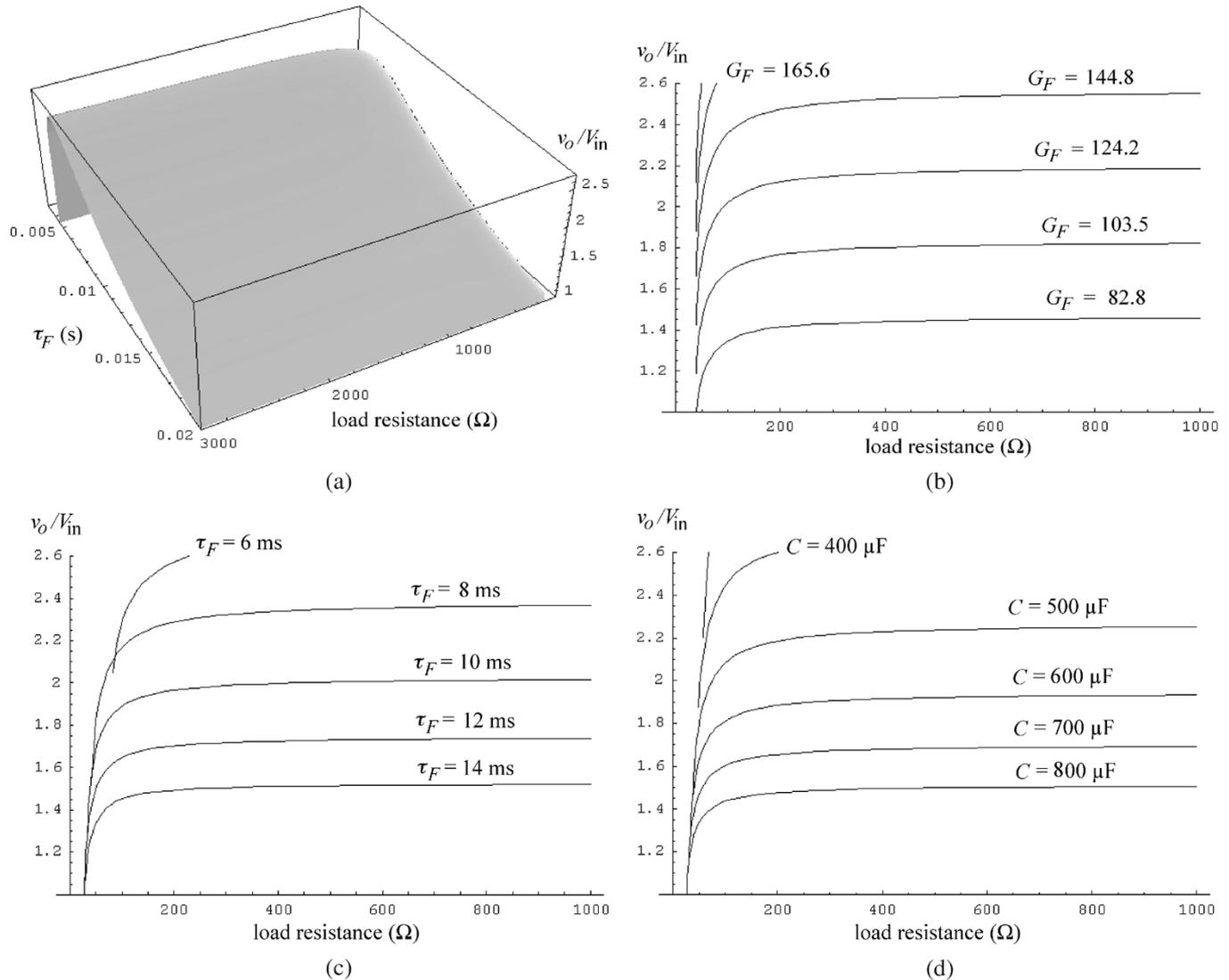


Fig. 4. Corresponding stability boundaries with: (a) $G_F = 128$ A, $C = 600$ μF , and $\omega_m = 2\pi(60)$ rad/s; (b) $\tau_F = 8.6$ ms, $C = 600$ μF , and $\omega_m = 2\pi(60)$ rad/s; (c) $G_F = 128$ A, $C = 600$ μF , and $\omega_m = 2\pi(60)$ rad/s; (d) $\tau_F = 8.6$ ms, $G_F = 128$ A, and $\omega_m = 2\pi(60)$ rad/s.

from which x_0 can be solved numerically. Note that x_0 is the steady-state dc component of the output voltage, which is close to V_{ref} .

V. APPLICATION OF DOUBLE AVERAGED MODEL TO STABILITY ANALYSIS

Our purpose in this section is to examine the stability of the system with respect to its operation at the forcing frequency $2\omega_m$, which is twice the mains frequency. This is because the input is a rectified sine wave repeating at $2\omega_m$, and all variables in the system are expected to repeat at this forcing frequency. Moreover, if the system is entrapped in an orbit of frequency ω_m , i.e., at half the expected repetition frequency,¹ the operation is considered undesirable for practical purposes since the device stresses would be drastically altered.

¹In the literature of dynamical systems, such a change in operation is known as period doubling, and the transition from the normal operation to a double-period operation is typically a sudden transition. The usual term used to describe this transition is *bifurcation* [20], [21].

In fact, such “period-doubling” has been observed experimentally when certain parameters are changed irrespective of the stability of the $2\omega_m$ -orbit [11], [12]. It is therefore of practical importance to locate the occurrence of period-doubling in the parameter space. In our double-averaged model, this is equivalent to finding the condition under which the system is entrapped into an ω_m -orbit.

Our approach to solve this stability problem is to consider the loop gain of x_1 (or y_1) in our double averaged model. Thus, from (22) and (25), and the steady-state values solved from (28) through (31), putting the time derivatives to zero, we get the loop gain for x_1 as shown in (34) at the bottom of the next page.

The system operates at the expected forcing frequency if it is not entrapped in the ω_m -orbit, i.e., $T_{x1} < 1$. (It turns out that the condition $T_{x1} > -1$ is irrelevant for practical sets of parameter values.) This condition is necessary for the system to operate in the expected regime. We may simplify (34) by putting $x_0 \approx V_{ref}$, giving (35), shown at the bottom of the next page. We note that the condition of $T_{x1} < 1$ is equivalent to (36), also shown at the bottom of the next page.

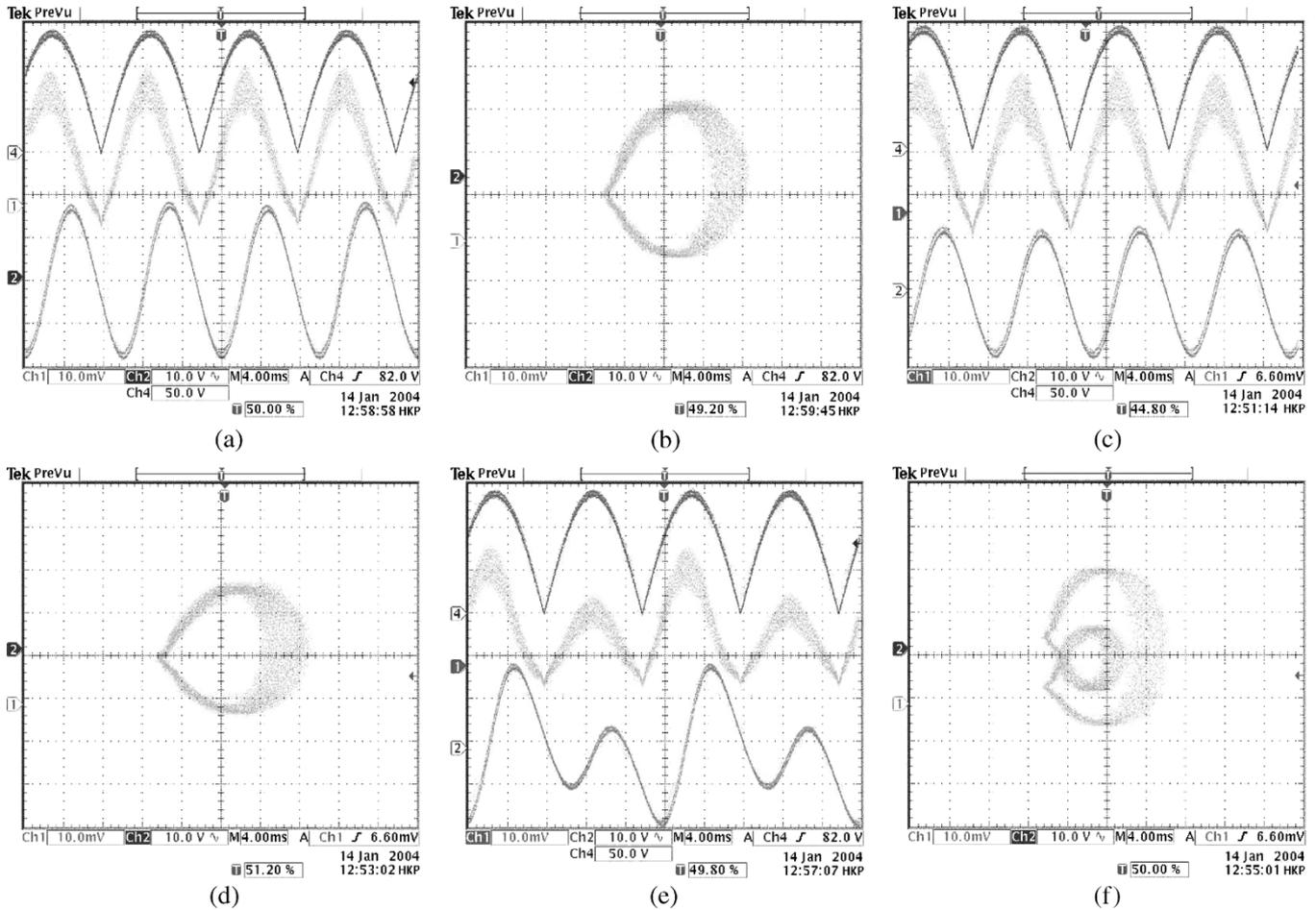


Fig. 6. Measured waveforms and phase portraits. (a), (b) Normal operation with $R = 454 \Omega$ and $V_o = 297 \text{ V}$; (c), (d) normal operation with $R = 645 \Omega$ and $V_o = 360 \text{ V}$; (e), (f) period-doubled (unstable) operation with $R = 645 \Omega$ and $V_o = 297 \text{ V}$. For all cases, $G_F \approx 20 \text{ A}$, input voltage is 100 V rms , $\tau_F = 8.46 \text{ ms}$, and line frequency is 50 Hz . For waveforms in (a), (c) and (e), upper trace is line voltage, middle trace is input current (1 A/div) and lower trace is output voltage ripple. For phase portraits in (b), (d) and (f), x axis is input current and y axis is output voltage.

ceases. Region above the boundary corresponds to normal operation, where the system is driven periodically at twice the mains frequency. These results are consistent with our observation. Specifically, the operation is stable up to a certain threshold set of parameter values. In Fig. 3, the capacitor values are relatively small and the output ripple is therefore relatively large, causing instability when such a large ripple is being fed back. Moreover, in Fig. 4, the capacitor values are large (normal in practical sense), and instability can only occur for relatively high feedback gain values. Our analysis provides a quantitative tool for evaluating the stability boundaries. We will verify this tool in the next section by experiments.

VII. EXPERIMENTAL VERIFICATIONS

An experimental PFC boost converter under average current mode control has been constructed for verification purposes. The complete schematic is shown in Fig. 5. The current-shaping control employs a standard 3854A integrated circuit controller. The line frequency is 50 Hz . The input voltage is fixed at 100 V rms .

Experiments have been performed to verify the stability boundaries predicted in Sections V and VI. First, for a quick illustration, typical waveforms and phase portraits for normal

operation and period-doubled (unstable) operation are shown in Fig. 6.

Our main objective is to locate the parameter points where period-doubling occurs. For convenience we replot the stability curves along with the experimental data in Fig. 7. We clearly see that the experimental data are in good agreement with the analytical results.

A few practical points are worth noting here.

- 1) One of the determining parameters is v_o/V_{in} , as shown in our analytical study. However, for measurement convenience, we fix the input voltage in our experiments and vary the output voltage.
- 2) In our analytical model, the feedback loop is defined by two parameters, namely G_F and τ_F . In order to find the corresponding value of G_F for the experimental circuit, we have to carefully identify all gains that may possibly contribute to G_F since the model assumes no gain in the multiplier-divider block. In our experiment, we find this gain by direct measurement.
- 3) Our definition of τ_F is the real time constant. Thus, its reciprocal should be compared directly with the angular mains frequency in rad/s instead of the frequency in

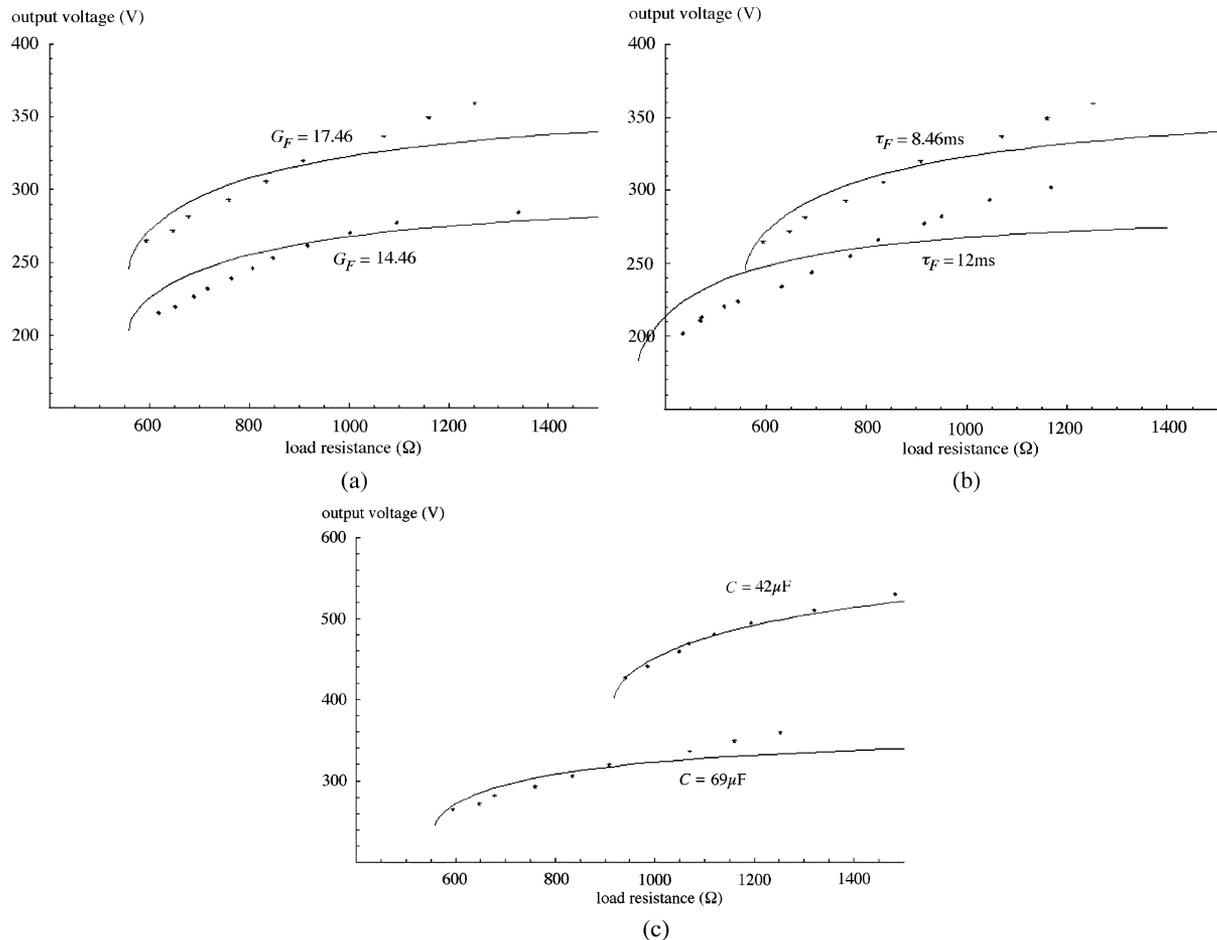


Fig. 7. Measured stability boundaries taken at $V_{in} = 100$ V (rms) and 50 Hz line frequency. Experimental data points are plotted as dots. Solid curves are from analytical expressions. (a) $C = 69 \mu\text{F}$ and $\tau_F = 8.46$ ms; (b) $C = 69 \mu\text{F}$ and $G_F = 17.46$ A and $\tau_F = 8.46$ ms; (c) $G_F = 17.46$ A and $\tau_F = 8.46$ ms. Note that the y axis would have been v_o/V_{in} for better generality, as shown in our analytical study. However, for measurement convenience, we have fixed the input voltage in our experiments and varied the output voltage.

Hertz. For instance, $\tau_F = 1/2\pi(50)$ s = 3.18 ms refers to a time constant equal to the period of a 50 Hz signal.

- 4) We have performed the experiments for relatively small output capacitor values in order to allow instability to be seen with the usual gain values readily implementable by the particular controller UC3854A. However, we should stress that the phenomenon remains invariant for practical capacitor values (e.g., hundreds of μF) provided that higher feedback gain values are used accordingly, as per Fig. 4 or (36).

VIII. CONCLUSION

At present, the main technique for the analysis of power converters is based on averaging and subsequent small-signal linearization. Since PFC converters are subject to a time-varying input voltage and the current-shaping process is nonlinear, the results obtained from small-signal linear analysis have limited validity and often provide inaccurate stability information, as has been explained and demonstrated by Orabi *et al.* [11]–[13] with computer simulations and laboratory measurements. Nonlinear analytical models would be desirable for stability evaluation. This paper has applied a method of double averaging for modeling PFC converters, leading to nonlinear models from

which closed-form stability conditions can be derived. This method is capable of revealing the *low-frequency stability* of the circuit. For instability near the switching frequency, moreover, we need a different type of models, such as those based on discrete-time mapping [15], [22], for revealing the fast-scale instability.

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