# Generating Functional Analysis for Iterative CDMA Multiuser Detectors

Kazushi Mimura, Member, IEEE, and Masato Okada

*Abstract*—We investigate the detection dynamics of a soft parallel interference canceller (soft-PIC), which includes a hard-PIC as a special case, for code-division multiple-access (CDMA) multiuser detection, applied to a randomly spread, fully synchronous base-band uncoded CDMA channel model with additive white Gaussian noise under perfect power control in the large-system limit. We analyze the detection dynamics of some iterative detectors, namely soft-PIC, the Onsager-reactioncancelling parallel interference canceller (ORC-PIC) and the belief-propagation-based detector (BP-based detector), by the generating functional analysis (GFA). The GFA allows us to study the asymptotic behavior of the dynamics in the infinitely large system without assuming the independence of messages. We study the detection dynamics and the stationary estimates of an iterative algorithm.

We also show the decoupling principle in iterative multiuser detection algorithms in the large-system limit. For a generic iterative multiuser detection algorithm with binary input, it is shown that the multiuser channel is equivalent to a bank of independent single-user additive non-Gaussian channels, whose signal-to-noise ratio degrades due to both the multiple-access interference and the Onsager reaction, at each stage of the algorithm. If an algorithm cancels the Onsager reaction, the equivalent single-user channels coincide with an additive white Gaussian noise channel. We also discuss ORC-PIC and the BPbased detector.

*Index Terms*—generating functional analysis, Code-Division Multiple-Access, iterative algorithms, detection dynamics

#### I. INTRODUCTION

**D**ETECTION DYNAMICS of generic iterative codedivision multiple-access (CDMA) multiuser detectors, which utilize soft-decision in the large system limit, is presented in this paper. The CDMA is a digital modulation system that employs spreading codes to enable access to a mobile communication system by multiple users [54], [55]. The statistical-mechanical approach has been applied to evaluate the performance of various wireless communication systems [51]. This kind of systems is widely used in communications and signal processing, such as the code-division multipleaccess, the multiple-input multiple-output channels [39], [20], and compressed sensing [5], [47], [14], [15].

Various types of multiuser detectors utilizing soft-decision have been proposed so far [55], [54], [13], [25]. Tanaka has first evaluated the properties of the maximum a posteriori detector and the marginal-posterior-mode detector by

M. Okada is with Graduate School of Frontier Sciences, University of Tokyo, Chiba 277-5861, Japan and Brain Science Institute, RIKEN. e-mail: okada@k.u-tokyo.ac.jp.

the statistical-mechanical analysis, which is called the *replica analysis* (or the *replica method*) [49], [51], [43]. The replica analysis has widely been applied to analyze communication systems [20], [21], [40] and other information theoretic problems. It can treat the properties of detection results, but cannot directly treat the detection dynamics of detectors. Since the optimal marginal-posterior-mode detector itself is known to be NP-hard [55], it is important to construct suboptimal methods. Iterative algorithms are generally useful as such methods from the viewpoint of computational cost. Various kinds of iterative detection algorithms have been developed to date [41], [13], [25], [3], [45], [53], [23]. The analysis of iterative multiuser detection algorithms is therefore expected to play an important role in developments and improvements of detectors.

Recently, the state evolution to evaluate the dynamics of the approximate belief propagation has been proposed by Bayati and Montanari [2]. The detection dynamics of iterative algorithms which are characterized by a dense graph has attracted a great deal of attention from theoretical and practical viewpoints up to now [25], [52], [2]. Kabashima has proposed the belief-propagation-based detector and analyzed its performance [25]. Tanaka and Okada have analyzed the detection dynamics [52] of the soft parallel interference canceller (soft-PIC) proposed by Divsalar et al [13] by means of a dynamical theory for the Hopfield model [44]. Bayati and Montanari have analyzed the reconstruction dynamics of approximate belief propagation algorithm for compressed sensing [2]. These existing studies [25], [52], [2] have succeeded in analyzing of various kinds of iterative algorithms including belief-propagation-based methods.

However, these analyses, such as density evolution [46], [52] and state evolution [2], are justified only for the cases that the correlation between present estimates and their past values can be neglected. In other words, these can only be applied to the case where there is not a retarded self-interaction, which is caused by iterations and this is often called the Onsager reaction, and their predictions systematically deviate from computer simulation results in general.

We have already applied the *generating functional analysis* (*GFA*) [11] to the hard-PIC, i.e., Varanasi and Aazhang's conventional PIC [32], [33] so far. This analysis can however treat the hard-decision only [32], [33]. The GFA, which uses the saddle-point method [8], [29], allows us to study the asymptotic behavior of the dynamics in the infinitely large system [11], [6], [7]. Since it is not based on the S/N analysis, it does not therefore need the Gaussian assumption of the noise part. In the S/N analysis, the signal part, that contains the user's information being estimated, is separated from the remaining noise part; besides the noise part is generally

K. Mimura is with Faculty of Information Sciences, Hiroshima City University, Hiroshima 731-3194, Japan. e-mail: mimura@hiroshima-cu.ac.jp.

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assumed to follow a given distribution such as the Gaussian distribution.

In this paper, we investigate the detection dynamics of some iterative algorithms for CDMA multiuser detection, applied to a randomly spread, fully synchronous base-band uncoded CDMA channel model with additive white Gaussian noise under perfect power control. We here treat soft-PIC, the *Onsagerreaction-cancelling parallel interference canceller (ORC-PIC)*, which is an analogue of soft-PIC, and the *belief-propagation-based detector (BP-based detector)*. These models have the retarded self-interaction. To confirm the validity of our analysis, we have performed computer simulations under some typical system loads and channel noise conditions.

This paper is organized as follows. The next section explains the system model. Sections III and IV introduce some multiuser detection algorithms and the generating functional analysis, respectively. Sections V - VII present analyses of soft-PIC, ORC-PIC, and the BP-based detector. Section VIII explains the stationary states of iterative algorithms and their stability. In section IX, we discuss the decoupling principle [20], [21] in iterative detectors. The final section is devoted to a summary.

## **II. SYSTEM MODEL AND NOTATIONS**

Let us focus on the basic fully synchronous *K*-user baseband direct-sequence / binary phase-shift-keying CDMA channel model with perfect power control as

$$y^{\mu} := \frac{1}{\sqrt{N}} \sum_{k=1}^{K} s_{k}^{\mu} b_{k} + \sigma_{0} n^{\mu}, \qquad (1)$$

where  $y^{\mu}$  is the received signal at chip interval  $\mu \in \{1, \dots, N\}$ , and where  $b_k \in \{-1, 1\}$  and  $s_k^{\mu} \in \{-1, 1\}$  are the binary phase-shift-keying-modulated information bit and the spreading code of user  $k \in \{1, \dots, K\}$  at chip interval  $\mu$ , respectively. The Gaussian random variable  $\sigma_0 n^{\mu}$ , where  $n^{\mu} \sim \mathcal{N}(0, 1)$ , represents channel noise whose variance is  $\sigma_0^2$ . The spreading codes are independently generated from the identical symmetric distribution  $P(s_k^{\mu} = 1) = P(s_k^{\mu} = -1) =$ 1/2. The factor  $1/\sqrt{N}$  is introduced in order to normalize the power per symbol to 1. The signal-to-noise ratio is obtained as  $E_b/N_0 = 1/(2\sigma_0^2)$  by using these normalizations. The ratio  $\beta := K/N$  is called system load.

In this paper, the letters k, k' denote indices in  $\{1, \dots, K\}$ and the letters  $\mu, \mu'$  denote indices in  $\{1, \dots, N\}$ . The (k, k')element of the matrix W, whose index is a pair of the user numbers, is indicated as  $W_{kk'}$ . The elements of vectors  $\boldsymbol{y} = (y_1, \dots, y_N)$  and  $\boldsymbol{b} = (b_1, \dots, b_K)$ , whose indices are the chip interval index or the user number, are indicated as  $y^{\mu}$  and  $b_k$ , respectively. The *s*-th element of vector  $\boldsymbol{x} = (x^{(-1)}, x^{(0)}, \dots, x^{(t)})^{\top}$ , whose index is the stage number, is indicated by  $(\boldsymbol{x})^{(s)}$ , e.g.,  $(\boldsymbol{x})^{(1)} = x^{(1)}$  (Note that  $(\boldsymbol{x})^{(1)}$  does not mean  $x^{(-1)}$ ). Here,  $\boldsymbol{X}^{\top}$  denotes the transpose of  $\boldsymbol{X}$ . The (s, s') element of the matrix

$$\boldsymbol{X} = \begin{pmatrix} X^{(-1,-1)} & X^{(-1,0)} & \cdots & X^{(-1,t)} \\ X^{(0,-1)} & X^{(0,0)} & \cdots & X^{(0,t)} \\ \vdots & \vdots & & \vdots \\ X^{(t,-1)} & X^{(t,0)} & \cdots & X^{(t,t)} \end{pmatrix},$$

whose index (s, s') is a pair of the stage numbers, is indicated as  $(\mathbf{X})^{(s,s')}$ , e.g.,  $(\mathbf{X})^{(1,1)} = X^{(1,1)}$  (Note that  $(\mathbf{X})^{(1,1)}$ does not mean  $X^{(-1,-1)}$ ). The notations are summarized in Appendix A.

## **III. DETECTION ALGORITHMS**

We discuss the detection dynamics of the following three kinds of iterative detection algorithms in this paper.

## A. Soft-PIC

The soft parallel interference canceller has been proposed by Divsalar et al [13]. Kaiser and Hagenauer have also proposed a similar algorithm [27].

Definition 1: (Soft-PIC) The updating rule for tentative decision  $\tilde{b}_k^{(t)} \in \mathbb{R}$  of bit signal  $b_k$  at stage t is

$$\tilde{b}_{k}^{(t)} = f\left(h_{k} - \sum_{k'=1, \neq k}^{K} W_{kk'} \tilde{b}_{k'}^{(t-1)}\right),$$
(2)

where  $f : \mathbb{R} \to \mathbb{R}$ , which is called a transfer function, that is arbitrary and  $h_k$  is the output of the matched filter for user k:

$$h_k := \frac{1}{\sqrt{N}} \sum_{\mu=1}^{N} s_k^{\mu} y^{\mu}, \tag{3}$$

and  $W_{kk'}$  is the kk'-element of sample correlation matrix W of the spreading code:

$$W_{kk'} := \frac{1}{N} \sum_{\mu=1}^{N} s_k^{\mu} s_{k'}^{\mu}.$$
 (4)

The initial condition of iteration is  $\tilde{b}_k^{(-1)} = 0$ . When the transfer function chooses  $f(x) = \tanh(x/\sigma^2)$ , this iterative detection algorithm is called soft-PIC. Here,  $\sigma^2$  is a control parameter representing the detector's estimate of channel noise variance.

Müller and Huber have improved Soft-PIC and have numerically evaluated its performance [38].

#### B. ORC-PIC

In soft-PIC, matched filter output has a very complex correlation between all estimates. The correlation due to iterative calculation worsens performance of detection. The ORC-PIC is an analogue of soft-PIC, which has a term to cancel such correlation. The updating rule is modified to

$$\tilde{b}_{k}^{(t)} = f \left( h_{k} - \sum_{k' \neq k} W_{kk'} \tilde{b}_{k'}^{(t-1)} - \hat{\Gamma}_{k}^{(t,t-1)} \tilde{b}_{k}^{(t-1)} - \dots - \hat{\Gamma}_{k}^{(t,t-1)} \tilde{b}_{k}^{(-1)} \right), \quad (5)$$

with  $\tilde{b}_k^{(-1)} = 0$ . If we properly choose coefficients  $\{\hat{\Gamma}_k^{(s,s')}\}$  of the term  $-\hat{\Gamma}_k^{(t,t-1)}\tilde{b}_k^{(t-1)} - \cdots - \hat{\Gamma}_k^{(t,-1)}\tilde{b}_k^{(-1)}$ , the correlation can be cancelled. Tanaka and Okada derived the preceding parameter  $\{\hat{\Gamma}_k^{(s,s')}\}$  by applying density evolution [52] based on the statistical neurodynamics [44].

*Definition 2:* The ORC-PIC is defined by updating rule (5) with

$$(\hat{\Gamma}_k \tilde{\boldsymbol{b}}_k)^{(t)} = \beta G^{(t,t-1)} [\tilde{b}_k^{(t-1)} - (\hat{\Gamma}_k \tilde{\boldsymbol{b}}_k)^{(t-1)}], \qquad (6)$$

and  $(\hat{\boldsymbol{\Gamma}}_k \tilde{\boldsymbol{b}}_k)^{(-1)} = (\hat{\boldsymbol{\Gamma}}_k \tilde{\boldsymbol{b}}_k)^{(0)} = 0$ . Here,  $\tilde{\boldsymbol{b}}_k^{(t)} = (\tilde{\boldsymbol{b}}_k^{(-1)}, \cdots, \tilde{\boldsymbol{b}}_k^{(t)})^\top$ ,  $\hat{\boldsymbol{\Gamma}}_k$  is a  $(t+1) \times (t+1)$  matrix whose (s,s') element is  $\hat{\boldsymbol{\Gamma}}_k^{(s,s')}$ , and  $G^{(t,t-1)}$  denotes the average single-user response function.

Detail of the parameter  $G^{(t,t-1)}$  is introduced later (in Section V-B). Note that  $G^{(t,t-1)}$  does not depend on the user index k.

## C. BP-based Detector

Assuming that information bits **b** are independently generated from the symmetric distribution, the posterior distribution from received signals **y** is given as  $p(\mathbf{b}|\mathbf{y}) = p(\mathbf{y}|\mathbf{b}) / \sum_{\mathbf{b} \in \{\pm 1\}^K} p(\mathbf{y}|\mathbf{b})$ , where

$$p(\boldsymbol{y}|\boldsymbol{b}) = \prod_{\mu=1}^{N} \frac{1}{\sqrt{2\pi\sigma}} \exp\left[-\frac{1}{2\sigma^2} \left(y^{\mu} - \frac{1}{\sqrt{N}} \sum_{k=1}^{K} s_k^{\mu} b_k\right)^2\right],\tag{7}$$

and  $\sigma^2$  is a control parameter when true noise level parameter  $\sigma_0^2$  is not known. The marginal-posterior-mode detector [51] is represented by

$$\tilde{b}_k = \operatorname*{argmax}_{b_k \in \{-1,1\}} \left( \sum_{\boldsymbol{b} \setminus b_k \in \{-1,1\}^{K-1}} p(\boldsymbol{b} | \boldsymbol{y}) \right).$$
(8)

The BP-based detector is an iterative algorithm that employs the belief propagation to approximately calculate the posterior marginal included in (8).

*Definition 3:* (BP-based detector) The BP-based detector [25] is given by the following iterative equations:

$$\tilde{\boldsymbol{b}}^{(t+1)} = \tanh(R^{(t)}\boldsymbol{h} - \boldsymbol{U}^{(t)} + A^{(t)}\tilde{\boldsymbol{b}}^{(t)}), \qquad (9)$$

$$\boldsymbol{U}^{(t)} = A^{(t)} \boldsymbol{W} \tilde{\boldsymbol{b}}^{(t)} + A^{(t)} \beta (1 - Q^{(t)}) \boldsymbol{U}^{(t-1)}, \qquad (10)$$

where

$$R^{(t)} = A^{(t)} + A^{(t)}\beta(1 - Q^{(t)})R^{(t-1)}, \qquad (11)$$

$$A^{(t)} = \frac{1}{\sigma^2 + \beta(1 - Q^{(t)})},$$
(12)

$$Q^{(t)} = \frac{1}{K} \sum_{k=1}^{K} (\tilde{b}_k^{(t)})^2,$$
(13)

with initial conditions:  $R^{(-1)} = A^{(-1)}$  and  $\tilde{b}_k^{(-1)} = 0$ . Function tanh is applied componentwise. From posterior average  $\tilde{b}_k^{(t)}$ , the tentative decision at the  $t^{\text{th}}$  update is evaluated as  $\hat{b}_k^{(t)} = \operatorname{sgn}(\tilde{b}_k^{(t)})$ , where function  $\operatorname{sgn}(x)$  denotes the sign function taking 1 for  $x \ge 0$  and -1 for x < 0.

This BP-based detector can be rewritten as

$$\tilde{b}_{k}^{(t+1)} = \tanh\left(R^{(t)}h_{k} + A^{(t)}\tilde{b}_{k}^{(t)} - \sum_{s=-1}^{t} J^{(t,s)}\sum_{k'=1}^{K} W_{kk'}\tilde{b}_{k}^{(s)}\right), \quad (14)$$

where

$$J^{(t,s)} = \begin{cases} 0, & s = -1 \\ A^{(s)} \prod_{s'=s+1}^{t} A^{(s')} \beta [1 - q^{(s')}], & 0 \le s \le t - 1 \\ A^{(t)}, & s = t. \end{cases}$$
(15)

## IV. GENERATING FUNCTIONAL ANALYSIS

In this section, we briefly summarize methods on GFA [11]. Some books that introduces GFA are available, e.g., the analyses of minority games [7] and spin glasses [18], [12].

First of all, let us compare GFA with the replica analysis [49], [51], [43], [30]. Both have been developed in the literature of statistical mechanics [17], [28] and have been applied to the analysis of problems in the field of information theory so far.

# A. Replica Analysis Versus Generating Functional Analysis

In the replica analysis, the goal is to understand macroscopic *static* properties in the large-system limit, i.e., the number of interacting bodies becomes large. One evaluate the *free energy* that is proportinal to a logarithm of the partition function. The partition function is a logarithm of a Gibbs distribution. In the static computation based on the replica analysis, the average over disorder generates a coupling between distinct *replicas*.

In GFA, the goal is, on the other hand, to understand macroscopic *dynamic* properties in the large-system limit. One evaluates not the free energy but the *generating functional* that is a kind of the characteristic function of the path probability. The generating functional is not a static object but contains dynamical information. In the dynamical computation, we do not have to introduce replicas. The effect of disorder is to generate a coupling between distinct *times*. Under some assumptions concerning stationary states, one can obtain static properties from GFA. We can, therefore, avoid the replica trick and obtain macroscopic static properties by GFA if the dynamics reaches a stationary state. From this point of view, GFA can be regarded as one of alternative methods for the replica analysis.

In the next two subsections, we briefly introduce both the replica analysis and GFA, and discuss the difference between them in more detail.

## B. Outline of Replica Analysis

Suppose that we want to infer the *Marginal-Posterior-Mode* (*MPM*) performance of a system. Let vectors  $\boldsymbol{x} = (x_k) \in \mathcal{X}^K$  be the states of the model, where  $\mathcal{X}$  denotes a set of values of each element. Let  $p(\boldsymbol{x}; \boldsymbol{w}) = [Z(\beta; \boldsymbol{w})]^{-1} e^{-\beta H(\boldsymbol{x}; \boldsymbol{w})}$  be a posterior distribution in the Bayesian framework, where

w denote parameters in it and  $\beta$  denotes a non-negative parameter that is called the *inverse temperature* and commonly corresponds to a noise amplitude, i.e., a larger  $\beta$  gives a smaller noise. (Note that  $\beta$  denotes the system load except in Section IV.) A distribution of this form is called the *Gibbs distribution*. The function is referred to as a cost function or the *Hamiltonian*. The normalization constant of the posterior distribution

$$Z(\beta; \boldsymbol{w}) = \sum_{\boldsymbol{x} \in \mathcal{X}^{K}} e^{-\beta H(\boldsymbol{x}; \boldsymbol{w})}$$
(16)

is called the partition function.

The Gibbs distribution is derived by extremization (maximization) of the entropy  $S(\beta; w) := -\sum_{x \in \mathcal{X}^K} p(x; w) \ln p(x; w)$  under an average cost  $U(\beta; w) := \sum_{x \in \mathcal{X}^K} H(x)p(x; w)$  and a parameter w are fixed. When we fix the parameter  $\beta$  instead of the average cost, this relationship can be rewritten as follows. Namely, for a fixed  $\beta$  and a fixed w, a functional  $\tilde{F}[q(x; w)] = \sum_{x \in \mathcal{X}^K} H(x; w)q(x; w) + \beta^{-1} \sum_{x \in \mathcal{X}^K} q(x; w) \ln q(x; w)$  with respect to q(x; w) takes an extremum (minimum) value at q(x; w) = p(x; w). The extremum (minimum) value of the functional is  $\tilde{F}[p(x; w)] = -\beta^{-1}Z(\beta; w) =: F(\beta; w)$ . This value is an important potential, called the (Helmholtz) free energy. From the free energy one can evaluate various quantities such as an average cost and the entropy for a fixed  $\beta$ , e.g.,

$$U(\beta; \boldsymbol{w}) = \frac{\partial [\beta F(\beta; \boldsymbol{w})]}{\partial \beta},$$
  
$$S(\beta; \boldsymbol{w}) = \beta^2 \frac{\partial F(\beta; \boldsymbol{w})}{\partial \beta}.$$

by taking derivatives with respect to  $\beta$ . Since the free energy  $F(\beta; w)$  is often proportional to K, the free energy density  $f_K(\beta; w) := F(\beta; w)/K$ , that is the free energy per an interacting body, is defined.

For a given realization  $\boldsymbol{w}$  it is, however, hard to calculate the free energy density, since the result must depend on detail of the realization. We therefore assume that the limit  $\lim_{K\to\infty} f_K(\beta; \boldsymbol{w}) =: \tilde{\mathcal{F}}$  exists and it is equal to its average  $\lim_{K\to\infty} K^{-1}\mathbb{E}[F(\beta; \boldsymbol{w})]$  for almost all realizations, where the  $\mathbb{E}_X$  denotes the expectation with respect to random variables X. Using the identity  $\ln x = \lim_{n\to 0} \frac{\partial x^n}{\partial n}$ , we have

$$\tilde{\mathcal{F}} = \lim_{K \to \infty} \left( \lim_{n \to 0} \frac{\partial \mathbb{E}_{\boldsymbol{w}}[Z(\beta; \boldsymbol{w})^n]}{\partial n} \right).$$
(17)

The idea of the replica analysis is to calculate  $\mathbb{E}_{\boldsymbol{w}}[Z(\beta; \boldsymbol{w})^n]$ as if n were in integer. The *n*-th moment is

$$Z(\beta; \boldsymbol{w})^{n} = \sum_{\boldsymbol{x}^{1}, \cdots, \boldsymbol{x}^{n} \in \mathcal{X}^{K}} \exp\left[-\beta \sum_{a=1}^{n} H(\boldsymbol{x}^{a}; \boldsymbol{w})\right], \quad (18)$$

where a is a replica index. We shall refer to copies  $x^1, \dots, x^n$  as *replicas*. It should be noted that the replicas are no longer statistically independent since the replicas have a common parameter w. The average over the parameter w generates a coupling between distinct replicas.

In the replica analysis, one therefore has to assume (i) the self-averaging property applies, (ii) the "replica trick" is valid,

and (iii) replica symmetry (or same kind of symmetries) holds, to keep tractability. The replica analysis, however, still does not have a rigorous justification.

## C. Outline of Generating Functional Analysis

The generating functional analysis, or the *path integral methods*, have been applied to the model which is described using realizations of random variables by de Dominicis [11]. In GFA, one can analyze the asymptotic behavior of the dynamics in the infinitely large system using the generating functional that can be regarded as a kind of the characteristic function in statistics.

We here consider the following model. Let vectors  $x^{(s)} = (x_k^{(s)}) \in \mathbb{R}^K$  be the states of the model at stage s and let the initial state be  $x^{(-1)}$ . Let the updating rule be

$$\boldsymbol{x}^{(s+1)} = \mathcal{F}(\boldsymbol{x}^{(s)}; \boldsymbol{w}) \tag{19}$$

for  $s \in \{-1, 0, \dots, t-1\}$ , where  $\mathcal{F} : \mathbb{R} \to \mathbb{R}$  denotes a function and w again denotes parameters in it.

If we know the probability of seeing a particular state at a given stage  $p_s(\boldsymbol{x}^{(s)})$ , we can evaluate the property of the system by using it. In GFA we consider the probability of observing a particular sequence or path of states, i.e.,  $p(\boldsymbol{x}^{(-1)}, \boldsymbol{x}^{(0)}, \dots, \boldsymbol{x}^{(t)})$  up to some finite time t instead of  $p_s(\boldsymbol{x}^{(s)})$ . The probability  $p(\boldsymbol{x}^{(-1)}, \boldsymbol{x}^{(0)}, \dots, \boldsymbol{x}^{(t)})$  is referred to as the *path probability*. The way to do this is to introduce a generating functional which is defined as

$$Z[\boldsymbol{\psi}] = \left\langle \exp\left[-i\sum_{s=-1}^{t}\sum_{k=1}^{K}\psi_{k}^{(s)}x_{k}^{(s)}\right]\right\rangle,$$
(20)

where the bracket  $\langle \cdot \rangle$  denotes the average over the path probability  $p(\boldsymbol{x}^{(-1)}, \boldsymbol{x}^{(0)}, \cdots, \boldsymbol{x}^{(t)})$  and we have introduced the dummy variables  $\boldsymbol{\psi}^{(s)} = (\boldsymbol{\psi}_k^{(s)}) \in \mathbb{R}^K$ . Note that the generating functional differs from the partition function in the previous subsection. Taking derivatives with respect to the dummy variables allow us to examine some averages, e.g.,

$$\langle x_k^{(s)} \rangle = \mathrm{i} \lim_{\psi \to \mathbf{0}} \frac{\partial Z[\psi]}{\partial \psi_k^{(s)}},\tag{21}$$

$$\langle x_k^{(s)} x_{k'}^{(s')} \rangle = -\lim_{\psi \to \mathbf{0}} \frac{\partial^2 Z[\psi]}{\partial \psi_k^{(s)} \partial \psi_k^{(s')}}.$$
 (22)

Since one does not have to average a ratio or a logarithm in this context, one can compute correlations by entirely avoiding the replica trick. In the replica analysis, macroscopic quantities are obtained from derivatives with respect to the "scalar" parameter  $\beta$ . In GFA, averages are given from derivatives with respect to not a scalar parameter but some elements of the dummy variables (vectors) with same dimension as the original vectors  $\boldsymbol{x}^{(-1)}, \boldsymbol{x}^{(0)}, \cdots, \boldsymbol{x}^{(t)}$ .

For a given realization w it is, again, hard to calculate the generating functional since the result must depend on detail of the realization. We, therefore, assume that the generating functional is concentrated to its average over the parameter  $\mathbb{E}_{w}(Z[\psi])$  in the large system limit. The effect of the parameter w is to generate a coupling between distinct times. Averaging over the random variables, we will move to a

saddle-point problem in the limit  $K \to \infty$ . It should be noted that the normalization relation  $Z[\mathbf{0}] = 1$  plays an important role in the elimination of spurious solutions to the saddle-point equations. (We see detail of this point in Section V-B.)

The terms in the averaged generating functional can be split into three related parts. The first one is a signal part. The second one is a static noise part due to the random variables within the model. The last one is retarded self-interaction due to the influence of the state at the previous stage, which may be able to affect the present state. The GFA allows us to treat the last part.

## V. GENERATING FUNCTIONAL ANALYSIS FOR SOFT-PIC

## A. Averaged Generating Functional

The goal of multiuser detection is to simultaneously infer information bits  $b_1, \dots, b_K$  after base-band signals  $y_1, \dots, y_N$ are received.

If the transfer function takes f(x) = sgn(x), this updating rule coincides with the *hard-PIC* [54]. Note that the hard-PIC, which means the Varanasi and Aazhang's conventional PIC, is also obtained by taking the limit  $\sigma \to 0$  in soft-PIC.

We assume the matched filter stage, i.e.,  $b_k(0) = f(h_k)$ , for initialization. This initialization is easily treated by formally assuming  $p(\tilde{b}_k^{(-1)}) = \delta(\tilde{b}_k^{(-1)})$  for all k, where  $\delta$  denotes the Dirac delta function. The widely used measure to determine the performance of a demodulator is the *bit error rate (BER)*. The BER  $P_b^{(t)}$  of hard decisions  $\hat{b}_k^{(t)} = \operatorname{sgn}(\tilde{b}_k^{(t)})$  at the  $t^{\text{th}}$ stage of soft PIC is given by  $P_b^{(t)} = (1 - m_h^{(t)})/2$ , where  $m_h^{(t)} = \frac{1}{K} \sum_{k=1}^K b_k \tilde{b}_k^{(t)}$  is the overlap between information vector  $\mathbf{b}^{(t)} = (b_1, \cdots, b_K)^{\top}$  and tentative hard-decision vector  $(\tilde{b}_1^{(t)}, \cdots, \tilde{b}_K^{(t)})^{\top}$ . Without loss of generality, we can assume that the true information bits are all 1, i.e.,  $b_k = 1$  for all k, because the spreading codes are symmetric.

Let us analyze the detection dynamics in the large system limit where  $K, N \to \infty$ , while the system load  $\beta$  is kept finite. We introduce inverse temperature  $\gamma$  for generating functional analysis. The stochastic updating rule for tentative decision  $\tilde{b}_k^{(t)} \in \mathbb{R}$  of bit signal  $b_k$  at stage t is given by

$$p(\tilde{b}_{k}^{(s+1)}|\tilde{b}^{(s)}) = \frac{\gamma}{\sqrt{2\pi}} e^{-\frac{\gamma^{2}}{2}\{\tilde{b}_{k}^{(s+1)} - f(u_{k}^{(s)})\}^{2}}, \quad (23)$$

with

$$u_k^{(t)} := h_k - \sum_{k'=1, \neq k}^K W_{kk'} \tilde{b}_{k'}^{(t)} + \theta_k^{(t)}, \qquad (24)$$

which is a summation over all messages from other tentative decisions. Note that this updating rule coincides with  $\tilde{b}_k^{(s+1)} = f(u_k^{(s)}) + \mathcal{N}(0, \gamma^{-2})$ . In the limit where  $\gamma \to \infty$ , this stochastic updating rule is equivalent to the deterministic rule (2). Term  $\theta_k^{(t)}$  is a stage-dependent external message that is introduced to define a response function. The inverse temperature and the external message are respectively set to  $\gamma \to \infty$  and  $\theta_k^{(t)} = 0$  at the end of analysis.

The stochastic updating rule for tentative decision vector  $\tilde{\boldsymbol{b}}^{(t)} = (\tilde{b}_k^{(t)}) \in \mathbb{R}^K$  at stage t is given by using (23), i.e.,  $p(\tilde{\boldsymbol{b}}^{(s+1)}|\tilde{\boldsymbol{b}}^{(s)}) = \prod_{k=1}^K p(\tilde{b}_k^{(s+1)}|\tilde{\boldsymbol{b}}^{(s)})$ . The dynamics is a Markov chain, since the present tentative decision depends

only on the past decisions. A path probability (density)  $p(\tilde{b}^{(-1)}, \dots, \tilde{b}^{(t)})$  is therefore simply given by the individual transition probability  $p(\tilde{b}^{(s+1)}|\tilde{b}^{(s)})$  of the chain:

$$p(\tilde{\boldsymbol{b}}^{(-1)}, \cdots, \tilde{\boldsymbol{b}}^{(t)}) = p(\tilde{\boldsymbol{b}}^{(-1)}) \prod_{s=-1}^{t-1} p(\tilde{\boldsymbol{b}}^{(s+1)} | \tilde{\boldsymbol{b}}^{(s)}).$$
 (25)

The initial state probability becomes  $p(\tilde{\boldsymbol{b}}^{(-1)}) = \prod_{k=1}^{K} \delta(\tilde{b}_{k}^{(-1)})$ . Therefore, we can calculate an expectation with respect to arbitrary function  $\mathcal{G} = \mathcal{G}(\tilde{\boldsymbol{b}}^{(-1)}, \dots, \tilde{\boldsymbol{b}}^{(t)})$  of tentative decisions as

$$\langle \mathcal{G} \rangle := \int_{\mathbb{R}^{(t+2)K}} \left( \prod_{s=-1}^{t} d\tilde{\boldsymbol{b}}^{(s)} \right) p(\tilde{\boldsymbol{b}}^{(-1)}, \cdots, \tilde{\boldsymbol{b}}^{(t)}) \mathcal{G}.$$
(26)

To analyze the detection dynamics of the system, we define a generating functional as  $Z[\psi] := \langle \exp[-i\sum_{s=-1}^{t} \tilde{\boldsymbol{b}}^{(s)} \cdot \boldsymbol{\psi}^{(s)}] \rangle$ , where  $\tilde{\boldsymbol{b}}^{(s)} = (\tilde{b}_{1}^{(s)}, \cdots, \tilde{b}_{K}^{(s)})^{\top}$  and  $\boldsymbol{\psi}^{(s)} = (\boldsymbol{\psi}_{1}^{(s)}, \cdots, \boldsymbol{\psi}_{K}^{(s)})^{\top}$ . The basic idea underlying generating functional formalism is very simple [11], [6], [7]. If the generating functional  $Z[\boldsymbol{\psi}]$  can be evaluated as a functional with respect to dummy functions  $\boldsymbol{\psi}^{(s)}$ , one can obtain all averages of interest by differentiation from  $Z[\boldsymbol{\psi}]$ . The generating functional includes the random variables  $\{s_k^{k}\}$  and  $\{n^{\mu}\}$ .

We here assume that the generating functional is selfaveraging, namely, in the large system limit, the generating functional is concentrated to its average over random variables  $\{s_k^{\mu}\}$  and  $\{n^{\mu}\}$  and the typical behavior of the system only depends on the statistical properties of the random variables. We therefore evaluate the averaged generating functional defined as follows.

Definition 4: (Average Generating functional) The average generating functional  $\overline{Z}[\psi]$  is defined by

$$\bar{Z}[\boldsymbol{\psi}] := \left\langle \exp\left[-i\sum_{s=-1}^{t} \tilde{\boldsymbol{b}}^{(s)} \cdot \boldsymbol{\psi}^{(s)}\right] \right\rangle, \quad (27)$$

where  $\overline{[\cdots]} := \mathbb{E}_{s_1, \cdots, s_K, n} [\cdots]$  denotes the average over spreading codes  $s_1, \cdots, s_K$  and the noise n.

From the averaged generating functional, we can obtain

$$\lim_{\boldsymbol{\psi}\to\mathbf{0}}\frac{\partial Z[\boldsymbol{\psi}]}{\partial \psi_k^{(s)}} = -\mathrm{i}\overline{\langle \tilde{b}_k^{(s)} \rangle},\tag{28}$$

$$\lim_{\boldsymbol{\psi}\to\boldsymbol{0}}\frac{\partial^2 \bar{Z}[\boldsymbol{\psi}]}{\partial\psi_k^{(s)}\partial\psi_{k'}^{(s')}} = -\overline{\langle \tilde{b}_k^{(s)}\tilde{b}_{k'}^{(s')}\rangle},\tag{29}$$

$$\lim_{\boldsymbol{\psi}\to\mathbf{0}}\frac{\partial^2 \bar{Z}[\boldsymbol{\psi}]}{\partial \psi_k^{(s)}\partial \theta_{k'}^{(s')}} = -\mathrm{i}\frac{\partial \langle \tilde{b}_k^{(s)} \rangle}{\partial \theta_{k'}^{(s')}}.$$
(30)

Calculating the average over spreading codes  $s_1, \dots, s_K$  and the noise n, we have the following result.

*Lemma 1:* The averaged generating functional is simplified to

$$\bar{Z}[\psi] = \int d\eta d\hat{\eta} d\mathbf{k} d\mathbf{k} d\mathbf{q} d\hat{\mathbf{q}} d\mathbf{Q} d\mathbf{Q} d\mathbf{L} d\hat{\mathbf{L}}$$
$$\times \exp\left[K(\Phi + \Psi + \Omega) + O(\ln K)\right], (31)$$

in which functions  $\Phi$ ,  $\Psi$ , and  $\Omega$  are given by

$$\Phi := i \sum_{s=-1}^{t-1} \{ \hat{\eta}^{(s)} \eta^{(s)} + \hat{k}^{(s)} k^{(s)} \}$$
$$+ i \sum_{s=-1}^{t-1} \sum_{s'=-1}^{t-1} \{ \hat{q}^{(s,s')} q^{(s,s')} + \hat{L}^{(s,s')} L^{(s,s')} \}, \quad (32)$$

$$\Psi := \frac{1}{K} \sum_{k=1}^{K} \ln \left\{ \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t} \mathrm{d}\tilde{b}^{(s)} \right) p[\tilde{b}^{(-1)}] \int \underline{\mathrm{d}}u \underline{\mathrm{d}}\hat{u} \\ \times \exp \left[ \sum_{s=-1}^{t-1} \left\{ \ln \frac{\gamma}{\sqrt{2\pi}} - \frac{\gamma^{2}}{2} [\tilde{b}^{(s+1)} - f(u^{(s)})]^{2} \right\} \\ -\mathrm{i} \sum_{s=-1}^{t-1} \sum_{s'=-1}^{t-1} \left\{ \hat{q}^{(s,s')} \tilde{b}^{(s)} \tilde{b}^{(s')} \\ + \hat{Q}^{(s,s')} \hat{u}^{(s)} \hat{u}^{(s')} + \hat{L}^{(s,s')} \tilde{b}^{(s)} \hat{u}^{(s')} \right\} \\ +\mathrm{i} \sum_{s=-1}^{t-1} \hat{u}^{(s)} \left\{ u^{(s)} - \tilde{b}^{(s)} - \theta_{k}^{(s)} - \hat{k}^{(s)} \right\} \\ -\mathrm{i} \sum_{s=-1}^{t-1} \tilde{b}^{(s)} \hat{\eta}^{(s)} - \mathrm{i} \sum_{s=-1}^{t} \tilde{b}^{(s)} \psi_{k}^{(s)} \right] \right\},$$
(33)

$$\begin{split} \Omega &:= \frac{1}{K} \ln \int \underline{d} \boldsymbol{v} \underline{d} \hat{\boldsymbol{v}} \underline{d} \boldsymbol{w} \underline{d} \hat{\boldsymbol{w}} \\ &\times \exp \left[ i \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \{ \hat{v}_{\mu}^{(s)} v_{\mu}^{(s)} + \hat{w}_{\mu}^{(s)} w_{\mu}^{(s)} - \beta v_{\mu}^{(s)} w_{\mu}^{(s)} \} \right. \\ &\left. - \frac{1}{2} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \sum_{s'=-1}^{t-1} \{ \beta \sigma_{0}^{2} v_{\mu}^{(s)} v_{\mu}^{(s')} + \hat{v}_{\mu}^{(s)} Q^{(s,s')} \hat{v}_{\mu}^{(s')} \} \right. \\ &\left. - \frac{1}{2} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \sum_{s'=-1}^{t-1} \{ \hat{v}_{\mu}^{(s)} [k^{(s)} - L^{(s',s)}] \hat{w}_{\mu}^{(s')} \right. \\ &\left. + \hat{w}_{\mu}^{(s)} [k^{(s')} - L^{(s,s')}] \hat{v}_{\mu}^{(s')} \} \right. \\ &\left. - \frac{1}{2} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \sum_{s'=-1}^{t-1} \{ \hat{w}_{\mu}^{(s)} \\ &\left. \times [1 - \eta^{(s)} - \eta^{(s')} + q^{(s,s')}] \hat{w}_{\mu}^{(s')} \} \right]. \end{split}$$
(34)

Details on the derivation and definitions of the notations are given in Appendix B.

# B. Saddle-Point Equations and Meaning of Macroscopic Parameters

One can deduce the meaning of macroscopic parameters by differentiating the averaged generating functional  $\bar{Z}[\psi]$  of (31) with respect to the external messages  $\theta_k^{(s)}$  and dummy functions  $\psi_k^{(s)}$ . The averaged generating functional  $\bar{Z}[\psi]$  is dominated by a saddle-point for  $K \to \infty$ . We can thus simplify the saddle-point problem in Lemma 1 in the large system limit. Using the shorthand  $\{d\} := d\eta d\hat{\eta} dk d\hat{k} dq d\hat{q} dQ d\hat{Q} dL d\hat{L}$ 

and the normalization identity  $\overline{Z}[\mathbf{0}] = \langle 1 \rangle = 1$ , we now find derivatives of the averaged generating functional:

$$\lim_{\psi \to \mathbf{0}} \frac{\partial Z[\psi]}{\partial \psi_k^{(s)}} = \lim_{\psi \to \mathbf{0}} \int \{d\} e^{K(\Phi + \Psi + \Omega)} \frac{\partial (K\Psi)}{\partial \psi_k^{(s)}} = -i \langle \tilde{b}^{(s)} \rangle_k, \quad (35)$$

$$\lim_{\psi \to \mathbf{0}} \frac{\partial^2 \bar{Z}[\psi]}{\partial \psi_k^{(s)} \partial \psi_{k'}^{(s')}} = \lim_{\psi \to \mathbf{0}} \int \{d\} e^{K(\Phi + \Psi + \Omega)} \times \left( \frac{\partial^2 (K\Psi)}{\partial \psi_k^{(s)} \partial \psi_{k'}^{(s')}} + \frac{\partial (K\Psi)}{\partial \psi_k^{(s)}} \cdot \frac{\partial (K\Psi)}{\partial \psi_{k'}^{(s')}} \right) = \delta_{k,k'} [-\langle \tilde{b}^{(s)} \tilde{b}^{(s')} \rangle_k + \langle \tilde{b}^{(s)} \rangle_k \langle \tilde{b}^{(s')} \rangle_{k'}] - \langle \tilde{b}^{(s)} \rangle_k \langle \tilde{b}^{(s')} \rangle_{k'} = -\delta_{k,k'} \langle \tilde{b}^{(s)} \tilde{b}^{(s')} \rangle_k - (1 - \delta_{k,k'}) \langle \tilde{b}^{(s)} \rangle_k \langle \tilde{b}^{(s')} \rangle_{k'} \quad (36)$$

$$\lim_{\psi \to \mathbf{0}} \frac{\partial \mathcal{L}[\Psi]}{\partial \psi_k^{(s)} \partial \theta_{k'}^{(s')}} = \lim_{\psi \to \mathbf{0}} \int \{d\} e^{K(\Phi + \Psi + \Omega)} \\ \times \left( \frac{\partial^2 (K\Psi)}{\partial \psi_k^{(s)} \partial \theta_{k'}^{(s')}} + \frac{\partial (K\Psi)}{\partial \psi_k^{(s)}} \cdot \frac{\partial (K\Psi)}{\partial \theta_{k'}^{(s')}} \right) \\ = -\delta_{k,k'} \langle \tilde{b}^{(s)} \hat{u}^{(s')} \rangle_k - (1 - \delta_{k,k'}) \langle \tilde{b}^{(s)} \rangle_k \langle \hat{u}^{(s')} \rangle_{k'}$$
(37)

$$\lim_{\psi \to \mathbf{0}} \frac{\partial Z[\psi]}{\partial \theta_k^{(s)}} = -i \langle \hat{u}^{(s)} \rangle_k \stackrel{(a)}{=} 0, \tag{38}$$

$$\lim_{\boldsymbol{\psi}\to\mathbf{0}}\frac{\partial^2 Z[\boldsymbol{\psi}]}{\partial\theta_k^{(s)}\partial\theta_{k'}^{(s')}} = -\delta_{k,k'} \langle \hat{u}^{(s)} \hat{u}^{(s')} \rangle_k \stackrel{(b)}{=} 0, \tag{39}$$

where  $\langle \rangle_k$  denotes the average as

$$\frac{\langle f(\lbrace b, u, \hat{u} \rbrace) \rangle_k :=}{\int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^t \mathrm{d}\tilde{b}^{(s)} \right) \int \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} w_k(\lbrace \tilde{b}, u, \hat{u} \rbrace) f(\lbrace \tilde{b}, u, \hat{u} \rbrace)}{\int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^t \mathrm{d}\tilde{b}^{(s)} \right) \int \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} w_k(\lbrace \tilde{b}, u, \hat{u} \rbrace)} \tag{40}$$

with

$$w_{k}(\{\tilde{b}, u, \hat{u}\})$$

$$:= \delta[\tilde{b}^{(-1)}] \exp\left[\sum_{s=-1}^{t-1} \{\ln \frac{\gamma}{\sqrt{2\pi}} - \frac{\gamma^{2}}{2} [\tilde{b}^{(s+1)} - f(u^{(s)})]^{2}\}\right]$$

$$-i \sum_{s=0}^{t-1} \sum_{s'=0}^{t-1} \{\hat{q}^{(s,s')} \tilde{b}^{(s)} \tilde{b}^{(s')} + \hat{Q}^{(s,s')} \hat{u}^{(s)} \hat{u}^{(s')}$$

$$+ \hat{L}^{(s,s')} \tilde{b}^{(s)} \hat{u}^{(s')}\}$$

$$+ i \sum_{s=0}^{t-1} \hat{u}^{(s)} \{u^{(s)} - \tilde{b}^{(s)} - \theta_{k}^{(s)} - \hat{k}^{(s)}\}$$

$$-i \sum_{s=0}^{t-1} \tilde{b}^{(s)} \hat{\eta}^{(s)} \Big]_{saddle}, \qquad (41)$$

which is referred to as a *single-user measure*. Here, evaluation  $f|_{\text{saddle}}$  denotes an evaluation of function f at the dominating saddle-point. To derive the  $=^{(a)}$  in (38) and  $=^{(b)}$  in (39), we use identities

$$\lim_{\boldsymbol{\psi}\to\mathbf{0}} \frac{\partial \bar{Z}[\boldsymbol{\psi}]}{\partial \theta_k^{(s)}} = \frac{\partial \bar{Z}[\mathbf{0}]}{\partial \theta_k^{(s)}} = 0, \tag{42}$$

$$\lim_{\boldsymbol{\psi}\to\boldsymbol{0}} \frac{\partial^2 \bar{Z}[\boldsymbol{\psi}]}{\partial \theta_k^{(s)} \partial \theta_{k'}^{(s')}} = \frac{\partial^2 \bar{Z}[\boldsymbol{0}]}{\partial \theta_k^{(s)} \partial \theta_{k'}^{(s')}} = 0, \qquad (43)$$

respectively. These are obtained from  $\overline{Z}[\mathbf{0}] = 1$ . Substituting (28) – (30) and (38) – (39) into (35) – (37), the spurious solutions, that depend on  $\langle \hat{u}^{(s)} \rangle_k$  or  $\langle \hat{u}^{(s)} \hat{u}^{(s')} \rangle_k$ , are eliminated and we then have

$$\overline{\langle \tilde{b}_{k}^{(s)} \rangle} = \langle \tilde{b}^{(s)} \rangle_{k}, \qquad (44)$$
$$\overline{\langle \tilde{b}^{(s)} \tilde{b}^{(s')} \rangle} = \delta_{k+1} \langle \tilde{b}^{(s)} \tilde{b}^{(s')} \rangle_{k}.$$

$$\frac{\partial \langle \tilde{b}_{k'}^{(s)} \rangle}{\partial \langle \tilde{b}_{k'}^{(s)} \rangle} = \delta_{k,k'} \langle \delta^{(s)} \rangle_{k'} \langle \tilde{b}^{(s)} \rangle_{k'} \langle \tilde{b}^{(s')} \rangle_{k'}, \quad (45)$$

$$\frac{\partial \langle \hat{b}_{k}^{(s')} \rangle}{\partial \theta_{k'}^{(s')}} = -\mathrm{i}\delta_{k,k'} \langle \tilde{b}^{(s)} \hat{u}^{(s')} \rangle_{k}. \tag{46}$$

In the large system limit  $K \to \infty$ , integral (31) will be evaluated by the dominating saddle-point of exponent  $\Phi + \Psi + \Omega$ . We can now derive the saddle-point equations by differentiation with respect to integral variables  $\eta^{(s)}$ ,  $\hat{\eta}^{(s)}$ ,  $k^{(s)}$ ,  $\hat{k}^{(s)}$ ,  $q^{(s,s')}$ ,  $\hat{q}^{(s,s')}$ ,  $Q^{(s,s')}$ ,  $L^{(s,s')}$ , and  $\hat{L}^{(s,s')}$ . These equations will involve the average overlap  $m^{(s)}$  (which measures the bit error rate), the average single-user correlation  $C^{(s,s')}$  and the average single-user response function  $G^{(s,s')}$ :

$$m^{(s)} := \lim_{K \to \infty} \frac{1}{K} \sum_{k=1}^{K} \overline{\langle \tilde{b}_k^{(s)} \rangle}, \tag{47}$$

$$C^{(s,s')} := \lim_{K \to \infty} \frac{1}{K} \sum_{k=1}^{K} \overline{\langle \tilde{b}_k^{(s)} \tilde{b}_k^{(s')} \rangle}, \qquad (48)$$

$$G^{(s,s')} := \lim_{K \to \infty} \frac{1}{K} \sum_{k=1}^{K} \frac{\partial \overline{\langle \tilde{b}_k^{(s)} \rangle}}{\partial \theta_k^{(s')}}.$$
 (49)

Using the identities (35) – (39) and (44) – (46), the straightforward differentiation of  $\Phi + \Psi + \Omega$  with respect to  $\eta^{(s)}$ ,  $\hat{\eta}^{(s)}$ ,  $k^{(s)}$ ,  $\hat{k}^{(s)}$ ,  $q^{(s,s')}$ ,  $\hat{q}^{(s,s')}$ ,  $Q^{(s,s')}$ ,  $\hat{Q}^{(s,s')}$ ,  $L^{(s,s')}$ , and  $\hat{L}^{(s,s')}$  leads us to the following saddle-point equations:

$$\hat{\eta}^{(s)} = i \left. \frac{\partial \Omega}{\partial \eta^{(s)}} \right|_{\text{saddle}},\tag{50}$$

$$\eta^{(s)} = \lim_{K \to \infty} \frac{1}{K} \sum_{k=1}^{K} \langle \tilde{b}^{(s)} \rangle_k, \tag{51}$$

$$\hat{k}^{(s)} = i \left. \frac{\partial \Omega}{\partial k^{(s)}} \right|_{\text{saddle}},\tag{52}$$

$$k^{(s)} = \frac{1}{K} \sum_{k=1}^{K} \langle \hat{u}^{(s)} \rangle_k = 0,$$
(53)

$$\hat{q}^{(s,s')} = i \left. \frac{\partial \Omega}{\partial q^{(s,s')}} \right|_{\text{saddle}},\tag{54}$$

$$q^{(s,s')} = \lim_{K \to \infty} \frac{1}{K} \sum_{k=1}^{K} \langle \tilde{b}^{(s)} \tilde{b}^{(s')} \rangle_k, \tag{55}$$

$$^{(s,s')} = \mathrm{i} \left. \frac{\partial \Omega}{\partial Q^{(s,s')}} \right|_{\mathrm{saddle}},$$
(56)

$$Q^{(s,s')} = \frac{1}{K} \sum_{k=1}^{K} \langle \hat{u}^{(s)} \hat{u}^{(s')} \rangle_k = 0,$$
 (57)

$$\hat{L}^{(s,s')} = i \left. \frac{\partial \Omega}{\partial L^{(s,s')}} \right|_{\text{saddle}},\tag{58}$$

$$L^{(s,s')} = \lim_{K \to \infty} \frac{1}{K} \sum_{k=1}^{K} \langle \tilde{b}^{(s)} \hat{u}^{(s')} \rangle_k,$$
(59)

for all s and s', respectively. Comparing them with (47) – (49), the following relationships are obtained:

$$\eta^{(s)} = m^{(s)}, \tag{60}$$

$$q^{(s,s')} = C^{(s,s')}, (61)$$

$$L^{(s,s')} = iG^{(s,s')}.$$
 (62)

Therefore, we hereafter make use of  $\{m^{(s)}, C^{(s,s')}, iG^{(s,s')}\}$  instead of  $\{\eta^{(s)}, q^{(s,s')}, L^{(s,s')}\}$ . It should be noted that causality

$$\frac{\partial \langle \hat{b}^{(s)} \rangle}{\partial \theta^{(s')}} = 0, \tag{63}$$

should hold for  $s \leq s'$ , therefore  $L^{(s,s')} = G^{(s,s')} = 0$  for  $s \leq s'$ . In the next subsection, we calculate the remaining derivatives in the above saddle-point equations.

## C. Derivation of Saddle-Point Equations

Q

The integral in  $\Omega$  with respect to  $\hat{v}$  and  $\hat{w}$  is calculated as

$$\Omega = \frac{1}{\beta} \int \frac{\mathrm{d}\hat{\boldsymbol{v}}}{(2\pi)^{(t+1)/2}} \frac{\mathrm{d}\hat{\boldsymbol{w}}}{(2\pi)^{(t+1)/2}} \\
\times e^{\mathrm{i}\hat{\boldsymbol{w}}^{\top}(\beta^{-1}\mathbf{1})\hat{\boldsymbol{v}} - \frac{1}{2}\hat{\boldsymbol{v}}^{\top}\boldsymbol{Q}\hat{\boldsymbol{v}} - \frac{1}{2}\hat{\boldsymbol{w}}^{\top}\boldsymbol{B}\hat{\boldsymbol{v}} - \frac{1}{2}\hat{\boldsymbol{w}}^{\top}\boldsymbol{B}\hat{\boldsymbol{v}} - \frac{1}{2}\hat{\boldsymbol{w}}^{\top}\hat{\boldsymbol{D}}\hat{\boldsymbol{w}}}{\left(2\pi\right)^{(t+1)/2}} \\
= \frac{1}{\beta} \int \frac{\mathrm{d}\hat{\boldsymbol{v}}}{(2\pi)^{(t+1)/2}} \\
\times e^{-\frac{1}{2}\hat{\boldsymbol{v}}^{\top}\boldsymbol{Q}\hat{\boldsymbol{v}}} |\hat{\boldsymbol{D}}|^{-1/2} e^{-\frac{1}{2}\hat{\boldsymbol{v}}^{\top}(\beta^{-1}\mathbf{1} - \boldsymbol{B})^{\top}\hat{\boldsymbol{D}}^{-1}(\beta^{-1}\mathbf{1} - \boldsymbol{B})\hat{\boldsymbol{v}}} \\
= -\frac{1}{2\beta} \{\ln|\hat{\boldsymbol{D}}| \\
+ \ln|\boldsymbol{Q} + (\beta^{-1}\mathbf{1} - \boldsymbol{B})^{\top}\hat{\boldsymbol{D}}^{-1}(\beta^{-1}\mathbf{1} - \boldsymbol{B})|\}, \quad (64)$$

where B and  $\hat{D}$  are matrices whose elements are defined by

$$B^{(s,s')} := -ik^{(s')} - G^{(s,s')}, \tag{65}$$

$$\hat{D}^{(s,s')} := \frac{\sigma_0}{\beta} + 1 - m^{(s)} - m^{(s')} + C^{(s,s')}, \quad (66)$$

respectively. The 1 denotes an identity matrix.

The saddle-point equations including  $\Omega$  are evaluated as follows. We find

$$\begin{split} \hat{\eta}^{(s)} &= \mathbf{i} \left. \frac{\partial \Omega}{\partial m^{(s)}} \right|_{\text{saddle}, \mathbf{k} = \mathbf{0}, \mathbf{Q} = \mathbf{0}} \\ &= \mathbf{i} \frac{\partial}{\partial m^{(s)}} \left( -\frac{1}{2\beta} \{ \ln |\hat{\boldsymbol{D}}| \\ &+ \ln |(\beta^{-1} \mathbf{1} + \boldsymbol{G})^{\top} \hat{\boldsymbol{D}}^{-1} (\beta^{-1} \mathbf{1} + \boldsymbol{G})| \} \right) \\ &= \mathbf{i} \frac{\partial}{\partial m^{(s)}} \left( -\frac{1}{\beta} \ln |\beta^{-1} \mathbf{1} + \boldsymbol{G}| \right) \\ &= 0, \end{split}$$
(67)

and find

$$\hat{q}^{(s,s')} = 0,$$
 (68)

in the same way. Similarly, the  $\hat{L}^{(s,s')}$  becomes

$$\hat{L}^{(s,s')} = \frac{\partial \Omega}{\partial G^{(s,s')}} \Big|_{\text{saddle}, \boldsymbol{k} = \boldsymbol{0}, \boldsymbol{Q} = \boldsymbol{O}} \\
= \frac{\partial}{\partial G^{(s,s')}} \left( -\frac{1}{\beta} \ln |\boldsymbol{U}| \right) \\
= -\frac{1}{\beta} \cdot \frac{1}{|\boldsymbol{U}|} \frac{\partial |\boldsymbol{U}|}{\partial G^{(s,s')}} \\
\stackrel{(a)}{=} -\frac{1}{\beta} \cdot \frac{1}{|\boldsymbol{U}|} \frac{\partial}{\partial G^{(s,s')}} \sum_{\tau=-1}^{t-1} U^{(s,\tau)} \tilde{U}^{(s,\tau)} \\
= -\frac{1}{\beta} \frac{\tilde{U}^{(s,s')}}{|\boldsymbol{U}|},$$
(69)

where we put  $U := \beta^{-1}\mathbf{1} + G$ , whose (s, s') element is  $U^{(s,s')} := \beta^{-1}\delta_{s,s'} + G^{(s,s')}$ , and  $\tilde{U}^{(s,s')}$  denotes a cofactor of the (s, s') element of U. We here use the cofactor expansion at  $=^{(a)}$ . We therefore have

$$\hat{\boldsymbol{L}} = -\frac{1}{\beta} (\boldsymbol{U}^{\top})^{-1}$$
$$= -(\mathbf{1} + \beta \boldsymbol{G}^{\top})^{-1}.$$
(70)

Since Q = O,  $\Omega$  can be expanded with respect to Q as  $\ln |A + Q| = \operatorname{tr} \ln A + \operatorname{tr} A^{-1}Q$ . The  $\hat{Q}^{(s,s')}$  can be evaluated as

$$\hat{Q}^{(s,s')} = \operatorname{i} \lim_{\boldsymbol{Q} \to \boldsymbol{O}} \left. \frac{\partial \Omega}{\partial Q^{(s,s')}} \right|_{\text{saddle},\boldsymbol{k}=\boldsymbol{0}} \\ = -\frac{\mathrm{i}}{2\beta} \lim_{\boldsymbol{Q} \to \boldsymbol{O}} \frac{\partial}{\partial Q^{(s,s')}} \operatorname{tr} \boldsymbol{M} \boldsymbol{Q} \\ = -\frac{\mathrm{i}}{2\beta} M^{(s',s)},$$
(71)

where we put

$$M := (\beta^{-1}\mathbf{1} + \mathbf{G})^{-1}\hat{\mathbf{D}}(\beta^{-1}\mathbf{1} + \mathbf{G}^{\top})^{-1}, \qquad (72)$$
$$D^{(s,s')} := \beta\hat{D}^{(s,s')} \tag{73}$$

$$D^{(s,s)} := \beta D^{(s,s)}$$
(73)  
=  $\sigma_0^2 + \beta [1 - m^{(s)} - m^{(s')} + C^{(s,s')}].$  (74)

We then have  $\hat{Q} = -i\frac{1}{2\beta}M^{\top} = -i\frac{1}{2}(1+\beta G)^{-1}D(1+\beta G^{\top})^{-1}$ . Finally, we turn to  $k^{(s)}$ :

$$\begin{split} \hat{k}^{(s)} &= \operatorname{i} \lim_{k \to 0} \left. \frac{\partial \Omega}{\partial k^{(s,s')}} \right|_{\operatorname{saddle}, \boldsymbol{Q} = \boldsymbol{O}} \\ &= -\frac{\mathrm{i}}{\beta} \lim_{k \to 0} \left. \frac{\partial}{\partial k^{(s)}} \ln \left| \mathbf{1} - \beta \boldsymbol{B}^{\top} \right| \\ \stackrel{(a)}{=} -\frac{\mathrm{i}}{\beta} \lim_{k \to 0} \frac{1}{\left| \mathbf{1} - \beta \boldsymbol{B}^{\top} \right|} \frac{\partial}{\partial k^{(s)}} \\ &\left\{ \begin{vmatrix} \cdots & \delta_{-1,s'} + \mathrm{i}\beta k^{(-1)} + \beta G^{(s',-1)} & \cdots \\ \vdots \\ \cdots & \delta_{s,s'} + \beta G^{(s',s)} & \cdots \\ \vdots \\ \cdots & \delta_{t-1,s'} + \mathrm{i}\beta k^{(t-1)} + \beta G^{(s',t-1)} & \cdots \\ &\vdots \\ \cdots & \delta_{-1,s'} + \mathrm{i}\beta k^{(-1)} + \beta G^{(s',-1)} & \cdots \\ &\vdots \\ \vdots \\ \cdots & \delta_{t-1,s'} + \mathrm{i}\beta k^{(t-1)} + \beta G^{(s',t-1)} & \cdots \\ &\vdots \\ \cdots & \delta_{t-1,s'} + \mathrm{i}\beta k^{(t-1)} + \beta G^{(s',t-1)} & \cdots \\ &\vdots \\ \cdots & \delta_{t-1,s'} + \mathrm{i}\beta k^{(t-1)} + \beta G^{(s',t-1)} & \cdots \\ &\vdots \\ \cdots & \delta_{t-1,s'} + \mathrm{i}\beta k^{(t-1)} + \beta G^{(s',t-1)} & \cdots \\ \end{vmatrix} \right\} \\ &= |\mathbf{A}_{[s]}|, \end{split}$$

where  $\Lambda_{[s]}$  denotes a  $(s+2) \times (s+2)$  matrix whose (s', s'') element is given by

$$\Lambda_{[s]}^{(s',s'')} = \begin{cases} \delta_{s',s''} + \beta G^{(s'',s')}, & \text{for } s' \neq s \\ 1, & \text{for } s' = s \end{cases} .$$
(76)

The  $\Lambda_{[s]}$  is a matrix whose elements in a row that represents values of stage *s*, i.e., the (s+2)th row, in  $1+\beta G$  are replaced to 1. Since  $|1 - \beta B^{\top}|$  contains  $k^{(s)}$  only in a single row,  $|1 - \beta B^{\top}|$  is expanded with respect to the row at  $=^{(a)}$  in (75).

*Example 1:* Parameter  $\Lambda_{[2]}$  has the following form:

$$\mathbf{\Lambda}_{[2]} = \begin{pmatrix} 1 & \beta G^{(0,-1)} & \beta G^{(1,-1)} & \beta G^{(2,-1)} \\ 0 & 1 & \beta G^{(1,0)} & \beta G^{(2,0)} \\ 0 & 0 & 1 & \beta G^{(2,1)} \\ 1 & 1 & 1 & 1 \end{pmatrix}, \quad (77)$$
when  $t = 2$ .

## D. Bit Error Rate

One can obtain the polynomial expressions of soft-bits  $\tilde{b}_k^{(s)}$ , which are averaged over the path probability (25), by GFA. Therefore, we can also evaluate the averaged value of analytic functions with respect to the soft-bits. The *n*-th order differentiation of the averaged generating functional, which has the forms of (27) and (31), with respect to  $\psi_k^{(s)}$  gives

$$\lim_{\psi \to \mathbf{0}} \left( \frac{\partial}{\partial \psi_k^{(s)}} \right)^n \bar{Z}[\psi] \stackrel{(a)}{=} (-\mathbf{i})^n \overline{\langle (\tilde{b}_k^{(s)})^n \rangle} \\ \stackrel{(b)}{=} (-\mathbf{i})^n \langle (\tilde{b}_k^{(s)})^n \rangle_k, \tag{78}$$

where  $=^{(a)}$  and  $=^{(b)}$  are derived by differentiation of (27) and (31), representing For arbitrary function F(x) that can be

expanded around x = 0, we thus have the identity:

$$\langle F(\tilde{b}_k^{(s)}) \rangle = \langle F(\tilde{b}_k^{(s)}) \rangle_k.$$
(79)

The Hamming distance between single letters is defined by

$$d(x, \hat{x}) := \begin{cases} 0, & \text{if } x = \hat{x} \\ 1, & \text{if } x \neq \hat{x} \end{cases}.$$
 (80)

The Hamming distance between K-bit sequences **b** and  $\hat{\mathbf{b}}$  is measured by the averaged single-letter distortion as  $d(\mathbf{x}, \hat{\mathbf{x}}) := K^{-1} \sum_{k=1}^{K} d(b_k, \hat{b}_k)$ . The hard-decision estimates can be represented by  $\hat{b}_k^{(t)} = \operatorname{sgn}(\tilde{b}_k^{(t)})$  If we therefore choose F(x) such as a function which approaches  $\operatorname{sgn}(x)$ , e.g.,  $\lim_{a\to\infty} \tanh(ax)$ , the bit error rate  $P_b^{(t)}$  can then be evaluated as

$$P_{b}^{(t)} := \overline{\langle d(\boldsymbol{b}, \hat{\boldsymbol{b}}^{(t)}) \rangle}$$
(81)

$$=\frac{1}{2}\left(1-\frac{1}{K}\sum_{k=1}^{K}\langle\operatorname{sgn}\left(\tilde{b}_{k}^{(t)}\right)\rangle_{k}\right),\tag{82}$$

where  $d(b, \hat{b}) = (1 - b\hat{b})/2$  for  $b, \hat{b} \in \{-1, 1\}$ .

## E. Effective Single-User Problem and Analytical Result

We here summarize our calculation. Some macroscopic parameters are found to vanish in the saddle-point:  $k^{(s)} = Q^{(s,s')} = 0$ . The remaining macroscopic parameters can all be expressed in terms of three observables: the average overlap  $m^{(s)}$ , the average single-user correlation  $C^{(s,s')}$ , and the average single-user response function  $G^{(s,s')}$ , which are defined by (47) – (49). The averaged generating functional  $\overline{Z}[\psi]$  is dominated by a saddle-point for  $K \to \infty$ . We can thus simplify the saddle-point problem of Lemma 1 in the large system limit.

We derive a single-user saddle-point problem. Note that we remained the user index k to deduce the meaning of macroscopic parameters with respect to external messages  $\{\theta_k^{(s)}\}$  and dummy functions  $\{\psi_k^{(s)}\}$ . We set these parameters as

$$\psi_k^{(s)} \to 0, \tag{83}$$

$$\theta_k^{(s)} \to 0. \tag{84}$$

Consequently, the single-user measure (41) becomes user independent. We then arrive at the following proposition.

Proposition 1: Setting  $\gamma \to \infty$  and  $\theta^{(s)} = 0$ , the dynamics of soft-PIC can be obtained as the following equations in the large system limit, i.e.,  $K \to \infty$ .

$$m^{(s)} = \langle \langle \tilde{b}^{(s)} \rangle \rangle, \tag{85}$$

$$C^{(s,s')} = \langle\!\langle \tilde{b}^{(s)} \tilde{b}^{(s')} \rangle\!\rangle.$$
(86)

$$G^{(s,s')} = \langle \tilde{b}^{(s)}(\mathbf{R}^{-1}\mathbf{v})^{(s')} \rangle, \tag{87}$$

with the causality:  $G^{(s,s')} = 0$  for  $s \leq s'$ . The bit error rate of hard decisions  $\{ \operatorname{sgn} [\tilde{b}_k^{(s)}] \}$  at the  $t^{\text{th}}$  stage of soft-PIC (2) is obtained by

$$P_{b}^{(s)} = \frac{1}{2} (1 - \langle\!\langle \operatorname{sgn}(\tilde{b}^{(s)}) \rangle\!\rangle).$$
(88)

The effective path measure is given by

$$\langle\!\langle g(\boldsymbol{b}, \boldsymbol{v}) \rangle\!\rangle$$

$$:= \int \mathcal{D}\boldsymbol{v} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t-1} \mathrm{d}\tilde{b}^{(s)} \right) g(\tilde{\boldsymbol{b}}, \boldsymbol{v}) \, \delta[\tilde{b}^{(-1)}]$$

$$\times \prod_{s=-1}^{t-1} \delta[\tilde{b}^{(s+1)} - f(\hat{k}^{(s)} + v^{(s)} + (\Gamma \tilde{\boldsymbol{b}})^{(s)})], \qquad (89)$$

where

$$\mathcal{D}\boldsymbol{v} := \mathrm{d}\boldsymbol{v} |2\pi\boldsymbol{R}|^{-1/2} e^{-\frac{1}{2}\boldsymbol{v}\cdot\boldsymbol{R}^{-1}\boldsymbol{v}}, \qquad (90)$$

$$\boldsymbol{R} = (\boldsymbol{1} + \beta \boldsymbol{G}^{\top})^{-1} \boldsymbol{D} (\boldsymbol{1} + \beta \boldsymbol{G})^{-1}, \qquad (91)$$

$$\Gamma = (\mathbf{1} + \beta \mathbf{G})^{-1} \beta \mathbf{G},\tag{92}$$

$$\hat{k}^{(s)} = |\mathbf{\Lambda}_{[s]}|. \tag{93}$$

with

$$D^{(s,s')} := \sigma_0^2 + \beta [1 - m^{(s)} - m^{(s')} + C^{(s,s')}], \qquad (94)$$

$$\Lambda_{[s]}^{(s',s'')} = \begin{cases} \delta_{s',s''} + \beta G^{(s',s')}, & \text{for } s' \neq s \\ 1, & \text{for } s' = s. \end{cases}$$
(95)

The terms  $(\mathbf{R}^{-1}\mathbf{v})^{(s)}$  and  $(\Gamma \tilde{\mathbf{b}})^{(s)}$  denote the  $s^{\text{th}}$  element of the vector  $\mathbf{R}^{-1}\mathbf{v}$  and  $\Gamma \tilde{\mathbf{b}}$ , respectively.

Details of derivation is available in Appendix C. Equations (85)-(95) entirely describe the dynamics of the system. Term  $(\Gamma \tilde{b})^{(s)}$  in (89) is called the *Onsager reaction term*.

## F. Numerical Analysis and Experiments

To validate the results obtained here, we performed numerical experiments in an N = 8,000 system. The system of N = 8,000 is too large for a practical system, but we are now concerned with the verification of the analytical result derived under the large system limit.

Figure 1(a) plots the first few stages of the detection dynamics of the hard-PIC ( $f(x) = \operatorname{sgn}(x)$ ) and soft-PIC ( $f(x) = \operatorname{tanh}(x/\sigma^2)$  with  $\sigma = \sigma_0$ ) predicted by GFA and density evolution [52] for  $E_b/N_0 = 8$  [dB]. The detailed derivation of the density evolution analysis is available in the reference [52]. The system load is  $\beta = 0.5 < \beta_c$ , where  $\beta_c$  is the critical system load defined as the minimum system load at which the dynamics fail to converge to the replica-symmetric solution of the marginal-posterior-mode detector. The critical load  $\beta_c$  of soft-PIC for  $E_b/N_0 = 8$  [dB] is about 0.6. The predictions of density evolution systematically deviate from computer simulation results at transients.

Figure 1(b) plots the first few stages of the detection dynamics of the hard-PIC and soft-PIC for  $E_b/N_0 = 8$  [dB], predicted by GFA and density evolution with the system load of  $\beta = 0.7 > \beta_c$ . Oscillation of the detection dynamics was observed, when  $\beta > \beta_c$ . The density evolution results have residual deviations in Fig. 1 due to the lack of the Onsager reaction term and the assumption that the summation over all messages follows a Gaussian distribution. In particular, the deviation between the density evolution predictions and the simulation results becomes large when  $\beta > \beta_c$ . In contrast, GFA exhibits good consistency with the simulation results.

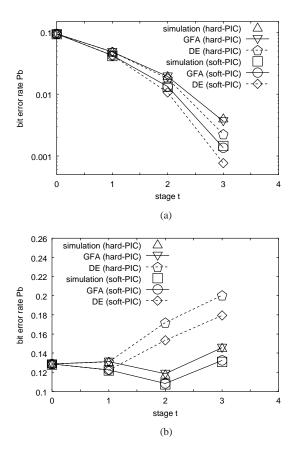


Fig. 1. First few stages of detection dynamics of the hard-PIC (upper) and soft-PIC with  $f(x) = \tanh(x/\sigma^2)$  with  $\sigma = \sigma_0$  (lower) predicted by GFA (solid lines) and density evolution (dashed lines). Computer simulations (triangles and squares) were evaluated with N = 8,000 for  $E_b/N_0 = 8$  [dB]. The system loads were (a)  $\beta = 0.5 < \beta_c$  and (b)  $\beta = 0.7 > \beta_c$ , respectively.

G. Derivation of Existing Results by Generating Functional Analysis

If we put  $\Gamma b = 0$ , viz., if we neglect the Onsager reaction term, the GFA recovers the density evolution framework based on the statistical neurodynamics [52].

The multiple integral including (89) can be partially performed, when we put  $\Gamma \tilde{b} = 0$ . Namely,  $m^{(s)}$ ,  $G^{(s')}$ , and  $P_b^{(s)}$ only depend on  $v^{(s-1)}$  among v; and  $C^{(s,s')}$  only depends on  $v^{(s-1)}$  and  $v^{(s'-1)}$ . We here separate v into two subvectors:  $v = (v_1, v_2)^{\top}$ . The correlation matrix R can then be represented as

$$\boldsymbol{R} = \begin{pmatrix} \boldsymbol{R}_{11} & \boldsymbol{R}_{12} \\ \boldsymbol{R}_{21} & \boldsymbol{R}_{22} \end{pmatrix}.$$
 (96)

For example,  $v_1$  and  $R_{11}$  are chosen as  $v_1 = (v^{(s-1)}, v^{(s'-1)})^\top$  and

$$\boldsymbol{R}_{11} = \begin{pmatrix} R^{(s-1,s-1)} & R^{(s-1,s'-1)} \\ R^{(s'-1,s-1)} & R^{(s'-1,s'-1)} \end{pmatrix}, \quad (97)$$

to calculate (86). For arbitrary function  $f(v_1)$ , the following identity holds:

$$\int \mathcal{D} \boldsymbol{v} f(\boldsymbol{v}_1) = \frac{1}{\sqrt{|2\pi \boldsymbol{R}_{11}|}} \int \mathrm{d} \boldsymbol{v}_1 e^{-\frac{1}{2} \boldsymbol{v}_1 \cdot \boldsymbol{R}_{11}^{-1} \boldsymbol{v}_1} f(\boldsymbol{v}_1).$$
(98)

Since it turns out that the response function becomes zero except for  $G^{(s,s-1)}$ , the  $\hat{k}^{(s)}$  has a simple form.

We then obtain the following result. The bit error rate  $P_b(t)$ of hard decisions  $\hat{b}_k(t) = \text{sgn}[u_k(t-1)]$  at the *t*th stage is given by

$$P_b(t) = \mathcal{Q}\bigg(\frac{\hat{k}^{(s-1)}}{\sqrt{R^{(s-1,s'-1)}}}\bigg),$$
(99)

where  $Q(z) := \int_{z}^{\infty} Dx$  denotes the error function and we here put  $Dz := dz(2\pi)^{-1/2}e^{-z^{2}/2}$  to simplify the notation. The  $m^{(s)}$  are to be evaluated with the following set of recursive equations:

$$\hat{k}^{(s)} = 1 - \beta G^{(s,s-1)} \hat{k}^{(s-1)}, \qquad (100)$$

$$R^{(s,s')} = D^{(s,s')} + \beta^2 G^{(s,s-1)} G^{(s',s'-1)} R^{(s-1,s'-1)} + \sum_{\lambda=-1}^{s-1} D^{(\lambda,s')} \prod_{\tau=\lambda+1}^{s} (-\beta G^{(\tau,\tau-1)}) + \sum_{\lambda=-1}^{s'-1} D^{(\lambda,s)} \prod_{\tau=\lambda+1}^{s'} (-\beta G^{(\tau,\tau-1)}), \quad (101)$$

$$D^{(s,s')} = \sigma_0^2 + \beta (1 - m^{(s)} - m^{(s')} + C^{(s,s')}), \qquad (102)$$

$$m^{(s+1)} = \int_{\mathbb{R}} \mathrm{D}z f(\hat{k}^{(s-1)} + z\sqrt{R^{(s-1,s-1)}}), \qquad (103)$$

$$G^{(s,s-1)} = \frac{1}{\sqrt{R^{(s-1,s-1)}}} \int_{\mathbb{R}} \mathrm{D}zz f(\hat{k}^{(s-1)} + z\sqrt{R^{(s-1,s-1)}}),$$
(104)

$$C^{(s,s')} = \int_{\mathbb{R}^3} \mathrm{D}z \mathrm{D}u \mathrm{D}v$$

$$\times f(\hat{k}^{(s-1)} + z\sqrt{R^{(s-1,s'-1)}} + u\sqrt{R^{(s-1,s-1)} - R^{(s-1,s'-1)}})$$

$$\times f(\hat{k}^{(s'-1)} + z\sqrt{R^{(s-1,s'-1)}} + u\sqrt{R^{(s'-1,s'-1)} - R^{(s-1,s'-1)}}). \quad (105)$$

The initialization condition is  $R^{(-1,-1)} = D^{(-1,-1)} = \sigma_0^2 + \beta$ ,  $\hat{k}^{(-1)} = 1$  and  $m^{(-1)} = G^{(-1,-1)} = C^{(-1,-1)} = 0$ ,

This result is identical to that of density evolution [52]. In the derivation by means of density evolution, it is assumed that the local field  $u_k(t)$  follows the Gaussian distribution with mean  $B_t$  and covariance  $C_{t,\tau}$ . Furthermore, the Onsager reaction term is ignored. The GFA, on the other hand, can treat the Onsager reaction term correctly.

## VI. GENERATING FUNCTIONAL ANALYSIS FOR ORC-PIC

This section discusses the dynamics of ORC-PIC.

*Proposition 2:* The dynamics of ORC-PIC is given by the following average over the effective path measure:

$$\langle\!\langle g(\boldsymbol{b}, \boldsymbol{v}) \rangle\!\rangle$$

$$:= \int \mathcal{D}\boldsymbol{v} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t-1} \mathrm{d}\tilde{b}^{(s)} \right) g(\tilde{\boldsymbol{b}}, \boldsymbol{v}) \, \delta[\tilde{b}^{(-1)}]$$

$$\times \prod_{s=-1}^{t-1} \delta[\tilde{b}^{(s+1)} - f(\hat{k}^{(s)} + v^{(s)} + (\Gamma\tilde{\boldsymbol{b}})^{(s)} - (\hat{\Gamma}\tilde{\boldsymbol{b}})^{(s)})].$$

$$(106)$$

All other parameters are identical to those in Proposition 1.  $\Box$ 

For ORC-PIC, the summation over all messages, which is the argument of  $f(\cdots)$  in (5) becomes

$$u_{k}^{(t)} = \sum_{k=1}^{K} W_{kk'}(b_{k'} - \tilde{b}_{k'}^{(t)}) + \tilde{b}_{k}^{(t)} - \sum_{s'=-1}^{s} \hat{\Gamma}_{k}^{(s,s')} \tilde{b}_{k}^{(s')} + \frac{\sigma_{0}}{\sqrt{N}} \sum_{\mu=1}^{N} s_{k}^{\mu} n^{\mu} + \theta_{k}^{(t)}, \quad (107)$$

with external messages  $\{\theta_k^{(t)}\}\$ . The averaged generating functional  $\bar{Z}[\psi]$  is represented as

$$\bar{Z}[\psi] = \mathbb{E}_{s_{1}, \cdots, s_{K}, n} \left[ \int_{\mathbb{R}^{(t+1)K}} \underline{d}u \\
\times p[\tilde{b}^{(-1)}] \left( \prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} e^{-\frac{\gamma^{2}}{2} \{\tilde{b}_{k}^{(s+1)} - f(u_{k}^{(s)})\}^{2}} \right) \\
\times \exp\left[ -i \sum_{s=-1}^{t} \tilde{b}^{(s)} \cdot \psi^{(s)} \right] \\
\times \prod_{s=0}^{t-1} \prod_{k=1}^{K} \delta\left( u_{k}^{(s)} - \sum_{k=1}^{K} W_{kk'}(b_{k'} - \tilde{b}_{k'}^{(s)}) - \tilde{b}_{k}^{(s)} \right) \\
+ \sum_{s'=-1}^{s} \hat{\Gamma}_{k}^{(s,s')} \tilde{b}_{k}^{(s')} - \frac{\sigma_{0}}{\sqrt{N}} \sum_{\mu=1}^{N} s_{k}^{\mu} n^{\mu} - \theta_{k}^{(s)} \right) \right].$$
(108)

The difference between (27) and (108) is only in the fourth term  $\sum_{s'=-1}^{s} \hat{\Gamma}_k^{(s,s')} \tilde{b}_k^{(s')}$  in  $\delta(\cdots)$  in (108). One can straightforwardly show Proposition 2 in the same manner as the derivation of Proposition 1. The difference in the analytical results only appears in (89) of Proposition 1.

We will now consider how to choose matrix  $\hat{\Gamma} = (\hat{\Gamma}^{(s,s')})$ in the Onsager reaction canceling term to cancel the Onsager reaction term.

*Proposition 3:* When the Onsager reaction term  $(\hat{\Gamma}\hat{b})^{(t)}$  is chosen as

$$(\hat{\boldsymbol{\Gamma}}\hat{\boldsymbol{b}})^{(t)} = \sum_{s=0}^{t-1} (-1)^{t-s-1} \beta^{t-s} \left(\prod_{\tau=s}^{t-1} G^{(\tau+1,\tau)}\right) \tilde{b}^{(s)}, \quad (109)$$

for t > 0 and  $(\hat{\Gamma}\hat{b})^{(-1)} = (\hat{\Gamma}\hat{b})^{(0)} = 0$ , then the Onsager reaction term  $\Gamma\hat{b}$  is canceled, viz.,  $\Gamma\hat{b} - \hat{\Gamma}\hat{b} = O$ , and the response functions become  $G^{(t,t')} = 0$  except for t = t' + 1.  $\Box$ 

Details of derivation is available in Appendix D. Note that (6) is a recursive formula of (109). One can confirm the ORC-PIC algorithm can correctly cancel the Onsager reaction term

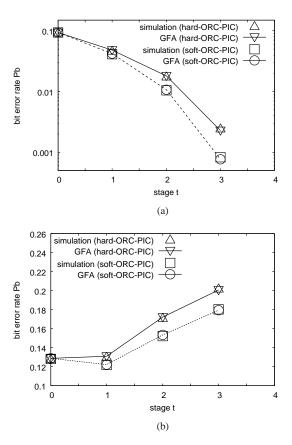


Fig. 2. First few stages of detection dynamics of the hard-ORC-PIC (upper) and soft-ORC-PIC with  $f(x) = \tanh(x/\sigma^2)$  with  $\sigma = \sigma_0$  (lower) predicted by GFA (solid lines). In this case, the density evolution result is identical to the GFA result. Computer simulations (triangles and squares) were evaluated with N = 8,000 for  $E_b/N_0 = 8$  [dB]. The system loads were (a)  $\beta = 0.5 < \beta_c$  and (b)  $\beta = 0.7 > \beta_c$ , respectively.

by GFA. The density evolution results completely agree with the GFA results [52]. The density evolution curves of Fig. 1 are equivalent to the performance of ORC-PIC.

To validate the results obtained here, we performed numerical experiments in an N = 8,000 system. Figure 2(a) plots the first few stages of the detection dynamics of the hard-ORC-PIC (f(x) = sgn(x)) and soft-ORC-PIC ( $f(x) = \tanh(x/\sigma^2)$  with  $\sigma = \sigma_0$ ) predicted by GFA [52] for  $E_b/N_0 = 8$  [dB]. Figure 2(b) plots the first few stages of the detection dynamics of the hard-ORC-PIC and soft-ORC-PIC for  $E_b/N_0 = 8$  [dB], predicted by GFA and density evolution with the system load of  $\beta = 0.7 > \beta_c$ . Oscillation of the detection dynamics was observed, when  $\beta > \beta_c$ .

# VII. GENERATING FUNCTIONAL ANALYSIS FOR BP-based detector

We next apply this scheme to the BP-based detector. One can obtain the following result in the same manner as soft-PIC.

Proposition 4: The dynamics of the BP-based detector represented by (14) is described by equations (85)-(89) and the following equations:  $\mathbf{R} = (\mathbf{1} - \beta \mathbf{B})^{-1} \mathbf{D} (\mathbf{1} - \beta \mathbf{B}^{\top})^{-1}, \mathbf{\Gamma} = \text{diag}(A^{(-1)}, \dots, A^{(t-1)}) - [\sum_{s=-1}^{t-1} \mathbf{J}_{s+1} \otimes \mathbf{\Delta}^{s+1} (\mathbf{U}^{\top})^{-1}]^{\top}$  and  $\hat{k}^{(s)} = |\mathbf{\Lambda}_{[s]}|$  where  $\mathbf{B}, \mathbf{U}$ , and  $\mathbf{D}$  are matrices having

matrix elements

$$B^{(s,s')} = -\sum_{\tau=-1}^{s} J^{(s,\tau)} G^{(\tau,s')},$$
(110)

$$U^{(s,s')} = \delta_{s,s'} + \beta \sum_{\tau=-1}^{s} J^{(s,\tau)} G^{(\tau,s')}, \qquad (111)$$

$$D^{(s,s')} = \sigma_0^2 R^{(s)} R^{(s')} + \beta \bigg[ R^{(s)} R^{(s')} - R^{(s)} \sum_{\tau'=-1}^{s'} J^{(s',\tau')} m^{(\tau')} - R^{(s')} \sum_{\tau=-1}^{s} J^{(s,\tau)} m^{(\tau)} + \sum_{\tau=-1}^{s} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} C^{(\tau,\tau')} \bigg],$$
(112)

and

$$\Lambda_{s}^{(s',s'')} = \begin{cases} \delta_{s',s''} + \beta \sum_{\tau'=-1}^{s'} J^{(s',\tau')} G^{(\tau',s')}, & s' \neq s \\ R^{(s')}, & s' = s. \end{cases}$$
(113)

and matrix  $J_s$  denotes

$$\boldsymbol{J}_{s} := \begin{pmatrix} J^{(s-1,-1)} & \cdots & J^{(s-1,-1)} \\ \vdots & & \vdots \\ J^{(t-1,t-s-1)} & \cdots & J^{(t-1,t-s-1)} \\ 0 & \cdots & 0 \\ \vdots & & \vdots \\ 0 & \cdots & 0 \end{pmatrix}.$$
(114)

Here,  $\Delta = (\delta_{s+1,s'})$  denotes a  $(t+1) \times (t+1)$  matrix whose (s, s') element is given by  $\delta_{s+1,s'}$ , and the operator  $\otimes$  denotes the Hadamard product  $A \otimes B = (a_{ij}b_{ij})$  for  $A = (a_{ij})$  and  $B = (b_{ij})$ .

For the BP-based detector, the summation over all messages, which is the argument of  $tanh(\cdots)$  of (14) becomes

$$u_k^{(t)} = R^{(t)}h_k + A^{(t)}\tilde{b}_k^{(t)} - \sum_{s=-1}^t J^{(t,s)}\sum_{k'=1}^K W_{kk'}\tilde{b}_k^{(s)}.$$
 (115)

The averaged generating functional  $\bar{Z}[\psi]$  is represented as

$$\bar{Z}[\psi] = \mathbb{E}_{s_1, \cdots, s_K, n} \left[ \int_{\mathbb{R}^{(t+1)_K}} \underline{d}u \\ \times p[\tilde{b}^{(-1)}] \left( \prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} e^{-\frac{\gamma^2}{2} \{\tilde{b}_k^{(s+1)} - f(u_k^{(s)})\}^2} \right) \\ \times \exp\left[ -i \sum_{s=-1}^t \tilde{b}^{(s)} \cdot \psi^{(s)} \right] \\ \times \prod_{s=0}^{t-1} \prod_{k=1}^K \delta\left( u_k^{(s)} - R^{(t)}h_k - A^{(t)}\tilde{b}_k^{(t)} \\ + \sum_{s=-1}^t J^{(t,s)} \sum_{k'=1}^K W_{kk'}\tilde{b}_k^{(s)} \right) \right].$$
(116)

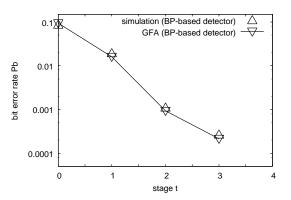


Fig. 3. First few stages of detection dynamics of the BP-based detector predicted by GFA (solid lines). Computer simulations (triangles and squares) were evaluated with N=8,000 for  $E_b/N_0=8$  [dB]. The control parameter and the system load were set as  $\sigma = \sigma_0$  and  $\beta = 0.5$ .

Details on the derivation are given in Appendix E. For the first few time-steps, we confirmed explicit expressions for the solutions to our dynamic equations as

$$m^{(t)} = \int Dz \tanh(\sqrt{R^{(t,t)}}z + \hat{k}^{(t)}), \qquad (117)$$

$$q^{(t,t)} = \int Dz \tanh^2(\sqrt{R^{(t,t)}}z + \hat{k}^{(t)}), \quad (118)$$

$$\hat{k}^{(t+1)} = \frac{1}{\sigma^2 + \beta [1 - q^{(t,t)}]},$$
(119)

$$R^{(t+1,t+1)} = \frac{\beta[1 - 2m^{(t)} + q^{(t,t)}]) + \sigma_0^2}{(\sigma^2 + \beta[1 - q^{(t,t)}])^2}, \quad (120)$$

where the BER of hard decisions at the  $t^{\rm th}$  stage is  $P_b^t = H(\hat{k}^{(t)}/\sqrt{R^{(t,t)}})$  with  $H(x) = \int_x^\infty Dz$ . In the BP-based detector, the Onsager reaction term vanished. These are identical to the results from density evolution [25]. Figure 3 plots the first few stages of the detection dynamics of the BP-based detector predicted by GFA [52] for  $E_b/N_0 = 8$  [dB]. The control parameter correct was set as  $\sigma = \sigma_0$ .

On the other hand, Kabashima analyzed the stability of the density evolution results and obtained an unstable condition for the fixed point solution [25]. This unstable condition coincided with the Almeida-Thouless (AT) instability condition [1] of replica analysis obtained by Tanaka [51]. The AT instability is the local instability of the RS saddle-point solution in the replica analysis. In GFA, it is known that the AT instability condition is derived by a condition that the response function diverges [48]. When the AT instability condition is satisfied, the relaxation times of GFA dynamic equations are expected to extend to infinity. In the system addressed here, this condition must be equivalent to the AT instability condition. In Section VIII-C, we will discuss the stability analysis that corresponds to the AT instability.

## VIII. STATIONARY ESTIMATES

This section explains how we approximately extract the stationary estimate, which is a stationary state of the iterative algorithm, under some assumptions. To simplify the problem, we restrict ourselves to the soft-PIC algorithm. We follow the method of Ref. [7].

## A. Analysis of Stationary Estimates

We recognize that the representation of the effective path measure given by (89) is fully equivalent to the measure corresponding to a single-user process of the form:

$$\tilde{b}^{(t+1)} = f(\hat{k}^{(t)} + v^{(t)} + \theta^{(t)} + (\Gamma \tilde{\boldsymbol{b}})^{(t)}).$$
(121)

Variable  $\boldsymbol{v} = (v^{(t)})$  can be regarded as a random variable that obeys Normal distribution  $\mathcal{N}(\mathbf{0}, \mathbf{R})$ , e.g.,  $\langle v^{(t)} \rangle_{\boldsymbol{v}} = 0$  and  $\langle v^{(t)}v^{(t')} \rangle_{\boldsymbol{v}} = R^{(t,t')}$ , from (91), where  $\langle \cdots \rangle_{\boldsymbol{v}}$  denotes the average over random variable  $\boldsymbol{v}$ . We here put

$$\phi^{(t+1)} := \hat{k}^{(t)} + v^{(t)} + \theta^{(t)} + (\Gamma \tilde{\boldsymbol{b}})^{(t)}, \qquad (122)$$

then the single-user process can be rewritten as  $\tilde{b}^{(t)} = f(\phi^{(t)})$ . We therefore have the following form.

$$\phi^{(t+1)} = \hat{k}^{(t)} + v^{(t)} + \theta^{(t)} + \sum_{t'=-1}^{t} \Gamma^{(t,t')} f(\phi^{(t')}), \quad (123)$$

The relationship of (123) is referred as a *single-user process*. Using (123), the overlap (85), the correlation function (87), the response function (87), and the bit error rate (88) are obtained as

$$m^{(t)} = \langle f(\phi^{(t)}) \rangle_{\boldsymbol{v}},\tag{124}$$

$$C^{(t,t')} = \langle f(\phi^{(t)}) f(\phi^{(t')}) \rangle_{v},$$
(125)

$$G^{(t,t')} = \langle f(\phi^{(t)})(\boldsymbol{R}^{-1}\boldsymbol{v})^{(t')} \rangle_{\boldsymbol{v}}, \qquad (126)$$

$$P_b^{(t)} = \frac{1}{2} (1 - \langle \operatorname{sgn}(f(\phi^{(t)})) \rangle_{\boldsymbol{v}}).$$
(127)

We make the following assumptions to evaluate the stationary estimates.

Assumption 1: (Time-translation invariance: TTI) The dynamics reaches a time-translation invariant estimate:

$$\lim_{t \to \infty} m^{(t)} = m, \tag{128}$$

$$\lim_{t \to \infty} C^{(t+\tau,\tau)} = C^{(\tau)},\tag{129}$$

$$\lim_{t \to \infty} G^{(t+\tau,\tau)} = G^{(\tau)}.$$
 (130)

If this property holds, the dynamics reaches stationary estimates.  $\hfill \Box$ 

Assumption 2: (Finite integrated response: FIR) The integrated response

$$\chi := \lim_{t \to \infty} \sum_{t' \le t} G^{(t,t')} \tag{131}$$

is a finite non-negative number, i.e.,  $\chi < \infty$ .

The integrated response is also called as a susceptibility.

Assumption 3: (Weak long-term memory: WLTM)

$$\lim_{t \to \infty} G^{(t,t')} = 0 \tag{132}$$

with

1

for any finite t' [9], [10].

Since the response function represents the memory which means what happened to the system, the weakness of the longterm memory implies that the system responds to its past in an averaged way. The details of what takes place during finite stages tend to be washed away.

We assume that the stationary estimate is unique and Assumptions 1 - 3 hold. Note that violation of WLTM

immediately implies  $\chi = \infty$ . Under Assumption 1, the matrices  $C = (C^{(s,s')} \text{ and } G = (G^{(s,s')})$  can be regarded as symmetric Toeplitz matrices, since we consider that the size of square matrices C, G, and D are sufficiently large, viz., t + 2 is sufficiently large. The matrix  $D = (D^{(s,s')})$  whose element is given by (94) is also regarded as a symmetric Toeplitz matrix from  $D^{(s',s)} = \sigma_0^2 + \beta[1 - m^{(s')} - m^{(s)} + C^{(s',s)}] = \sigma_0^2 + \beta[1 - m^{(s')} - m^{(s)} + C^{(s,s')}] = D^{(s,s')}$  and  $\lim_{t\to\infty} D^{(t+\tau,t)} = \sigma_0^2 + \beta[1 - 2m + C(\tau)] =: D^{(\tau)}$ , where we again use Assumption 1. Therefore, we hereafter make the following assumption.

Assumption 4: The matrices  $\{C, G, D\}$  and their powers are Toeplitz matrices. All pairs of matrices commute.

In this subsection, we derive the following proposition under some assumptions.

*Proposition 5:* Let Assumptions 1 - 4 hold, then soft-PIC converges to an unique stationary estimate whose BER is given by

$$P_b = \frac{1}{2} \left( 1 - \int_{\mathbb{R}} \mathrm{D}z \operatorname{sgn}\left(g(z)\right) \right), \tag{133}$$

where g(z) is to be determined by equations that describe the stationary estimate:

$$m = \int_{\mathbb{R}} \mathrm{D}z g(z), \tag{134}$$

$$c = \int_{\mathbb{R}} \mathrm{D}z g(z)^2, \tag{135}$$

$$\chi = \frac{1}{\sqrt{\mathsf{F}}} \int_{\mathbb{R}} \mathrm{D}z z g(z), \tag{136}$$

$$g(z) = f(\mathsf{E} + z\sqrt{\mathsf{F}} + \mathsf{G}g(z) + \bar{\theta}), \tag{137}$$

$$\mathsf{E} = \frac{1}{1 + \beta \chi},\tag{138}$$

$$\mathsf{F} = \frac{\sigma_0^2 + \beta(1 - 2m + c)}{(1 + \beta\chi)^2},$$
(139)

$$\mathsf{G} = \frac{\beta \chi}{1 + \beta \chi}.\tag{140}$$

This result is identical to that from the statistical-mechanical analysis [50], which is called the naive mean-field theory [4]. If we put  $f(x) = \tanh(x/\sigma^2)$ ,  $\bar{\theta} = 0$  and G = 0, the BER of Proposition 5 can then be rewritten as

$$P_b = \mathcal{Q}\bigg(\frac{\tilde{\mathsf{E}}}{\sqrt{\tilde{\mathsf{F}}}}\bigg),\tag{141}$$

$$n = \int_{\mathbb{R}} \mathrm{D}z \tanh(\tilde{\mathsf{E}} + z\sqrt{\tilde{\mathsf{F}}}), \qquad (142)$$

$$c = \int_{\mathbb{R}} \mathrm{D}z \tanh^2(\tilde{\mathsf{E}} + z\sqrt{\tilde{\mathsf{F}}}), \qquad (143)$$

$$\tilde{\mathsf{E}} = \frac{1}{\sigma^2 + \beta(1-c)},\tag{144}$$

$$\tilde{\mathsf{F}} = \frac{\sigma_0^2 + \beta(1 - 2m + c)}{[\sigma^2 + \beta(1 - c)]^2},$$
(145)

where the identity,

$$\chi = \frac{1-c}{\sigma^2},\tag{146}$$

holds in this parameter settings. This result recovers that of the marginal-posterior-mode detectors [51]. It may mean that the relationship G = 0, which corresponds to the case where the self-coupling term vanishes, is required to achieve the individually optimal performance.

#### B. Derivation of Proposition 5

Since the covariance matrix  $\mathbf{R}$  of (91) becomes

$$\boldsymbol{R} = \left(\sum_{m=0}^{\infty} (-\beta \boldsymbol{G}^{\top})^{m}\right) \boldsymbol{D} \left(\sum_{n=0}^{\infty} (-\beta \boldsymbol{G})^{n}\right)$$
(147)  
$$= \boldsymbol{D} (\boldsymbol{1} + \beta \boldsymbol{G}^{\top} + \beta \boldsymbol{G} + \beta^{2} \boldsymbol{G}^{\top} \boldsymbol{G})^{-1}.$$
(148)

$$-D(\mathbf{I} + p\mathbf{G} + p\mathbf{G} + p\mathbf{G} + p\mathbf{G} + p\mathbf{G}), \quad (\mathbf{I} + \mathbf{0})$$

by using Assumption 4, we then have  $R^{+} = R$ . Similarly, the reaction matrix  $\Gamma$  is represented as

$$\boldsymbol{\Gamma} = (\mathbf{1} + \beta \boldsymbol{G})^{-1} \beta \boldsymbol{G} \tag{149}$$

$$=-\sum_{n=1}^{\infty}(-\beta)^{n}\boldsymbol{G}^{n}.$$
(150)

Since G is a lower triangular Toeplitz matrix, the  $\Gamma$  is also a lower triangular Toeplitz matrix, i.e.,  $\Gamma^{(s,s')} := \Gamma^{(s-s')}$  for s > s' and  $\Gamma^{(s,s')} = 0$  for  $s \le s'$ .

Under Assumption 3, it can be considered that the system responds to its past in an averaged way. We therefore consider the average of  $\phi^{(t)}$  instead of  $\lim_{t\to\infty} \phi^{(t+1)}$ , which gives

$$\bar{\phi} := \frac{1}{t+2} \sum_{t'=-1}^{t} \phi^{(t'+1)}$$
(151)

$$=\bar{\hat{k}} + \bar{v} + \bar{\theta} + \frac{1}{t+2} \sum_{t'=-1}^{t} \sum_{s=-1}^{t'} \Gamma^{(t',s)} f(\phi^{(s)}), \quad (152)$$

where we put  $\bar{\hat{k}} := \frac{1}{t+2} \sum_{t'=-1}^{t} \hat{k}^{(s)}$ ,  $\bar{v} := \frac{1}{t+2} \sum_{t'=-1}^{t} v^{(s)}$ and  $\bar{\theta} := \frac{1}{t+2} \sum_{t'=-1}^{t} \theta^{(s)}$ . Under Assumptions 1 and 2,  $\sum_{t'=-1}^{\infty} (G^{(t,t')})^n$  is written by

$$\lim_{t \to \infty} \sum_{t'=-1}^{t} (G^{(t,t')})^n$$

$$= \lim_{t \to \infty} \sum_{t'=-1}^{t} \left( \sum_{t_1=-1}^{t} \cdots \sum_{t_{n-1}=-1}^{t} G^{(t,t_1)} G^{(t,t_1)} \times \cdots \times G^{(t_{n-2},t_{n-1})} G^{(t_{n-1},t')} \right)$$
(153)
(154)

$$=\chi \lim_{t \to \infty} \sum_{t_1=-1}^{t} \cdots \sum_{t_{n-1}=-1}^{t} G^{(t,t_1)} G^{(t,t_1)} \times \cdots \times G^{(t_{n-2},t_{n-1})}$$
(155)  
$$=\chi^n.$$
(156)

$$\lim_{t \to \infty} \frac{1}{t+2} \sum_{t'=-1}^{t} \sum_{s=-1}^{t'} \Gamma^{(t',s)} f(\phi^{(s)})$$
(157)

$$=\bar{f}\sum_{n=1}^{\infty} (-\beta)^n \left(\lim_{t \to \infty} \sum_{s=-1}^t (G^{(t,s)})^n\right)$$
(158)

$$=\frac{\beta\chi}{1+\beta\chi}\bar{f},\tag{159}$$

where

=

$$\bar{f} := \frac{1}{t+2} \sum_{t'=-1}^{t} f(\phi^{(s)}).$$
(160)

Taking the  $t \to \infty$  limit and using (159), the average singleuser process  $\bar{\phi}$  of (152) becomes

$$\bar{\phi} = \bar{\hat{k}} + \bar{v} + \bar{\theta} + \frac{\beta\chi}{1 + \beta\chi}\bar{f}.$$
(161)

First, we evaluate  $\hat{k}$ , which corresponds to the signal part. If series  $\hat{k}^{(-1)}, \hat{k}^{(0)}, \cdots$  converges, then  $\hat{k} = \hat{k}^{(\infty)}$ . Since  $\hat{k}$  is obtained as  $\bar{k} = 1 - \beta \hat{k}^{(\infty)} \sum_{t'=-1}^{\infty} G^{(t')} = 1 - \beta \chi \bar{k}$ , we then have

$$\bar{\hat{k}} = \frac{1}{1+\beta\chi} =: \mathsf{E}.$$
 (162)

Next, we consider the variance of Gaussian random variable  $\bar{v}$ . The average of  $\bar{v}$  is obviously  $\langle \bar{v} \rangle_{\boldsymbol{v}} = \lim_{t \to \infty} \frac{1}{t+2} \sum_{t'=-1}^{t} \langle v^{(t')} \rangle_{\boldsymbol{v}} = 0$ . Using the definition of the persistent correlation  $c := \lim_{t \to \infty} C(t)$ , the variance of  $\bar{v}$  is given by

$$\langle \bar{v}^2 \rangle_{\boldsymbol{v}} = \lim_{\tau \to \infty} \frac{1}{(\tau+2)^2} \sum_{t=-1}^{\tau} \sum_{t'=-1}^{\tau} ((\mathbf{1} + \beta \boldsymbol{G}^{\top})^{-1} \boldsymbol{D} (\mathbf{1} + \beta \boldsymbol{G})^{-1})^{(t,t')}$$
(163)  
$$= \sum_{s=0}^{\infty} \sum_{s'=0}^{\infty} ((\mathbf{1} + \beta \boldsymbol{G}^{\top})^{-1})^{(s)} ((\mathbf{1} + \beta \boldsymbol{G}^{\top})^{-1})^{(s')}$$
$$\times \lim_{\tau \to \infty} \frac{1}{(\tau+2)^2} \sum_{t=s-1}^{\tau} \sum_{t'=s'-1}^{\tau} \boldsymbol{D}^{(t-s-t'+s')}$$
(164)  
$$= \frac{\tau^2 + \beta(1-2m+s)}{\tau^2 + \beta(1-2m+s)}$$

$$=\frac{\sigma_0^2 + \beta(1-2m+c)}{(1+\beta\chi)^2},$$
(165)

where we here write  $((\mathbf{1} + \beta \mathbf{G}^{\top})^{-1})^{(t+\tau,t)}$  as  $((\mathbf{1} + \beta \mathbf{G}^{\top})^{-1})^{(\tau)}$ , since  $(\mathbf{1}+\beta \mathbf{G}^{\top})^{-1}$  is a Toeplitz matrix. Namely,  $\bar{v}$  obeys a Gaussian distribution with mean zero and variance F:  $\bar{v} \sim \mathcal{N}(0, \mathsf{F})$ . The Gaussian random variable  $\bar{v}$  can be represented as  $\bar{v} = \sqrt{\mathsf{F}}z$  by using standard Gaussian random variable  $z \sim \mathcal{N}(0, 1)$ .

Limit  $f(\phi^{(t)})$  as  $t \to \infty$  converges by using Assumptions 1, 2 and 3. We therefore have  $\bar{\phi} = \lim_{t\to\infty} \frac{1}{t+2} \sum_{t'=-1}^{t} \phi^{(t')} = \phi^{(\infty)}$  and  $\bar{f} = \lim_{t\to\infty} \frac{1}{t+2} \sum_{t'=-1}^{t} f(\phi^{(t')}) = f(\phi^{(\infty)})$ . These relationships give  $\bar{f} = f(\bar{\phi})$ . Applying function f to both sides of (161) and letting  $g(z) := \bar{f}$ , we obtain

$$g(z) = f(\mathsf{E} + z\sqrt{\mathsf{F}} + \bar{\theta} + \mathsf{G}g(z)), \qquad (167)$$

where  $G := \beta \chi / (1 + \beta \chi)$ . The persistent overlap  $m := \lim_{t\to\infty} m^{(t)}$ , which is identical to that of (128), then becomes

$$m = \lim_{t \to \infty} \langle f(\phi^{(t)}) \rangle_{\boldsymbol{v}} \tag{168}$$

$$= \langle \bar{f} \rangle_{\bar{v}} = \int_{\mathbb{R}} \mathrm{D}z g(z). \tag{169}$$

We then have the persistent correlation *c*:

$$c = \lim_{\tau \to \infty} \lim_{t \to \infty} \langle f(\phi^{(t+\tau)}) f(\phi^{(t)}) \rangle_{\boldsymbol{v}}$$
(170)

$$= \langle \bar{f}^2 \rangle_{\bar{v}} = \int_{\mathbb{R}} \mathrm{D}z g(z)^2, \qquad (171)$$

in a similar way. Since we assume Assumption 2, the integrated response  $\chi$  of (131) is a finite non-negative number and is then given by

$$\chi = \lim_{t \to \infty} \sum_{t'=-1}^{t} \frac{\partial \langle f(\phi^{(t)}) \rangle_{\boldsymbol{v}}}{\partial \theta^{(t')}}$$
(172)

$$= \langle \frac{\partial f}{\partial \bar{\theta}} \rangle_{\bar{v}} = \int_{\mathbb{R}} \mathrm{D}z \frac{\partial g(z)}{\partial \bar{\theta}} = \frac{1}{\sqrt{\mathsf{F}}} \int_{\mathbb{R}} \mathrm{D}z z g(z).$$
(173)

For arbitrary function F(x) that can be expanded around x = 0, the bit error rate  $P_b := \lim_{t \to \infty} P_b^{(t)}$  can be evaluated as

$$\lim_{t \to \infty} \langle F(f(\phi^{(t)})) \rangle_{\boldsymbol{v}} = \langle F(\bar{f}) \rangle_{\bar{\boldsymbol{v}}} = \int_{\mathbb{R}} \mathrm{D}z F(g(z)).$$
(174)

Therefore, the bit error rate  $P_b$  can be evaluated as (133).

Setting  $\theta^{(s)} = 0$ , i.e.,  $\bar{\theta} = 0$ , we have then arrived at Proposition 5.

## C. Stability Analysis for Stationary Estimates

If Assumption 3 (WLTM) is violated, the dynamics do not achieve the stationary estimate. For example, oscillatory behavior in the dynamics corresponds to this situation. In this subsection, we investigate where Assumption 3 is violated and obtain the following proposition.

*Proposition 6:* The condition where WLTM is violated is given by

$$\frac{\beta}{(1+\beta\chi)^2} \int_{\mathbb{R}} \mathrm{D}z \left(\frac{\partial g(z)}{\partial \bar{\theta}}\right)^2 > 1, \tag{175}$$

where g(z) and all parameters are identical to the set of (134) – (140).

If we put  $f(x) = \tanh(x/\sigma^2)$  and  $\bar{\theta} = 0$  and G = 0, condition (175) can be simplified as

$$\frac{\beta}{[\sigma^2 + \beta(1-c)]^2} \int_{\mathbb{R}} \mathrm{D}z [1 - \tanh^2(\tilde{\mathsf{E}} + z\sqrt{\tilde{\mathsf{F}}})]^2 > 1, \quad (176)$$

where  $\tilde{E}$ ,  $\tilde{F}$ , and c are the same as (144), (145), and (146), respectively. Equation (176) coincides with the AT instability [1] of the marginal-posterior-mode detector [51]. This condition (176) is also identical to the microscopic instability, which is derived by the belief update of the BP-based detector [25].

In the rest of this subsection, we derive Proposition 6. Let us separate the response function  $G^{(t,t')}$  and the Onsager

reaction term  $\Gamma^{(t,t')}$  explicitly into TTI parts and further small contributions:

$$G^{(t,t')} = \tilde{G}^{(t-t')} + \hat{G}^{(t,t')}, \qquad (177)$$

$$\Gamma^{(t,t')} = \tilde{\Gamma}^{(t-t')} + \hat{\Gamma}^{(t,t')}, \qquad (178)$$

with  $\lim_{t\to\infty} G^{(t+\tau,t)} = \tilde{G}^{(\tau)}$ , and  $\lim_{t\to\infty} \Gamma^{(t+\tau,t)} = \tilde{\Gamma}^{(\tau)}$ , where  $\tilde{G}^{(t-t')}$  and  $\tilde{\Gamma}^{(t-t')}$  denote TTI parts that are referred to as *persistent parts*. The  $\hat{G}^{(t,t')}$  and  $\hat{\Gamma}^{(t,t')}$  denote small contribution terms that are referred to as *non-persistent parts*. If weak long-term memory holds, then  $\hat{G}^{(t,t')}$  and  $\hat{\Gamma}^{(t,t')}$ vanish for  $t \to \infty$ . We here assume that some small longterm memory, but such that the limits  $\hat{G}^{(t')} := \lim_{t\to\infty} \hat{G}^{(t,t')}$ and  $\hat{\Gamma}^{(t')} := \lim_{t\to\infty} \hat{\Gamma}^{(t,t')}$  exist and also both  $\hat{\chi} :=$  $\lim_{t\to\infty} \sum_{t'=-1}^t \hat{G}(t')$ ,  $\hat{\gamma} := \lim_{t\to\infty} \sum_{t'=-1}^t \hat{\Gamma}(t')$ ,  $\tilde{\chi} :=$  $\lim_{t\to\infty} \sum_{t'=-1}^t \tilde{G}(t')$  and  $\tilde{\gamma} := \lim_{t\to\infty} \sum_{t'=-1}^t \tilde{\Gamma}(t')$  exist. When Assumption 3 holds, the identities  $\lim_{t\to\infty} \hat{G}^{(t,t')} = 0$ and  $\lim_{t\to\infty} \hat{\Gamma}^{(t,t')} = 0$  hold.

We expand (92) for small  $\hat{G}$ , and find

$$\boldsymbol{\Gamma} = \sum_{n=1}^{\infty} (-\beta)^n \sum_{m=0}^n {n \choose m} \tilde{\boldsymbol{G}}^{n-m} \hat{\boldsymbol{G}}^m$$
(179)  
= $(\mathbf{1} + \beta \tilde{\boldsymbol{G}})^{-1} \beta \tilde{\boldsymbol{G}}$   
+ $\sum_{n=1}^{\infty} \sum_{m=0}^{n-1} (-\beta \tilde{\boldsymbol{G}})^m (\beta \hat{\boldsymbol{G}}) (-\beta \tilde{\boldsymbol{G}})^{n-1-m} + O(\tilde{\boldsymbol{G}}^2),$ (180)

where the (t, t') elements of the  $(t+2) \times (t+2)$  matrices  $\tilde{G}$ ,  $\hat{G}$ ,  $\tilde{\Gamma}$ , and  $\hat{\Gamma}$  are  $\tilde{G}^{(t,t')}$ ,  $\hat{G}^{(t,t')}$ ,  $\tilde{\Gamma}^{(t,t')}$ , and  $\hat{\Gamma}^{(t,t')}$ , respectively. From (180), the persistent and the non-parsistent parts of  $\Gamma$  is given by  $\tilde{\Gamma} = (\mathbf{1} + \beta \tilde{G})^{-1} \beta \tilde{G}$  and  $\hat{\Gamma}$ . Then,  $\hat{\gamma}$  becomes

$$\hat{\gamma} = \beta \lim_{t \to \infty} \sum_{t'=-1}^{t} \sum_{s=-1}^{t} \sum_{n=0}^{\infty} (n+1) (-\beta \tilde{G}^{(t,s)})^n \hat{G}^{(s,t')}$$
(181)  
$$= \frac{\beta \hat{\chi}}{(1+\beta \tilde{\chi})^2}.$$
(182)

We next substitute (177) into (123) and take the average with respect to iterative steps t, i.e.,  $\frac{1}{t+2}\sum_{t'=-1}^{t}$ , to both-sides of the substituted equation, which gives

$$\bar{\phi} = \bar{\hat{k}} + \bar{v} + \bar{\theta} + \frac{\beta \tilde{\chi}}{1 + \beta \tilde{\chi}} + \frac{1}{t+2} \sum_{t'=-1}^{t} \sum_{s=-1}^{t} \Gamma^{(t',s)} f(\phi^{(s)}).$$
(183)

As  $\lim_{t\to\infty} \tilde{G}^{(t,t')}$ , since we assume Assumption 2 (finite integrated response), we find

$$\hat{G}^{(t)} = \lim_{t \to \infty} G^{(t,t')}$$
 (184)

$$=\lim_{t\to\infty}\frac{\partial\langle f(\phi^{(t)})\rangle_{\boldsymbol{v}}}{\partial\theta^{(t')}}$$
(185)

$$= \lim_{t \to \infty} \left\langle \left( \frac{\partial \bar{f}}{\partial \bar{\theta}} \right)^2 \frac{1}{t+2} \sum_{s=-1}^t \hat{\Gamma}^{(s)} \right\rangle_{\bar{v}}.$$
 (186)

Substituting this equation into the definition of  $\hat{\chi}$ , the following identity is found:

$$\hat{\chi} = \hat{\gamma} \left\langle \left( \frac{\partial \bar{f}}{\partial \bar{\theta}} \right)^2 \right\rangle_{\bar{v}}.$$
(187)

Although  $\hat{\chi} = \hat{\gamma} = 0$  always solves this equation, one finds another solution when  $(\hat{\gamma}/\hat{\chi})\langle(\partial\bar{f}/\partial\bar{\theta})^2\rangle_{\bar{v}} = 1$  holds, which is equivalent to the AT line [1]. Since  $\hat{\chi} = 0$  is required to hold Assumption 2 and (187) represents an evaluation of  $\hat{\chi}$ , the evaluated value should not be greater than  $\hat{\chi}$ , i.e.,  $\hat{\chi} < \hat{\gamma}\langle(\partial\bar{f}/\partial\bar{\theta})^2\rangle_{\bar{v}}$ , to diverge stationary estimates. We therefore have an instability condition of Assumption 2 as

$$\frac{\beta}{(1+\beta\tilde{\chi})^2} \left\langle \left(\frac{\partial\bar{f}}{\partial\bar{\theta}}\right)^2 \right\rangle_{\bar{v}} > 1, \tag{188}$$

by substituting (182). Using  $\chi = \tilde{\chi} + \hat{\chi} \approx \tilde{\chi}$  and (137), we have then arrived at Proposition 6.

## IX. DECOUPLING PRINCIPLE

With the statistical-mechanical approach [51], [20], [21], a complex *many-user problem* such as an inference problem where many users' transmitted signals in the CDMA system is reduced to an equivalent *single-user problem*. This approach brings a significant interpretation in the communication theory. This is known as the *decoupling principle* found by Guo and Verdú [20]. It claims that the vector channel concatenated with optimal detection is equivalent to a bank of independent singleuser additive white Gaussian noise channels, whose signal-tonoise ratio degrades due to the multiple-access interference, in the large-system limit under a certain randomness assumption of the channel.

The decoupling principle has recently attracted a great deal of attention and has been investigated in detail [24], [42]. Especially, by applying density evolution, Ikehara and Tanaka have found that the decoupling principle holds not only at equilibrium but also at each stage of the BP-based detector [24]. Their analysis was however based on an assumption of independence of messages. The GFA allows us to study algorithms without assuming the independence of messages [6], [7], [32], [33], [34]. In this section, we investigate the decoupling principle via detection dynamics using GFA by considering an arbitrary fading model.

## A. Generic Detection Algorithm

In this section, we focus on a generic system model and a generic iterative multiuser detection algorithm.

We now consider a K-input N-output vector channel, which can be regarded as a fully-synchronous K-user CDMA system with spreading factor N or a multiple-input multiple-output (MIMO) system. Let  $\boldsymbol{b} = (b_1, \dots, b_K)^\top \in \mathbb{R}^K$  denote the input vector, and  $\boldsymbol{y} = (y_1, \dots, y_N)^\top \in \mathbb{R}^N$  denote the output vector of the channel. The output  $\boldsymbol{y}$  is given by a linear transform as

$$\boldsymbol{y} = \boldsymbol{S}\boldsymbol{b} + \boldsymbol{n},\tag{189}$$

where  $S = (a_1 s_1, \dots, a_K s_K)$  denotes an  $N \times K$  channel state matrix, which includes the received amplitude  $a_k$ .  $s_k = \frac{1}{\sqrt{N}} (s_{1k}, \dots, s_{Nk})^{\top}$  denote the channel parameters for the spreading code sequence of user k, and n is a vector consisting of i.i.d. zero-mean Gaussian random variables with variance  $\sigma_0^2$ . We use control parameter  $\sigma^2$  as the estimate of the channel noise variance instead of true noise variance. Each element  $s_{\mu k} \in \{-1, 1\}$  is an i.i.d. random variable with equal probability. Also, let  $p_{b}(b) = \prod_{k=1}^{K} p_{b}(b_{k})$  denote the joint prior distribution of input  $b \in \{-1, 1\}^{K}$ . Input distribution  $p_{b}(b)$  is arbitrary. The fading model is represented by arbitrary probability distribution  $p_{a}(a) = \prod_{k=1}^{K} p_{a}(a_{k})$ . The distribution  $p_{a}(a)$  is assumed to have finite moments. Especially it has the second-order moment  $\gamma$ . We analyze the system in the large-system limit where  $K, N \to \infty$ , while the ratio  $\beta = K/N$  is kept finite.

Definition 5: (Generic iterative algorithm) The updating rule for tentative decision  $\tilde{b}_k^{(t)} \in \mathbb{R}$  of bit signal  $b_k$  at stage t is

$$\tilde{b}_k^{(t+1)} = f(u_k^{(t)}), \tag{190}$$

with

$$u_{k}^{(t)} = R^{(t)}h_{k} - \sum_{s=-1}^{t} J^{(t,s)} \sum_{k=1}^{K} W_{kk'} \tilde{b}_{k'}(s) - \sum_{s=-1}^{t} \hat{\Gamma}_{k}^{(t,s)} \tilde{b}_{k}^{(s)} + \theta_{k}^{(t)}, \qquad (191)$$

where  $\boldsymbol{h} = (h_1, \dots, h_K)^{\top}$ ,  $\boldsymbol{W} = (W_{kk'})$ ,  $R^{(t)}$ ,  $J^{(t,s)}$ ,  $A^{(s)}$ ,  $\hat{\Gamma}_k^{(t,s)}$ , and  $\theta^{(t)}$  are parameters. Transfer function  $f : \mathbb{R} \to \mathbb{R}$  is arbitrary and applied componentwise.

The  $u_k^{(t)}$  is a summation of messages from tentative decisions. The  $\theta_k^{(t)}$  is again introduced to define a response function. Let  $h_k$  be an output of the matched filter for user k, i.e.,  $h_k = s_k \cdot y$ , and let W be a  $K \times K$  correlation matrix which is defined by  $W = \tilde{S}^{\top}\tilde{S}$  with  $\tilde{S} = (s_1, \dots, s_K)$ . For initialization, we use  $\tilde{b}_k^{(-1)} = 0$ . This generic algorithm includes various types of iterative multiuser detectors as in the following examples.

*Example 2:* If we set the transfer function f and the parameters as

$$f(x) = \tanh(x), \tag{192}$$

$$R^{(s)} = \sigma^{-2}, \tag{193}$$

$$J^{(s,s')} = \sigma^{-2} \delta_{s,s'}, \tag{194}$$

$$A^{(s)} = \theta_k^{(s)} = \hat{\Gamma}_k^{(s,s')} = 0, \qquad (195)$$

the algorithm of (190) is identical to soft-PIC [27]:  $\tilde{b}_k^{(t+1)} = \tanh[\frac{1}{\sigma^2}(h_k - \sum_{k'=1}^K W_{kk'}\tilde{b}_k^{(t)})].$ 

*Example 3:* All parameters except  $\hat{\Gamma}_k^{(s,s')}$  use those of soft-PIC. If the  $\hat{\Gamma}_k^{(s,s')}$  are chosen to cancel the Onsager reaction, then algorithm (190) gives ORC-PIC [52].

Example 4: We next consider the following parameters and

function:

$$f(x) = \tanh(x),\tag{196}$$

$$A^{(t)} = \{\sigma^2 + \beta [1 - Q^{(t)}]\}^{-1}, \tag{197}$$

$$R^{(t)} = A^{(t)} + A^{(t)}\beta[1 - Q^{(t)}]R^{(t-1)},$$
(198)

$$\theta_k^{(s)} = 0, \tag{199}$$

$$\hat{\Gamma}_{k}^{(s,s')} = -\delta_{s,s'} A^{(s)},$$
(200)

$$J^{(t,s)} = \mathbb{I}(0 \le s < t)A^{(s)} \prod_{s'=s+1}^{t} A^{(s')}\beta[1-q^{(s')}] + \mathbb{I}(s=t)A^{(t)},$$
(201)

with  $Q^{(t)} = K^{-1} \sum_{k=1}^{K} [\tilde{b}_{k}^{(t)}]^{2}$ . Here,  $\mathbb{I}(\mathcal{P})$  denotes an indicator function, which takes value 1 if  $\mathcal{P}$  is true, or 0 otherwise. We assume that Q(t) can be regarded as a constant with respect to  $\tilde{\boldsymbol{b}}$  by using the central limit theorem. Then (190) is equivalent to the BP-based detector [25].

As previously mentioned, the iterative detection algorithm (190) includes various types of iterative algorithms. It should be noted that in all these detectors,  $\tilde{b}_k^{(t)}$  gives an approximate value for the posterior-mean at the  $t^{\rm th}$  stage [52], [25]. We obtain the analytical result, which involves the following measure:

$$\langle g(\hat{\boldsymbol{b}}_{k}, \boldsymbol{v}_{k}) \rangle_{*_{k}} = \int \mathcal{D}\boldsymbol{v}_{k} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t-1} \mathrm{d}\hat{b}_{k}^{(s)} \right) g(\hat{\boldsymbol{b}}_{k}, \boldsymbol{v}_{k})$$
$$\times \delta [\tilde{b}_{k}^{(-1)}] \prod_{s=-1}^{t-1} \delta [\hat{b}_{k}^{(s+1)} - f(u_{k}^{(s)})], \quad (202)$$

with

$$u_k^{(s)} = a_k \hat{k}^{(s)} b_k + v_k^{(s)} + \sum_{s'=-1}^s [\Gamma^{(s,s')} - \hat{\Gamma}_k^{(s,s')}] \tilde{b}_k^{(s')}, \quad (203)$$

where  $g(\hat{\boldsymbol{b}}_k, \boldsymbol{v}_k)$  denotes an arbitrary function of  $\tilde{\boldsymbol{b}}_k = (\hat{b}_k^{(-1)}, \cdots, \tilde{b}_k^{(t)})^\top$  and  $\boldsymbol{v}_k = (v_k^{(-1)}, \cdots, v_k^{(t)})^\top$ . The parameters are as follows:

$$\boldsymbol{R} = (\boldsymbol{1} - \beta \boldsymbol{B})^{-1} \boldsymbol{D} (\boldsymbol{1} - \beta \boldsymbol{B}^{\top})^{-1}, \qquad (204)$$

$$\boldsymbol{\Gamma} = -[\sum_{s=-1}^{i-1} \boldsymbol{J}_{s+1} \otimes \boldsymbol{\Delta}^{s+1} (\boldsymbol{U}^{\top})^{-1}]^{\top}$$
(205)

$$\hat{k}^{(s)} = |\mathbf{\Lambda}_s|,\tag{206}$$

where B, U, and D are matrices with elements

$$B^{(s,s')} = -\sum_{\tau=-1}^{s} J^{(s,\tau)} G^{(\tau,s')},$$
(207)

$$U^{(s,s')} = \delta_{s,s'} + \beta \sum_{\tau=-1}^{s} J^{(s,\tau)} G^{(\tau,s')}$$
(208)

$$D^{(s,s')} = \sigma_0^2 R^{(s)} R^{(s')} + \beta [R^{(s)} R^{(s')} - R^{(s)} \sum_{\tau'=-1}^{s'} J^{(s',\tau')} m^{(\tau')} - R^{(s')} \sum_{\tau=-1}^{s} J^{(s,\tau)} m^{(\tau)} + \sum_{\tau=-1}^{s} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} C^{(\tau,\tau')}], \quad (209)$$
$$\Lambda_s^{(s',s'')} = \mathbb{I}(s' \neq s) \left( \delta_{s',s''} + \beta \sum_{\tau'=-1}^{s'} J^{(s',\tau')} G^{(\tau',s')} \right)$$

$$+ \mathbb{I}(s'=s)R^{(s')}.$$
 (210)

with  $\mathcal{D}\boldsymbol{v} = |2\pi\boldsymbol{R}|^{-1/2} d\boldsymbol{v} \exp[-\frac{1}{2}\boldsymbol{v} \cdot \boldsymbol{R}^{-1}\boldsymbol{v}]$ . All these parameters mentioned above are obtained from the following three kinds of quantities:

$$m^{(s)} = \lim_{K \to \infty} \frac{1}{K} \sum_{k=1}^{K} \langle a_k b_k \tilde{b}_k^{(s)} \rangle_{*_k}$$
$$= \langle a_k b_k \tilde{b}_k^{(s)} \rangle_{*_k, b_k, a_k},$$
(211)

$$C^{(s,s')} = \lim_{K \to \infty} \frac{1}{K} \sum_{k=1}^{K} \langle \tilde{b}_{k}^{(s)} \tilde{b}_{k}^{(s')} \rangle_{*_{k}}$$
  
=  $\langle \tilde{b}_{k}^{(s)} \tilde{b}_{k}^{(s')} \rangle_{*_{k}, b_{k}, a_{k}},$  (212)  
$$G^{(s,s')} = \mathbb{I}(s \leq s') \lim_{K \to \infty} \frac{1}{K} \sum_{k=1}^{K} \langle \tilde{b}_{k}^{(s)} (\mathbf{R}^{-1} \mathbf{v}_{k})^{(s')} \rangle_{*_{k}}$$

$$= \mathbb{I}(s \le s') \langle \tilde{b}_k^{(s)} (\boldsymbol{R}^{-1} \boldsymbol{v}_k)^{(s')} \rangle_{*_k, b_k, a_k}, \qquad (213)$$

where  $G^{(s,s')} = 0$  for s < s' due to causality. Here,  $\Delta = (\delta_{s+1,s'})$  is a  $(t+1) \times (t+1)$  matrix. Operator  $\langle \cdot \rangle_U$  denotes the expectation with respect to random variable U and operator  $\otimes$  denotes the Hadamard product, i.e.,  $A \otimes B = (a_{ij}b_{ij})$  for  $A = (a_{ij})$  and  $B = (b_{ij})$ . Terms  $(\mathbf{R}^{-1}\mathbf{v})^{(s)}$  denote the  $s^{\text{th}}$  element of the vector  $\mathbf{R}^{-1}\mathbf{v}$ . Term  $\sum_{s'=-1}^{s} \Gamma^{(s,s')}\tilde{b}^{(s')}$  in (203) represents the retarded self-interaction, which is called the Onsager reaction.

## B. Decoupling Principle

From the average of (202) we find the effective single-user process:  $\tilde{b}_k^{(s+1)} = f(u_k^{(s)})$  and  $\tilde{b}_k^{(-1)} = 0$ . Variable  $v_k$  can be regarded as the (t + 1)-dimensional Gaussian variable with mean vector **0** and covariance matrix **R**. Using (203), the effective single-user iterative process can be written in the

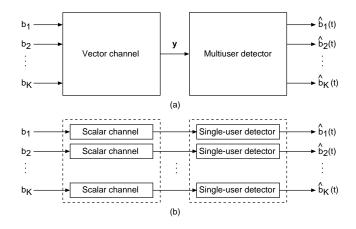


Fig. 4. Schematic of decoupling principle. (a) Vector channel followed by multiuser detection and estimates at stage t. (b) Equivalent scalar channels followed by single-user detectors and estimates at stage t.

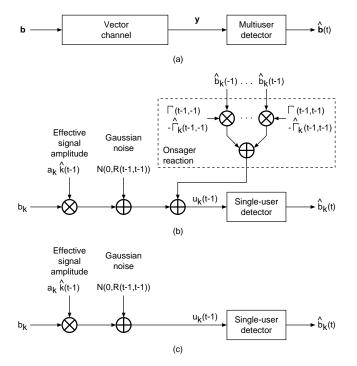


Fig. 5. Vector channel and equivalent single-user channels. (a) Vector channel, multiuser detection, and tentative decision. (b) Equivalent single-user non-Gaussian channel, single-user detection, and tentative decision. (c) Equivalent single-user Gaussian channel, single-user detection, and tentative decision.

following simple form.

$$u_{k}^{(s)} = a_{k} \hat{k}^{(s)} b_{k} + v_{k}^{(s)} + \sum_{s'=-1}^{s} [\Gamma^{(s,s')} - \hat{\Gamma}_{k}^{(s,s')}] f(u_{k}^{(s'-1)}), \quad (214)$$

with  $v_k^{(s)} \sim \mathcal{N}(0, R^{(s-1,s-1)})$ , where  $\tilde{b}_k^{(s)} = f(u_k^{(s-1)})$ ,  $\tilde{b}_k^{(-1)} = f(u_k^{(-2)}) = 0$ . Introducing the average over  $v_k$  as  $\langle \cdot \rangle_{v_k} = \int \mathcal{D}v_k \cdot$  and using (214), then  $m^{(s)}$ ,  $C^{(s,s')}$ , and  $G^{(s,s')}$  are obtained as  $m^{(s)} = \langle b_k \tilde{b}_k^{(s)} \rangle_{v,b_k,a_k}$ ,  $C^{(s,s')} = \langle \tilde{b}_k^{(s)} \tilde{b}_k^{(s')} \rangle_{v,b_k,a_k}$ , and  $G^{(s,s')} = \mathbb{I}(s \leq 1)$   $s'\rangle \langle \tilde{b}_k^{(s)}(\mathbf{R}^{-1}\boldsymbol{v}_k)^{(s')} \rangle_{\boldsymbol{v}_k, b_k, a_k}$ , respectively. The macroscopic parameters contain the effects of all users as averages; therefore, these can be represented by each user's distribution. Hence, (214) holds for any user k. All the parameters in (214) can only be represented by each user's distribution.

This result (214) means that each user experiences an equivalent single-user additive noise channel whose signal-to-noise ratio degrades due to both the multiple-access interference and the Onsager reaction, at each stage of the algorithm. This property is called the decoupling principle [20]. The additive channel noise generally becomes non-Gaussian due to the existence of the Onsager reaction  $\sum_{t'=-1}^{t-1} \Gamma^{(t-1,t')} f(u_k^{(t'-1)})$ . Figure 4 has a schematic of the decoupling principle.

Figure 5 shows the results of the analysis. Approximate *posterior-mean estimator (PME)*  $\tilde{b}_k^{(t)} = f(u_k^{(t-1)})$  is used as the single-user detector. It should be noted that (214) implies that equivalent channel noise becomes Gaussian if the Onsager reaction is cancelled. However, for soft-PIC, the Onsager reaction is not cancelled (Fig. 5 (b)). If the Onsager reaction does not vanish, the noise distribution is not centered at zero. Therefore, soft-PIC must be suboptimal.

Under the assumption of the large-system limit, the performance of algorithm (190) at stage t can be evaluated by using the equivalent scalar channel determined from three types of macroscopic parameters  $(\boldsymbol{m}, \boldsymbol{C}, \boldsymbol{G})$ , where  $\boldsymbol{m}, \boldsymbol{C}$ , and  $\boldsymbol{G}$  are a (t+1)-dimensional vector  $(m^{(-1)}, \cdots, m^{(t-1)})^{\top}$ , a  $(t+1) \times (t+1)$  matrix  $(C^{(s,s')})$  and a  $(t+1) \times (t+1)$  matrix  $(G^{(s,s')})$ , respectively. The bit error rate  $P_b^{(t)}$  of hard decisions sgn  $[\tilde{b}_k^{(t)}]$  at the t<sup>th</sup> stage is given by  $P_b^{(t)} = \frac{1}{2}(1 - \langle b_k \operatorname{sgn}[\tilde{b}_k^{(t)}] \rangle_{\boldsymbol{v}, b_k, a_k})$ .

We next consider the case where we set the parameters in (190) to those of ORC-PIC. The reaction term in the effective single-user process is determined by two matrices  $\Gamma$  and  $\hat{\Gamma}_k$ . It is considered that the parameter  $\Gamma$  represents a retarded self-interaction caused by iterative calculation. The parameter  $\hat{\Gamma}_k$ , on the other hand, is arbitrary; therefore we can choose  $\hat{\Gamma}_k$  that cancel the Onsager reaction. Using the inductive method, we can show that the Onsager reaction is entirely cancelled at each stage t, when the parameter  $\hat{\Gamma}^{(s,s')}$  is set to  $(\hat{\Gamma}\tilde{b})^{(t)} = \beta G^{(t,t-1)}[\tilde{b}^{(t-1)} - (\hat{\Gamma}\tilde{b})^{(t-1)}]$  and  $(\hat{\Gamma}\tilde{b})^{(-1)} = (\hat{\Gamma}\tilde{b})^{(0)} = 0$ . This parameter  $\hat{\Gamma}$  coincides with the parameter derived by density evolution [52]. Density evolution cannot treat the Onsager reaction, however density evolution can be applied to derive an algorithm that can correctly cancel the Onsager reaction.

In the case where we set the parameters in (190) to those of the BP-based detector, we confirmed the Onsager reaction is cancelled at least at stage 8 through straightforward calculations. In this case,  $\Gamma$  becomes diag $(A^{(-1)}, 0, \dots, 0)$ . Therefore, for both ORC-PIC and the BP-based detector (at least at stage 8), the equivalent single-user channels are given as an additive white Gaussian noise channel (Fig. 5 (c)).

# X. CONCLUSION

We evaluated the detection dynamics of soft-PIC, ORC-PIC and the BP-based detector by applying GFA in the large-system limit. We also showed that GFA could treat the dynamics of belief propagation. From the practical point of view, iterative algorithms are generally utilized as suboptimal methods. It is important to understand the detection dynamics in detail to improve detectors. We confirmed that the results from the density evolution analysis could be obtained from those of the GFA analysis by omitting the Onsager reaction term. The GFA could correctly treat the Onsager reaction, which is a retarded self-interaction due to iterations, and gives density evolution a basic theory.

We also studied the decoupling principle in iterative multiuser detection algorithms using GFA. and found that the decoupling principle holds. The reconstruction algorithm of compressed sensing could also be analyzed by GFA [37].

# APPENDIX A NOTATIONS

We here summarize notations, which are used in this paper, in Table I. It should be noted that notations to describe the GFA result don't contain the user number k, since a manyuser problem is reduced to an equivalent single-user problem in GFA. In Table I 'BP' denotes BP-based detector.

# APPENDIX B Derivation of Lemma 1

Evaluating the averaged generating functional straightforwardly, one can obtain Lemma 1. Substituting (1) and (3) into (24), the summation over all messages becomes

$$u_{k}^{(t)} = \sum_{k'=1}^{K} W_{kk'}(b_{k'} - \tilde{b}_{k'}^{(t)}) + \tilde{b}_{k}^{(t)} + \frac{\sigma_{0}}{\sqrt{N}} \sum_{\mu=1}^{N} s_{k}^{\mu} n^{\mu} + \theta_{k}^{(t)}.$$
(215)

TABLE I			
SUMMARY OF NOTATIONS			

SUMMART OF NOTATIONS				
variable	the first or top-left element	description		
(notations for models) $\begin{split} &K \in \mathbb{N}^+ \\ &N \in \mathbb{N}^+ \\ &N, k' \in \{1, \cdots, K\} \\ &\mu, \mu' \in \{1, \cdots, N\} \\ &t \in \mathbb{N}^+ \\ &s, s', t' \in \{-1, 0, \cdots, t\} \\ &s_k = (s_k^{n}) \in \{\pm 1\}^N \\ &b = (b_k) \in \{\pm 1\}^K \\ &g = (y^{\mu}) \in \mathbb{R}^N \\ &n = (n^{\mu}) \in \mathbb{R}^N \\ &n = (n^{\mu}) \in \mathbb{R}^N \\ &n = (m^{\mu}) \in \mathbb{R}^N \\ &m = (W_{kk'}) \in \mathbb{R}^K \\ &b = (b_k) \in \mathbb{R}^K \\ &b = (b_k) \in \mathbb{R}^K \\ &b = (h_k) \in \mathbb{R}^K \\ &f : \mathbb{R} \to \mathbb{R} \\ &\hat{\Gamma}_k = (\hat{\Gamma}_k^{(s,s')}) \in \mathbb{R}^{(t+1) \times (t+1)} \\ &R^{(t)} \in \mathbb{R}^+ \\ &A^{(t)} \in \mathbb{R}^+ \\ &J^{(t,s)} \in \mathbb{R}^+ \\ &J \in \mathbb{R}^+ \\ &S = (a_k s_k) \in \mathbb{R}^{N \times K} \\ &a_k \in \mathbb{R}^+ \end{split}$	$\begin{array}{c} - & - & - & - & - & - & - & - & - & - $	the number of users spreading code length user number chip interval index stage stage spreading code of user k information bits tentative soft decisions tentative hard decisions received signals noise vector true noise variance control parameter correlation matrix dummy functions external message matched filter outputs transfer function parameter for BP parameter for BP parameter for BP parameter for BP inverse temperature channel state matrix received amplitude		
$P_b^{(t)} \in \mathbb{R}^+$	_	bit error rate		
$ \begin{array}{l} (\text{notations for GFA}) \\ \tilde{\boldsymbol{b}} = (\tilde{\boldsymbol{b}}^{(t)}) \in \mathbb{R}^{t+1} \\ \tilde{\boldsymbol{b}} = (\tilde{\boldsymbol{b}}^{(t)}) \in \{\pm 1\}^{t+1} \\ \boldsymbol{\psi} = (\boldsymbol{\psi}^{(t)}) \in \mathbb{R}^{t+1} \\ \boldsymbol{\theta} = (\boldsymbol{\theta}^{(t)}) \in \mathbb{R}^{t+1} \\ \boldsymbol{\eta} = (\boldsymbol{\eta}^{(s)}) \in \mathbb{R}^{t+1} \\ \tilde{\boldsymbol{\eta}} = (\tilde{\boldsymbol{\eta}}^{(s)}) \in \mathbb{R}^{t+1} \\ \tilde{\boldsymbol{k}} = (k^{(s)}) \in \mathbb{R}^{t+1} \\ \tilde{\boldsymbol{k}} = (\tilde{\boldsymbol{k}}^{(s)}) \in \mathbb{R}^{t+1} \\ \boldsymbol{q} = (\boldsymbol{q}^{(s,s')}) \in \mathbb{R}^{(t+1)\times(t+1)} \\ \boldsymbol{q} = (\boldsymbol{q}^{(s,s')}) \in \mathbb{R}^{(t+1)\times(t+1)} \\ \boldsymbol{Q} = (Q^{(s,s')}) \in \mathbb{R}^{(t+1)\times(t+1)} \\ \boldsymbol{Q} = (Q^{(s,s')}) \in \mathbb{R}^{(t+1)\times(t+1)} \\ \boldsymbol{L} = (L^{(s,s')}) \in \mathbb{R}^{(t+1)\times(t+1)} \\ \boldsymbol{L} = (L^{(s,s')}) \in \mathbb{R}^{(t+1)\times(t+1)} \\ \boldsymbol{m} = (\boldsymbol{m}^{(s)}) \in \mathbb{R}^{t+1} \\ \boldsymbol{C} = (C^{(s,s')}) \in \mathbb{R}^{(t+1)\times(t+1)} \\ \boldsymbol{G} = (G^{(s,s')}) \in \mathbb{R}^{(t+1)\times(t+1)} \\ \end{array} $	$\begin{array}{l} \tilde{b}(-1) \\ \hat{b}(-1) \\ \psi(-1) \\ \eta(-1) \\ \eta(-1) \\ \hat{\eta}(-1) \\ \hat{k}(-1) \\ \hat{k}(-1) \\ \hat{q}(-1,-1) \\ \hat{q}(-1,-1) \\ Q(-1,-1) \\ \hat{Q}(-1,-1) \\ \hat{L}(-1,-1) \\ \hat{L}(-1,-1) \\ m(-1) \\ C(-1,-1) \\ G(-1,-1) \end{array}$	tentative soft decisions tentative hard decisions dummy functions external message macroscopic parameter macroscopic parameter overlap correlation function response function		

The average generating functional  $\bar{Z}[\psi]$  is represented using  $\boldsymbol{u}^{(s)} = (u_1^{(s)}, \cdots, u_K^{(s)})^{\top}$  whose elements are defined by (215). Using Dirac's delta function  $\delta$ , we first introduce the definition of  $\boldsymbol{u}^{(-1)}, \cdots, \boldsymbol{u}^{(t-1)}$  into  $Z[\psi]$ , which is written in terms of them, as

$$\bar{Z}[\psi] = \mathbb{E}_{\boldsymbol{s}_{1}, \cdots, \boldsymbol{s}_{K}, \boldsymbol{n}} \left[ \int_{\mathbb{R}^{(t+1)K}} \underline{d}\boldsymbol{u} \\
\times p[\tilde{\boldsymbol{b}}^{(-1)}] \left( \prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} e^{-\frac{\gamma^{2}}{2} \{\tilde{\boldsymbol{b}}_{k}^{(s+1)} - f(\boldsymbol{u}_{k}^{(s)})\}^{2}} \right) \\
\times \exp\left[ -i \sum_{s=-1}^{t} \tilde{\boldsymbol{b}}^{(s)} \cdot \boldsymbol{\psi}^{(s)} \right] \\
\times \prod_{s=-1}^{t-1} \prod_{k=1}^{K} \delta\left( \boldsymbol{u}_{k}^{(s)} - \sum_{k'=1}^{K} W_{kk'}(\boldsymbol{b}_{k'} - \tilde{\boldsymbol{b}}_{k'}^{(s)}) - \tilde{\boldsymbol{b}}_{k}^{(s)} - \frac{\sigma_{0}}{\sqrt{N}} \sum_{\mu=1}^{N} s_{k}^{\mu} \boldsymbol{n}^{\mu} - \boldsymbol{\theta}_{k}^{(s)} \right) \right],$$
(216)

where  $\underline{\mathrm{d}} u{:=}\prod_{s=-1}^{t-1}\prod_{k=1}^{K}\frac{\mathrm{d} u_{k}^{(s)}}{\sqrt{2\pi}}$ . The arguments of Dirac's delta functions  $\delta(\cdot)$  represent the definition of  $\{u_{k}^{(t)}\}$ . The term of  $\delta$  is rewritten by applying the Fourier integral form of Dirac's delta function  $\delta(x) = \int_{\mathbb{R}} \mathrm{d} \hat{x} e^{\mathrm{i} \hat{x} x}$  as  $\delta(u_{k}^{(s)} - [\cdots]) = \int_{\mathbb{R}} \mathrm{d} \hat{u}_{k}^{(s)} \exp\{\mathrm{i} \hat{u}_{k}^{(s)}(u_{k}^{(s)} - [\cdots])\}$ . We then have

$$\bar{Z}[\psi] = \sum_{\tilde{b}^{(-1)}, \dots, \tilde{b}^{(t)}} p[\tilde{b}^{(-1)}] \int_{\mathbb{R}^{2(t+1)K}} \underline{d}u \underline{d}\hat{u} \\
\times \exp\left[i \sum_{s=-1}^{t-1} \sum_{k=1}^{K} \hat{u}_{k}^{(s)} \{u_{k}^{(s)} - \tilde{b}_{k}^{(s)} - \theta_{k}^{(s)}\}\right] \\
-i \sum_{s=-1}^{t} \sum_{k=1}^{K} \tilde{b}_{k}^{(s)} \psi_{k}^{(s)} \\
+ \sum_{s=-1}^{t} \sum_{k=1}^{K} \{\ln \frac{\gamma}{\sqrt{2\pi}} - \frac{\gamma^{2}}{2} [\tilde{b}_{k}^{(t+1)} - f(u_{k}^{(s)})]^{2}\}\right] \\
\times \mathbb{E}_{s_{1}, \dots, s_{K}, n} \left\{ \exp\left[-i\sqrt{\beta}\sigma_{0} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \left(\frac{1}{\sqrt{K}} \sum_{k=1}^{K} s_{k}^{\mu} \hat{u}_{k}^{(s)}\right) n^{\mu} \\
-i\beta \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \left(\frac{1}{\sqrt{K}} \sum_{k=1}^{K} s_{k}^{\mu} \hat{u}_{k}^{(s)}\right) \\
\times \left(\frac{1}{\sqrt{K}} \sum_{k'=1}^{K} s_{k'}^{\mu} \{b_{k'} - \tilde{b}_{k'}^{(s)}\}\right) \right] \right\}, \quad (217)$$

where  $\underline{d}\hat{u} := \prod_{s=-1}^{t-1} \prod_{k=1}^{K} \frac{d\hat{u}_{k}^{(s)}}{\sqrt{2\pi}}$ . Without loss of generality, we can set to  $b_{k} = 1$  ( $\forall k$ ).

To take the average of  $s_1, \dots, s_K$ , we introduce variables

 $v^{(s)}_{\mu}$  and  $w^{(s)}_{\mu}$  which are defined as:

$$v_k^{(s)} := \frac{1}{\sqrt{K}} \sum_{k=1}^K s_k^\mu \hat{u}_k^{(s)}, \qquad (218)$$

$$w_k^{(s)} := \frac{1}{\sqrt{K}} \sum_{k=1}^K s_k^{\mu} \{ 1 - \tilde{b}_k^{(s)} \}.$$
 (219)

It should be noted that the averaged generating functional  $\overline{Z}[\psi]$  of (217) includes  $s_1, \dots, s_K$  only in terms of these variables  $v_{\mu}^{(s)}$  and  $w_{\mu}^{(s)}$ . Due to this, random variables  $s_1, \dots, s_K$  are isolated and are able to be averaged. Introducing  $v_{\mu}^{(s)}$  and  $w_{\mu}^{(s)}$  into term  $\mathbb{E}_{s_1,\dots,s_K,n}\{\dots\}$  in (217), one obtains

$$\int_{\mathbb{R}^{2tN}} \underline{\mathrm{d}} \boldsymbol{v} \underline{\mathrm{d}} \boldsymbol{w} \mathbb{E}_{\boldsymbol{s}_{1}, \cdots, \boldsymbol{s}_{K}, \boldsymbol{n}} \left\{ \exp\left[-\mathrm{i}\sqrt{\beta}\sigma_{0} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} v_{\mu}^{(s)} n^{\mu} - \mathrm{i}\beta \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} v_{\mu}^{(s)} w_{\mu}^{(s)}\right] \right\} \\
\times \delta\left(v_{k}^{(s)} - \frac{1}{\sqrt{K}} \sum_{k=1}^{K} s_{k}^{\mu} \hat{u}_{k}^{(s)}\right) \\
\times \delta\left(w_{k}^{(s)} - \frac{1}{\sqrt{K}} \sum_{k=1}^{K} s_{k}^{\mu} \{1 - \tilde{b}_{k}^{(s)}\}\right) \\
\int_{\mathbb{R}^{4tN}} \underline{\mathrm{d}} v \underline{\mathrm{d}} \hat{v} \underline{\mathrm{d}} w \underline{\mathrm{d}} \hat{w} \\
\times \exp\left[\mathrm{i} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \{\hat{v}_{\mu}^{(s)} v_{\mu}^{(s)} + \hat{w}_{\mu}^{(s)} w_{\mu}^{(s)} - \beta v_{\mu}^{(s)} w_{\mu}^{(s)}\}\right] \\
\times \mathbb{E}_{\boldsymbol{n}} \left\{ \exp\left[-\mathrm{i}\sqrt{\beta}\sigma_{0} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} v_{\mu}^{(s)} n^{\mu}\right] \right\} \\
\times \mathbb{E}_{\boldsymbol{s}_{1}, \cdots, \boldsymbol{s}_{K}} \left\{ \exp\left[-\mathrm{i} \frac{1}{\sqrt{K}} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \{\hat{v}_{\mu}^{(s)} \sum_{k=1}^{K} s_{k}^{\mu} \hat{u}_{k}^{(s)} + \hat{w}_{\mu}^{(s)} \sum_{k=1}^{K} s_{k}^{\mu} (1 - \tilde{b}_{k}^{(s)}) \right\} \right] \right\}. \tag{220}$$

=

We here again used the Fourier integral form of Dirac's delta. The term  $\mathbb{E}_n\{\cdots\}$  in (220) becomes

$$\mathbb{E}_{n} \left\{ \exp\left[-i\sqrt{\beta}\sigma_{0}\sum_{\mu=1}^{N}\sum_{s=-1}^{t-1}v_{\mu}^{(s)}n^{\mu}\right] \right\} \\ = \prod_{\mu=1}^{N} \exp\left[-\frac{1}{2}\beta\sigma_{0}^{2}\sum_{s=-1}^{t-1}\sum_{s'=-1}^{t-1}v_{\mu}^{(s)}v_{\mu}^{(s')}\right]. \quad (221)$$

Since t is finite, term  $\mathbb{E}_{s_1,\dots,s_K}\{\cdots\}$  in (220) is given by

$$\mathbb{E}_{s_{1},\cdots,s_{K}} \left\{ \exp\left[-i\frac{1}{\sqrt{K}}\sum_{\mu=1}^{N}\sum_{s=-1}^{t-1} \left\{\hat{v}_{\mu}^{(s)}\sum_{k=1}^{K}s_{k}^{\mu}\hat{u}_{k}^{(s)}\right. \\ \left. +\hat{w}_{\mu}^{(s)}\sum_{k=1}^{K}s_{k}^{\mu}(1-\tilde{b}_{k}^{(s)})\right\} \right] \right\} \\ = \prod_{\mu=1}^{N}\prod_{k=1}^{K}\cos\left[-\frac{1}{\sqrt{K}}\sum_{s=-1}^{t-1} \left\{\hat{v}_{\mu}^{(s)}\hat{u}_{k}^{(s)} + \hat{w}_{\mu}^{(s)}(1-\tilde{b}_{k}^{(s)})\right\}\right] \\ = \prod_{\mu=1}^{N}\prod_{k=1}^{K}\exp\left[-\frac{1}{2}\left(\frac{1}{\sqrt{K}}\sum_{s=-1}^{t-1} \left\{\hat{v}_{\mu}^{(s)}\hat{u}_{k}^{(s)} + \hat{w}_{\mu}^{(s)}(1-\tilde{b}_{k}^{(s)})\right\}\right)^{2} + O(K^{-2})\right] \\ = \prod_{\mu=1}^{N}\exp\left[-\frac{1}{2}\sum_{s=-1}^{t-1}\sum_{s'=-1}^{t-1} \left\{\hat{v}_{\mu}^{(s)}\left(\frac{1}{K}\sum_{k=1}^{K}\hat{u}_{k}^{(s)}\hat{u}_{k}^{(s')}\right)\hat{v}_{\mu}^{(s')} + \hat{v}_{\mu}^{(s)}\left(\frac{1}{K}\sum_{k=1}^{K}\hat{u}_{k}^{(s)} - \frac{1}{K}\sum_{k=1}^{K}\tilde{b}_{k}^{(s')}\hat{u}_{k}^{(s)}\right)\hat{v}_{\mu}^{(s')} \\ + \hat{w}_{\mu}^{(s)}\left(\frac{1}{K}\sum_{k=1}^{K}\hat{u}_{k}^{(s')} - \frac{1}{K}\sum_{k=1}^{K}\tilde{b}_{k}^{(s)}\hat{u}_{k}^{(s')}\right)\hat{v}_{\mu}^{(s')} \\ + \hat{w}_{\mu}^{(s)}\left(1 - \frac{1}{K}\sum_{k=1}^{K}\tilde{b}_{k}^{(s)} - \frac{1}{K}\sum_{k=1}^{K}\tilde{b}_{k}^{(s')}\right)\hat{v}_{\mu}^{(s')} \\ + \frac{1}{K}\sum_{k=1}^{K}\tilde{b}_{k}^{(s)}\hat{b}_{k}^{(s')}\right)\hat{v}_{\mu}^{(s')}\right\} + O(K^{-1})\right]. \tag{222}$$

We next separate the relevant one-stage and two-stage values which are appeared in (222):

$$\eta^{(s)} := \frac{1}{\sqrt{K}} \sum_{k=1}^{K} \tilde{b}_k^{(s)}, \qquad (223)$$

$$k^{(s)} := \frac{1}{\sqrt{K}} \sum_{k=1}^{K} \hat{u}_k^{(s)}, \qquad (224)$$

$$q^{(s,s')} := \frac{1}{\sqrt{K}} \sum_{k=1}^{K} \tilde{b}_k^{(s)} \tilde{b}_k^{(s')}, \qquad (225)$$

$$Q^{(s,s')} := \frac{1}{\sqrt{K}} \sum_{k=1}^{K} \hat{u}_k^{(s)} \hat{u}_k^{(s')}, \qquad (226)$$

$$L^{(s,s')} := \frac{1}{\sqrt{K}} \sum_{k=1}^{K} \tilde{b}_k^{(s)} \hat{u}_k^{(s')}.$$
 (227)

Equation (222) is only written in terms of these variables  $\eta^{(s)}$ ,  $k^{(s)}$ ,  $q^{(s,s')}$ ,  $Q^{(s,s')}$  and  $L^{(s,s')}$ . We hereafter call these variables macroscopic parameters, which describe the nature of the system and are also often called the order parameters. Similar to the way to derive (220), one can introduce  $m^{(s)}$ ,  $k^{(s)}$ ,  $q^{(s,s')}$ ,  $Q^{(s,s')}$  and  $L^{(s,s')}$  into (222). The term of  $\delta$  in (222) is again rewritten by applying the Fourier integral form of Dirac's delta such as  $\delta(\eta^{(s)} - [\cdots]) = \int_{\mathbb{R}} d\hat{\eta}^{(s)} \exp\{i\hat{\eta}^{(s)}(\eta^{(s)} - [\cdots])\}, \,\delta(k^{(s)} - [\cdots]) = \int_{\mathbb{R}} d\hat{q}^{(s,s')} \exp\{i\hat{q}^{(s,s')} - [\cdots])\}, \,\delta(Q^{(s,s')} - [\cdots]) = \int_{\mathbb{R}} d\hat{Q}^{(s,s')}$ 

$$\begin{split} & \exp\{\mathrm{i}\hat{Q}^{(s,s')}(Q^{(s,s')}-[\cdots])\} \ \text{and} \ \delta(L^{(s,s')}-[\cdots]) = \int_{\mathbb{R}} \\ & d\hat{L}^{(s,s')} \ \exp\{\mathrm{i}\hat{L}^{(s,s')}(L^{(s,s')}-[\cdots])\}. \ \text{We here introduce} \\ & \text{the notations} \ \underline{d}u \ := \ \prod_{s=-1}^{t-1} \frac{\mathrm{d}u^{(s)}}{\sqrt{2\pi}} \ \underline{d}\hat{u} \ := \ \prod_{s=-1}^{t-1} \frac{\mathrm{d}\hat{u}^{(s)}}{\sqrt{2\pi}}, \\ & \mathrm{d}\boldsymbol{\eta} \ := \ \prod_{s=-1}^{t-1} \mathrm{d}\boldsymbol{\eta}^{(s)}, \ \mathrm{d}\hat{\boldsymbol{\eta}} \ := \ \prod_{s=-1}^{t-1} \mathrm{d}\hat{\boldsymbol{\eta}}^{(s)}, \\ & \mathrm{d}\boldsymbol{k} \ := \ \prod_{s=-1}^{t-1} \mathrm{d}\boldsymbol{k}^{(s)}, \ \mathrm{d}\hat{\boldsymbol{\mu}} \ := \ \prod_{s=-1}^{t-1} \mathrm{d}\hat{\boldsymbol{\eta}}^{(s)}, \\ & \mathrm{d}\boldsymbol{k} \ := \ \prod_{s=-1}^{t-1} \mathrm{d}\boldsymbol{q}^{(s,s')}, \ \mathrm{d}\hat{\boldsymbol{q}} \ := \ \prod_{s=-1}^{t-1} \mathrm{d}\hat{\boldsymbol{k}}^{(s)}, \ \mathrm{d}\boldsymbol{q} \ := \\ & \prod_{s=-1}^{t-1} \prod_{s'=-1}^{t-1} \mathrm{d}\boldsymbol{q}^{(s,s')}, \ \mathrm{d}\hat{\boldsymbol{q}} \ := \ \prod_{s=-1}^{t-1} \prod_{s'=-1}^{t-1} \mathrm{d}\hat{\boldsymbol{q}}^{(s,s')}, \\ & \mathrm{d}\boldsymbol{Q} \ := \ \prod_{s=-1}^{t-1} \prod_{s'=-1}^{t-1} \mathrm{d}\boldsymbol{Q}^{(s,s')}, \ \mathrm{d}\boldsymbol{L} \ := \ \prod_{s=-1}^{t-1} \prod_{s'=-1}^{t-1} \mathrm{d}\boldsymbol{L}^{(s,s')} \\ & \mathrm{d}\hat{\boldsymbol{L}} \ := \ \prod_{s=-1}^{t-1} \prod_{s'=-1}^{t-1} \mathrm{d}\hat{\boldsymbol{L}}^{(s,s')}, \ \mathrm{d}\boldsymbol{v} \ := \ \prod_{\mu=1}^{N} \prod_{s=-1}^{t-1} \frac{\mathrm{d}\boldsymbol{v}^{(s)}}{\sqrt{2\pi}}, \\ & \mathrm{d}\hat{\boldsymbol{v}} \ := \ \prod_{\mu=1}^{N} \prod_{s=-1}^{t-1} \mathrm{d}\hat{\boldsymbol{L}}^{(s,s')}. \ \mathrm{d}\boldsymbol{v} \ := \ \prod_{\mu=1}^{N} \prod_{s=-1}^{t-1} \frac{\mathrm{d}\boldsymbol{w}^{(s)}}{\sqrt{2\pi}}, \\ & \mathrm{d}\hat{\boldsymbol{w}} \ := \ \prod_{\mu=1}^{N} \prod_{s=-1}^{t-1} \frac{\mathrm{d}\hat{\boldsymbol{v}}^{(s)}}{\sqrt{2\pi}}, \ \mathrm{d}\boldsymbol{w} \ := \ \prod_{\mu=1}^{N} \prod_{s=-1}^{t-1} \frac{\mathrm{d}\boldsymbol{w}^{(s)}}{\sqrt{2\pi}}, \ \mathrm{d}\hat{\boldsymbol{w}} \ := \ \prod_{\mu=1}^{K} \delta(\tilde{\boldsymbol{b}^{(-1)}_k)), \ \text{is factorizable}, \\ & \mathrm{the averaged generating functional} \ Z[\boldsymbol{\psi}] \ factorizes into \\ & \mathrm{single-user contributions, i.e., with respect to user index \ k. \\ & \text{We have then arrived at Lemma 1. \\ \end{array}$$

# APPENDIX C DERIVATION OF PROPOSITION 1

Taking the limit of (83) and (84) for all k and s, the single-user measure (40) becomes user independent. It can be, therefore, represented without the user index k. Equation (40) becomes

$$\lim_{\{\psi_{k}^{(s)} \to 0\}} \lim_{\{\theta_{k}^{(s)} \to 0\}} \langle f(\boldsymbol{b}, \boldsymbol{u}, \hat{\boldsymbol{u}}) \rangle_{k}$$

$$= \frac{\int \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t} \mathrm{d} \tilde{b}^{(s)} \right) w_{*}(\tilde{\boldsymbol{b}}, \boldsymbol{u}, \hat{\boldsymbol{u}}) f(\tilde{\boldsymbol{b}}, \boldsymbol{u}, \hat{\boldsymbol{u}})}{\int \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t} \mathrm{d} \tilde{b}^{(s)} \right) w_{*}(\tilde{\boldsymbol{b}}, \boldsymbol{u}, \hat{\boldsymbol{u}})} \quad (228)$$

$$=: \langle f(\tilde{\boldsymbol{b}}, \boldsymbol{u}, \hat{\boldsymbol{u}}) \rangle_{*}, \quad (229)$$

where

$$w_{*}(\tilde{\boldsymbol{b}}, \boldsymbol{u}, \hat{\boldsymbol{u}})$$

$$:= \delta[\tilde{b}^{(-1)}] \exp\left[\sum_{s=-1}^{t-1} \{\ln \frac{\gamma}{\sqrt{2\pi}} - \frac{\gamma^{2}}{2} [\tilde{b}^{(s+1)} - f(\boldsymbol{u}^{(s)})]^{2} \}\right]$$

$$- i \sum_{s=0}^{t-1} \sum_{s'=0}^{t-1} \{\hat{Q}^{(s,s')} \hat{\boldsymbol{u}}^{(s)} \hat{\boldsymbol{u}}^{(s')} + \hat{L}^{(s,s')} \tilde{b}^{(s)} \hat{\boldsymbol{u}}^{(s')} \}$$

$$+ i \sum_{s=0}^{t-1} \hat{\boldsymbol{u}}^{(s)} \{\boldsymbol{u}^{(s)} - \tilde{b}^{(s)} - \hat{k}^{(s)} \} \right]_{\text{saddle}}.$$
(230)

We consider two kind of functions as  $f(\hat{b}, u, \hat{u})$ . First, we treat an arbitrary function  $f(\tilde{b})$  that does not include u and  $\hat{u}$ . In this case, one can, therefore, perform the integral with

respect to u and  $\hat{u}$ . The numerator of (228) becomes

$$\int \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t} \mathrm{d} \tilde{b}^{(s)} \right) w_*(\tilde{\boldsymbol{b}}, \boldsymbol{u}, \hat{\boldsymbol{u}}) f(\tilde{\boldsymbol{b}})$$

$$= \int \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t} \mathrm{d} \tilde{b}^{(s)} \right) \delta[\tilde{b}^{(-1)}]$$

$$\times \left( \prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} \mathrm{e}^{-\frac{\gamma^2}{2} [\tilde{b}^{(s+1)} - f(\boldsymbol{u}^{(s)})]} \right)$$

$$\times \mathrm{e}^{-\frac{1}{2} \hat{\boldsymbol{u}}^{\top} (\mathbf{1} + \beta \boldsymbol{G})^{-1} \boldsymbol{D} (\mathbf{1} + \beta \boldsymbol{G})^{-1} \hat{\boldsymbol{u}}}$$

$$\times \mathrm{e}^{\mathrm{i} \tilde{\boldsymbol{b}}^{\top} (\mathbf{1} + \beta \boldsymbol{G}^{\top})^{-1} \hat{\boldsymbol{u}} + \mathrm{i} \hat{\boldsymbol{u}} \cdot (\boldsymbol{u} - \tilde{\boldsymbol{b}} - \hat{\boldsymbol{k}}) f(\tilde{\boldsymbol{b}})}$$

$$= \int \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t} \mathrm{d} \tilde{\boldsymbol{b}}^{(s)} \right) \delta[\tilde{b}^{(-1)}]$$

$$\times \left( \prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} \mathrm{e}^{-\frac{\gamma^2}{2} [\tilde{b}^{(s+1)} - f(\boldsymbol{u}^{(s)})]} \right)$$

$$\times \mathrm{e}^{-\frac{1}{2} \hat{\boldsymbol{u}}^{\top} \boldsymbol{R} \hat{\boldsymbol{u}} + \mathrm{i} \hat{\boldsymbol{u}} \cdot (\boldsymbol{u} - \hat{\boldsymbol{k}} - \boldsymbol{\Gamma} \tilde{\boldsymbol{b}})} f(\tilde{\boldsymbol{b}}), \qquad (231)$$

where  $\mathbf{R} := (\mathbf{1} + \beta \mathbf{G})^{-1} \mathbf{D} (\mathbf{1} + \beta \mathbf{G})^{-1}$  and  $\Gamma := (\mathbf{1} + \beta \mathbf{G})^{-1} \beta \mathbf{G}$ . We use variable transformations  $\mathbf{v} = \mathbf{u} - \hat{\mathbf{k}} - \Gamma \tilde{\mathbf{b}}$ and  $\mathbf{w} = \hat{\mathbf{u}}$ , which gives

$$\begin{split} &\int \left(\prod_{s=-1}^{t-1} \frac{\partial(u^{(s)}, \hat{u}^{(s)})}{\partial(v^{(s)}, w^{(s)})} \frac{\mathrm{d}v^{(s)} \mathrm{d}w^{(s)}}{2\pi}\right) \mathrm{e}^{\mathrm{i}\boldsymbol{v}\cdot\boldsymbol{w} - \frac{1}{2}\boldsymbol{w}\cdot\boldsymbol{R}\boldsymbol{w}} \\ &\times \int_{\mathbb{R}^{t+2}} \left(\prod_{s=-1}^{t} \mathrm{d}\tilde{b}^{(s)}\right) f(\tilde{\boldsymbol{b}}) \delta[\tilde{b}^{(-1)}] \\ &\times \left(\prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} \mathrm{e}^{-\frac{\gamma^2}{2}[\tilde{b}^{(s+1)} - f(\hat{k}^{(s)} + v^{(s)} + (\Gamma\tilde{\boldsymbol{b}})^{(s)})]^2}\right) \\ &= \int \left(\prod_{s=-1}^{t-1} \frac{\mathrm{d}v^{(s)}}{\sqrt{2\pi}}\right) \sqrt{|\boldsymbol{R}^{-1}|} \mathrm{e}^{-\frac{1}{2}\boldsymbol{v}\cdot\boldsymbol{R}\boldsymbol{v}} \\ &\times \int_{\mathbb{R}^{t+2}} \left(\prod_{s=-1}^{t} \mathrm{d}\tilde{b}^{(s)}\right) f(\tilde{\boldsymbol{b}}) \delta[\tilde{b}^{(-1)}] \\ &\times \left(\prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} \mathrm{e}^{-\frac{\gamma^2}{2}[\tilde{b}^{(s+1)} - f(\hat{k}^{(s)} + v^{(s)} + (\Gamma\tilde{\boldsymbol{b}})^{(s)})]^2}\right). \end{split}$$
(232)

Taking the limit  $\gamma \rightarrow \infty,$  the numerator of (40), finally, arrives at

$$\int \mathcal{D}\boldsymbol{v} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t-1} \mathrm{d}\tilde{b}^{(s)} \right) f(\tilde{\boldsymbol{b}}) \, \delta[\tilde{b}^{(-1)}] \\ \times \prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} \mathrm{e}^{-\frac{\gamma^2}{2} (\tilde{b}^{(s+1)} - f(\hat{k}^{(s)} + v^{(s)} + (\boldsymbol{\Gamma}\tilde{\boldsymbol{b}})^{(s)}))^2} =: \langle\!\langle f(\tilde{\boldsymbol{b}}) \rangle\!\rangle,$$
(233)

where  $\mathcal{D}\boldsymbol{v} := \mathrm{d}\boldsymbol{v} |2\pi\boldsymbol{R}|^{-1/2} e^{-\frac{1}{2}\boldsymbol{v}\cdot\boldsymbol{R}^{-1}\boldsymbol{v}}$ . Since  $\langle\!\langle 1 \rangle\!\rangle = 1$ , the effective path measure (40) in the limit  $\gamma \to \infty$  is given by

$$\lim_{\gamma \to \infty} \langle f(\tilde{\boldsymbol{b}}) \rangle_* = \frac{\langle\!\langle f(\boldsymbol{b}) \rangle\!\rangle}{\langle\!\langle 1 \rangle\!\rangle} \\ = \langle\!\langle f(\tilde{\boldsymbol{b}}) \rangle\!\rangle.$$
(234)

Equation (85) can be obtained from (234), (51) and (60). Equation (86) is given by (234), (55) and (61).

Next, we consider a function  $\tilde{b}^{(s)}\hat{u}$  that corresponds to (59). In a similar way to have (233), one obtain

$$\int \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t} \mathrm{d} \tilde{b}^{(s)} \right) w_{*}(\tilde{b}, u, \hat{u}) \tilde{b}^{(s)} \hat{u}$$

$$= \int \left( \prod_{s=-1}^{t-1} \frac{\partial(u^{(s)}, \hat{u}^{(s)})}{\partial(v^{(s)}, w^{(s)})} \frac{\mathrm{d} v^{(s)} \mathrm{d} w^{(s)}}{2\pi} \right) w \mathrm{e}^{\mathrm{i} \boldsymbol{v} \cdot \boldsymbol{w} - \frac{1}{2} \boldsymbol{w} \cdot \boldsymbol{R} \boldsymbol{w}}$$

$$\times \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t} \mathrm{d} \tilde{b}^{(s)} \right) \tilde{b}^{(s)} \delta[\tilde{b}^{(-1)}]$$

$$\times \left( \prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} \mathrm{e}^{-\frac{\gamma^{2}}{2} [\tilde{b}^{(s+1)} - f(\hat{k}^{(s)} + v^{(s)} + (\Gamma \tilde{b})^{(s)})]^{2}} \right)$$

$$= \int \mathcal{D} \boldsymbol{v} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t-1} \mathrm{d} \tilde{b}^{(s)} \right) \mathrm{i} \tilde{b}^{(s)} \boldsymbol{R}^{-1} \boldsymbol{v} \delta[\tilde{b}^{(-1)}]$$

$$\times \prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} \mathrm{e}^{-\frac{\gamma^{2}}{2} (\tilde{b}^{(s+1)} - f(\hat{k}^{(s)} + v^{(s)} + (\Gamma \tilde{b})^{(s)}))^{2}}$$

$$= \langle \mathrm{i} \tilde{b}^{(s)} \boldsymbol{R}^{-1} \boldsymbol{v} \rangle, \qquad (235)$$

by using the notation of (233). We therefore have

$$\int \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t} \mathrm{d} \tilde{b}^{(s)} \right) w_*(\tilde{\boldsymbol{b}}, \boldsymbol{u}, \hat{\boldsymbol{u}}) \tilde{b}^{(s)} \hat{u}^{(s')}$$
$$= \langle\!\langle \mathrm{i} \tilde{b}^{(s)} (\boldsymbol{R}^{-1} \boldsymbol{v})^{(s')} \rangle\!\rangle.$$
(236)

Equation (87) can be obtained from (236), (234), (59) and (62). We have then arrived at Proposition 1.

# APPENDIX D DERIVATION OF PROPOSITION 3

The inductive method is applied.

(i) The case of t = -1. This is an initial stage, therefore the Onsager reaction term does not exist. To cancel the Onsager reaction term, the matrix  $\hat{\Gamma}\hat{\boldsymbol{b}}$  is simply chosen as  $(\hat{\Gamma}\hat{\boldsymbol{b}})^{(-1)} = 0$ . This is automatically held by the initialization  $\tilde{\boldsymbol{b}}^{(-1)} = 0$ . The matrix G is G = 0 because of causality.

(ii) The case of t = 0. In this case, G is a  $1 \times 1$  zero matrix O and  $\Gamma$  becomes  $\Gamma = (\mathbf{1} + \beta G)^{-1}\beta G = O$ . The Onsager reaction term is therefore  $\Gamma \hat{\boldsymbol{b}} = O$ . To cancel the Onsager reaction term, we only have to choose matrix  $\hat{\Gamma}$  as  $(\hat{\Gamma} \hat{\boldsymbol{b}})^{(-1)} = 0$ . Matrix G is

$$G = \begin{pmatrix} 0 & 0 \\ G^{(0,-1)} & 0 \end{pmatrix}.$$
 (237)

until here.

(iii) The case of stage t = 1. Since  $G \in \mathbb{R}^{2 \times 2}$  is a nilpotent matrix, i.e.,  $G^2 = O$ , then the Onsager reaction term  $\Gamma \hat{\boldsymbol{b}}$  becomes

$$\Gamma \hat{\boldsymbol{b}} = [(\mathbf{1} + \beta G)^{-1} \beta G] \hat{\boldsymbol{b}}$$
  
=  $\beta G \hat{\boldsymbol{b}}$   
=  $\begin{pmatrix} 0 & 0 \\ \beta G^{(0,-1)} & 0 \end{pmatrix} \begin{pmatrix} 0 \\ \tilde{\boldsymbol{b}}^{(0)} \end{pmatrix}$   
=  $\begin{pmatrix} 0 \\ 0 \end{pmatrix}.$  (238)

To cancel the Onsager reaction term, one needs  $(\hat{\Gamma}\hat{b})^{(0)} = 0$ . If  $(\hat{\Gamma}\hat{b})^{(0)}$  is chosen as this, the average  $\langle\!\langle \tilde{b}^{(1)} \rangle\!\rangle = \int \mathcal{D} \boldsymbol{v}$  sgn  $[\hat{k}^{(0)} + v^{(0)} + \theta^{(0)}]$  only includes  $\theta^{(0)}$  and does not depend on  $\theta^{(-1)}$ . Therefore,  $G^{(1,-1)} = \partial\langle\!\langle \tilde{b}^{(1)} \rangle\!\rangle / \partial \theta^{(-1)} = 0$  holds. The matrix G is

$$G = \begin{pmatrix} 0 & 0 & 0 \\ G^{(0,-1)} & 0 & 0 \\ 0 & G^{(1,0)} & 0 \end{pmatrix}$$
(239)

until here.

(iv) The case of stage t = 2. Since  $G \in \mathbb{R}^{3 \times 3}$  is a nilpotent matrix, i.e.,  $G^3 = O$ , then the ORCT  $\Gamma \hat{b}$  is

$$\Gamma \hat{\boldsymbol{b}} = [(\mathbf{1} + \beta G)^{-1} \beta G] \hat{\boldsymbol{b}} 
 = [\beta G - (\beta G)^{2}] \hat{\boldsymbol{b}} 
 = \begin{pmatrix} 0 & 0 & 0 \\ \beta G^{(0,-1)} & 0 & 0 \\ -\beta^{2} G^{(1,0)} G^{(0,-1)} & \beta G^{(1,0)} & 0 \end{pmatrix} \begin{pmatrix} 0 \\ \tilde{\boldsymbol{b}}^{(0)} \\ \tilde{\boldsymbol{b}}^{(1)} \end{pmatrix} 
 = \begin{pmatrix} 0 \\ \beta G^{(1,0)} \tilde{\boldsymbol{b}}^{(0)} \end{pmatrix}.$$
(240)

Therefore, if  $\hat{\Gamma}\hat{\boldsymbol{b}}$  is chosen as

$$(\hat{\Gamma}\hat{\boldsymbol{b}})^{(1)} = \beta G^{(1,0)}\tilde{b}^{(0)}, \qquad (241)$$

then it cancels the Onsager reaction term. When we choose  $(\hat{\Gamma}\hat{\boldsymbol{b}})^{(1)}$  as this, the average  $\langle\!\langle \tilde{b}^{(2)} \rangle\!\rangle = \int \mathcal{D}\boldsymbol{v} \operatorname{sgn} [\hat{k}^{(1)} + v^{(1)} + \theta^{(1)}]$  only includes  $\theta(1)$  and does not depend on  $\theta^{(-1)}$  and  $\theta(0)$ . Therefore,  $G^{(2,s)} = \partial \langle\!\langle \tilde{b}^{(2)} \rangle\!\rangle / \partial \theta^{(-1)} = 0$  holds for  $s \in \{-1, 0\}$ . The matrix G is

$$G = \begin{pmatrix} 0 & 0 & 0 & 0 \\ G^{(0,-1)} & 0 & 0 & 0 \\ 0 & G^{(1,0)} & 0 & 0 \\ 0 & 0 & G^{(2,1)} & 0 \end{pmatrix}$$
(242)

until here.

(v) The case of stage t = 3. Since  $G \in \mathbb{R}^{4 \times 4}$  is a nilpotent matrix, i.e.,  $G^4 = O$ , then the ORCT  $\Gamma \hat{b}$  is

$$\Gamma \hat{\boldsymbol{b}} = [(\mathbf{1} + \beta G)^{-1} \beta G] \hat{\boldsymbol{b}} 
= [\beta G - (\beta G)^2 + (\beta G)^3] \hat{\boldsymbol{b}} 
= \begin{pmatrix} 0 \\ 0 \\ \beta G^{(1,0)} \tilde{\boldsymbol{b}}^{(0)} \\ \beta G^{(2,1)} \tilde{\boldsymbol{b}}^{(1)} - \beta^2 G^{(2,1)} G^{(1,0)} \tilde{\boldsymbol{b}}^{(0)} \end{pmatrix}.$$
(243)

Therefore, if  $\hat{\Gamma}\hat{\boldsymbol{b}}$  is chosen as

$$(\hat{\Gamma}\hat{\boldsymbol{b}})^{(2)} = \beta G^{(2,1)} \tilde{b}^{(1)} - \beta^2 G^{(2,1)} G^{(1,0)} \tilde{b}^{(0)}, \qquad (244)$$

then it cancels the Onsager reaction term. When we choose  $(\hat{\Gamma}\hat{\boldsymbol{b}})^{(2)}$  as this, the average  $\langle\!\langle \tilde{b}^{(3)} \rangle\!\rangle = \int \mathcal{D}\boldsymbol{v} \operatorname{sgn}[\hat{k}^{(2)} + v^{(2)} + \theta^{(2)}]$  only includes  $\theta(1)$  and does not depend on  $\theta^{(-1)}, \dots, \theta(1)$ . Therefore,  $G^{(3,s)} = \partial\langle\!\langle \tilde{b}^{(3)} \rangle\!\rangle / \partial \theta^{(-1)} = 0$ 

holds for  $s \in \{-1, 0, 1\}$ . The matrix G is

$$G = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ G^{(0,-1)} & 0 & 0 & 0 & 0 \\ 0 & G^{(1,0)} & 0 & 0 & 0 \\ 0 & 0 & G^{(2,1)} & 0 & 0 \\ 0 & 0 & 0 & G^{(3,2)} & 0 \end{pmatrix}$$
(245)

until here. Note that it is expected that

$$(\hat{\Gamma}\hat{\boldsymbol{b}})^{(t)} = \beta G^{(t,t-1)} [\tilde{\boldsymbol{b}}^{(t-1)} - (\hat{\Gamma}\hat{\boldsymbol{b}})^{(t-1)}] = -\sum_{s=0}^{t-1} (-\beta)^{t-s} \left(\prod_{\tau=s}^{t-1} G^{(\tau+1,\tau)}\right) \tilde{\boldsymbol{b}}^{(s)}$$
(246)

will holds for any  $t \ge 1$  in an analogy with (244).

(vi) Stage  $t \ge 3$ 

We here assume that the Theorem holds for stage  $t \in \{-1, 0, 1, \dots, t\}$ . Namely, we assume that  $(\hat{\Gamma}\hat{\boldsymbol{b}})^{(-1)}, \dots, (\hat{\Gamma}\hat{\boldsymbol{b}})^{(t)}$  are given by

$$\begin{split} (\hat{\Gamma}\boldsymbol{b})^{(-1)} &= 0\\ (\hat{\Gamma}\boldsymbol{b})^{(0)} &= 0\\ (\hat{\Gamma}\boldsymbol{b})^{(1)} &= \beta G^{(1,0)} \tilde{b}^{(0)}\\ \vdots\\ (\hat{\Gamma}\boldsymbol{b})^{(t)} &= -\sum_{s=0}^{t-1} (-\beta)^{t-s} \left(\prod_{\tau=s}^{t-1} G^{(\tau+1,\tau)}\right) \tilde{b}^{(s)}, (247) \end{split}$$

and G has the following form:

$$G = \begin{pmatrix} 0 & & & 0 \\ G^{(0,-1)} & 0 & & & \\ & G^{(1,0)} & 0 & & \\ & & \ddots & \ddots & \\ 0 & & & G^{(t+1,t)} & 0 \end{pmatrix}.$$
 (248)

Since  $G \in \mathbb{R}^{(t+3)\times(t+3)}$  is a nilpotent matrix, i.e.,  $G^{t+3} = O$ , then Onsager reaction term  $\Gamma \hat{\boldsymbol{b}} \in \mathbb{R}^{t+3}$  can be calculated as

$$\Gamma \boldsymbol{b} = [(\mathbf{1} + \beta G)^{-1} \beta G] \boldsymbol{b} \\
 = [\beta G - (\beta G)^{2} + \dots + (-1)^{t+1} (\beta G)^{t+2}] \hat{\boldsymbol{b}} \\
 = \begin{pmatrix} 0 \\ 0 \\ \beta G^{(1,0)} \tilde{\boldsymbol{b}}^{(0)} \\ \beta G^{(2,1)} \tilde{\boldsymbol{b}}^{(1)} - \beta^{2} G^{(2,1)} G^{(1,0)} \tilde{\boldsymbol{b}}^{(0)} \\ \vdots \\ - \sum_{s=0}^{t-1} (-\beta)^{t-s} \left(\prod_{\tau=s}^{t-1} G^{(\tau+1,\tau)}\right) \tilde{\boldsymbol{b}}^{(s)} \\ - \sum_{s=0}^{t} (-\beta)^{t+1-s} \left(\prod_{\tau=s}^{t} G^{(\tau+1,\tau)}\right) \tilde{\boldsymbol{b}}^{(s)} \end{pmatrix}.$$
(249)

If  $\hat{\Gamma}\hat{b}$  is chosen as

$$(\hat{\Gamma}\hat{\boldsymbol{b}})^{(t+1)} = -\sum_{s=0}^{t} (-\beta)^{t+1-s} \left(\prod_{\tau=s}^{t} G^{(\tau+1,\tau)}\right) \tilde{\boldsymbol{b}}^{(s)}, \quad (250)$$

the Onsager reaction term is cancelled. When we choose  $(\hat{\Gamma}\hat{\boldsymbol{b}})^{(t+1)}$  as this, the average  $\langle\!\langle \tilde{\boldsymbol{b}}^{(t+2)} \rangle\!\rangle = \int \mathcal{D}\boldsymbol{v} \operatorname{sgn}[\hat{k}^{(t+1)} +$ 

 $v^{(t+1)} + \theta^{(t+1)}$ ] only includes  $\theta^{(t+1)}$  and does not depend on  $\theta^{(-1)}, \cdots, \theta^{(t)}$ . Therefore,  $G^{(t+2,s)} = \partial \langle \langle \tilde{b}^{(3)} \rangle / \partial \theta^{(-1)} = 0$  holds for  $s \in \{-1, 0, 1, \cdots, t\}$ .

If the claim holds for stage  $s \in \{-1, 0, 1, \dots, t\}$ , it holds for stage t + 1. This proves Proposition 3.

# APPENDIX E Derivation of Proposition 4

One can evaluate the BP-based detector in the same manner as soft-PIC. The summation over all messages, which is the argument of  $tanh(\cdots)$  of (14) becomes

$$u_{k}^{(t)} = R^{(t)} \sum_{k'=1}^{K} W_{kk'} b_{k'} + A^{(t)} \tilde{b}_{k}^{(t)} - \sum_{s=-1}^{t} J^{(t,s)} \sum_{k'=1}^{K} W_{kk'} \tilde{b}_{k}^{(s)} + R^{(s)} \frac{\sigma_{0}}{\sqrt{N}} \sum_{\mu=1}^{N} s_{k}^{\mu} n^{\mu} + \theta_{k}^{(s)}$$
(251)

Substituting (1) and (3) into (24), the averaged generating functional  $\bar{Z}[\psi]$  is represented as

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$$\begin{split} [\psi] &= \sum_{\hat{b}^{(0)}, \cdots, \hat{b}^{(t)}} p[\hat{b}(0)] \int_{\mathbb{R}^{2(t+1)K}} \underline{\mathrm{d}} u \underline{\mathrm{d}} \hat{u} \\ &\times \exp\left[\mathrm{i} \sum_{s=0}^{t-1} \sum_{k=1}^{K} \hat{u}_{k}^{(s)} \{u_{k}^{(s)} - A^{(s)} \tilde{b}_{k}^{(s)} - \theta_{k}^{(s)}\} \right] \\ &+ \sum_{s=0}^{t} \sum_{k=1}^{K} \{\ln \frac{\gamma}{\sqrt{2\pi}} - \frac{\gamma^{2}}{2} [\tilde{b}_{k}^{(t+1)} - f(u_{k}^{(s)})]^{2}\} \right] \\ &\times \mathbb{E}_{s_{1}, \cdots, s_{K}, n} \left\{ \exp\left[ -\mathrm{i} \sqrt{\beta} \sigma_{0} \sum_{\mu=1}^{N} \sum_{s=0}^{t-1} R^{(s)} \left( \frac{1}{\sqrt{K}} \sum_{k=1}^{K} s_{k}^{\mu} \hat{u}_{k}^{(s)} \right) n^{\mu} \right. \\ &- \mathrm{i} \beta \sum_{\mu=1}^{N} \sum_{s=0}^{t-1} \left( \frac{1}{\sqrt{K}} \sum_{k=1}^{K} s_{k}^{\mu} \hat{u}_{k}^{(s)} \right) \\ &\times \left( \frac{1}{\sqrt{K}} \sum_{k'=1}^{K} s_{k'}^{\mu} \left\{ R^{(s)} - \sum_{s'=-1}^{s} J^{(s,s')} \tilde{b}_{k}^{(s')} \right\} \right) \right] \right\}, \end{split}$$

$$(253)$$

where  $\underline{d}\hat{\boldsymbol{u}} := \prod_{s=-1}^{t-1} \prod_{k=1}^{K} \frac{d\hat{u}_{k}^{(s)}}{\sqrt{2\pi}}$ . To take the average of  $\boldsymbol{s}_{1}, \cdots, \boldsymbol{s}_{K}$ , we introduce the following variables:

$$\begin{split} \bar{Z}[\psi] = &\mathbb{E}_{\boldsymbol{s}_{1},\cdots,\boldsymbol{s}_{K},\boldsymbol{n}} \left[ \int_{\mathbb{R}^{(t+1)K}} \underline{\mathrm{d}}\boldsymbol{u} \\ &\times p[\tilde{\boldsymbol{b}}^{(-1)}] \left( \prod_{s=-1}^{t-1} \frac{\gamma}{\sqrt{2\pi}} e^{-\frac{\gamma^{2}}{2}} \{\tilde{\boldsymbol{b}}_{k}^{(s+1)} - f(\boldsymbol{u}_{k}^{(s)})\}^{2} \right) \\ &\times \exp\left[ -\mathrm{i} \sum_{s=-1}^{t} \tilde{\boldsymbol{b}}^{(s)} \cdot \boldsymbol{\psi}^{(s)} \right] \\ &\times \prod_{s=0}^{t-1} \prod_{k=1}^{K} \delta\left( \boldsymbol{u}_{k}^{(s)} - R^{(t)} \sum_{k'=1}^{K} W_{kk'} \boldsymbol{b}_{k'} - A^{(t)} \tilde{\boldsymbol{b}}_{k}^{(t)} \\ &+ \sum_{s=-1}^{t} J^{(t,s)} \sum_{k'=1}^{K} W_{kk'} \tilde{\boldsymbol{b}}_{k}^{(s)} - R^{(s)} \frac{\sigma_{0}}{\sqrt{N}} \sum_{\mu=1}^{N} s_{k}^{\mu} \boldsymbol{n}^{\mu} \\ &- \theta_{k}^{(s)} \right) \bigg], \end{split}$$
(252)

where  $\underline{d}u := \prod_{s=0}^{t-1} \prod_{k=1}^{K} \frac{du_k^{(s)}}{\sqrt{2\pi}}$ . Without loss of generality we can put  $b_k = 1$  ( $\forall k$ ). We apply the Fourier integral form of

$$v_k^{(s)} := \frac{1}{\sqrt{K}} \sum_{k=1}^K s_k^\mu \hat{u}_k^{(s)}, \tag{254}$$

$$w_k^{(s)} := \frac{1}{\sqrt{K}} \sum_{k=1}^K s_k^{\mu} \bigg\{ R^{(s)} - \sum_{s'=-1}^s J^{(s,s')} \tilde{b}_k^{(s')} \bigg\}.$$
 (255)

It should be noted that the averaged generating functional  $\bar{Z}[\psi]$  of (27) only includes  $s_1, \dots, s_K$  in terms of these variables  $v_{\mu}^{(s)}$  and  $w_{\mu}^{(s)}$ . Due to this, the random variables  $s_1, \dots, s_K$  are isolated and their average is able to be taken. Introducing  $v_{\mu}^{(s)}$  and  $w_{\mu}^{(s)}$  into the term  $\mathbb{E}_{s_1,\dots,s_K,n}\{\dots\}$  in (27), one

Dirac's delta function, which gives

 $\bar{Z}$ 

obtains

$$\int_{\mathbb{R}^{2tN}} \underline{\mathrm{d}} \boldsymbol{v} \underline{\mathrm{d}} \boldsymbol{w} \mathbb{E}_{\boldsymbol{s}_{1}, \cdots, \boldsymbol{s}_{K}, \boldsymbol{n}} \left\{ \exp\left[-\mathrm{i}\sqrt{\beta}\sigma_{0} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} R^{(s)} v_{\mu}^{(s)} n^{\mu} - \mathrm{i}\beta \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} v_{\mu}^{(s)} w_{\mu}^{(s)}\right] \right] \\
\times \delta\left(v_{k}^{(s)} - \frac{1}{\sqrt{K}} \sum_{k=1}^{K} s_{k}^{\mu} \hat{u}_{k}^{(s)}\right) \\
\times \delta\left(w_{k}^{(s)} - \frac{1}{\sqrt{K}} \sum_{k=1}^{K} s_{k}^{\mu} \left\{R^{(s)} - \sum_{s'=-1}^{s} J^{(s,s')} \tilde{b}_{k}^{(s')}\right\}\right) \\
= \int_{\mathbb{R}^{4tN}} \underline{\mathrm{d}} v \underline{\mathrm{d}} \hat{v} \underline{\mathrm{d}} w \underline{\mathrm{d}} \hat{w} \\
\times \exp\left[\mathrm{i} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \left\{\hat{v}_{\mu}^{(s)} v_{\mu}^{(s)} + \hat{w}_{\mu}^{(s)} w_{\mu}^{(s)} - \beta v_{\mu}^{(s)} w_{\mu}^{(s)}\right\}\right] \\
\times \mathbb{E}_{\boldsymbol{n}} \left\{\exp\left[-\mathrm{i}\sqrt{\beta}\sigma_{0} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} R^{(s)} v_{\mu}^{(s)} n^{\mu}\right]\right\} \\
\times \mathbb{E}_{\boldsymbol{s}_{1}, \cdots, \boldsymbol{s}_{K}} \left\{\exp\left[-\mathrm{i} \frac{1}{\sqrt{K}} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \left(\hat{v}_{\mu}^{(s)} \sum_{k=1}^{K} s_{\mu}^{\mu} \hat{u}_{k}^{(s)} + \hat{w}_{\mu}^{(s)} \sum_{k=1}^{K} s_{\mu}^{\mu} \hat{u}_{k}^{(s)} + \hat{w}_{\mu}^{(s)} \sum_{k=1}^{K} s_{\mu}^{\mu} (R^{(s)} - \sum_{s'=-1}^{s} J^{(s,s')} \tilde{b}_{k}^{(s)})\right)\right]\right\}, \quad (256)$$

Term  $\mathbb{E}_{\boldsymbol{s}_1,\cdots,\boldsymbol{s}_K}\{\cdots\}$  in (256) becomes

=

=

$$\begin{split} \mathbb{E}_{s_{1},\cdots,s_{K}} \left\{ \exp\left[-\mathrm{i}\frac{1}{\sqrt{K}} \sum_{\mu=1}^{N} \sum_{s=0}^{t-1} \left\{ \hat{v}_{\mu}^{(s)} \sum_{k=1}^{K} s_{k}^{\mu} \hat{u}_{k}^{(s)} \right. \\ &+ \hat{w}_{\mu}^{(s)} \sum_{k=1}^{K} s_{k}^{\mu} (R^{(s)} - \sum_{s'=-1}^{s} J^{(s,s')} \tilde{b}_{k}^{(s')}) \right\} \right] \right\} \\ = \prod_{\mu=1}^{N} \prod_{k=1}^{K} \exp\left[-\frac{1}{2} \left(\frac{1}{\sqrt{K}} \sum_{s=0}^{t-1} \left\{ \hat{v}_{\mu}^{(s)} \hat{u}_{k}^{(s)} \right. \\ &+ \hat{w}_{\mu}^{(s)} (R^{(s)} - \sum_{s'=-1}^{s} J^{(s,s')} \tilde{b}_{k}^{(s')}) \right\} \right)^{2} \right]$$
(258)  
$$= \prod_{\mu=1}^{N} \exp\left[-\frac{1}{2} \sum_{s=0}^{t-1} \sum_{s'=0}^{t-1} \left\{ \hat{v}_{\mu}^{(s)} \left(\frac{1}{K} \sum_{k=1}^{K} \hat{u}_{k}^{(s)} \hat{u}_{k}^{(s')}\right) \hat{v}_{\mu}(s') \right. \\ &+ \hat{v}_{\mu}^{(s)} \left( R^{(s')} \left[\frac{1}{K} \sum_{k=1}^{K} \hat{u}_{k}^{(s)}\right] \right] \\ &- \sum_{\tau'=-1}^{s'} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{K} \hat{u}_{k}^{(s')} \hat{u}_{k}^{(s)}\right] \right) \hat{v}_{\mu}(s') \\ &+ \hat{w}_{\mu}^{(s)} \left( R^{(s)} \left[\frac{1}{K} \sum_{k=1}^{K} \hat{u}_{k}^{(s)} \hat{u}_{k}^{(s')}\right] \right) \hat{v}_{\mu}(s') \\ &+ \hat{w}_{\mu}^{(s)} \left( R^{(s)} R^{(s')} - R^{(s)} \sum_{\tau'=-1}^{s'} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{K} \tilde{b}_{k}^{(s')}\right] \right) \\ &- R^{(s')} \sum_{\tau=-1}^{s} J^{(s,\tau)} \left[\frac{1}{K} \sum_{k=1}^{K} \tilde{b}_{k}^{(s')}\right] \\ &+ \sum_{\tau=-1}^{s} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{K} \tilde{b}_{k}^{(s)}\right] \right) \hat{v}_{\mu}(s') \\ &+ \sum_{\tau=-1}^{s} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{K} \tilde{b}_{k}^{(s)} \hat{b}_{k}^{(s')}\right] \hat{v}_{\mu}(s') \\ &+ \sum_{\tau=-1}^{s} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{K} \tilde{b}_{k}^{(s)} \hat{b}_{k}^{(s')}\right] \hat{v}_{\mu}(s') \\ &+ \sum_{\tau=-1}^{s} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{K} \tilde{b}_{k}^{(s)} \hat{b}_{k}^{(s')}\right] \hat{v}_{\mu}(s') \\ &+ \sum_{\tau=-1}^{s} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{K} \tilde{b}_{k}^{(s)} \hat{b}_{k}^{(s')}\right] \hat{v}_{\mu}(s') \\ &+ \sum_{\tau'=-1}^{s'} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{s'} \tilde{b}_{k}^{(s)} \hat{b}_{k}^{(s')}\right] \hat{v}_{\mu}(s') \\ &+ \sum_{\tau'=-1}^{s'} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{s'} \tilde{b}_{k}^{(s)} \hat{b}_{k}^{(s')}\right] \hat{v}_{\mu}(s') \\ &+ \sum_{\tau'=-1}^{s'} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} \left[\frac{1}{K} \sum_{k=1}^{s'} \tilde{b}_{k}^{(s)} \hat{b}_{k}^{(s')} \right] \\ &+ \sum_{\tau'=-1}^{s'} \sum_{\tau'=$$

where  $\underline{d}\boldsymbol{v} := \prod_{\mu=1}^{N} \prod_{s=-1}^{t-1} \frac{dv_{\mu}^{(s)}}{\sqrt{2\pi}}, \underline{d}\hat{\boldsymbol{v}} := \prod_{\mu=1}^{N} \prod_{s=-1}^{t-1} \frac{d\hat{v}_{\mu}^{(s)}}{\sqrt{2\pi}}, \\ \underline{d}\boldsymbol{w} := \prod_{\mu=1}^{N} \prod_{s=-1}^{t-1} \frac{dw_{\mu}^{(s)}}{\sqrt{2\pi}} \text{ and } \underline{d}\hat{\boldsymbol{w}} := \prod_{\mu=1}^{N} \prod_{s=-1}^{t-1} \frac{d\hat{w}_{\mu}^{(s)}}{\sqrt{2\pi}}. \\ \text{We here again use the Fourier integral form of Dirac's delta.} \\ \text{Term } \mathbb{E}_{\boldsymbol{n}} \{\cdots\} \text{ in (256) is given by}$ 

 $\mathbb{E}_{\boldsymbol{n}}\left\{\exp\left[-\mathrm{i}\sqrt{\beta}\sigma_{0}\sum_{\mu=1}^{N}\sum_{s=0}^{t-1}R^{(s)}v_{\mu}^{(s)}n^{\mu}\right]\right\}$ 

 $= \prod_{\mu=1}^{N} \exp\left[-\frac{1}{2}\beta\sigma_0^2 \sum_{s=0}^{t-1} \sum_{s'=0}^{t-1} R^{(s)} R^{(s')} v_{\mu}^{(s)} v_{\mu}^{(s')}\right].$ 

(257)

for finite t. The relevant one-stage and two-stage values, which are equal to those of soft-PIC's analysis, appeared in (259) are separated by introducing (223) – (227). The term of  $\delta$  in (222) is again rewritten by applying the Fourier integral form of Dirac's delta function as well as the case of soft-PIC's analysis.

Since the initial probability, which is given as  $p[\hat{b}^{(-1)}] := \prod_{k=1}^{K} \delta(\tilde{b}_{k}^{(-1)})$ , is factorized, the averaged generating functional  $\bar{Z}[\psi]$  factorizes into single-user contributions. The averaged generating functional is therefore simplified to  $\bar{Z}[\psi] = \int d\eta d\hat{\eta} dk d\hat{k} dq d\hat{q} dQ d\hat{Q} dL d\hat{L} \exp[K(\Phi + \Psi + \Omega) + O(\ln K)]$ , in which functions  $\Phi$ ,  $\Psi$ ,  $\Omega$  are given by

$$\Phi := i \sum_{s=-1}^{t-1} \{ \hat{\eta}^{(s)} \eta^{(s)} + \hat{k}^{(s)} k^{(s)} \} + i \sum_{s=-1}^{t-1} \sum_{s'=-1}^{t-1} \{ \hat{q}^{(s,s')} q^{(s,s')} + \hat{Q}^{(s,s')} Q^{(s,s')} + \hat{L}^{(s,s')} L^{(s,s')} \}, \qquad (260)$$

$$\Psi := \frac{1}{K} \sum_{k=1}^{K} \ln \left\{ \int_{\mathbb{R}^{t+2}} \left( \prod_{s=-1}^{t-1} d\tilde{b}^{(s)} \right) p[\tilde{b}^{(-1)}] \int \underline{d}u \underline{d}\hat{u} \\ \times \exp \left[ \sum_{s=-1}^{t-1} \left\{ \ln \frac{\gamma}{\sqrt{2\pi}} - \frac{\gamma^2}{2} [\tilde{b}^{(s+1)} - f(u^{(s)})]^2 \right\} \\ - \mathrm{i} \sum_{s=-1}^{t-1} \sum_{s'=-1}^{t-1} \left\{ \hat{q}^{(s,s')} \tilde{b}^{(s)} \tilde{b}^{(s')} \\ + \hat{Q}^{(s,s')} \hat{u}^{(s)} \hat{u}^{(s')} + \hat{L}^{(s,s')} \tilde{b}^{(s)} \hat{u}^{(s')} \right\} \\ + \mathrm{i} \sum_{s=-1}^{t-1} \hat{u}^{(s)} \left\{ u^{(s)} - A^{(s)} \tilde{b}^{(s)} - \theta_k^{(s)} - \hat{k}^{(s)} \right\} \\ - \mathrm{i} \sum_{s=-1}^{t-1} \tilde{b}^{(s)} \hat{\eta}^{(s)} - \mathrm{i} \sum_{s=-1}^{t} \tilde{b}^{(s)} \psi_k^{(s)} \right],$$
(261)

the single-user measure

$$w_{k}(\{\tilde{b}, u, \hat{u}\})$$

$$:= \delta[\tilde{b}^{(-1)}] \exp\left[\sum_{s=-1}^{t-1} \{\ln \frac{\gamma}{\sqrt{2\pi}} - \frac{\gamma^{2}}{2} [\tilde{b}^{(s+1)} - f(u^{(s)})]^{2} \}\right]$$

$$-i \sum_{s=0}^{t-1} \sum_{s'=0}^{t-1} \{\hat{q}^{(s,s')} \tilde{b}^{(s)} \tilde{b}^{(s')} + \hat{Q}^{(s,s')} \hat{u}^{(s)} \hat{u}^{(s')}$$

$$+ \hat{L}^{(s,s')} \tilde{b}^{(s)} \hat{u}^{(s')} \}$$

$$+ i \sum_{s=0}^{t-1} \hat{u}^{(s)} \{u^{(s)} - A^{(s)} \tilde{b}^{(s)} - \theta_{k}^{(s)} - \hat{k}^{(s)} \}$$

$$- i \sum_{s=0}^{t-1} \tilde{b}^{(s)} \hat{\eta}^{(s)} \Big]_{\text{saddle}}, \qquad (263)$$

The integral in the averaged generating functional  $\overline{Z}[\psi]$  will

instead of the single-use measure (40) for soft-PIC.

be evaluated by the dominating saddle-point of the exponent  $\Phi + \Psi + \Omega \text{ in the large system limit } K \to \infty. \text{ Using the identities } (35) - (35) \text{ and } (44) - (46) \text{ with } (263), \text{ the differentiation of } \Phi + \Psi + \Omega \text{ with respect to } \eta^{(s)}, \hat{\eta}^{(s)}, k^{(s)}, \hat{k}^{(s)}, \hat{q}^{(s,s')}, \hat{q}^{(s,s')}, \hat{Q}^{(s,s')}, \hat{Q}^{(s,s')}, L^{(s,s')}, \text{ and } \hat{L}^{(s,s')} \text{ leads us to the following saddle-point equations } (50) - (59) \text{ with } (263), \text{ respectively. Comparing these saddle-point equations with } (47) - (49), we again find the following relationships: <math>\eta^{(s)} = m^{(s)}, q^{(s,s')} = C^{(s,s')}, \text{ and } L^{(s,s')} = iG^{(s,s')}.$ The integral in  $\Omega$  with respect to  $\hat{v}$  and  $\hat{w}$  is given by  $\Omega = -\frac{1}{2\beta} \{\ln |\hat{D}| + \ln |\mathbf{Q} + (\beta^{-1}\mathbf{1} - \hat{B})^{\top}\hat{D}^{-1}(\beta^{-1}\mathbf{1} - \hat{B})|\},$  where matrices  $\hat{B}$  and  $\hat{D}$  whose (s, s') elements are given by  $\Sigma^{(s',s)} \hat{w}_{\mu}^{(s')}$ 

$$\hat{B}^{(s,s')} := -iR^{(s)}k^{(s)} - \sum_{\tau=-1} J^{(s,\tau)}G^{(\tau,s)},$$
(264)
$$\hat{D}^{(s,s')} := \frac{\sigma_0^2}{\beta}R^{(s)}R^{(s')} + R^{(s)}R^{(s')} - R^{(s)}\sum_{\tau'=-1}^{s'} J^{(s',\tau')}m^{(\tau')} - R^{(s')}\sum_{\tau'=-1}^{s} J^{(s,\tau)}m^{(\tau)} + \sum_{\tau=-1}^{s}\sum_{\tau'=-1}^{s'} J^{(s,\tau)}J^{(s',\tau')}C^{(\tau,\tau')},$$
(265)

respectively.

The saddle-point equations including  $\Omega$  are evaluated as follows. One finds  $\hat{\eta}^{(s)} = 0$  and  $\hat{q}^{(s)} = 0$ . We put

$$\hat{U} := \beta^{-1} \mathbf{1} + \hat{B},$$
 (266)

Applying the same way as the derivation of soft-PIC, we again obtain (44) – (46), with the average (40)  $\langle \rangle_k$  which has

differentiation of the averaged generating functional  $\overline{Z}[\psi]$  with respect to  $\theta_k^{(s)}$  and  $\psi_k^{(s)}$ . The averaged generating functional  $\overline{Z}[\psi]$  is dominated by a saddle-point for  $K \to \infty$ .

One can deduce the meaning of macroscopic parameters by

for light notations. Noting that  $\hat{U}$  includes G via  $\hat{B}$  and using

$$\begin{split} \Omega &:= \frac{1}{K} \ln \int \underline{d} v \underline{d} \hat{v} \underline{d} w \underline{d} \hat{w} \\ &\times \exp \left[ i \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \{ \hat{v}_{\mu}^{(s)} v_{\mu}^{(s)} + \hat{w}_{\mu}^{(s)} w_{\mu}^{(s)} - \beta v_{\mu}^{(s)} w_{\mu}^{(s)} \} \right. \\ &- \frac{1}{2} \sum_{\mu=1}^{N} \sum_{s=-1}^{t-1} \sum_{s'=-1}^{t-1} \left\{ \beta \sigma_{0}^{2} R^{(s)} R^{(s')} v_{\mu}^{(s)} v_{\mu}^{(s')} \\ &+ \hat{v}_{\mu}^{(s)} Q^{(s,s')} \hat{v}_{\mu}^{(s')} \right. \\ &+ \hat{v}_{\mu}^{(s)} \left( R^{(s')} k^{(s)} - \sum_{\tau'=-1}^{s'} J^{(s',\tau')} L^{(s',s)} \right) \hat{w}_{\mu}^{(s')} \\ &+ \hat{w}_{\mu}^{(s)} \left( R^{(s)} k^{(s')} - \sum_{\tau=-1}^{s} J^{(s,\tau)} L^{(s,s')} \right) \hat{v}_{\mu}^{(s')} \\ &+ \hat{w}_{\mu}^{(s)} \left( R^{(s)} R^{(s')} \\ &- R^{(s)} \sum_{\tau'=-1}^{s'} J^{(s',\tau')} \eta^{(\tau')} - R^{(s')} \sum_{\tau=-1}^{s} J^{(s,\tau)} \eta^{(\tau)} \\ &+ \sum_{\tau=-1}^{s} \sum_{\tau'=-1}^{s'} J^{(s,\tau)} J^{(s',\tau')} q(\tau,\tau') \right) \hat{w}_{\mu}^{(s')} \bigg\} \bigg]. \tag{262}$$

the cofactor expansion of  $|\hat{U}|$ , we find

$$\hat{L}^{(s,s')} = \frac{\partial\Omega}{\partial G^{(s,s')}}$$

$$= -\frac{1}{\beta} \frac{\partial}{\partial G^{(s,s')}} \ln |U|$$

$$= -\frac{1}{\beta |U|} \frac{\partial}{\partial G^{(s,s')}} \sum_{s=-1}^{t-1} U^{(s,s')} \tilde{U}^{(s,s')}$$

$$= -\frac{1}{\beta |U|} \sum_{s=-1}^{t-1} J^{(s,s')} \tilde{U}^{(s,s')}, \qquad (267)$$

where  $\hat{U}^{(s,s')}$  denotes a cofactor of the (s,s') element of the matrix  $\hat{U}$ , which does not include  $G^{(s,s')}$ . The  $(t+1) \times (t+1)$  matrix  $\hat{L} = (\hat{L}^{(s,s')})$  then becomes

$$\begin{split} \hat{\boldsymbol{L}} &= -\frac{1}{\beta|\hat{\boldsymbol{U}}|} \begin{bmatrix} & J^{(-1,-1)}\tilde{\boldsymbol{U}}^{(-1,-1)} & \cdots & J^{(-1,-1)}\tilde{\boldsymbol{U}}^{(-1,t-1)} \\ J^{(0,0)}\tilde{\boldsymbol{U}}^{(0,-1)} & \cdots & J^{(0,0)}\tilde{\boldsymbol{U}}^{(0,t-1)} \\ \vdots & \vdots & \vdots \\ J^{(t-1,t-1)}\tilde{\boldsymbol{U}}^{(t-1,-1)} & \cdots & J^{(t-1,t-1)}\tilde{\boldsymbol{U}}^{(t-1,t-1)} \end{bmatrix} \\ &+ \cdots + & \begin{pmatrix} J^{(t-1,t-1)}\tilde{\boldsymbol{U}}^{(t-1,-1)} & \cdots & J^{(t-1,t-1)}\tilde{\boldsymbol{U}}^{(t-1,t-1)} \\ 0 & \cdots & 0 \\ \vdots & & \vdots \\ 0 & \cdots & 0 \end{pmatrix} \end{bmatrix} \\ &= -\frac{1}{\beta|\hat{\boldsymbol{U}}|} \begin{bmatrix} \begin{pmatrix} J^{(-1,-1)} & \cdots & J^{(-1,-1)} \\ J^{(0,0)} & \cdots & J^{(0,0)} \\ \vdots & & \vdots \\ J^{(t-1,t-1)} & \cdots & J^{(t-1,t-1)} \end{pmatrix} \\ &+ \cdots \\ &+ \begin{pmatrix} J^{(t-1,t-1)} & \cdots & J^{(t-1,t-1)} \\ 0 & \cdots & 0 \\ \vdots & & \vdots \\ 0 & \cdots & 0 \end{pmatrix} \otimes \boldsymbol{\Delta}^{t} \mathrm{adj} \hat{\boldsymbol{U}}^{\top} \end{bmatrix} \\ &= -J_{0} \otimes (\boldsymbol{U}^{\top})^{-1} - J_{1} \otimes \boldsymbol{\Delta} (\boldsymbol{U}^{\top})^{-1} \\ &- \cdots - J_{t} \otimes \boldsymbol{\Delta}^{t} (\boldsymbol{U}^{\top})^{-1}, \end{split}$$
(268)

where  $U := \beta \hat{U}$  and  $\operatorname{adj} A$  denotes an adjoint matrix of A. The definition of  $J_s$  is given by (114).

Due to  $\boldsymbol{Q} = \boldsymbol{O}$ ,  $\Omega$  can be expanded as  $\ln |\boldsymbol{A} + \boldsymbol{Q}| = \operatorname{Tr} \ln \boldsymbol{A} + \operatorname{tr} \boldsymbol{A}^{-1} \boldsymbol{Q}$ . The  $\hat{Q}^{(s,s')}$  is obtained as  $\hat{Q}^{(s,s')} = -\frac{\mathrm{i}}{2\beta} M^{(s',s)}$ , and we then have  $\hat{\boldsymbol{Q}} = -\mathrm{i}\frac{1}{2\beta} M^{\top} = -\mathrm{i}\frac{1}{2}(\mathbf{1} + \beta G)^{-1}D(\mathbf{1} + \beta G^{\top})^{-1}$ , where  $\boldsymbol{M} := (\beta^{-1}\mathbf{1} + \boldsymbol{G})^{-1}\hat{\boldsymbol{D}}(\beta^{-1}\mathbf{1} + \boldsymbol{G}^{\top})^{-1}$  and  $D^{(s,s')} := \beta \hat{\boldsymbol{D}}$ . Noting that  $|\mathbf{1} - \beta \hat{\boldsymbol{B}}^{\top}|$  only contains  $k^{(s)}$  in a single row,  $k^{(s)}$  is obtained as  $\hat{k}^{(s)} = |\Lambda_{[s]}|$ , where  $\boldsymbol{B} := \hat{\boldsymbol{B}}|_{\boldsymbol{k}=\mathbf{0}} = -\sum_{\tau=-1}^{s} J^{(s,\tau)}G^{(\tau,s)}$  and the definition of  $\Lambda_{[s]}$  is given by (113).

Some macroscopic parameters are found to vanish in the saddle-point:  $k^{(s)} = Q^{(s,s')} = 0$ . In the case of the BP-base detector, the remaining macroscopic parameters can all be expressed in terms of three observables, i.e., the average

overlap  $m^{(s)}$ , the average single-user correlation  $C^{(s,s')}$ , and the average single-user response function  $G^{(s,s')}$ , which are defined by (47) – (49). The averaged generating functional  $\overline{Z}[\psi]$  is dominated by a saddle-point in the large system limit. We then arrive at Proposition 4.

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