STATE-CONSTRAINT STATIC HAMILTON-JACOBI EQUATIONS IN NESTED DOMAINS

YEONEUNG KIM, HUNG V. TRAN, AND SON N. T. TU

ABSTRACT. We study state-constraint static Hamilton–Jacobi equations in a sequence of domains $\{\Omega_k\}_{k\in\mathbb{N}}$ in \mathbb{R}^n such that $\Omega_k \subset \Omega_{k+1}$ for all $k \in \mathbb{N}$. We obtain rates of convergence of \mathfrak{u}_k , the solution to the state-constraint problem in Ω_k , to \mathfrak{u} , the solution to the corresponding problem in $\Omega = \bigcup_{k\in\mathbb{N}} \Omega_k$. In many cases, the rates obtained are proven to be optimal. Various new examples and discussions are provided at the end of the paper.

1. INTRODUCTION

Let $\{\Omega_k\}_{k\in\mathbb{N}}$ be a sequence of domains in \mathbb{R}^n such that $\Omega_k \subset \Omega_{k+1}$ for all $k \in \mathbb{N}$. We say that $\{\Omega_k\}_{k\in\mathbb{N}}$ is a sequence of nested domains. Then, $\Omega = \bigcup_{k\in\mathbb{N}} \Omega_k$ is also a domain in \mathbb{R}^n . Let $H : \overline{\Omega} \times \mathbb{R}^n \to \mathbb{R}$ be a given continuous Hamiltonian. In this paper, we are interested in studying state-constraint solutions to the following static Hamilton-Jacobi equations:

$$u_k(x) + H(x, Du_k(x)) = 0 \qquad \text{in } \Omega_k, \tag{HJ}_k$$

and

$$u(x) + H(x, Du(x)) = 0$$
 in Ω . (HJ)

The precise definition of state-constraint viscosity solutions is given in Section 2. Under some appropriate conditions, (HJ_k) has a unique state-constraint viscosity solution $\mathbf{u}_k \in C(\overline{\Omega}_k)$ for each $k \in \mathbb{N}$, and (HJ) has a unique state-constraint viscosity solution $\mathbf{u} \in C(\overline{\Omega})$. Furthermore, by a priori estimates and the stability results of viscosity solutions, we have that $\mathbf{u}_k \to \mathbf{u}$ locally uniformly on $\overline{\Omega}$. Our main focus here is to study how fast this convergence is in two different types of nested domains.

1.1. Assumptions. In the paper, we consider the following two prototypes of nested domains, which are

- (P1) $\Omega_k = B(0, k)$ and $\Omega = \bigcup_{k \in \mathbb{N}} B(0, k) = \mathbb{R}^n$,
- (P2) $\Omega_k = B(0, 1 \frac{1}{k})$, and $\Omega = B(0, 1)$.

We list the main assumptions that will be used throughout the paper.

(H1) There exists $C_1 > 0$ such that

$$H(x,p) \ge -C_1$$
 for all $(x,p) \in \Omega \times \mathbb{R}^n$. (H1)

2010 Mathematics Subject Classification. 35B40, 35D40, 49J20, 49L25, 70H20.

Key words and phrases. first-order Hamilton–Jacobi equations; state-constraint problems; optimal control theory; rate of convergence; viscosity solutions.

The authors were supported in part by NSF grant DMS-1664424. The second and the third authors are supported in part by NSF CAREER grant DMS-1843320.

(H2) There exists $C_2 > 0$ such that

$$\sup_{\mathbf{x}\in\overline{\Omega}}|\mathsf{H}(\mathbf{x},\mathbf{0})|\leqslant C_2.\tag{H2}$$

(H3a) There exists a modulus $\omega_{\rm H}$: $[0,\infty) \to [0,\infty)$, which is a nondecreasing function satisfying $\omega_{\rm H}(0^+) = 0$ and

$$\begin{cases} |\mathsf{H}(x,p) - \mathsf{H}(y,p)| \leq \omega_{\mathsf{H}}(|x - y|(1 + |p|)), \\ |\mathsf{H}(x,p) - \mathsf{H}(x,q)| \leq \omega_{\mathsf{H}}(|p - q|), \end{cases}$$
(H3a)

for $\mathbf{x}, \mathbf{y} \in \overline{\Omega}$ and $\mathbf{p}, \mathbf{q} \in \mathbb{R}^{n}$.

(H3b) For every R > 0, there exists a modulus $\omega_R : [0, +\infty) \to [0, +\infty)$, which is nondecreasing with $\omega_R(0^+) = 0$ and

$$\begin{cases} |\mathsf{H}(\mathbf{x},\mathbf{p}) - \mathsf{H}(\mathbf{y},\mathbf{p})| \leq \omega_{\mathsf{R}}(|\mathbf{x} - \mathbf{y}|), \\ |\mathsf{H}(\mathbf{x},\mathbf{p}) - \mathsf{H}(\mathbf{x},\mathbf{q})| \leq \omega_{\mathsf{R}}(|\mathbf{p} - \mathbf{q}|), \end{cases}$$
(H3b)

 $\mathrm{for}\; x,y\in\overline{\Omega}\;\mathrm{and}\; p,q\in\mathbb{R}^n\;\mathrm{with}\; |p|,|q|\leqslant R.$

(H3c) For each R>0 there exists a constant C_R such that

$$\begin{cases} |\mathsf{H}(\mathbf{x},\mathbf{p}) - \mathsf{H}(\mathbf{y},\mathbf{p})| \leqslant C_{\mathsf{R}}|\mathbf{x} - \mathbf{y}|, \\ |\mathsf{H}(\mathbf{x},\mathbf{p}) - \mathsf{H}(\mathbf{x},\mathbf{q})| \leqslant C_{\mathsf{R}}|\mathbf{p} - \mathbf{q}|, \end{cases}$$
(H3c)

for $x, y \in \overline{\Omega}$ and $p, q \in \mathbb{R}^n$ with $|p|, |q| \leq R$.

(H4) H satisfies the coercivity assumption

$$\lim_{|\mathfrak{p}|\to\infty} \left(\inf_{\mathbf{x}\in\overline{\Omega}} \mathsf{H}(\mathbf{x},\mathfrak{p})\right) = +\infty. \tag{H4}$$

(H5) $\mathbf{p} \mapsto \mathbf{H}(\mathbf{x}, \mathbf{p})$ is convex for each $\mathbf{x} \in \Omega$.

Let us give some quick comments on the assumptions here. Assumption (H1) is necessary to ask a meaningful question about the rate of convergence of \mathbf{u}_k to \mathbf{u} . See the discussion in Section 7 in case where (H1) fails to hold. Besides, it is clear that (H3b) is weaker than both (H3a) and (H3c).

1.2. Main results. There have been many works in the literature on the well-posedness of state-constraint Hamilton-Jacobi equations and fully nonlinear elliptic equations. The state-constraint problem for first-order convex Hamilton-Jacobi equations using optimal control frameworks was first studied in [27, 28]. The general nonconvex, coercive first-order equations was then discussed in [10]. For further developments in using optimal control formulation and obtaining optimal paths, we refer the readers to [17, 19, 12, 11, 6, 1, 14, 26, 25] for the finite dimensional cases, and [8, 20] for the infinite dimensional cases. See [7] for discrete numerical schemes, and [24] for large time behavior results. We also refer to the classical books [5, 4] and the references therein.

The state-constraint problem for second-order equations was first studied in [21] for the Laplacian, and in [3] for the general possibly degenerate diffusion matrices. Boundary behavior of blow-up solutions was discussed in [21, 23, 3]. Convex solutions with state-constraint boundary were constructed in [2, 13]. The convergence of solutions to the vanishing discount problems was proved in [18].

In terms of state-constraint problems in nested domains, up to our knowledge, there are only qualitative results in the literature in [10, 3] where certain approximations were needed for the analysis of solutions. We provide here some first quantitative results on the rate of convergence of the solutions to (HJ_k) as k goes to infinity in two different types ((P1) or (P2)) of nested domains.

First of all, we show that the rate of convergence is $O\left(\frac{1}{k^2}\right)$ for the prototype (P1) for general nonconvex Hamiltonians.

Theorem 1.1. Under the assumptions (P1), (H1), (H2), (H3c), and (H4), we have

- (i) $u(x) \leq u_k(x)$ for every $x \in \overline{B(0,k)}$,
- (ii) there exists a constant C > 0 depending only on H such that

$$0 \leq u_k(x) - u(x) \leq \frac{C(1+|x|^2)}{k^2}$$

for $k \in \mathbb{N}$ and $x \in \overline{B(0,k)}$.

In particular, for any fixed R > 0 and $|x| \leq R$,

$$0 \leq u_k(x) - u(x) \leq \frac{C(1+R^2)}{k^2}.$$

The condition that $|\mathbf{x}| \leq \mathbf{R}$ is important since there are examples where the estimate above fails at the boundary of Ω_k . In Proposition 5.10, we have, for each $\mathbf{k} \in \mathbb{N}$, $|\mathbf{u}_k(\mathbf{x}) - \mathbf{u}(\mathbf{x})| = 1$ for some $\mathbf{x} \in \partial \Omega_k$.

Theorem 1.2. Assume (P1). Assume further that H(x, p) = a(x)K(p) for $(x, p) \in \mathbb{R}^n \times \mathbb{R}^n$. Here, K(0) = 0, K is locally Lipschitz and coercive in \mathbb{R}^n , and $a \in BUC(\mathbb{R}^n)$ satisfies $\alpha \leq a(\cdot) \leq \beta$ for some given $\alpha, \beta > 0$. Then, $u \equiv 0$, and for every $x \in \overline{B(0, k)}$, we have

$$0 \leq u_k(x) \leq \left(Ce^{\frac{|x|}{C}}\right)e^{-\frac{k}{C}},$$

where C is a constant depending only on H. In particular, for any fixed R > 0, we have

$$0 \leqslant u_k(x) \leqslant \left(Ce^{\frac{R}{C}}\right)e^{-\frac{k}{C}}$$

for every $x \in \overline{B(0, R)}$ and $k \ge R$. In addition to that, this exponential rate is optimal.

It is quite interesting to observe that we obtain the exponential rate of convergence for this particular class of nonconvex Hamiltonians and the rate is indeed optimal. When a(x) is a positive constant, the assumption K(0) = 0 in the theorem above can be removed.

Corollary 1.3. Assume (P1). Assume further that H(x, p) = H(p) for $(x, p) \in \mathbb{R}^n \times \mathbb{R}^n$. Here, H is locally Lipschitz and coercive in \mathbb{R}^n . Then, $u \equiv -H(0)$, and for every $x \in \overline{B(0, k)}$, we have

$$0 \leq u_k(x) - u(x) \leq \left(Ce^{\frac{|x|}{C}}\right)e^{-\frac{k}{C}},$$

where C is a constant depending only on H. In particular, for any fixed R > 0, we have

$$0 \leq u_k(x) - u(x) \leq \left(Ce^{\frac{R}{C}}\right)e^{-\frac{k}{C}}$$

for every $x \in \overline{B(0, R)}$ and $k \ge R$. In addition to that, this rate is optimal.

When H(x,p) = K(p) + V(x), the analysis becomes much more complicated due to the interaction between K and V. We provide an example where the exponential rate of convergence is obtained in Example 2.

For convex Hamiltonians, we are able to establish the exponential rate of convergence using optimal control theory. Some examples for which the exponential rate is obtained are given in Proposition 5.10 and Proposition 5.11.

Theorem 1.4. Under the assumptions (P1), (H1), (H2), (H3b), (H4), and (H5), we have

- (i) $\mathfrak{u}(\mathbf{x}) \leq \mathfrak{u}_k(\mathbf{x})$ for every $\mathbf{x} \in \overline{\mathsf{B}(\mathbf{0},\mathbf{k})}$,
- (ii) for each fixed $x \in B(0, k)$ we have

$$\mathfrak{u}_{k}(\mathbf{x}) \leq \mathfrak{u}(\mathbf{x}) + \left(Ce^{\frac{|\mathbf{x}|}{C}}\right)e^{-\frac{k}{C}},\tag{1.1}$$

where C is a constant depending only on the growth of H. In particular, for any fixed R > 0, we have

$$0 \leq u_k(x) - u(x) \leq \left(Ce^{\frac{R}{C}}\right)e^{-\frac{k}{C}}$$

for all $x \in B(0, R)$ and k > R.

As a byproduct, we prove the existence of a minimizer η with bounded velocity to the minimizing problem (5.2) for each given $\mathbf{x} \in \mathbb{R}^n$, which is a key element in the proof of Theorem 1.4. Moreover, the bound on the velocity of η only depends on the growth of H and not on its smoothness. We believe that this bound (Theorem 5.7 and Lemma 5.9) is new in the literature. See Remark 6 for further discussions.

For the second prototype (P2), we establish the rate $\mathcal{O}\left(\frac{1}{k}\right)$ for a quite general class of Hamiltonians. The rate is also optimal, as pointed out in Remark 9.

Theorem 1.5. Under assumptions (P2), (H1), (H2), (H3c) and (H4), for any $k \ge 2$,

$$0 \leq u_k(x) - u(x) \leq \frac{C}{k}$$

for every $x \in \overline{B\left(0, 1-\frac{1}{k}\right)}$ where C is a constant depending only on H. Moreover, this rate is optimal.

Although we only deal with two prototype cases (P1) and (P2) in this paper, the obtained results can be extended to more general domains in a similar fashion under some appropriate conditions. See Remarks 8 and 10 for example.

1.3. Organization of the paper. The paper is organized in the following way. We first provide some results on state-constraint Hamilton-Jacobi equations needed throughout the paper in Section 2. Section 3 and 4 are devoted to proving Theorem 1.1 and Theorem 1.2, respectively. In the following section, we deal with the rate of convergence for convex Hamiltonians (Theorem 1.5). In Section 6, the second prototype case is considered. We provide some examples and further discussion in Section 7. The proofs for some results concerning minimizers of the corresponding optimal control problem are provided in Appendix.

Acknowledgement. We would like to thank the two referees very much for carefully reading our manuscript and giving very helpful comments to improve the presentation of the paper.

2. Preliminaries

For an open subset $\Omega \subset \mathbb{R}^n$, we denote the space of bounded uniformly continuous functions defined in Ω by BUC(Ω ; \mathbb{R}).

Definition 1. We say

- (i) $\mathbf{v} \in \text{BUC}(\Omega; \mathbb{R})$ is a viscosity subsolution of (HJ) in Ω if for every $\mathbf{x} \in \Omega$ and $\varphi \in C^1(\Omega)$ such that $\mathbf{v} \varphi$ has a local maximum over Ω at \mathbf{x} then $\mathbf{v}(\mathbf{x}) + H(\mathbf{x}, \mathsf{D}\varphi(\mathbf{x})) \leq 0$.
- (ii) $\nu \in BUC(\overline{\Omega}; \mathbb{R})$ is a viscosity supersolution of (HJ) on $\overline{\Omega}$ if for every $\mathbf{x} \in \overline{\Omega}$ and $\varphi \in C^1(\overline{\Omega})$ such that $\nu \varphi$ has a local minimum over $\overline{\Omega}$ at \mathbf{x} then $\nu(\mathbf{x}) + H(\mathbf{x}, D\varphi(\mathbf{x})) \ge 0$.

If \mathbf{v} is a viscosity subsolution to (HJ) in Ω , and is a viscosity supersolution to (HJ) on Ω , that is,

$$\begin{cases} \nu(\mathbf{x}) + \mathbf{H}(\mathbf{x}, \mathsf{D}\nu(\mathbf{x})) &\leq 0 & \text{ in } \Omega, \\ \nu(\mathbf{x}) + \mathbf{H}(\mathbf{x}, \mathsf{D}\nu(\mathbf{x})) &\geq 0 & \text{ on } \overline{\Omega} \end{cases}$$
(2.1)

in the viscosity sense, then we say that ν is a state-constraint viscosity solution of (HJ).

Remark 1. As pointed out in [27], the state-constraint implicitly imposes a boundary condition to solutions. Indeed, when $\partial\Omega$ is smooth, we can define an outward normal vector $\vec{v}(x)$ at $x \in \partial\Omega$. Moreover, if the state-constraint solution $\nu \in C^1(\overline{\Omega})$, then ν solves $\nu(x) + H(x, D\nu(x)) = 0$ in Ω and satisfies

$$H(x, Dv(x)) \leq H(x, Dv(x) + \beta \vec{v}(x))$$
 for any $\beta \ge 0, x \in \partial \Omega$.

If H is differentiable in p, the above condition can also be phrased as a constraint on the normal derivative on the boundary as

$$\mathsf{D}_{\mathfrak{p}}\mathsf{H}(\mathfrak{x},\mathsf{D}\mathfrak{v}(\mathfrak{x}))\cdot\vec{\mathfrak{v}}(\mathfrak{x}) \ge 0 \qquad \text{for any } \mathfrak{x}\in\partial\Omega. \tag{2.2}$$

Now we construct a state-constraint viscosity solution to (HJ) based on Perron's method. It is a variant of the classical result in [15] but we include the proof here for the sake of the readers' convenience. Note that the Lipschitz regularity of subsolutions is encoded directly into the admissible class \mathcal{F} .

Definition 2. For a real valued function w(x) define for $x \in \Omega$, we define the superdifferential and sub-differential of w at x as

$$D^+w(x) = \left\{ p \in \mathbb{R}^n : \limsup_{y \to x} \frac{w(y) - w(x) - p \cdot (y - x)}{|y - x|} \leq 0 \right\},$$
$$D^-w(x) = \left\{ p \in \mathbb{R}^n : \liminf_{y \to x} \frac{w(y) - w(x) - p \cdot (y - x)}{|y - x|} \geq 0 \right\}.$$

Theorem 2.1. Assume (H1), (H2) and (H4). Then, there exists a state-constrained viscosity solution $\mathbf{u} \in C(\overline{\Omega}) \cap W^{1,\infty}(\Omega)$ to (2.1).

Proof of Theorem 2.1. Under (H1) and (H2), C_1 and $-C_2$ are a supersolution on $\overline{\Omega}$ and a subsolution in Ω of (HJ), respectively. By the coercivity assumption (H4), we can find a constant $C_3 > 0$ such that

 $\mathsf{H}(x,p) \leqslant \max\{C_1,C_2\} \ \ \mathrm{for \ some} \ x \in \overline{\Omega} \qquad \Longrightarrow \qquad |p| \leqslant C_3.$

Let us define

$$\mathcal{F} = \left\{ w \in \mathrm{C}(\overline{\Omega}) \cap \mathrm{W}^{1,\infty}(\Omega) : -C_2 \leqslant w(x) \leqslant C_1, \|\mathrm{D}w\|_{\mathrm{L}^{\infty}(\overline{\Omega})} \leqslant C_3, \right.$$

and w is a viscosity subsolution to $w(x) + H(x, Dw(x)) \leq 0$ in Ω

and for each $x \in \overline{\Omega}$, we define

$$\mathfrak{u}(\mathbf{x}) := \sup \left\{ w(\mathbf{x}) : w \in \mathcal{F} \right\}.$$

By the stability of viscosity subsolutions, we have that u is a viscosity subsolution to (HJ) in Ω . Thus, $u \in \mathcal{F}$ as well.

We now check that \mathfrak{u} is a viscosity supersolution to (HJ) on $\overline{\Omega}$. Assume that \mathfrak{u} is not a supersolution on $\overline{\Omega}$. Then, there exists $\mathfrak{x}_0 \in \overline{\Omega}$, $\varphi \in C^1(\overline{\Omega})$ with $\|D\varphi\|_{L^{\infty}(B(\mathfrak{x}_0,\mathfrak{r}))} \leq C_3$ and $\mathfrak{r} > 0$ such that $\mathfrak{u}(\mathfrak{x}_0) = \varphi(\mathfrak{x}_0)$ and $(\mathfrak{u} - \varphi)(\mathfrak{x}) \geq |\mathfrak{x} - \mathfrak{x}_0|^2$ for all $\mathfrak{x} \in B(\mathfrak{x}_0, \mathfrak{r}) \cap \overline{\Omega}$, and

$$\varphi(\mathbf{x}_0) + H(\mathbf{x}_0, D\varphi(\mathbf{x}_0)) < 0.$$
 (2.3)

From (H1) and (2.3), we obtain $\varphi(x_0) = u(x_0) < C_1$. By continuity of φ and H, one can choose $\delta, \varepsilon \in (0, \frac{r}{2})$ small enough so that $\varepsilon < \delta^2$ and

$$\begin{cases} \phi(x) + \delta^2 < C_1, \\ \phi(x) + \delta^2 + H(x, D\phi(x)) < 0 \end{cases} \implies \begin{cases} \phi(x) + \epsilon^2 < C_1, \\ \phi(x) + \epsilon^2 + H(x, D\phi(x)) < 0 \end{cases}$$

for all $x \in B(x_0, 2\epsilon) \cap \overline{\Omega}$. Clearly, $x \mapsto \phi(x) + \epsilon^2$ is a viscosity subsolution to (HJ) in $B(x_0, 2\epsilon) \cap \overline{\Omega}$ and $u(x) \ge \phi(x) + \epsilon^2$ for $x \in B(x, 2\epsilon) \setminus B(x_0, \epsilon)$. Let us define $w : \overline{\Omega} \to \mathbb{R}$ by

$$w(\mathbf{x}) = \begin{cases} \max \left\{ u(\mathbf{x}), \varphi(\mathbf{x}) + \varepsilon^2 \right\} & \mathbf{x} \in B(\mathbf{x}_0, \varepsilon) \cap \overline{\Omega}, \\ u(\mathbf{x}) & \mathbf{x} \in \overline{\Omega} \setminus B(\mathbf{x}_0, \varepsilon). \end{cases}$$

Then, $w(x) = \max \{ u(x), \phi(x) + \varepsilon^2 \}$ in $B(x_0, 2\varepsilon) \cap \overline{\Omega}$ belongs to \mathcal{F} . Therefore, w(x) is a viscosity subsolution to (HJ). However, $w(x_0) = \phi(x_0) + \varepsilon^2 = u(x_0) + \varepsilon^2 > u(x_0)$, which is a contradiction to the definition of u.

The argument used in the proof of Perron's method implies the following corollary as well (see also [10]).

Corollary 2.2. Let $\mathbf{u} \in C(\overline{\Omega})$ be a viscosity subsolution to (HJ) in Ω . Assume further that $\mathbf{v} \leq \mathbf{u}$ on $\overline{\Omega}$ for all viscosity subsolutions $\mathbf{v} \in C(\overline{\Omega})$ of (HJ) in Ω . Then, \mathbf{u} is a viscosity supersolution to (HJ) on $\overline{\Omega}$.

The uniqueness of (2.1) follows from the comparison principle. It was first studied by M. Soner in [27] under the following assumption on $\partial\Omega$:

(A) There exists a universal pair $(\mathbf{r}, \mathbf{h}) \in (0, \infty) \times (0, \infty)$ and a uniformly bounded continuous function $\eta \in \text{BUC}(\overline{\Omega}; \mathbb{R}^n)$ such that

$$B(x + t\eta(x), rt) \subset \Omega \quad \text{for all } x \in \partial\Omega, \ t \in (0, h].$$
 (A)

See also [10] for other conditions to establish the comparison principle.

Theorem 2.3. Assume (A). If $v_1 \in BUC(\Omega; \mathbb{R})$ is a viscosity subsolution of (HJ) in Ω , and $v_2 \in BUC(\overline{\Omega}; \mathbb{R})$ is a viscosity supersolution of (HJ) on $\overline{\Omega}$. If either

• (H3a) holds, or

• (H3b) holds and v_2 is Lipschitz,

then $v_1(x) \leq v_2(x)$ for all $x \in \overline{\Omega}$.

When the uniqueness of (2.1) is guaranteed, the unique viscosity solution to (2.1) is the maximal viscosity subsolution of (HJ). This property will play a crucial role in dealing with the second prototype (P2).

3. A RATE OF CONVERGENCE FOR GENERAL HAMILTONIANS IN UNBOUNDED DOMAIN

In this section, we consider the first prototype (P1). The assumptions (H1), (H2), (H3c) and (H4) are enforced throughout the section. By Theorems 2.1 and 2.3, there exists $u_k \in \text{Lip}(\overline{B(0, k)})$ which is the unique solution to

$$\begin{cases} u_k(x) + H(x, Du_k(x)) \leq 0 & \text{ in } B(0, k), \\ u_k(x) + H(x, Du_k(x)) \geq 0 & \text{ on } \overline{B(0, k)} \end{cases}$$
(3.1)

in the viscosity sense. Based on the construction of solutions via Perron's method together with the coercivity of H, we have the following a priori estimate:

$$|\mathfrak{u}_k(x)| + |\mathsf{D}\mathfrak{u}_k(x)| \leqslant C_H$$

for all $x \in B(0, k)$ in the viscosity sense. Here, C_H is a positive constant depending only on H (one can take $C_H = \max\{C_1, C_2, C_3\}$ from Theorem 2.1). By the Arzelà–Ascoli theorem, there is a subsequence $\{k_m\} \to \infty$, and a function $u \in \operatorname{Lip}(\mathbb{R}^n)$ such that

$$\mathfrak{u}_{k_{\mathfrak{m}}} \to \mathfrak{u}$$
 locally uniformly in $\mathbb{R}^{\mathfrak{n}}$. (3.2)

Theorem 3.1. The function u defined in (3.2) is a viscosity solution to

$$u(\mathbf{x}) + H(\mathbf{x}, \mathsf{D}u(\mathbf{x})) = 0 \qquad in \,\mathbb{R}^n. \tag{3.3}$$

Moreover, $u_k \rightarrow u$ locally uniformly in \mathbb{R}^n as k grows to infinity.

Proof. It is clear from the stability of viscosity solutions that \mathbf{u} is a solution to (3.3). The fact that $\mathbf{u}_k \to \mathbf{u}$ locally uniformly in \mathbb{R}^n follows from the uniqueness of solutions to (3.3). \Box

Now we are ready to give a proof for Theorem 1.1 using the doubling variables method.

Proof of Theorem 1.1. We first note that u_k solves $u_k(x) + H(x, Du_k(x)) \ge 0$ on B(0, k), and u solves $u(x) + H(x, Du(x)) \le 0$ in B(0, k) in viscosity sense. By the comparison principle, we get $u_k(x) \ge u(x)$ for all $x \in \overline{B(0, k)}$.

For the upper bound of $u_k - u$, we define the following auxiliary function

$$\Phi^{k}(x,y) = u_{k}(x) - u(y) - 2C_{H}k^{2}|x-y|^{2} - \frac{8C_{H}}{k^{2}}|y|^{2}$$

for $(x, y) \in \overline{B(0, k)} \times \mathbb{R}^n$. It is clear that Φ^k is bounded above by $2C_H$ independent of $k \in \mathbb{N}$. If $|y| > \frac{k}{2}$, then we have

$$\Phi^{k}(0,0) - \Phi^{k}(x,y) \ge -u_{k}(x) + u_{k}(0) - u(0) + u(y) + 2C_{H}k^{2}|x-y|^{2} + \frac{8C_{H}}{k^{2}}|y|^{2} > 0,$$

which implies that for each $k \in \mathbb{N}$, $\Phi^{k}(x,y)$ achieves a global maximum over $\overline{B(0,k)} \times \mathbb{R}^{n}$ at $(x_{k},y_{k}) \in \overline{B(0,k)} \times \overline{B(0,\frac{k}{2})}$. Of course, $|y_{k}| \leq \frac{k}{2}$. Now we use $\Phi^{k}(x_{k},y_{k}) \geq \Phi^{k}(y_{k},y_{k})$ to get

$$2C_{H}k^{2}|x_{k}-y_{k}|^{2} \leq u_{k}(x_{k})-u_{k}(y_{k}) \leq C_{H}|x_{k}-y_{k}|$$

Therefore, we deduce that

$$|\mathbf{x}_{k}| \leq |\mathbf{y}_{k}| + \frac{1}{2k^{2}} < k \tag{3.4}$$

for all $k \ge 1$ since $|y_k| \le \frac{k}{2}$. Observing that $x \mapsto \Phi^k(x, y_k)$ obtains a maximum at x_k with $|x_k| < k$, we have

$$\mathbf{u}_{\mathbf{k}}(\mathbf{x}_{\mathbf{k}}) + \mathbf{H}\left(\mathbf{x}_{\mathbf{k}}, \mathbf{p}_{\mathbf{k}}\right) \leqslant \mathbf{0},\tag{3.5}$$

where $p_k = 4C_H k^2 (x_k - y_k)$ by the definition of viscosity subsolutions. We also observe that $y \mapsto \Phi^k(x_k, y)$ obtains a maximum at y_k , which implies that

$$u(y) - \left(-2C_{H}k^{2}|x_{k} - y_{k}|^{2} - \frac{8C_{H}}{k^{2}}|y|^{2}\right)$$

has a minimum at y_k . By the definition of viscosity supersolutions, we get

$$\mathfrak{u}(\mathfrak{y}_k) + \mathsf{H}(\mathfrak{y}_k, \mathfrak{p}_k + \mathfrak{q}_k) \ge 0 \tag{3.6}$$

where $q_k = -\frac{16C_H}{k^2}y_k$. Here, it needs to be noted that

$$|\mathbf{p}_{k}|, |\mathbf{p}_{k} + \mathbf{q}_{k}| \leq C_{H},$$

which comes from Lipschitz continuity of u_k . Using (3.5), (3.6) and assumption (H3c), there exists a constant \tilde{C}_H such that

$$\begin{aligned} u_{k}(x_{k}) - u(y_{k}) &\leq H(y_{k}, p_{k} + q_{k}) - H(x_{k}, p_{k}) \\ &= H(y_{k}, p_{k} + q_{k}) - H(y_{k}, p_{k}) + H(y_{k}, p_{k}) - H(x_{k}, p_{k}) \\ &\leq \tilde{C}_{H} |q_{k}| + \tilde{C}_{H} |x_{k} - y_{k}| \\ &\leq \frac{16\tilde{C}_{H}C_{H}}{k^{2}} |y_{k}| + \frac{\tilde{C}_{H}}{k^{2}} \leq \frac{8\tilde{C}_{H}C_{H}}{k} + \frac{\tilde{C}_{H}}{k^{2}}. \end{aligned}$$
(3.7)

If we stop here, the fact that $\Phi^k(x_k, y_k) \ge \Phi^k(x, x)$ for $x \in B(0, k)$ gives

$$u_k(x) - u(x) \le u_k(x_k) - u(y_k) + \frac{8C_H}{k^2} |x|^2 \le \frac{C}{k} + \frac{C(1 + |x|^2)}{k^2}$$

for all $k \ge 2$. This gives us the rate of convergence of u_k to u is $O(\frac{1}{k})$ for $x \in B(0, R)$, which is typically the case in light of the doubling variables method.

Nevertheless, a key new point here is to bootstrap once more to improve this rate. The monotonicity of $\{u_k\}$ allows us to bound $|y_k|$ better. We use that $\Phi^k(x_k, y_k) \ge \Phi^k(0, 0)$ together with (3.7) and $u_k \ge u$ to yield

$$\begin{split} 2C_{\mathrm{H}}k^{2}|x_{k}-y_{k}|^{2} + \frac{8C_{\mathrm{H}}}{k^{2}}|y_{k}|^{2} &\leqslant \mathfrak{u}_{k}(x_{k}) - \mathfrak{u}_{k}(0) + \mathfrak{u}(0) - \mathfrak{u}(y_{k}) \\ &\leqslant \mathfrak{u}_{k}(x_{k}) - \mathfrak{u}(y_{k}) \\ &\leqslant \frac{16\tilde{C}_{\mathrm{H}}C_{\mathrm{H}}}{k^{2}}|y_{k}| + \frac{\tilde{C}_{\mathrm{H}}}{k^{2}}. \end{split}$$

Therefore,

$$|y_{k}|^{2} \leq 2\tilde{C}_{H}|y_{k}| + \frac{\tilde{C}_{H}}{8C_{H}} \leq \frac{1}{2}|y_{k}|^{2} + 2\tilde{C}_{H}^{2} + \frac{\tilde{C}_{H}}{8C_{H}} = \frac{1}{2}|y_{k}|^{2} + C.$$

In particular, $|y_k| \leq C$. This bound is much better than the earlier bound that $|y_k| \leq \frac{k}{2}$.

Now for any $x \in \overline{B(0,k)}$, clearly we have that $\Phi^k(x_k, y_k) \ge \Phi^k(x, x)$. This, together with (3.7) and $|y_k| \le C$, implies

$$u_k(x) - u(x) \le u_k(x_k) - u(y_k) + \frac{8C_H}{k^2} |x|^2 \le \frac{C(1 + |x|^2)}{k^2}$$

for all $k \ge 2$. If $|x| \le R$, then

$$0 \leq u_k(x) - u(x) \leq \frac{C(1+R^2)}{k^2},$$

which gives the desired result.

Remark 2. In the general setting, one only has that $\Phi^k(x, y)$ achieves a global maximum over $\overline{B(0, k)} \times \mathbb{R}^n$ at (x_k, y_k) where $|y_k| \leq \frac{k}{2}$ and $|x_k| < k$. In our current situation, the monotonicity of $\{u_k\}$ allows us to bootstrap once more to deduce further that $|y_k| \leq C$, which helps to obtain $\mathcal{O}\left(\frac{1}{k^2}\right)$ rate of convergence. This seems to be the best convergence rate that one is able to get through the doubling variables method here as it is unlikely that $|y_k|$ vanishes as $k \to \infty$.

We do not know yet whether the $O\left(\frac{1}{k^2}\right)$ rate of convergence is optimal or not in the general nonconvex setting. See Questions 1 and 2 in Section 7 below.

4. An optimal rate for a class of nonconvex Hamiltonians on unbounded domain

In this section, we show that the rate of convergence $u_k \to u$ is of order $\mathcal{O}(e^{-Ck})$ for a class of possibly nonconvex Hamiltonians which are written as H(x,p) = a(x)K(p) with K(0) = 0and $0 < \alpha \leq a(x) \leq \beta$. The aforementioned rate is indeed optimal.

A brief idea for the proof is that we construct a supersolution to (3.1) by finding a symmetric Hamiltonian \tilde{H} such that $\tilde{H}(0) = 0$ and $\tilde{H} \leq H$. The following proposition is needed as a building block.

Proposition 4.1. Let $H : \mathbb{R}^n \to \mathbb{R}$ be defined by

$$H(\mathbf{p}) = \begin{cases} -\alpha |\mathbf{p}| & \text{for } |\mathbf{p}| \leq \beta, \\ f(\mathbf{p}) & \text{for } |\mathbf{p}| \geq \beta, \end{cases}$$

where $\alpha, \beta > 0$ and $f : \mathbb{R}^n \to \mathbb{R}$ is a coercive continuous function such that $f(p) = -\alpha\beta$ for $|p| = \beta$ and $\min_{\mathbb{R}^n} f = -\alpha\beta$. Then,

$$\mathfrak{u}_k(\mathfrak{x}) = \alpha \beta e^{\frac{|\mathfrak{x}|-k}{\alpha}}$$

for $\mathbf{x} \in \overline{B(\mathbf{0}, \mathbf{k})}$ is the unique solution to the state-constraint problem (3.1).

Proof. It is clear that $u_k(x) + H(Du_k(x)) = 0$ in $B(0,k) \setminus \{0\}$ in classical sense. For $x \in \partial B(0,k)$ and $\varphi \in C^1(\overline{B(0,k)})$ such that $u_k - \varphi$ has a local minimum over $\overline{B(0,k)}$ at x, we have $u_k(x) + H(D\varphi(x)) \ge 0$ since $u_k(x) = \alpha\beta = -\min H$. We only need to check if u_k is a viscosity supersolution at x = 0.

Let $\varphi \in C^1(\mathbb{R}^n)$ such that $\varphi(0) = u_k(0)$ and $u_k - \varphi$ has a local minimum over $\overline{B(0,k)}$ at x = 0. Since u_k is convex, we can replace φ by an affine function $\varphi(x) = \xi \cdot x + u_k(0)$ for some $\xi \in \mathbb{R}^n$. Without loss of generality, it suffices to consider $\xi \neq 0$. For |x| sufficiently small, we have $u_k(x) - \varphi(x) \ge u_k(0) - \varphi(0)$, which implies that

$$\alpha\beta e^{-\frac{k}{\alpha}} \left(e^{\frac{|\mathbf{x}|}{\alpha}} - 1 \right) \geqslant \xi \cdot \mathbf{x}. \tag{4.1}$$

Now we choose $\mathbf{x} = \mathbf{t} \frac{\xi}{|\xi|}$ for $\mathbf{t} > 0$ small, then (4.1) implies that $\alpha \beta e^{-\frac{\mathbf{k}}{\alpha}} \left(e^{\frac{\mathbf{t}}{\alpha}} - 1 \right) \ge \mathbf{t} |\xi|$ for all $\mathbf{t} > 0$ sufficiently small. Dividing both sides by \mathbf{t} and sending \mathbf{t} to 0, we deduce that $|\mathsf{D}\varphi(0)| = |\xi| \le \beta e^{-\frac{\mathbf{k}}{\alpha}}$. Therefore,

$$\mathfrak{u}_{k}(\mathfrak{0}) + \mathsf{H}(\mathsf{D}\varphi(\mathfrak{0})) = \alpha\beta e^{-\frac{k}{\alpha}} - \alpha|\mathsf{D}\varphi(\mathfrak{0})| \ge 0.$$

Consequently, u_k is the unique viscosity solution to (3.1).

Proof of Theorem 1.2. Since K(0) = 0, $u \equiv 0$ is the unique solution to (3.3). Recalling the a priori estimate $\|u_k\|_{L^{\infty}(B(0,k))} + \|Du_k\|_{L^{\infty}(B(0,k))} \leq C_H$, condition (H3c) gives

$$|\mathsf{K}(\mathsf{p}) - \mathsf{K}(\mathsf{q})| \leq \mathsf{L}|\mathsf{p} - \mathsf{q}|$$

for all $p, q \in \overline{B(0, C_H)}$. Let $K(p_0) = \min K \leq 0$ for some $p_0 \in \mathbb{R}^n$. Let f(p) be a coercive, continuous function such that $f(p) = -L|p_0|$ for $|p| \leq |p_0|$, $\min_{\mathbb{R}^n} f = -L|p_0|$, and $f(p) \leq K(p)$ for $|p| \geq |p_0|$. Now we consider

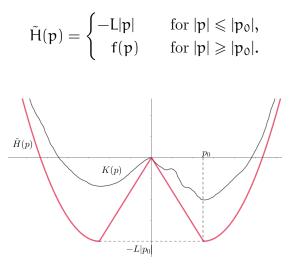


FIGURE 4.1. The graph of H(p) and K(p).

The graph of \tilde{H} is described in Figure 4.1. It is clear that $\tilde{H}(p) \leq K(p)$ for all $p \in \mathbb{R}^n$. Moreover, using Proposition 4.1, the unique viscosity solution to the state-constraint problem $\tilde{u}_k(x) + \beta \tilde{H}(D\tilde{u}_k(x)) = 0$ in B(0,k) is given by $\tilde{u}_k(x) = \beta L |p_0| e^{\frac{|x|-k}{\beta L}}$ for $x \in \overline{B(0,k)}$. It is clear that \tilde{u}_k is also the unique viscosity solution to $\frac{1}{\beta}\tilde{u}_k(x) + \tilde{H}(D\tilde{u}_k(x)) = 0$ in B(0,k). Since $\beta \ge a(x) \ge \alpha > 0$ and $\tilde{H} \le K$, we deduce that

$$\frac{1}{\beta}\tilde{u}_{k}(x) + \tilde{H}(D\tilde{u}_{k}(x)) \leqslant \frac{1}{a(x)}\tilde{u}_{k}(x) + K(D\tilde{u}_{k}(x))$$

on B(0,k). Therefore, $\tilde{u}_k(x) + a(x)K(D\tilde{u}_k(x)) \ge 0$ on B(0,k). By the comparison principle, one gets

$$0 \leq u_k(x) \leq \beta L |p_0| e^{\frac{|x|-k}{\beta L}}$$

for all $x \in \overline{B(0, k)}$. The conclusion for $|x| \leq R$ follows immediately.

In case that H(x,p) = K(p) for $(x,p) \in \mathbb{R}^n \times \mathbb{R}^n$, where K is locally Lipschitz continuous and coercive in \mathbb{R}^n , we have the unique viscosity solution to (3.3) is $u \equiv -K(0)$. Therefore, we can assume that K(0) = 0, and Corollary 1.3 follows without assuming that K(0) = 0.

It should be noted that the local Lipschitz continuity of Hamiltonians is important when it comes to getting an exponential rate of convergence. If a Hamiltonian is only Hölder continuous around 0, we get a slower rate of convergence depending on the regularity of H as described in the following proposition.

Proposition 4.2. Let $H : \mathbb{R}^n \to \mathbb{R}$ defined by

$$H(\mathbf{p}) = \begin{cases} -|\mathbf{p}|^{\gamma} & \text{if } |\mathbf{p}| \leq 1, \\ f(\mathbf{p}) & \text{if } |\mathbf{p}| \geq 1, \end{cases}$$

where $\gamma \in (0,1)$ and $f : \mathbb{R}^n \to \mathbb{R}$ is a continuous, coercive function with f(x) = -1 for $|x| \leq 1$, and $\min_{\mathbb{R}^n} f = -1$. Then, the solution to (3.1) is given by

$$u_{k}(\mathbf{x}) = \left[\frac{1-\gamma}{\gamma}\left(k + \frac{\gamma}{1-\gamma} - |\mathbf{x}|\right)\right]^{\frac{\gamma}{\gamma-1}}, \qquad \mathbf{x} \in \overline{B(\mathbf{0}, \mathbf{k})}.$$
(4.2)

As a consequence, $\mathbf{u}_k \to \mathbf{0}$ with the rate $\mathbb{O}\left(\frac{1}{k^{\frac{\gamma}{1-\gamma}}}\right)$.

Proof. Let us first consider the one dimensional case. The higher dimensional setting can be done in a same manner. Let $\mu = \gamma^{-1}$, we look for a nonnegative solution to $u(x)^{\mu} = u'(x)$ where $x \in (0, k)$. We have

$$u(x)^{1-\mu} = (1-\mu)x - C_k \implies u(x) = (\mu-1)^{\frac{1}{1-\mu}}(C_k-x)^{\frac{1}{1-\mu}}.$$

We want to choose C_k such that $u'(x) \in [0, 1]$ for $x \in (0, k)$. Equivalently,

$$\mathfrak{u}'(x) = \mathfrak{u}(x)^{\mu} = (\mu - 1)^{\frac{\mu}{1-\mu}} (C_k - x)^{\frac{\mu}{1-\mu}} \in [0, 1]$$

for $x \in (0, k)$. Since it is an increasing function, $C_k = k + \frac{1}{\mu - 1}$. Using symmetry, we guess that u_k is written as

$$u_k(x) = (\mu - 1)^{\frac{1}{1-\mu}} \left(k + \frac{1}{\mu - 1} - |x|\right)^{\frac{1}{1-\mu}}.$$

It is straightforward to see that u_k satisfies the equation in the classical sense $u_k(x) - |u'_k(x)|^{\gamma} = 0$ in $(-k, k) \setminus \{0\}$. Since $|u'_k(x)| \leq 1$ on $(-k, k) \setminus \{0\}$, we have $u_k(x) + H(u'_k(x)) = 0$ in the classical sense in $(-k, k) \setminus \{0\}$. At |x| = k, we have $u_k(x) = 1 \geq -\min H$. Therefore,

the supersolution test at these points are satisfied. Finally, at x = 0 we only need to verify the supersolution test, which is simple since if $p \in D^-u_k(0)$ then

$$|\mathfrak{p}| \leqslant (\mu-1)^{\frac{\mu}{1-\mu}} \left(k + \frac{\mu}{\mu-1} \right)^{\frac{\mu}{1-\mu}} \implies \mathfrak{u}_k(\mathfrak{0}) + \mathfrak{H}(\mathfrak{p}) \geqslant \mathfrak{u}_k(\mathfrak{0}) - |\mathfrak{p}|^{\mu} \geqslant \mathfrak{0}.$$

Thus, $\mathbf{u}_{\mathbf{k}}$ defined above is the unique viscosity solution to the constraint problem (3.1). Using a similar argument as in the proof of Proposition 4.1, this formula of $\mathbf{u}_{\mathbf{k}}$ can be extended naturally to the **n**-dimensional case, as given in (4.2) and the conclusion follows.

Remark 3. From Proposition 4.2 we see that the optimal rate of convergence can be as slow as we wish as the Hölder exponent $\gamma \to 0^+$. This shows that the required condition (H3c) is really essential in this section.

When Hamiltonians are of the form H(x,p) = K(p) + V(x), the situation becomes much more complicated. See Example 2 for a situation where we get the optimal exponential rate of convergence with nonconvex K.

5. An optimal rate for convex Hamiltonians

In this section, the assumptions (H1), (H2), (H3b), (H5) are always in force. The stateconstraint problem was studied in the context of optimal control for convex Hamiltonians (see [27, 10, 4] for instance). When H is convex, we are able to obtain a representation formula for the viscosity solution based on the optimal control theory. Let us assume the following superlinear property (see Remark 4 where we can remove this assumption), which is

(H6) $\mathbf{p} \mapsto \mathbf{H}(\mathbf{x}, \mathbf{p})$ is superlinear uniformly for $\mathbf{x} \in \overline{\Omega}$, that is,

$$\lim_{|\mathfrak{p}|\to\infty} \left(\inf_{\mathbf{x}\in\Omega} \frac{\mathsf{H}(\mathbf{x},\mathfrak{p})}{|\mathfrak{p}|}\right) = +\infty.$$
(H6)

If (H5) and (H6) hold, then the Legendre transform $L: \overline{\Omega} \times \mathbb{R}^n$ of H is defined as

$$L(x,\nu) := \sup_{p \in \mathbb{R}^n} \{ p \cdot \nu - H(x,p) \}, \qquad (x,\nu) \in \overline{\Omega} \times \mathbb{R}^n.$$

Lemma 5.1. Assume (H5) and (H6). Then, $L: \overline{\Omega} \times \mathbb{R}^n \to \mathbb{R}$ is continuous satisfying:

- (L1) If (H1) holds, then $L(x, 0) \leq C_1$ for all $x \in \overline{\Omega}$;
- (L2) If (H2 holds, then $L(x, v) \ge -C_2$ for all $(x, v) \in \overline{\Omega} \times \mathbb{R}^n$;
- (L3) If (H3b) holds, then for each R > 0 there exists a modulus $\tilde{\omega}_{R}(\cdot)$ such that

$$|L(x,\nu) - L(y,\nu)| \leqslant \tilde{\omega}_{\mathsf{R}}(|x-y|) \qquad \textit{for all } x,y \in \overline{\Omega}, |\nu| \leqslant \mathsf{R}$$

- (L5) $\mathbf{v} \mapsto \mathbf{L}(\mathbf{x}, \mathbf{v})$ is convex for each $\mathbf{x} \in \overline{\Omega}$;
- (L6) $\mathbf{v} \mapsto \mathbf{L}(\mathbf{x}, \mathbf{v})$ is superlinear uniformly in $\mathbf{x} \in \overline{\Omega}$, *i.e.*,

$$\lim_{|\mathfrak{p}|\to\infty} \left(\inf_{\mathbf{x}\in\overline{\Omega}} \frac{\mathsf{L}(\mathbf{x},\mathbf{\nu})}{|\mathbf{\nu}|}\right) = +\infty.$$
 (L6)

We omit the proof of this lemma and refer the interested readers to [9]. For each $\mathbf{x} \in \overline{\Omega}$, we define the admissible set of paths as

$$\mathcal{A}_{x} = \left\{ \eta \in \mathrm{AC}\big([0,\infty); \mathbb{R}^{n}\big) : \eta(0) = x \text{ and } \eta(s) \in \overline{\Omega} \text{ for all } s \ge 0 \right\}$$

where $AC([0,\infty); \mathbb{R}^n)$ denotes the set of absolutely continuous curves from $[0,\infty)$ to \mathbb{R}^n . Note that $\mathcal{A}_x \neq \emptyset$ since $\eta(s) \equiv x$ for all $s \in [0,\infty)$ is an admissible path. From this, define the value function as

$$\mathfrak{u}(\mathbf{x}) := \inf_{\eta \in \mathcal{A}_{\mathbf{x}}} J\left[\mathbf{x}, \eta\right] \tag{5.1}$$

where the cost functional is defined as

$$J[x,\eta] = \int_0^\infty e^{-s} L(\eta(s), -\dot{\eta}(s)) ds$$

for $(x,\eta) \in \overline{\Omega} \times \mathcal{A}_x$. Now we have the following classical dynamic programming principle.

Theorem 5.2 (Dynamic Programming Principle). For any t > 0, we have

$$\mathfrak{u}(\mathbf{x}) = \inf_{\eta \in \mathcal{A}_{\mathbf{x}}} \left\{ \int_{0}^{t} e^{-s} L(\eta(s), -\dot{\eta}(s)) \, \mathrm{d}s + e^{-t} \mathfrak{u}(\eta(t)) \right\}.$$

Using the Dynamic Programming Principle, one can prove that $\mathfrak{u} \in BUC(\overline{\Omega})$ and indeed a viscosity solution to (2.1) as stated in the following theorems.

Theorem 5.3. Assume (H1), (H2), (H3c), (H5) and (H6), then the function $\mathbf{u}(\mathbf{x})$ defined by (5.1) is bounded and is uniformly continuous up to the boundary, which is $\mathbf{u} \in \text{BUC}(\overline{\Omega})$.

Theorem 5.4. The value function $\mathbf{u} \in \text{BUC}(\overline{\Omega})$ defined in (5.1) is a viscosity solution to the state-constraint Hamilton-Jacobi equation $\mathbf{u}(\mathbf{x}) + \mathbf{H}(\mathbf{x}, \mathbf{Du}(\mathbf{x})) = \mathbf{0}$ in Ω , i.e.,

$$\begin{cases} u(x) + H(x, Du(x)) \leq 0 & in \Omega, \\ u(x) + H(x, Du(x)) \geq 0 & on \overline{\Omega}. \end{cases}$$

We omit the proofs of Theorems 5.2, 5.3 and 5.4. We refer to [27, 10, 4] for those who are interested.

On the other hand, when $\Omega = \mathbb{R}^n$, it is known that the function u(x) defined in (5.1) satisfies the Hamilton-Jacobi equation (3.3) in viscosity sense (see [4, 22] for instance).

Theorem 5.5. For each $x \in \mathbb{R}^n$, we define

$$u(\mathbf{x}) = \inf_{\eta \in \mathcal{A}_{\mathbf{x}}} \int_{0}^{\infty} e^{-s} L\left(\eta(s), -\dot{\eta}(s)\right) \, \mathrm{d}s \tag{5.2}$$

subject to $\mathcal{A}_{\mathbf{x}} = \{ \eta \in \mathrm{AC}([0,\infty); \mathbb{R}^n) : \eta(0) = \mathbf{x} \}$. Then, $\mathbf{u} \in \mathrm{BUC}(\mathbb{R}^n)$ is a viscosity solution to (3.3) and we have the following priori estimate:

$$\|\mathbf{u}\|_{\mathbf{L}^{\infty}(\mathbb{R}^{n})} + |\mathbf{D}\mathbf{u}\|_{\mathbf{L}^{\infty}(\mathbb{R}^{n})} \leqslant C_{\mathbf{H}}.$$
(5.3)

Remark 4. We may assume that H is just coercive rather than superlinear. When a Hamiltonian is coercive, we still have that (5.3) holds for some $C = C_H > 0$. Therefore, for $|\mathbf{p}| \ge C$, we can modify H so that (H6) holds. Furthermore, we can impose a quadratic growth rate on H as following.

(H7) There exist some positive constants A, B such that

$$A^{-1}|v|^2 - B \leqslant H(x,p) \leqslant A|v|^2 + B \quad \text{for } (x,p) \in \mathbb{R}^n \times \mathbb{R}^n.$$
(H7)

It is easy to see from (H7) that we have $(4A)^{-1}|\nu|^2 - B \leq L(x,\nu) \leq 4A|\nu|^2 + B$ for all $(x,\nu) \in \mathbb{R}^n \times \mathbb{R}^n$. By making A bigger, we can assume the following.

(L7)

$$A^{-1}|\nu|^2 - B \leqslant L(x,\nu) \leqslant A|\nu|^2 + B \qquad \text{for } (x,\nu) \in \mathbb{R}^n \times \mathbb{R}^n. \tag{L7}$$

We give a proof for the existence of a minimizer with bounded velocity to (5.2) for the sake of readers' convenience in Appendix. This is an extremely important fact in our analysis and is a key element in the proof of Theorem 1.4 (see Remark 6 for further discussions). To establish this point, the following lemma on the subdifferentials of L(x, v) in v is needed. For continuously differentiable Lagrangians, it is obvious, but we state here a slightly more general version.

Lemma 5.6. Let $L : \mathbb{R}^n \times \mathbb{R}^n \to \mathbb{R}$ be continuous and satisfy (L5) and (L7). There exists $C_L > 0$ such that for all $\nu \in \mathbb{R}^n$, we have

$$|\xi| \leqslant C_{L}(1+|\nu|) \qquad whenever \ \xi \in D_{\nu}^{-}L(x,\nu).$$
(5.4)

For simplicity, let us assume further that

(L8) $(\mathbf{x}, \mathbf{v}) \mapsto \mathbf{L}(\mathbf{x}, \mathbf{v})$ is continuously differentiable on $\mathbb{R}^n \times \mathbb{R}^n$.

This assumption can be removed in the proof of Theorem 1.4 due to the fact that the estimate (1.1) does not depend on the regularity of H, hence, we can approximate H by convex, smooth Hamiltonians.

Theorem 5.7 (Existence of a minimizer). Let L(x, v) be a continuous Lagrangian satisfying (L5), (L7) and (L8). Then, for each $x \in \mathbb{R}^n$, there exists $\eta \in \mathcal{A}_x$ such that $J[x, \eta] = u(x)$ and also $\|e^{-s/2}\dot{\eta}(s)\|_{L^2(0,\infty)} \leq C_4$ where C_4 depends only on C_H, A, B .

The existence of minimizers of smooth Lagrangian is sufficient for our proof of Theorem 1.4 since the last estimate does not depend on the smoothness of L or H. Clearly, a minimizer for a general continuous Lagrangian can be obtained via approximation of smooth Lagrangians (see Appendix).

A minimizer to (5.2) satisfies the following properties.

Lemma 5.8. Let $x \in \mathbb{R}^n$ and η be a corresponding minimizer. For any t > 0, we have

$$u(x) = \int_0^t e^{-s} L(\eta(s), -\dot{\eta}(s)) \, ds + e^{-t} u(\eta(t)) \,.$$
 (5.5)

Furthermore, for every t, h > 0, we have

$$\mathfrak{u}(\eta(t)) = e^t \int_t^\infty e^{-s} L\big(\eta(s), -\dot{\eta}(s)\big) \, ds \tag{5.6}$$

and

$$e^{-t}\mathfrak{u}(\eta(t)) = \int_{t}^{t+h} e^{-s} L(\eta(s), -\dot{\eta}(s)) \, ds + e^{-(t+h)}\mathfrak{u}(\eta(t+h)).$$
 (5.7)

Lemma 5.9. Let $x \in \mathbb{R}^n$ and η be a minimizer to (5.2) associated with it. Then, there exists a constant $C_5 > 0$ depending only on C_H, A, B such that $|\dot{\eta}(s)| \leq C_5$ for a.e. $s \in (0, \infty)$.

Remark 5. We provide here a connection between a minimizer η of $u(x) = J[x, \eta]$ and some properties in the view of the method of characteristics. If H is assumed to be C^2 , then $L \in C^2$ and η is a weak solution to the Euler-Lagrange equation

$$D_{x}L(\eta(s), -\dot{\eta}(s)) - D_{\nu}L(\eta(s), -\dot{\eta}(s)) + \frac{d}{ds} \Big(D_{\nu}L(\eta(s), -\dot{\eta}(s)) \Big) = 0.$$
(5.8)

14

Assume that $\eta \in C^2$ (it holds if, for instance $L \in C^{2,\alpha}$ for some $\alpha \in (0,1)$). Then, one can define the momentum $\mathbf{p}(s) = D_{\nu}L(\eta(s), -\dot{\eta}(s))$ and show that

$$\mathbf{u}(\mathbf{\eta}(\mathbf{t})) + \mathbf{H}(\mathbf{\eta}(\mathbf{t}), \mathbf{p}(\mathbf{t})) = \mathbf{0}$$
(5.9)

for t > 0. Indeed, for every fixed $x \in \mathbb{R}^n$, we recall that

$$\nu \in D_p^- H(x, p) \quad \Leftrightarrow \quad p \in D_\nu^- L(x, \nu) \quad \Leftrightarrow \quad H(x, p) + L(x, \nu) = p \cdot \nu.$$
 (5.10)

Using (5.10) we can deduce that

$$\begin{cases} \frac{\mathrm{d}}{\mathrm{d}s} \left(e^{-s} \mathbf{p}(s) \right) &= e^{-s} \mathsf{D}_{\mathsf{x}} \mathsf{L} \left(\eta(s), \dot{\eta}(s) \right), \\ \frac{\mathrm{d}}{\mathrm{d}s} \left(\mathsf{H} \left(\eta(s), \mathbf{p}(s) \right) \right) &= -\dot{\eta}(s) \cdot \mathbf{p}(s), \\ \frac{\mathrm{d}}{\mathrm{d}s} \left(e^{-s} \mathsf{H} \left(\eta(s), \mathbf{p}(s) \right) \right) &= e^{-s} \mathsf{L} \left(\eta(s), -\dot{\eta}(s) \right). \end{cases}$$

From that, we can derive the characteristic ODEs for s > 0, which are

$$\begin{cases} -\dot{\eta}(s) &= \mathsf{D}_{\mathsf{p}}\mathsf{H}(\eta(s),\mathbf{p}(s)), \\ \dot{\mathbf{p}}(s) &= \mathbf{p}(s) - \mathsf{D}_{\mathsf{x}}\mathsf{L}\big(\eta(s),-\dot{\eta}(s)\big). \end{cases}$$

This together with (5.6) yields that

$$\mathfrak{u}(\eta(t)) + H\big(\eta(t),\mathbf{p}(t)\big) = Ce^t \qquad \text{where} \qquad C = \lim_{a\to\infty} e^{-\mathfrak{a}} H\big(\eta(a),\mathbf{p}(a)\big).$$

Lemma 5.6 together with Lemma 5.9 gives us a uniform bound on \mathbf{p} , thus $\mathbf{C} = \mathbf{0}$. Hence, (5.9) follows.

Now we give a proof for Theorem 1.4. Recall that we have the value function

$$u_{k}(\mathbf{x}) = \inf_{\eta \in \mathcal{A}_{\mathbf{x}}^{k}} \int_{0}^{\infty} e^{-s} L\left(\eta(s), -\dot{\eta}(s)\right) \, \mathrm{d}s, \tag{5.11}$$

where $\mathcal{A}_{x}^{k} = \left\{ \eta \in \mathrm{AC}([0,\infty); \mathbb{R}^{n}) : \eta(0) = x \text{ and } \eta(s) \in \overline{B(0,k)} \text{ for } s \ge 0 \right\}$. Then, u_{k} solves the state-constraint problem (3.1).

Proof of Theorem 1.4. Let $k \in \mathbb{N}$ be given. We may assume that H satisfies (H6) and (H7) up to modification for |p| large enough. Also, since the final estimate does not depend on the smoothness of L, we can assume H is smooth and thus L is smooth without any loss of generality. Clearly, $\mathcal{A}_x^k \subset \mathcal{A}_x$ for any $x \in \overline{B(0,k)}$, which implies that $u_k(x) \ge u(x)$.

For $x \in B(0, k)$, let $\eta \in A_x$ be a minimizer to (5.2), if $\eta(s) \in \overline{B(0, k)}$ for all s > 0, then $\eta \in A_x^k$ as well, hence $u(x) = u_k(x)$. Otherwise, there exists t > 0 such that $\eta(t) \in \partial B(0, k)$ and $\eta(s) \in B(0, k)$ for all $s \in (0, t)$. By Lemma 5.9, we have

$$k = |\eta(t)| \leqslant |\eta(0)| + \int_0^t |\dot{\eta}(s)| \, ds \leqslant |x| + C_5 t,$$

which implies that $t \ge \frac{k-|x|}{C_5}$. Let us define

$$\gamma(s) = egin{cases} \eta(s) & ext{if } s \in [0,t], \ \eta(t) & ext{if } s \in [t,\infty), \end{cases}$$

so that $\gamma \in \mathcal{A}_{\mathbf{x}}^{\mathbf{k}}$. Using Lemma 5.8, we have

$$\begin{split} u(x) &= \int_{0}^{t} e^{-s} L(\eta(s), -\dot{\eta}(s)) \, ds + e^{-t} u(\eta(t)) \\ &\geq \int_{0}^{t} e^{-s} L(\gamma(s), -\dot{\gamma}(s)) \, ds - C_{H} e^{-t} \\ &\geq \int_{0}^{\infty} e^{-s} L(\gamma(s), -\dot{\gamma}(s)) \, ds - C_{1} e^{-t} - C_{H} e^{-t} \\ &\geq u_{k}(x) - \left((C_{1} + C_{H}) e^{\frac{|x|}{C_{5}}} \right) e^{-\frac{k}{C_{5}}}. \end{split}$$

Consequently, we obtain (1.1). The conclusion for $|\mathbf{x}| \leq \mathbf{R}$ follows immediately.

Remark 6. Here, we note that the constants in the proof above do not depend on the regularity of the Lagrangian. As long as a minimizer exists, we get the same exponential rate of convergence. See Appendix for a discussion on the existence of minimizers. It is worth noting here that, for each $x \in \overline{B(0,k)}$, the existence of a minimizer $\eta \in \mathcal{A}_x$ to (5.2) with bounded velocity is a nontrivial fact and plays an essential role in the proof above. Moreover, the bound on the velocity of η only depends on C_H, A, B .

In the rest of this section, we provide two explicit examples to show that the rate $O\left(e^{-\frac{k}{C}}\right)$ is indeed optimal.

5.1. Examples with exponential rate of convergence.

Proposition 5.10. Let $H(p) : \mathbb{R} \to \mathbb{R}$ be defined by H(p) = |p - 1| - 1 for $p \in [0, 2]$ and $H(p) \ge 0$ elsewhere such that H is continuous and coercive. Let u_k be the solution to (3.1) on [-k, k], then $u_k \to 0$ locally uniformly on \mathbb{R} as $k \to \infty$. Here, $u \equiv 0$ is the unique solution to (3.3). Furthermore, we have $u_k(k) = 1$ for all $k \in \mathbb{N}$ and

$$u_k(x) \ge e^{-2k}$$
 on $[-k,k]$.

Proof. It is clear that $u_k(x) = e^{x-k}$ solves v(x) + H(v'(x)) = 0 in (-k, k) in the classical sense, and indeed, in viscosity sense. We need to verify that u_k is a viscosity supersolution on [-k, k]. Let $u_k - \phi$ has a local minimum at x = -k for $\phi(x) \in C^1(\mathbb{R})$. Clearly, we can see that

$$\varphi'(-k) \leqslant \mathfrak{u}'_k(-k) = e^{-2k},$$

which implies $e^{-2k} + H(\varphi'(-k)) \ge 0$. On the other hand, at x = k, one has

$$\mathfrak{u}_{k}(k) + H(\varphi'(k)) = 1 + H(\varphi'(k)) \ge 0$$

since by definition of H, it is bounded below by -1. Therefore, $u_k(x) = e^{x-k}$ is the unique viscosity solution to (3.1), and furthermore $e^{-2k} \leq u_k(x) \leq (e^{|x|}) e^{-k}$ for all $x \in [-k, k]$. In addition to that, we have $u_k(k) = 1$ for all $k \in \mathbb{N}$, hence, the convergence fails when x = k. 5.2. Optimal control formulations. We give another example from the optimal control theory point of view (see [27]). Let us recall briefly the setting of optimal control as follows. Let U be a compact metric space. We regard a control as a Borel measurable map α : $[0, \infty) \mapsto U$. Let Ω be an open subset of \mathbb{R}^n with the connected boundary satisfying (A). We also assume that $\mathbf{b} = \mathbf{b}(\mathbf{x}, \mathbf{a}) : \overline{\Omega} \times \mathbf{U} \to \mathbb{R}^n$, $\mathbf{f} = \mathbf{f}(\mathbf{x}, \mathbf{a}) : \overline{\Omega} \times \mathbf{U} \to \mathbb{R}$ satisfy

$$\begin{split} \sup_{a \in U} |b(x, a) - b(y, a)| &\leq L(b)|x - y| & \text{for all } x, y \in \overline{\Omega}, \\ \sup_{a \in U} |b(x, a)| &\leq K(b) & \text{for all } x \in \overline{\Omega}, \\ \sup_{a \in U} |f(x, a) - f(y, a)| &\leq \omega_f(|x - y|) & \text{for all } x, y \in \overline{\Omega}, \\ \sup_{a \in U} |f(x, a)| &\leq K(f) & \text{for all } x \in \overline{\Omega}, \end{split}$$

where K(b), L(b), K(f) are positive constants and ω_f is a nondecreasing continuous function with $\omega_f(0^+) = 0$.

For each $x \in \overline{\Omega}$ and a given control $\alpha(\cdot) : [0, \infty) \to U$, let $y^{x,\alpha}(t)$ be a controlled process (we will write α instead of $\alpha(\cdot)$ as a control for simplicity), which is a solution to

$$\begin{cases} \frac{\mathrm{d}}{\mathrm{d}t}y^{x,\alpha}(t) = b\left(y^{x,\alpha}(t),\alpha(t)\right) & \quad \mathrm{for} \ t > 0, \\ y^{x,\alpha}(0) = x. \end{cases}$$

We denote the set of controls (strategies) α where $y^{x,\alpha}(t) \in \overline{\Omega}$ for all $t \ge 0$ and $y^{x,\alpha}$ solves the ODE above by \mathcal{A}_x . The value function is defined by

$$\mathfrak{u}(\mathbf{x}) = \inf_{\alpha \in \mathcal{A}_{\mathbf{x}}} \int_{0}^{\infty} e^{-t} f(\mathbf{y}^{\mathbf{x},\alpha}(t), \alpha(t)) dt.$$

Here, one can define the Hamiltonian associated with b and f as

$$H(x,p) := \sup_{a \in U} \{-b(x,a) \cdot p - f(x,a)\} \in C\left(\overline{\Omega} \times \mathbb{R}^n; \mathbb{R}\right).$$

It was proved in [27] that **u** is a viscosity solution to (2.1).

Proposition 5.11. Let n = 1 and U = [-1, 1]. Let us consider the following Hamiltonian defined as

$$H(\mathbf{x},\mathbf{p}) = \sup_{\mathbf{a}\in[-1,1]} \left\{ -\mathbf{a}\cdot\mathbf{p} - e^{-|\mathbf{x}|} \right\} = |\mathbf{p}| - e^{-|\mathbf{x}|}, \qquad (\mathbf{x},\mathbf{p}) \in \mathbb{R} \times \mathbb{R}.$$

Then, the solution to (3.1) is given by $u_k(x) = \frac{e^{-|x|}}{2} + \frac{e^{|x|-2k}}{2}$ for $x \in [-k, k]$, while the solution to (3.3) is $u(x) = \frac{e^{-|x|}}{2}$. Hence, the exponential rate of convergence is obtained.

Proof. In the optimal control setting, the Hamiltonian above is obtained by considering U = [-1, 1], b(x, a) = a and $f(x, a) = e^{-|x|}$. To find $u_k(x_0)$ and $u(x_0)$, one needs to find a control $\alpha(t)$ that minimizes

$$\int_0^\infty e^{-s-|y(s)|} \, ds \qquad \mathrm{subject \ to} \qquad \begin{cases} \dot{y}(t) &= \alpha(t) \in [-1,1], \\ y(0) &= x_0. \end{cases}$$

It is easy to see the following points:

- (i) An optimal control for the unconstrained problem with $x_0 \ge 0$ ($x_0 < 0$) is $\alpha(t) \equiv 1$ ($\alpha(t) \equiv -1$, respectively).
- (ii) An optimal control for the constrained problem on [-k, k] with $x_0 \ge 0$ ($x_0 < 0$) is $\alpha(t) \equiv 1$ on $[0, k x_0]$ and 0 elsewhere ($\alpha(t) \equiv -1$ on $[0, k + x_0]$ and 0 elsewhere, respectively).

Once we have the optimal controls, we can easily compute the value function and the result follows. In conclusion, for all $x \in [-k, k]$ we have

$$0 \leq u_k(x) - u(x) = \left(\frac{e^{|x|}}{2}\right)e^{-2k}.$$

In this example, the convergence holds everywhere in [-k, k] with the rate $\mathcal{O}(e^{-k})$.

Remark 7. One interesting fact to point out here is that the optimal path starting from x_0 for the state-constraint problem on [-k, k] in Proposition 5.11 stays on the boundary $\pm k$ for all $t \ge k - |x_0|$.

6. The case of bounded domain

The second prototype case is considered in this section. Let us assume that (P2), (H1), (H2), (H3c) and (H4) are enforced. Recall that $\Omega_k = B\left(0, 1 - \frac{1}{k}\right)$ for $k \in \mathbb{N}$, and $\Omega = B(0, 1)$. Let $u_k \in \text{Lip}\left(\overline{\Omega}_k\right)$ be the unique viscosity solution to

$$\begin{cases} u_{k}(x) + H(x, Du_{k}(x)) &\leq 0 & \text{ in } \Omega_{k}, \\ u_{k}(x) + H(x, Du_{k}(x)) &\geq 0 & \text{ on } \overline{\Omega}_{k}. \end{cases}$$
(6.1)

It is clear that we still have the following priori estimate

$$\|\mathbf{u}_{k}\|_{\mathbf{L}^{\infty}(\overline{\Omega}_{k})} + \|\mathbf{D}\mathbf{u}_{k}\|_{\mathbf{L}^{\infty}(\overline{\Omega}_{k})} \leqslant C_{\mathsf{H}}.$$
(6.2)

Proposition 6.1. For each $k \in \mathbb{N}$, let u_k be the unique solution to (6.1). Then, there exists $u \in BUC(\overline{\Omega})$ such that $u_k \to u$ locally uniformly on $\overline{\Omega}$ as k grows to infinity. Moreover, u has the same bounds as in (6.2) and solves

$$\begin{cases} u(x) + H(x, Du(x)) &\leq 0 & in \Omega, \\ u(x) + H(x, Du(x)) &\geq 0 & on \overline{\Omega} \end{cases}$$
(6.3)

in viscosity sense.

Proof. From a priori estimate (6.2), by Arzelà-Ascoli's theorem and a diagonal argument we can extract a subsequence such that $\mathbf{u}_{\mathbf{k}_m} \to \mathbf{u}$ uniformly on compact subsets of Ω . By the stability of viscosity solutions we obtain that $\mathbf{u} \in C(\Omega)$ is a viscosity solution to

$$\mathbf{u}(\mathbf{x}) + \mathbf{H}(\mathbf{x}, \mathbf{D}\mathbf{u}(\mathbf{x})) = \mathbf{0} \quad \text{in } \Omega.$$
(6.4)

We deduce that $|\mathfrak{u}(\mathbf{x})| \leq C_B$ and $|\mathfrak{u}(\mathbf{x}) - \mathfrak{u}(\mathbf{y})| \leq C_H |\mathbf{x} - \mathbf{y}|$ for $\mathbf{x}, \mathbf{y} \in \Omega$. We can extend $\mathfrak{u} \in \operatorname{Lip}(\overline{\Omega})$ with the same priori bound as in (6.2). We need to show that \mathfrak{u} is a viscosity supersolution to $\mathfrak{u}(\mathbf{x}) + H(\mathbf{x}, \operatorname{Du}(\mathbf{x})) = \mathfrak{0}$ on $\overline{\Omega}$.

We can verify it using Corollary 2.2. Indeed, let $v \in C(\overline{\Omega})$ be a viscosity subsolution to (6.4) in Ω . Applying the comparison principle to $u_k(x) + H(x, Du_k(x)) \ge 0$ on $\overline{\Omega}_k$, we have that $v(x) \le u_k(x)$ for $x \in \overline{\Omega}_k$. Now fixing $r \in (0, 1)$, we have $v(x) \le u_k(x)$ for all

 $x \in \overline{B(0,r)}$ and $r \leq 1 - \frac{1}{k}$ if k is large enough. Letting $k \to \infty$, we deduce that $\nu(x) \leq u(x)$ for $x \in \overline{B(0,r)}$. Since we have $u, \nu \in C(\overline{B(0,1)})$, the inequality $\nu \leq u$ on $\overline{B(0,1)}$ follows. Hence, u is a viscosity supersolution to (6.4) by Corollary 2.2.

Now we are ready to give a proof for Theorem 1.5. We note that star-shaped and scaling properties of $\{\Omega_k\}$ play an important role.

Proof of Theorem 1.5. The fact that $u_k(x) \ge u(x)$ on $\overline{\Omega}_k$ is clear by the comparison principle. For $k \ge 2$, let us define

$$\tilde{\mathfrak{u}}_k(\mathbf{x}) := \frac{k}{k-1} \mathfrak{u}_k\left(\frac{k-1}{k}\mathbf{x}\right) \qquad \text{for } \mathbf{x} \in \overline{\mathsf{B}}(\mathbf{0},\mathbf{1}).$$

It is clear that \tilde{u}_k is a viscosity subsolution to

$$\frac{k-1}{k}\tilde{u}_{k}(x) + H\left(\frac{k-1}{k}x, D\tilde{u}_{k}(x)\right) = 0 \quad \text{in } B(0,1).$$
(6.5)

From (6.2) and (H3c), there exists \tilde{C}_H such that $|H(x,p) - H(x,p)| \leq \tilde{C}_H |x-y|$ for all $x, y \in \overline{\Omega}$ and $|p| \leq C_H$. Therefore, by using (6.5) we have

$$\tilde{u}_{k}(x) + H(x, D\tilde{u}_{k}(x)) \leq \frac{1}{k}\tilde{u}_{k}(x) + H(x, D\tilde{u}_{k}(x)) - H\left(\frac{k-1}{k}x, D\tilde{u}_{k}(x)\right) \leq \frac{C_{H} + \tilde{C}_{H}}{k}$$

for all $x \in B(0, 1)$. By the comparison principle and the fact that u solves (6.3) in the viscosity sense, we deduce that

$$\tilde{\mathfrak{u}}_k(x) - \frac{C_H + \tilde{C}_H}{k} \leqslant \mathfrak{u}(x) \quad \text{ for all } x \in \overline{B(0,1)}.$$

Consequently, we obtain the conclusion $u_k(x) \leq u(x) + \frac{C}{k}$ for $x \in \overline{\Omega}_k$ where the constant C can be chosen as $C = 2C_H + \tilde{C}_H$.

Remark 8.

- (i) It is clear from the proof that prototype condition (P2) can be relaxed as following.
 - (P2') Assume $0 \in \Omega_k$ for all $k \in \mathbb{N}$, $\Omega = \bigcup_{k \in \mathbb{N}} \Omega_k$ is bounded, and the comparison principle for the state-constraint problem holds on Ω_k, Ω . Assume further that, for $k \in \mathbb{N}$,

$$\left(1-\frac{1}{k}\right)\Omega\subset\Omega_k.$$

(ii) Theorem 1.5 can also be proved using the doubling variable method with the following auxiliary function (see [10])

$$\Phi^{k}(\mathbf{x},\mathbf{y}) := \frac{k+1}{k-1} \mathfrak{u}_{k}\left(\frac{k-1}{k+1}\mathbf{x}\right) - \mathfrak{u}(\mathbf{y}) - C_{H}k^{2}|\mathbf{x}-\mathbf{y}|^{2}$$

for $(x, y) \in (1 + \frac{1}{k}) \overline{\Omega} \times \overline{\Omega}$.

The following remark shows that the rate $O\left(\frac{1}{k}\right)$ is indeed optimal.

Remark 9. Let H be defined as in Proposition 5.10, we see that $u_k(x) = e^{x - (1 - \frac{1}{k})}$ solves (6.1) and $u(x) = e^{x-1}$ solves (6.3), therefore

$$0 \leq u_k(x) - u(x) = e^{x-1} \left(e^{\frac{1}{k}} - 1 \right) \leq \frac{2}{k}$$

for $x \in \left[-\left(1-\frac{1}{k}\right), 1-\frac{1}{k}\right]$. Besides, $e^{\frac{1}{k}}-1 \ge \frac{1}{k}$, and so, $O\left(\frac{1}{k}\right)$ is optimal.

7. DISCUSSIONS

We give here some further discussions along the line with the topics considered in the paper. Firstly, when our Hamiltonian is given as H(x,p) = a(x)K(p) in the first prototype (P1), we get an exponential rate of convergence provided that the assumption (H1) is enforced (Theorem 1.2). Without this assumption, we have an example with a polynomial rate of convergence whose power can be increased or decreased as much as we want.

Example 1. Let us consider n = 1, $H(x,p) = \left(\frac{1+|x|}{m}\right)K(p)$ for m > 1 and $K : \mathbb{R} \to \mathbb{R}$ defined by

$$K(\mathbf{p}) = \begin{cases} -|\mathbf{p}| & \text{for } |\mathbf{p}| \leq 1, \\ |\mathbf{p}| - 2 & \text{for } |\mathbf{p}| \geq 1. \end{cases}$$
(7.1)

The unique viscosity solution to (3.1) is

$$u_k(x) = \frac{(1+|x|)^m}{m(1+k)^{m-1}} \qquad \textit{for } x \in [-k,k].$$

Clearly, $\mathbf{u}_{\mathbf{k}}(\mathbf{x}) \to \mathbf{0}$ locally uniformly with rate $\mathcal{O}\left(\frac{1}{\mathbf{k}^{m-1}}\right)$ for any given $\mathbf{m} > 1$. We should note that the limit $\mathbf{0}$ is not a unique solution to (3.3). Another solution to (3.3) is $\mathbf{u}(\mathbf{x}) = \mathbf{m}^{-1}(1+|\mathbf{x}|)^{\mathbf{m}}$, but it does not belong to BUC (\mathbb{R}).

Example 2. Assume n = 1, H(x, p) = K(p) + V(x) where $V(x) = e^{-|x|}$ and $K : \mathbb{R} \to \mathbb{R}$ defined by

$$\mathsf{K}(\mathsf{p}) = \begin{cases} -|\mathsf{p}| & \text{for } |\mathsf{p}| \leq 1, \\ |\mathsf{p}| - 2 & \text{for } |\mathsf{p}| \geq 1. \end{cases}$$

The unique state-constraint viscosity solution to (3.1) is

$$u_k(x) = -\frac{1}{2}e^{-|x|} + \left(e^{-k} - \frac{1}{2}e^{-2k}\right)e^{|x|}, \qquad x \in [-k,k],$$

and the unique viscosity solution to (3.3) is

$$\mathfrak{u}(\mathbf{x}) = -\frac{1}{2}e^{-|\mathbf{x}|}, \qquad \mathbf{x} \in \mathbb{R}.$$

We have $u_k \to u$ locally uniformly in \mathbb{R} with rate $O(e^{-k})$.

Secondly, prototype condition (P1) can be relaxed as follows.

Remark 10. It is clear from the proofs of our main results (Theorems 1.1, 1.2, 1.4, and Corollary 1.3) that prototype condition (P1) can be relaxed as following.

(P1') Assume Ω_k is bounded, $B(0, k) \subset \Omega_k$, and the comparison principle for the stateconstraint problem holds on Ω_k for all $k \in \mathbb{N}$. Of course, $\Omega = \bigcup_{k \in \mathbb{N}} \Omega_k = \mathbb{R}^n$ here.

Thirdly, there are some open questions we are not able to answer yet.

Question 1. In the first prototype (P1) case, what is the optimal rate of convergence of u_k to u in the general nonconvex setting?

A more specific question is as follows.

Question 2. Assume (P1), and H(x,p) = K(p) + V(x), where $K \in Lip(\mathbb{R}^n)$ is coercive and nonconvex, and $V \in BUC(\mathbb{R}^n)$. Is it true that we always have an exponential rate of convergence of u_k to u?

8. Appendix

8.1. Proofs of some lemmas.

Proof of Lemma 5.6. We first prove the result for all (x, v) with $|v| \leq 1$, then by scaling we get the result for all (x, v). Using (L7) we have $-B \leq L(x, v) \leq 4A + B$ for all (x, v) with $|v| \leq 2$. For $u, v \in B(0, 1)$ with $u \neq v$, let $w = v + |v - u|^{-1}(v - u)$. Then, |w| < |v| + 1 < 2. Let $\lambda = (1 + |u - v|)^{-1} \in (0, 1)$, we have $v = \lambda u + (1 - \lambda)w$. By the convexity, one obtains

$$L(x,\nu) - L(x,u) \leq (1-\lambda) (L(x,w) - L(x,u)) \leq (4A + 2B)|u - \nu|.$$

By symmetry, we deduce that $|L(x, u) - L(x, v)| \leq (4A + 2B)|u - v|$ for all (x, v) with $|v| \leq 1$. In other words, we have that $|\xi| \leq 4A + 2B$ whenever $\xi \in D_v^-L(x, v)$ for $(x, v) \in \mathbb{R}^n \times \overline{B(0, 1)}$. Now for r > 1, we define $L_r(x, v) = r^{-2}L(x, rv)$ for $(x, v) \in \mathbb{R}^n \times \mathbb{R}^n$. We observe that

$$|A^{-1}|v|^2 - B \leq |A^{-1}|v|^2 - Br^{-2} \leq L_r(x,v) \leq |A|v|^2 + Br^{-2} \leq |A|v|^2 + Br^{-2}$$

for all (x, v). For $v \in \mathbb{R}^n$ with $|v| \ge 1$, let r = 2|v| > 1 and $u = \frac{v}{2|v|} \in B(0, 1)$ so that v = ru. Since $\xi \in D_v^-L(x, v)$ implies $\frac{\xi}{r} \in D_u^-L_r(x, u)$, we have $|\xi| \le (4A + 2B)|r| = (8A + 4B)|v|$. \Box

Proof of Theorem 5.7. Let $\{\eta_k\}_{k=1}^{\infty} \subset \mathcal{A}_x$ be a minimizing sequence in $AC([0, \infty))$ such that $\lim_{k\to\infty} J[x,\eta_k] = u(x)$. From the uniform boundedness of u and the quadratic bounds of $L(x,\nu)$, we have

$$\left\|e^{-\frac{3}{2}}\dot{\eta}_{k}(s)\right\|_{L^{2}((0,\infty);\mathbb{R}^{n})} \leqslant C_{4}.$$

Here, C_4 can be chosen as $(A(2C_H + B))^{\frac{1}{2}}$. By the weak compactness of L^2 , there exists g such that $e^{-\frac{s}{2}}g(s) \in L^2((0,\infty);\mathbb{R}^n)$ and a subsequence $\{k_j\} \to \infty$ such that $e^{-\frac{s}{2}}\dot{\eta}_{k_j} \rightharpoonup e^{-\frac{s}{2}}g$ weakly in $L^2((0,\infty);\mathbb{R}^n)$ as $j \to \infty$.

Writing g as $e^{\frac{s}{2}} g \cdot e^{-\frac{s}{2}}$ and using the Cauchy-Schwartz inequality, we get $g \in L^1_{loc}((0,\infty); \mathbb{R}^n)$. For t > 0, we let $\eta(t) = x + \int_0^t g(s) \, ds$. Clearly, $\eta \in \mathcal{A}_x$ and one obtains that $\eta_{k_j} \to \eta$ pointwise with $\dot{\eta} = g$ almost everywhere. On the other hand, the convexity of L implies

$$L\left(\eta_{k_{j}}(s),-\dot{\eta}_{k_{j}}(s)\right) \geqslant L\left(\eta_{k_{j}}(s),-\dot{\eta}(s)\right) - D_{\nu}L\left(\eta_{k_{j}}(s),-\dot{\eta}(s)\right) \cdot \left(\dot{\eta}_{k_{j}}(s)-\dot{\eta}(s)\right).$$

Therefore,

$$\begin{split} \int_0^\infty e^{-s} L\left(\eta_{k_j}(s), -\dot{\eta}_{k_j}(s)\right) ds & \ge \int_0^\infty e^{-s} L\left(\eta_{k_j}(s), -\dot{\eta}(s)\right) ds \\ & + \int_0^\infty e^{-s/2} D_\nu L\left(\eta_{k_j}(s), -\dot{\eta}(s)\right) \cdot e^{-s/2} \left(\dot{\eta}_{k_j}(s) - \dot{\eta}(s)\right) ds. \end{split}$$

Since $|D_{\nu}L(\eta_{k_j}(s), -\dot{\eta}(s))| \leq C_L(1+|\dot{\eta}(s)|)$ for a.e. $s \in (0, \infty)$, it is clear that

$$e^{-s/2} D_{\nu} L\left(\eta_{k_j}(s), -\dot{\eta}(s)\right) \to e^{-s/2} D_{\nu} L\left(\eta(s), -\dot{\eta}(s)\right)$$

in $L^2((0,\infty);\mathbb{R}^n)$ and thus

$$\int_0^\infty e^{-s} \mathsf{D}_{\nu} \mathsf{L}\left(\eta_{k_j}(s), -\dot{\eta}(s)\right) \cdot \left(\dot{\eta}_{k_j}(s) - \dot{\eta}(s)\right) \, \mathrm{d}s$$

converges to 0 as k goes to infinity, which yields that $J[x, \eta] \leq u(x)$. Hence $J[x, \eta] = u(x)$. \Box *Proof of Lemma 5.8.* By the definition of u in (5.1), we have

$$\mathfrak{u}\big(\eta(t)\big) \leqslant \int_0^\infty e^{-s} L\big(\gamma(s), -\dot{\gamma}(s)\big) \, ds = e^t \int_t^\infty e^{-\xi} L\left(\eta(\xi), -\dot{\eta}(\xi)\right) \, d\xi,$$

where $\gamma(s) = \eta(t + s)$ for $s \ge 0$. Thus,

$$e^{-t}\mathfrak{u}(\eta(t)) \leqslant \int_{t}^{\infty} e^{-\xi} L(\eta(\xi), -\dot{\eta}(\xi)) d\xi.$$
(8.1)

By the dynamic programming principle and (8.1), we have

$$\mathfrak{u}(\mathfrak{\eta}(0)) \leqslant \int_0^t e^{-s} L(\mathfrak{\eta}(s), -\dot{\mathfrak{\eta}}(s)) \, \mathrm{d}s + e^{-t} \mathfrak{u}(\mathfrak{\eta}(t)) \leqslant \int_0^\infty e^{-s} L(\mathfrak{\eta}(s), -\dot{\mathfrak{\eta}}(s)) \, \mathrm{d}s = \mathfrak{u}(\mathfrak{\eta}(0)).$$

herefore, (5.5), (5.6) and (5.7) follow.

Therefore, (5.5), (5.6) and (5.7) follow.

Proof of Lemma 5.9. For every t, h > 0, by Lemma 5.8 we have that

$$\frac{e^{-t}u(\eta(t))-e^{-(t+h)}u(\eta(t+h))}{h}=\frac{1}{h}\int_{t}^{t+h}e^{-s}L(\eta(s),-\dot{\eta}(s))\ ds.$$

Let $\varphi \in C^1(\mathbb{R})$ such that $u - \varphi$ has a local min at $\eta(t)$ and $u(\eta(t)) = \varphi(\eta(t))$, then

$$\frac{e^{-t}u\big(\eta(t)\big)-e^{-(t+h)}u\big(\eta(t+h)\big)}{h}\leqslant \frac{e^{-t}\phi\big(\eta(t)\big)-e^{-(t+h)}\phi\big(\eta(t+h)\big)}{h}.$$

Therefore,

$$\frac{1}{h}\int_{t}^{t+h} e^{-s}L\Big(\eta(s),-\dot{\eta}(s)\Big) \ ds \leqslant \frac{e^{-t}\phi\big(\eta(t)\big)-e^{-(t+h)}\phi\big(\eta(t+h)\big)}{h}.$$

Since $\eta(t)$ is differentiable a.e. in $(0, \infty)$, at those t where $\eta(t)$ is differentiable, let $h \to 0^+$ we deduce that

$$e^{-t}L\big(\eta(t),-\dot{\eta}(t)\big) \leqslant -\frac{d}{dt}\Big(e^{-t}\phi\big(\eta(t)\big)\Big) = e^{-t}\phi\big(\eta(t)\big) - e^{-t}D\phi\big(\eta(t)\big) \cdot \dot{\eta}(t).$$

Thus, for a.e. t > 0 where η is differentiable, we have

 $L(\eta(t), -\dot{\eta}(t)) \leq \varphi(\eta(t)) - D\varphi(\eta(t)) \cdot \dot{\eta}(t).$

By (L7) and a priori estimate (5.3) for a.e. $t \in (0, \infty)$ we have that

$$A^{-1}|\dot{\eta}(t)|^2 - B \leqslant \phi\big(\eta(t)\big) - D\phi\big(\eta(t)\big) \cdot \dot{\eta}(t) \leqslant C_H + C_H |\dot{\eta}(t)|.$$

This shows that $|\dot{\eta}(t)| \leq C_5$ for a.e. $t \in (0, \infty)$, and C_5 only depends on C_H, A, B .

It is worth emphasizing again here that the bound C_5 on the velocity of η only depends on C_H , A, B, which can be seen clearly from the last chain of inequalities in the above proof. In fact, one can choose explicitly that $C_5 = (2AB + 2AC_H + A^2C_H^2)^{1/2}$.

8.2. Existence of minimizers in the general case. We show that one can remove the smoothness of L in Theorem 5.7 under the assumption (L3).

Let us consider mollifiers in \mathbb{R}^{2n} defined as $\{\eta_{\varepsilon}\}_{\varepsilon>0}$ such that $\eta_{\varepsilon}(x) = \frac{1}{\varepsilon^{2n}}\eta\left(\frac{x}{\varepsilon}\right)$ for $x \in \mathbb{R}^{2n}$ where $\eta \in C_{c}^{\infty}(\mathbb{R}^{2n})$ satisfying $0 \leq \eta \leq 1$, supp $(\eta) \subset B_{\mathbb{R}^{2n}}(0,1)$ and $\int_{\mathbb{R}^{2n}} \eta(x) dx = 1$.

For each $\varepsilon > 0$ we define the convolution $L^{\varepsilon} = \eta_{\varepsilon} * L \in C^{\infty}(\mathbb{R}^n \times \mathbb{R}^n)$. It is easy to see that L^{ε} is bounded below, (L5), (L6) are preserved to L^{ε} and (L7) now becomes:

(L7^{ε}) There exist positive constants $A_{\varepsilon}, B_{\varepsilon}$ such that $A_{\varepsilon}^{-1} - B_{\varepsilon}^{-1} \leq L^{\varepsilon}(x, \nu) \leq A_{\varepsilon}|\nu|^2 + B_{\varepsilon}$ for all $(x, \nu) \in \mathbb{R}^n \times \mathbb{R}^n$, and $A_{\varepsilon} \to A, B_{\varepsilon} \to B$ as $\varepsilon \to 0$.

By Theorem 5.7, there exists a minimizer γ_{ε} in \mathcal{A}_{χ} such that

$$\mathfrak{u}^{\varepsilon}(\mathfrak{x}) := \inf_{\zeta \in \mathcal{A}_{\mathfrak{x}}} \int_{0}^{\infty} e^{-s} L^{\varepsilon} \big(\zeta(s), -\dot{\zeta}(s) \big) \, \mathrm{d}s = \int_{0}^{\infty} e^{-s} L^{\varepsilon} \big(\gamma_{\varepsilon}(s), -\dot{\gamma}_{\varepsilon}(s) \big) \, \mathrm{d}s.$$

Let H^{ε} be the Legendre transform of L^{ε} . Then, we can show that u^{ε} is the unique solution to $u^{\varepsilon}(x) + H^{\varepsilon}(x, Du^{\varepsilon}(x)) = 0$ in \mathbb{R}^n . It is easy to see that $H^{\varepsilon} \to H$ locally uniformly in $\mathbb{R}^n \times \mathbb{R}^n$, therefore by stability of viscosity solutions, $u^{\varepsilon} \to u$ locally uniformly in \mathbb{R}^n as $\varepsilon \to 0$.

We indeed have that γ^{ε} is smooth according to Remark 5. Furthermore, Theorem 5.7 yields that $\|e^{-\frac{s}{2}}\dot{\gamma}_{\varepsilon}(s)\|_{L^{2}} \leq C$ and $|\dot{\gamma}_{\varepsilon}| \leq C$ pointwise in $(0,\infty)$. Therefore, we can define $\gamma \in \mathcal{A}_{x}$ such that (up to subsequence) $\gamma_{\varepsilon} \to \gamma$ locally uniformly on $[0,\infty)$ and $e^{-\frac{s}{2}}\dot{\gamma}_{\varepsilon} \rightharpoonup e^{-\frac{s}{2}}\dot{\gamma}$ weakly in L². Since $L^{\varepsilon} \to L$ uniformly on a compact set and $\{-\dot{\gamma}_{\varepsilon}(s)\}_{\varepsilon>0}$ is bounded, we obtain that

$$L^{\varepsilon}(\gamma_{\varepsilon}(s), -\dot{\gamma}_{\varepsilon}(s)) = L(\gamma(s), -\dot{\gamma}_{\varepsilon}(s)) + \tilde{\omega}_{C_{5}}(\varepsilon) + \tilde{\omega}_{C_{5}}(|\gamma_{\varepsilon}(s) - \gamma(s)|)$$

using (L3). Therefore, it suffices to show that

$$\int_{0}^{\infty} e^{-s} \mathsf{L}(\gamma(s), -\dot{\gamma}(s)) \, \mathrm{d}s \leqslant \liminf_{\varepsilon \to 0} \int_{0}^{\infty} e^{-s} \mathsf{L}(\gamma(s), -\dot{\gamma}_{\varepsilon}(s)) \, \mathrm{d}s. \tag{8.2}$$

For simplicity, let $d\mu = e^{-s} ds$ be a probability measure on $[0, \infty)$. It is easy to see that the functional $I : L^2(\mu) \to \mathbb{R}$ maps $f \mapsto \int_0^\infty L(\gamma(s), f(s)) d\mu(s)$ is convex and lower semicontinuous, thus it is also weakly lower semicontinuous. Now since $\dot{\gamma}_{\varepsilon} \rightharpoonup \dot{\gamma}$ weakly in $L^2(d\mu)$, we obtain (8.2) and thus γ is a minimizer for u(x).

Remark 11. Inequality (8.2) for the Cauchy problem (finite time horizon) is proved using a different argument by H. Ishii in [16] under more general assumptions. Such inequalities are crucial for the analysis of large time behavior of solutions to the time-dependent problems.

References

- ALTAROVICI, A., BOKANOWSKI, O., AND ZIDANI, H. A general Hamilton–Jacobi framework for nonlinear state-constrained control problems. *ESAIM: COCV 19*, 2 (2013), 337–357.
- [2] ALVAREZ, O., LASRY, J.-M., AND LIONS, P.-L. Convex viscosity solutions and state constraints. Journal de Mathématiques Pures et Appliquées 76, 3 (1997), 265 – 288.
- [3] ARMSTRONG, S. N., AND TRAN, H. V. Viscosity solutions of general viscous Hamilton-Jacobi equations. *Mathematische Annalen 361*, 3 (Apr 2015), 647–687.
- [4] BARDI, M., AND CAPUZZO-DOLCETTA, I. Optimal Control and Viscosity Solutions of Hamilton-Jacobi-Bellman Equations. Modern Birkhäuser Classics. Birkhäuser Basel, 1997.
- [5] BARLES, G. Solutions de viscosité des équations de Hamilton-Jacobi. Mathématiques et Applications. Springer-Verlag Berlin Heidelberg, 1994.
- [6] BOKANOWSKI, O., FORCADEL, N., AND ZIDANI, H. Deterministic state-constrained optimal control problems without controllability assumptions. ESAIM: COCV 17, 4 (2011), 995–1015.
- [7] CAMILLI, F., AND FALCONE, M. Approximation of optimal control problems with state constraints: Estimates and applications. In *Nonsmooth Analysis and Geometric Methods in Deterministic Optimal Control* (New York, NY, 1996), B. S. Mordukhovich and H. J. Sussmann, Eds., Springer New York, pp. 23–57.
- [8] CANNARSA, P., GOZZI, F., AND SONER, H. M. A boundary-value problem for Hamilton–Jacobi equations in Hilbert spaces. Applied Mathematics and Optimization 24, 1 (Jul 1991), 197–220.
- [9] CANNARSA, P., AND SINESTRARI, C. Semiconcave Functions, Hamilton-Jacobi Equations, and Optimal Control. Birkhäuser Boston, Boston, MA, 2004.
- [10] CAPUZZO-DOLCETTA, I., AND LIONS, P.-L. Hamilton-Jacobi equations with state constraints. Transactions of the American Mathematical Society 318, 2 (1990), 643–683.
- [11] CARLIER, G., AND TAHRAOUI, R. Hamilton–Jacobi–Bellman equations for the optimal control of a state equation with memory. *ESAIM: COCV 16*, 3 (2010), 744–763.
- [12] CLARKE, F. H., RIFFORD, L., AND STERN, R. J. Feedback in state constrained optimal control. ESAIM: COCV 7 (2002), 97–133.
- [13] FELDMAN, M., AND MCCUAN, J. Constructing convex solutions via Perron's method. Annali Dell'Universita' Di Ferrara 53, 1 (May 2007), 65–94.
- [14] FRANKOWSKA, H., AND MAZZOLA, M. Discontinuous solutions of Hamilton-Jacobi-Bellman equation under state constraints. *Calculus of Variations and Partial Differential Equations* 46, 3 (Mar 2013), 725–747.
- [15] ISHII, H. Perron's method for Hamilton–Jacobi equations. Duke Math. J. 55, 2 (06 1987), 369–384.
- [16] ISHII, H. Asymptotic solutions for large time of Hamilton–Jacobi equations in Euclidean n space. Ann. Inst. H. Poincaré Anal. Non Linéaire 25, 2 (2008), 231–266.
- [17] ISHII, H., AND KOIKE, S. A new formulation of state constraint problems for first-order PDEs. SIAM Journal on Control and Optimization 34, 2 (1996), 554–571.
- [18] ISHII, H., MITAKE, H., AND TRAN, H. V. The vanishing discount problem and viscosity mather measures. part 2: Boundary value problems. *Journal de Mathématiques Pures et Appliquées 108*, 3 (2017), 261 – 305.
- [19] KIM, K.-E. Hamilton-Jacobi-Bellman equation under states constraints. Journal of Mathematical Analysis and Applications 244, 1 (2000), 57 – 76.
- [20] KOCAN, M., AND SORAVIA, P. A viscosity approach to infinite-dimensional Hamilton-Jacobi equations arising in optimal control with state constraints. SIAM Journal on Control and Optimization 36, 4 (1998), 1348–1375.
- [21] LASRY, J. M., AND LIONS, P.-L. Nonlinear elliptic equations with singular boundary conditions and stochastic control with state constraints. I. the model problem. *Mathematische Annalen 283*, 4 (1989), 583–630.
- [22] LE, N. Q., MITAKE, H., AND TRAN, H. V. Dynamical and Geometric Aspects of Hamilton-Jacobi and Linearized Monge-Ampère Equations, vol. 2183 of Lecture Notes in Mathematics. Springer International Publishing, 2017.

- [23] LEONORI, T., AND PORRETTA, A. The boundary behavior of blow-up solutions related to a stochastic control problem with state constraint. *SIAM Journal on Mathematical Analysis 39*, 4 (2008), 1295–1327.
- [24] MITAKE, H. Asymptotic solutions of Hamilton-Jacobi equations with state constraints. Applied Mathematics and Optimization 58, 3 (Dec 2008), 393–410.
- [25] SANDHOLM, W. H., TRAN, H. V., AND ARIGAPUDI, S. Hamilton–Jacobi–Bellman equations with semilinear costs and state constraints, with applications to large deviations. *preprint* (July 2018).
- [26] SEDRAKYAN, H. Stability of solutions to Hamilton–Jacobi equations under state constraints. Journal of Optimization Theory and Applications 168, 1 (Jan 2016), 63–91.
- [27] SONER, H. Optimal control with state-space constraint I. SIAM Journal on Control and Optimization 24, 3 (1986), 552–561.
- [28] SONER, H. Optimal control with state-space constraint. II. SIAM Journal on Control and Optimization 24, 6 (1986), 1110–1122.

(Y. Kim) DEPARTMENT OF MATHEMATICS, UNIVERSITY OF WISCONSIN MADISON, 480 LINCOLN DRIVE, MADISON, WI 53706, USA

Email address: yeonkim@math.wisc.edu

(H. V. Tran) DEPARTMENT OF MATHEMATICS, UNIVERSITY OF WISCONSIN MADISON, 480 LINCOLN DRIVE, MADISON, WI 53706, USA

Email address: hung@math.wisc.edu

(S. N.T. Tu) DEPARTMENT OF MATHEMATICS, UNIVERSITY OF WISCONSIN MADISON, 480 LINCOLN DRIVE, MADISON, WI 53706, USA

Email address: thaison@math.wisc.edu