# Uncertainty product for Vilenkin groups 

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#### Abstract

We study a localization of functions defined on Vilenkin groups. To measure the localization we introduce two uncertainty products $U P_{\lambda}$ and $U P_{G}$ that are similar to the Heisenberg uncertainty product. $U P_{\lambda}$ and $U P_{G}$ differ from each other by the metric used for the Vilenkin group $G$. We discuss analogs of a quantitative uncertainty principle. Representations for $U P_{\lambda}$ and $U P_{G}$ in terms of Walsh and Haar basis are given.


Keywords Vilenkin group; uncertainty product; Haar wavelet; modified Gibbs derivative; generalized Walsh function.

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## 1 Introduction

An uncertainty product for a function characterizes how concentrated is the function in time and frequency domain. Initially the notion of uncertainty product was introduced for $f \in L_{2}(\mathbb{R})$ by W. Heisenberg [6] and E. Schrödinger [12]. Later on extensions of this notion appeared for various algebraic and topological structures. For periodic functions, it was suggested by E. Breitenberger [1]. For some particular cases of locally compact groups (namely a euclidean motion groups, non-compact semisimple Lie groups, Heisenberg groups) the counterpart was derived in [11]. Uncertainty products on compact Riemannian manifolds was discussed in [4]. In [8], this concept was introduced for functions defined on the Cantor group. In this paper, we discuss localization of functions defined on Vilenkin groups.

To measure the localization we introduce a functional that is similar to the Heisenberg uncertainty product (see Definition 11). It depends on the metric used for the Vilenkin group $G$. Two equivalent metrics are in common use for the group $G$. So we discuss two uncertainty

[^0]products $U P_{\lambda}$ and $U P_{G}$. The first one is a strict counterpart of "dyadic uncertainty constant" introduced in 8] (see Theorems 11 and 2). Usage of another metric in the second uncertainty product allows for exploitation of a modified Gibbs derivative that plays a role of usual derivative for the Heisenberg uncertainty product. At the same time it turns out that usage of Haar basis is a good approach for evaluation of $U P_{G}$ (see Theorem 3). In particular, it allows for an estimate of Fourier-Haar coefficients for functions defined on the Vilenkin group (see Corollary 2). The connection between $U P_{\lambda}$ and $U P_{G}$ is showed in Lemma 1.

## 2 Auxiliary results

We recall necessary facts about the Vilenkin group. More details can be found in [3, 13. The Vilenkin group $G=G_{p}, p \in \mathbb{N}, p \neq 1$, is a set of the sequences

$$
x=\left(x_{j}\right)=\left(\ldots, 0,0, x_{-k}, x_{-k+1}, x_{-k+2}, \ldots\right),
$$

where $x_{j} \in\{0, \ldots, p-1\}$ for $j \in \mathbb{Z}$. The operation on $G$ is denoted by $\oplus$ and defined as the coordinatewise addition modulo $p$ :

$$
\left(z_{j}\right)=\left(x_{j}\right) \oplus\left(y_{j}\right) \Longleftrightarrow z_{j}=x_{j}+y_{j}(\bmod p) \quad \text { for } \quad j \in \mathbb{Z}
$$

The inverse operation of $\oplus$ is denoted by $\ominus$. The symbol $\ominus x$ denotes the inverse element of $x \in G$. The sequence $\mathbf{0}=(\ldots, 0,0, \ldots)$ is a neutral element of $G$. If $x \neq \mathbf{0}$, then there exists a unique number $N=N(x)$ such that $x_{N} \neq 0$ and $x_{j}=0$ for $j<N$. The Vilenkin group $G_{p}$, where $p=2$ is called the Cantor group. In this case the inverse operation $\ominus$ coicides with the group operation $\oplus$.

Define a map $\lambda: G \rightarrow[0,+\infty)$

$$
\lambda(x)=\sum_{j \in \mathbb{Z}} x_{j} p^{-j-1}, \quad x=\left(x_{j}\right) \in G
$$

The mapping $x \mapsto \lambda(x)$ is a bijection taking $G \backslash \mathbb{Q}_{0}$ onto $[0, \infty)$, where $\mathbb{Q}_{0}$ is a set of all elements terminating with $p-1$ 's.

Two equivalent metrics are in common use for the group $G$. One metric is defined by $d_{1}(x, y):=\lambda(x \ominus y)$ for $x, y \in G$. To define another one $d_{2}$ we consider a map $\|\cdot\|_{G}: G \rightarrow$ $[0, \infty)$, where $\|\mathbf{0}\|_{G}:=0$ and $\|x\|_{G}:=p^{-N(x)}$ for $x \neq \mathbf{0}$. Then $d_{2}(x, y):=\|x \ominus y\|_{G}, x, y \in G$. Given $n \in \mathbb{Z}$ and $x \in G$, denote by $I_{n}(x)$ the ball of radius $2^{-n}$ with the center at $x$, i.e.

$$
I_{n}(x)=\left\{y \in G: d(x, y)<2^{-n}\right\}
$$

For brevity we set $I_{j}:=I_{j}(\mathbf{0})$ and $I:=I_{0}$.
We denote dilation on $G$ by $D: G \rightarrow G$, and set $(D x)_{k}=x_{k+1}$ for $x \in G$. Then $D^{-1}: G \rightarrow G$ is the inverse mapping $\left(D^{-1} x\right)_{k}=x_{k-1}$. Set $D^{k}=D \circ \cdots \circ D(k$ times $)$ if $k>0$, and $D^{k}=D^{-1} \circ \cdots \circ D^{-1}(-k$ times $)$ if $k<0 ; D^{0}$ is the identity mapping.

We deal with functions taking $G$ to $\mathbb{C}$. Denote $\mathbb{1}_{E}$ the characteristic function of a set $E \subset G$. Given a function $f: G \rightarrow \mathbb{C}$ and a number $h \geq 0$, for every $x \in G$ we define $f_{0, h}(x)=f\left(x \oplus \lambda^{-1}(h)\right)$. Finally, we set for $j \in \mathbb{Z}$

$$
f_{j, h}(x)=p^{j / 2} f_{0, h}\left(D^{j} x\right), \quad x \in G .
$$

The functional spaces $L_{q}(G)$ and $L_{q}(E)$, where $E$ is a measurable subset of $G$, are derived using the Haar measure (see [7]).

Given $\xi \in G$, a group character of $G$ is defined by

$$
\chi_{\xi}(x)=\chi(x, \xi):=\exp \left(\frac{2 \pi i}{p} \sum_{j \in \mathbb{Z}} x_{j} \xi_{-1-j}\right) .
$$

The functions $\mathrm{w}_{n}(x):=\chi\left(\lambda^{-1}(n), x\right)$ are called the generalized Walsh functions. If $p=2$, than $\mathrm{w}_{n}$ are called the Walsh functions.

The Fourier transform of a function $f \in L^{1}(G)$ is defined by

$$
\begin{equation*}
F f(\omega)=\int_{G} f(x) \overline{\chi(x, \omega)} d \mu(x), \quad \omega \in G \tag{1}
\end{equation*}
$$

The Fourier transform is extended to $L_{2}(G)$ in a standard way, and the Plancherel equality takes place

$$
\langle f, g\rangle:=\int_{G} f(x) \overline{g(x)} d x=\int_{G} F f(\xi) \overline{F g(\xi)} d \xi=\langle F f, F g\rangle, \quad f, g \in L_{2}(G)
$$

The inversion formula is valid for any $f \in L_{2}(G)$

$$
F^{-1} F f(x)=\int_{G} F f(\omega) \chi(x, \omega) d \mu(\omega)=f(x) .
$$

It is straightforward to see that

$$
\begin{equation*}
F\left(f_{j, n}\right)(\xi)=p^{-j / 2} \chi\left(k, D^{-j} \xi\right) F f\left(D^{-j} \xi\right), \quad n \in \mathbb{Z}_{+}, j \in \mathbb{Z} \tag{2}
\end{equation*}
$$

The discrete Vilenkin-Chrestenson transform of a vector $x=\left(x_{k}\right)_{k=0, p^{n}-1} \in \mathbb{C}^{p^{n}}$ is a vector $y=\left(y_{k}\right)_{k=\overline{0, p^{n}-1}} \in \mathbb{C}^{p^{n}}$, where

$$
\begin{equation*}
y_{k}=p^{-n} \sum_{s=0}^{p^{n}-1} x_{s} \mathrm{w}_{k}\left(\lambda^{-1}\left(s / p^{n}\right)\right), \quad 0 \leq k \leq p^{n}-1 \tag{3}
\end{equation*}
$$

The inverse transform is

$$
\begin{equation*}
x_{k}=\sum_{s=0}^{p^{n}-1} y_{s} \overline{\mathrm{w}_{k}\left(\lambda^{-1}\left(s / p^{n}\right)\right)} . \quad 0 \leq k \leq p^{n}-1 \tag{4}
\end{equation*}
$$

Given $f: G_{2} \rightarrow \mathbb{C}$, the function

$$
f^{[1]}(x):=\lim _{n \rightarrow \infty} \sum_{j=-n}^{n} 2^{j-1}\left(f(x)-f_{0,2^{-j-1}}(x)\right)
$$

is called the Gibbs derivative of a function $f$. The following properties hold true

$$
\begin{equation*}
F f^{[1]}(\xi)=\lambda(\xi) F f(\xi), \quad \mathrm{w}_{n}^{[1]}(x)=n \mathrm{w}_{n}(x) \tag{5}
\end{equation*}
$$

Set $\varphi=\mathbb{1}_{I}$. The Haar functions $\psi^{\nu}, \nu=1, \ldots, p-1$ are defined by

$$
\begin{equation*}
\psi^{\nu}(x)=\sum_{n=0}^{p-1} \exp \left(\frac{2 \pi i \nu n}{p}\right) \varphi\left(D x \oplus \lambda^{-1}(n)\right) . \tag{6}
\end{equation*}
$$

The system $\psi_{j, k}^{\nu}, \nu=1, \ldots, p-1, j \in \mathbb{Z}, k \in \mathbb{Z}_{+}$, forms an orthonormal basis (Haar basis) for $L_{2}(G)$, see [5, 9].

It follows from (11) that $F \varphi=\varphi=\mathbb{1}_{I}$ and $F \psi=\mathbb{1}_{I_{0} \oplus \lambda^{-1}(p-\nu)}$. Taking into account (2), we get

$$
\begin{equation*}
F \psi_{j, k}^{\nu}(\xi)=p^{-j / 2} \chi\left(k, D^{-j} \xi\right) \mathbb{1}_{I_{-j} \oplus \lambda^{-1}\left((p-\nu) p^{j}\right)} . \tag{7}
\end{equation*}
$$

Given $f \in L_{1}(G)$, the modified Gibbs derivative $\mathcal{D}$ is defined by

$$
\begin{equation*}
F \mathcal{D} f=\|\cdot\|_{G} F f \tag{8}
\end{equation*}
$$

It was introduced in [2] for $L_{1}\left(G_{2}\right)$. Such kind of operators are often called pseudo-differential.
Proposition 1. Suppose $g, F g,\|\cdot\|_{G} F g$ are locally integrable on $G, j \in \mathbb{Z}$. Then the assertion $\operatorname{supp} \widehat{g} \subset I_{-j-1} \backslash I_{-j}$ is necessary and sufficient for $g$ to be an eigenfunction of $\mathcal{D}$ corresponding to the eigenvalue $p^{j}$.

The proof can be rewritten from Proposition 1 [10], where it is proved for the Cantor group.

Corollary 1. Any Haar function $\psi_{j, k}^{\nu}$ is an eigenfunction of $\mathcal{D}^{\alpha}$ corresponding to the eigenvalue $p^{j}$.

Proof. The statement follows from Proposition 1 and (7).

## 3 Uncertainty product and metrics

Originally, the concept of an uncertainty product was introduced for the real line case in 1927. The Heisenberg uncertainty product of $f \in L_{2}(\mathbb{R})$ is the functional $U C_{H}(f):=\Delta_{f} \Delta_{\hat{f}}$ such that

$$
\Delta_{f}^{2}:=\|f\|_{L^{2}(\mathbb{R})}^{-2} \int_{\mathbb{R}}\left(x-x_{f}\right)^{2}|f(x)|^{2} d x, \quad \Delta_{\widehat{f}}^{2}:=\|\widehat{f}\|_{L^{2}(\mathbb{R})}^{-2} \int_{\mathbb{R}}\left(t-t_{\widehat{f}}\right)^{2}|\widehat{f}(t)|^{2} d t
$$

$$
x_{f}:=\|f\|_{L^{2}(\mathbb{R})}^{-2} \int_{\mathbb{R}} x|f(x)|^{2} d x, \quad t_{\widehat{f}}:=\left.\|\widehat{f}\|_{L^{2}(\mathbb{R})}^{-2} \int_{\mathbb{R}}| | \widehat{f}(t)\right|^{2} d t,
$$

where $\widehat{f}$ denotes the Fourier transform of $f \in L_{2}(\mathbb{R})$. It is well known that $U C_{H}(f) \geq 1 / 2$ for a function $f \in L_{2}(\mathbb{R})$ and the minimum is attained on the Gaussian. To motivate the definition of a localization characteristic for the Vilenkin group we note that on one hand $x_{f}$ is the solution of the minimization problem

$$
\min _{\tilde{x}} \int_{\mathbb{R}}(x-\tilde{x})^{2}|f(x)|^{2} d x
$$

and on another hand the sense of the sign "-" in the definition of $\Delta_{f}$ is the distance between $x$ and $x_{f}$. So we come to the main definition.

Definition 1. Suppose $f: G \rightarrow \mathbb{C}, f \in L_{2}(G)$, and $d$ is a metric on $G$, then a functional

$$
\begin{gathered}
U P(f):=V(f) V(F f), \quad \text { where } \\
V(f):=\frac{1}{\|f\|_{L_{2}(G)}^{2}} \min _{\widetilde{x}} \int_{G}(d(x, \tilde{x}))^{2}|f(x)|^{2} d x
\end{gathered}
$$

is called the uncertainty product of a function $f$ defined on the Vilenkin group.
Thus, we study two uncertainty products $U P_{\lambda}$ and $U P_{G}$ that corresponds to the metric $d_{1}(x, y):=\lambda(x \ominus y)$ and $d_{2}(x, y):=\|x \ominus y\|_{G}$. More precisely,

$$
\begin{gathered}
U P_{\lambda}(f):=V_{\lambda}(f) V_{\lambda}(F f), \quad \text { where } \\
V_{\lambda}(f):=\frac{1}{\|f\|_{L_{2}(G)}^{2}} \min _{\tilde{x}} \int_{G}(\lambda(x \ominus \tilde{x}))^{2}|f(x)|^{2} d x .
\end{gathered}
$$

The functional $U P_{G}$ is defined as

$$
\begin{gathered}
U P_{G}(f):=V_{G}(f) V_{G}(F f), \quad \text { where } \\
V_{G}(f):=\frac{1}{\|f\|_{L_{2}(G)}^{2}} \min _{\tilde{x}} \int_{G}\|x \ominus \tilde{x}\|_{G}^{2}|f(x)|^{2} d x
\end{gathered}
$$

The functional $U P_{\lambda}$ for functions defined on the Cantor group was introduced and studied in [8]. The following results are extended from the Cantor group to the Vilenkin group without any essential changes. So we omit the proofs.

Theorem 1. Suppose $f: G \rightarrow \mathbb{C}, f \in L_{2}(G)$. Then the following inequality holds true

$$
U P_{\lambda}(f) \geq C, \text { where } C \simeq 8.5 \times 10^{-5}
$$

Theorem 2. Let $f(x)=\mathbb{1}_{\lambda^{-1}[0,1)}(x) \sum_{k=0}^{\infty} a_{k} \mathrm{w}_{k}(x)$ be a uniformly convergent series. Denote

$$
f_{n}(x)=\mathbb{1}_{\lambda^{-1}[0,1)}(x) \sum_{k=0}^{p^{n}-1} a_{k} \mathrm{w}_{k}(x)
$$

Let $V_{\lambda}(f)<+\infty, V_{\lambda}(F f)<+\infty$. Then $U P_{\lambda}(f)=\lim _{n \rightarrow \infty} V_{\lambda}\left(f_{n}\right) V_{\lambda}\left(F f_{n}\right)$, where

$$
\begin{aligned}
& V_{\lambda}\left(f_{n}\right)=\frac{\min _{k_{0}=\overline{0, p^{n}-1}} \sum_{k=0}^{p^{n}-1} p^{-n}\left|b_{\lambda\left(\lambda \lambda^{-1}(k) \oplus \lambda^{-1}\left(k_{0}\right)\right)}\right|^{2}\left((k+1)^{3}-k^{3}\right) / 3}{\sum_{k=0}^{p^{n}-1}\left|a_{k}\right|^{2}}, \\
& V_{\lambda}\left(F f_{n}\right)=\frac{\min _{k_{1}=\overline{0, p^{n}-1}} \sum_{k=0}^{p^{n}-1}\left|a_{\lambda\left(\lambda^{-1}(k) \oplus \lambda^{-1}\left(k_{1}\right)\right)}\right|^{2}\left((k+1)^{3}-k^{3}\right) / 3}{\sum_{k=0}^{p^{n}-1}\left|a_{k}\right|^{2}}
\end{aligned}
$$

and $b_{k}, 0 \leq k \leq p^{n}-1$, is the inverse discrete Vilenkin-Chrestenson transform (4).
The following Lemma shows that the functionals $U P_{\lambda}$ and $U P_{G}$ have the same order.
Lemma 1. Suppose $f \in L_{2}(G)$, then $p^{-4} U P_{G}(f) \leq U P_{\lambda}(f)<U P_{G}(f)$.
Proof. It is sufficient to note that $p^{-1}\|x\|_{G} \leq \lambda(x)<\|x\|_{G}$.
Taking into account Theorem 1, we conclude that $U P_{G}$ has a positive lower bound. So, $U P_{G}$ satisfies the uncertainty principle.
Example 1. Let us illustrate a definition of $U P_{G}$ for $p=2$ using functions $f_{1}, g_{1}, f_{2}$, and $g_{2}$ taken from [8, Example 1]. Recall $f_{1}(x)=\mathbb{1}_{\lambda^{-1}[0,1 / 4)}(x), g_{1}(x)=\mathbb{1}_{\lambda^{-1}[3 / 4,1)}(x)$, $f_{2}(x)=\mathbb{1}_{\lambda^{-1}[0,3 / 8)}(x)$, and $g_{2}(x)=\mathbb{1}_{\lambda^{-1}[3 / 4,9 / 8)}(x)$. Their Walsh-Fourier transforms are $F f_{1}=\mathbb{1}_{\lambda^{-1}[0,4)} / 4, F g_{1}=\mathrm{w}_{3}(\cdot / 4) \mathbb{1}_{\lambda^{-1}[0,4)} / 4, F f_{2}=\mathbb{1}_{\lambda^{-1}[0,4)} / 4+\mathrm{w}_{1}(\cdot / 4) \mathbb{1}_{\lambda^{-1}[0,8)} / 8$, and $F g_{2}=\mathrm{w}_{3}(\cdot / 4) \mathbb{1}_{\lambda^{-1}[0,4)} / 4+\mathrm{w}_{1}(\cdot) \mathbb{1}_{\lambda^{-1}[0,8)} / 8$. Given $\alpha \in[0, \infty)$, since the mapping $\alpha \mapsto$ $\left\|\lambda^{-1}(\alpha)\right\|_{G}$ is increasing and a measure of the set $\lambda^{-1}[a, b) \ominus \tilde{x}$ does not depend on $\tilde{x}$, it follows that

$$
\min _{\tilde{x}} \int_{\lambda^{-1}\left[0, \frac{1}{4}\right)}\|x \ominus \tilde{x}\|_{G} d x=\min _{\tilde{x}} \int_{\lambda^{-1}\left[0, \frac{1}{4}\right) \ominus \tilde{x}}\|\tau\|_{G} d \tau=\int_{\lambda^{-1}\left[0, \frac{1}{4}\right)}\|\tau\|_{G} d \tau
$$

and $\lambda^{-1}[0,1 / 4)$ is a set of minimizing $\tilde{x}$ 's as well. So, taking into account $\left\|f_{1}\right\|_{L_{2}(G)}^{2}=$ $\left\|F f_{1}\right\|_{L_{2}(G)}^{2}=1 / 4$, we get

$$
\begin{aligned}
& V_{G}\left(f_{1}\right)=\frac{1}{\left\|f_{1}\right\|_{L_{2}(G)}^{2}} \min _{\widetilde{x}} \int_{G}\|x \ominus \tilde{x}\|_{G}^{2}\left|f_{1}(x)\right|^{2} d x=4 \min _{\widetilde{x}} \int_{\lambda^{-1}\left[0, \frac{1}{4}\right)}\|x \ominus \tilde{x}\|_{G}^{2} d x \\
= & 4 \int_{\lambda^{-1}\left[0, \frac{1}{4}\right)}\|\tau\|_{G}^{2} d \tau=4 \sum_{i=2}^{\infty} \int_{\lambda^{-1}\left[\frac{1}{2^{i+1}}, \frac{1}{2^{2}}\right)}\|\tau\|_{G}^{2} d \tau=4 \sum_{i=2}^{\infty}\left(\frac{1}{2^{i}}-\frac{1}{2^{i+1}}\right) 2^{-2 i}=\frac{1}{28} .
\end{aligned}
$$

Analogously, we obtain

$$
\begin{aligned}
& V_{G}\left(F f_{1}\right)=\frac{1}{\left\|F f_{1}\right\|_{L_{2}(G)}^{2}} \min _{\widetilde{x}} \int_{G}\|x \ominus \tilde{x}\|_{G}^{2}\left|F f_{1}(x)\right|^{2} d x=\frac{1}{4} \min _{\widetilde{x}} \int_{\lambda^{-1}[0,4)}\|x \ominus \widetilde{x}\|_{G}^{2} d x \\
& =\frac{1}{4} \int_{\lambda^{-1}[0,4)}\|\tau\|_{G}^{2} d \tau=\frac{1}{4} \sum_{i=-2}^{\infty} \int_{\lambda^{-1}\left[\frac{1}{\left.2^{i+1}, \frac{1}{2^{i}}\right)}\right.}\|\tau\|_{G}^{2} d \tau=\frac{1}{4} \sum_{i=-2}^{\infty}\left(\frac{1}{2^{i}}-\frac{1}{2^{i+1}}\right) 2^{-2 i}=\frac{64}{7} .
\end{aligned}
$$

Thus, $U P_{G}\left(f_{1}\right)=16 / 49$. Using the same arguments, we calculate $U P_{G}$ for the remaining functions. We collect all the information in Table 1. Values of $U P_{\lambda}$ we extract from [8,

Example 1]. Columns named $\tilde{x}_{0}(f)$ and $\tilde{t}_{0}(f)$ contain sets of $\tilde{x}$ and $\tilde{t}$ minimizing the functionals $V_{\lambda}(f), V_{G}(f)$ and $V_{\lambda}(F f), V_{G}(F f)$ respectively. With respect both uncertainty products $U P_{G}$ and $U P_{\lambda}$, functions $f_{1}$ and $g_{1}$ have the same localization, while function $f_{2}$ is more localized then $g_{2}$, that is adjusted with a naive idea of localization as a characteristic of a measure for a function support.

Table 1: $U P_{G}$ and $U P_{\lambda}$ : Example 1.

| $f$ | $\tilde{x}_{0}(f)$ | $\tilde{t}_{0}(f)$ | $V_{\lambda}(f)$ | $V_{\lambda}(F f)$ | $U P_{\lambda}(f)$ | $V_{G}(f)$ | $V_{G}(F f)$ | $U P_{G}(f)$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $f_{1}$ | $[0,1 / 4)$ | $[0,4)$ | $1 / 48$ | $16 / 3$ | $1 / 9$ | $1 / 28$ | $64 / 7$ | $16 / 49$ |
| $g_{1}$ | $[3 / 4,1)$ | $[0,4)$ | $1 / 48$ | $16 / 3$ | $1 / 9$ | $1 / 28$ | $64 / 7$ | $16 / 49$ |
| $f_{2}$ | $[0,1 / 8)$ | $[0,2)$ | $3 / 64$ | 8 | $3 / 8$ | $4 / 21$ | $96 / 7$ | $128 / 49$ |
| $g_{2}$ | $[3 / 4,7 / 8)$ | $[0,4)$ | $71 / 64$ | $32 / 3$ | $71 / 6$ | $19 / 14$ | $255 / 14$ | $4845 / 196$ |

Example 2. Here we discuss a dependence of a localization for a fixed function on a parameter $p$ of the Vilenkin group $G_{p}$. Let us consider a function $f_{1}(x)=\mathbb{1}_{\lambda^{-1}[0,1 / 4)}(x)$ and $p=2^{k}, k \in \mathbb{N}$. We calculate $U P_{G}\left(f_{1}\right)$.
(1) If $k=1$, then $U P_{G}\left(f_{1}\right)=\frac{16}{49}$ (see Example 1.);
(2) If $k=2$, then

$$
\begin{aligned}
& V_{G}\left(f_{1}\right)=\frac{1}{\left\|f_{1}\right\|_{L_{2}(G)}^{2}} \min _{\widetilde{x}} \int_{G}\|x \ominus \tilde{x}\|_{G}^{2}\left|f_{1}(x)\right|^{2} d x=4 \min _{\widetilde{x}} \int_{\lambda^{-1}\left[0, \frac{1}{4}\right)}\|x \ominus \tilde{x}\|_{G}^{2} d x \\
& =4 \int_{\lambda^{-1}\left[0, \frac{1}{4}\right)}\|\tau\|_{G}^{2} d \tau=4 \sum_{i=1}^{\infty} \int_{\lambda^{-1}\left[\frac{1}{\left.4^{2+1}, \frac{1}{4^{2}}\right)}\right.}\|\tau\|_{G}^{2} d \tau=4 \sum_{i=1}^{\infty}\left(\frac{1}{4^{i}}-\frac{1}{4^{i+1}}\right) 4^{-2 i}=\frac{1}{21} . \\
& V_{G}\left(F f_{1}\right)=\frac{1}{\left\|F f_{1}\right\|_{L_{2}(G)}^{2}} \min _{\tilde{x}} \int_{G}\|x \ominus \tilde{x}\|_{G}^{2}\left|F f_{1}(x)\right|^{2} d x=\frac{1}{4} \min _{\tilde{x}} \int_{\lambda^{-1}[0,4)}\|x \ominus \widetilde{x}\|_{G}^{2} d x \\
& =\frac{1}{4} \int_{\lambda^{-1}[0,4)}\|\tau\|_{G}^{2} d \tau=\frac{1}{4} \sum_{i=-1}^{\infty} \int_{\lambda^{-1}\left[\frac{1}{4^{2+1}}, \frac{1}{4^{2}}\right)}\|\tau\|_{G}^{2} d \tau=\frac{1}{4} \sum_{i=-1}^{\infty}\left(\frac{1}{4^{i}}-\frac{1}{4^{i+1}}\right) 4^{-2 i}=\frac{256}{21} . \\
& \text { Hence, } U P_{G}\left(f_{1}\right)=\frac{256}{441} . \\
& \text { (3) If } k>2, \text { then }
\end{aligned}
$$

$$
\begin{gathered}
V_{G}\left(f_{1}\right)=\frac{1}{\left\|f_{1}\right\|_{L_{2}(G)}^{2}} \min _{\widetilde{x}} \int_{G}\|x \ominus \tilde{x}\|_{G}^{2}\left|f_{1}(x)\right|^{2} d x=4 \min _{\widetilde{x}} \int_{\lambda^{-1}\left[0, \frac{1}{4}\right)}\|x \ominus \tilde{x}\|_{G}^{2} d x \\
=4 \int_{\lambda^{-1}\left[0, \frac{1}{2^{k}}\right) \oplus\left[\frac{1}{2^{k}}, \frac{1}{4}\right)}\|\tau\|_{G}^{2} d \tau=4\left(\sum_{i=1}^{\infty}\left(\frac{1}{\left(2^{k}\right)^{i}}-\frac{1}{\left(2^{k}\right)^{i+1}}\right)\left(2^{k}\right)^{-2 i}+\left(\frac{1}{4}-\frac{1}{2^{k}}\right)\right) \\
=1-\frac{4}{2^{k}}+\frac{4}{2^{k}\left(2^{2 k}+2^{k}+1\right)} .
\end{gathered}
$$

$$
\begin{gathered}
V_{G}\left(F f_{1}\right)=\frac{1}{\left\|F f_{1}\right\|_{L_{2}(G)}^{2}} \min _{\widetilde{x}} \int_{G}\|x \ominus \tilde{x}\|_{G}^{2}\left|f_{1}(x)\right|^{2} d x=\frac{1}{4} \min _{\widetilde{x}} \int_{\lambda^{-1}[0,4)}\|x \ominus \tilde{x}\|_{G}^{2} d x \\
=\frac{1}{4} \int_{\lambda^{-1}[0,1) \oplus[1,4)}\|\tau\|_{G}^{2} d \tau= \\
=\frac{1}{4}\left(\sum_{i=0}^{\infty}\left(\frac{1}{\left(2^{k}\right)^{i}}-\frac{1}{\left(2^{k}\right)^{i+1}}\right)\left(2^{k}\right)^{-2 i}+(4-1) \cdot 2^{2 k}\right) \\
=\frac{3}{4} \cdot 2^{2 k}+\frac{1}{4} \cdot \frac{2^{2 k}}{2^{2 k}+2^{k}+1} .
\end{gathered}
$$

Therefore, $U P_{G}\left(f_{1}\right)=\left(1-\frac{4}{2^{k}}+\frac{4}{2^{k}\left(2^{2 k}+2^{k}+1\right)}\right)\left(\frac{3}{4} \cdot 2^{2 k}+\frac{1}{4} \cdot \frac{2^{2 k}}{2^{2 k}+2^{k}+1}\right)$.
It is easy to see that time variance $V_{G}\left(f_{1}\right)$ goes to 1 , and frequency variance $V_{G}\left(F f_{1}\right)$ goes to infinity as $k \rightarrow \infty$.

## 4 Uncertainty product $U P_{G}$.

In this section we concentrate on the uncertainty product corresponding to the metric $d_{2}$. It turns out that the modified Gibbs derivative $\mathcal{D}$ plays a role of a usual derivative in this case. And since the Haar functions are the eigenfunctions of $\mathcal{D}$, it is possible to get representation for $U P_{G}$ using the Haar coefficients.

Theorem 3. Suppose $f \in L_{2}(G) \cap L_{1}(G),\|\cdot\|_{G} f \in L_{2}(G)$, where "dot" • means the argument $x \in G$ of a function $f$, and $f(x)=\sum_{\nu=1}^{p-1} \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}_{+}} c_{j, k}^{\nu} \psi_{j, k}^{\nu}(x)$. Then

$$
\begin{gather*}
\int_{G}\|t\|_{G}^{2}|F f(t)|^{2} d t=\int_{G}|\mathcal{D} f(t)|^{2} d t=\sum_{\nu=1}^{p-1} \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}_{+}}\left|p^{j} c_{j, k}^{\nu}\right|^{2}  \tag{9}\\
\int_{G}\|x\|_{G}^{2}|f(x)|^{2} d x=\int_{G}|\mathcal{D} F f(x)|^{2} d x=\sum_{\nu=1}^{p-1} \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}_{+}}\left|p^{j} d_{j, k}^{\nu}\right|^{2} \tag{10}
\end{gather*}
$$

where $d_{j, k}^{\nu}, j \in \mathbb{Z}, k \in \mathbb{Z}_{+}, \nu=1, \ldots, p-1$, are the coefficients in the Haar series for the function $F f$, that is $F f(t)=\sum_{\nu=1}^{p-1} \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}_{+}} d_{j, k}^{\nu} \psi_{j, k}^{\nu}(t)$.

Proof. By the definition of the modified Gibbs derivative and the Plancherel equality we get

$$
\int_{G}\|t\|_{G}^{2}|F f(t)|^{2} d t=\int_{G}|F \mathcal{D} f(t)|^{2} d t=\int_{G}|\mathcal{D} f(t)|^{2} d t
$$

Expanding a function in the Haar series and applying Corollary [1 we get

$$
\int_{G}|\mathcal{D} f(t)|^{2} d t=\int_{G}\left|\sum_{\nu=1}^{p-1} \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}_{+}} c_{j, k}^{\nu} \mathcal{D} \psi_{j, k}^{\nu}(t)\right|^{2} d t
$$

$$
=\int_{G}\left|\sum_{\nu=1}^{p-1} \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}_{+}} c_{j, k}^{\nu} p^{j} \psi_{j, k}^{\nu}(t)\right|^{2} d t=\sum_{\nu=1}^{p-1} \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}_{+}}\left|p^{j} c_{j, k}^{\nu}\right|^{2}
$$

The last equality follows from the orthonormality of the Haar system. Equality (10) is proved analogously to (9).

Remark 1. Formally, it is possible to write $\int_{G} \lambda^{2}(x)|F f(x)|^{2} d x=\int_{G}\left|f^{[1]}(x)\right|^{2} d x$ and to try to represent $U C_{\lambda}$ in terms of eigenfunctions of the Gibbs derivative $f^{[1]}$ in the case of the Cantor group. (The Gibbs derivative is defined for functions defined on the Cantor group only.) However, the Gibbs differentiation is not a local operation, that is $\left(f \mathbb{1}_{E}\right)^{[1]} \neq f^{[1]} \mathbb{1}_{E}$, see also discussion in [10]. So, usage of Walsh functions instead of Haar basis might give interesting results for periodic functions only.

We did not found in the literature a formula expressing $d_{j, k}^{\mu}$ in terms of $c_{j, k}^{\nu}$. So we obtain this formula in the following lemma.

Lemma 2. Suppose $f \in L_{2}(G)$ and the coefficients $c_{j, k}^{\nu}, d_{j, k}^{\mu}, j \in \mathbb{Z}, k \in \mathbb{Z}_{+}, \nu, \mu=1, \ldots, p-1$, are defined in Theorem 图. Then

$$
\begin{equation*}
d_{j, k}^{\mu}=\sum_{\nu=1}^{p-1} p^{q_{0} / 2} b_{k}^{\nu}+p^{j / 2} \sum_{\nu=1}^{p-1} c_{-j-1,0}^{\nu} \exp \left(-\frac{2 \pi i \nu \mu}{p}\right) \delta_{k, 0}+p^{j / 2} \sum_{i=-\infty}^{-j-2} \sum_{\nu=1}^{p-1} c_{i, 0}^{\nu} \delta_{k, 0} \tag{11}
\end{equation*}
$$

where $b_{k}^{\nu}=p^{-q_{0}} \sum_{n=0}^{p^{q_{0}}-1} c_{q_{0}-j, n+(p-\mu) p^{q_{0}}}^{\nu} \chi\left(\lambda^{-1}(n), D^{-q_{0}} \lambda^{-1}(k)\right)$ is the $k$-th term of the discrete Vilenkin-Chrestenson transform of $\left(c_{q_{0}-j, n+(p-\mu) p_{0}^{q}}^{\nu}{ }_{n=0}^{p_{0}^{q}-1}, q_{0}=\left[\log _{p} \frac{k}{p-\nu}\right]\right.$, and $\delta_{0,0}=1$, and $\delta_{k, 0}=0$, if $k \neq 0$.

Proof. Using the Plancherel equality and (7), we get

$$
\begin{gathered}
d_{j, k}^{\mu}=\int_{G} F f(x) \overline{\psi_{j, k}^{\mu}(x)} d x=\int_{G} f(x) \overline{F \psi_{j, k}^{\mu}(x)} d x=\sum_{\nu=1}^{p-1} \sum_{i \in \mathbb{Z}} \sum_{n \in \mathbb{Z}_{+}} c_{i, n}^{\nu} \int_{G} \psi_{i, n}^{\nu}(x) \overline{F \psi_{j, k}^{\mu}(x)} d x \\
=\sum_{\nu=1}^{p-1} \sum_{i \in \mathbb{Z}} \sum_{n \in \mathbb{Z}_{+}} c_{i, n}^{\nu} \int_{G} \psi_{i, n}^{\nu}(x) p^{-j / 2} \overline{\chi\left(\lambda^{-1}(k), D^{-j} x\right)} \mathbb{1}_{I_{-j} \oplus \lambda^{-1}\left((p-\mu) p^{j}\right)}(x) d x .
\end{gathered}
$$

Since supp $\psi_{i, n}^{\nu}=\lambda^{-1}\left(\left[n p^{-i},(n+1) p^{-i}\right)\right)$, it follows that the last expression takes the form

$$
\begin{aligned}
& \sum_{\nu=1}^{p-1} \sum_{i=-j}^{\infty} \sum_{n=(p-\mu) p^{i+j}}^{(p-\mu+1) p^{i+j}-1} c_{i, n}^{\nu} \int_{G} \psi_{i, n}^{\nu}(x) p^{-j / 2} \overline{\chi\left(\lambda^{-1}(k), D^{-j} x\right)} d x \\
& \quad+p^{-j / 2}\left(\sum_{\nu=1}^{p-1} c_{-j-1,0}^{\nu} \exp \left(-\frac{2 \pi i \nu \mu}{p}\right)+\sum_{\nu=1}^{p-1} \sum_{i=-\infty}^{-j-2} c_{i, 0}^{\nu}\right)
\end{aligned}
$$

$$
\times \int_{G} \overline{\chi\left(\lambda^{-1}(k), D^{-j} x\right)} \mathbb{1}_{I_{-j} \oplus \lambda^{-1}\left((p-\mu) p^{j}\right)}(x) d x=: S_{1}+S_{2} .
$$

For the first sum by (7) we note that

$$
\begin{aligned}
& \int_{G} \psi_{i, n}^{\nu}(x) \overline{\chi\left(\lambda^{-1}(k), D^{-j} x\right)} d x=F \psi_{i, n}^{\nu}\left(D^{-j} \lambda^{-1}(k)\right) \\
= & p^{-i / 2} \chi\left(n, D^{-i-j} \lambda^{-1}(k)\right) \mathbb{1}_{I_{-i} \oplus \lambda^{-1}\left((p-\nu) p^{i}\right)}\left(D^{-j} \lambda^{-1}(k)\right) .
\end{aligned}
$$

Therefore, the first sum takes the form

$$
\begin{aligned}
S_{1}= & \sum_{\nu=1}^{p-1} \sum_{i=-j}^{\infty} \sum_{n=(p-\mu) p^{i+j}}^{(p-\mu+1) p^{i+j}-1} p^{-(j+i) / 2} c_{i, n}^{\nu} \chi\left(\lambda^{-1}(n), D^{-i-j} \lambda^{-1}(k)\right) \mathbb{1}_{I_{-i-j} \oplus \lambda^{-1}\left((p-\nu) p^{(i+j))}\right.}\left(\lambda^{-1}(k)\right) \\
& =\sum_{\nu=1}^{p-1} \sum_{q=0}^{\infty} p^{-q / 2} \sum_{n=0}^{p^{q}-1} c_{q-j, n+(p-\mu) p^{q}}^{\nu} \chi\left(\lambda^{-1}(n), D^{-q} \lambda^{-1}(k)\right) \mathbb{1}_{I_{-q} \oplus \lambda^{-1}\left((p-\nu) p^{q}\right)}\left(\lambda^{-1}(k)\right) .
\end{aligned}
$$

Since $\mathbb{1}_{I_{-q} \oplus \lambda^{-1}\left((p-\nu) p^{q}\right)}\left(\lambda^{-1}(k)\right)=1$ for $(p-\nu) p^{q} \leq k<(p-\nu+1) p^{q}$ and $\mathbb{1}_{I_{-q} \oplus \lambda^{-1}\left((p-\nu) p^{q}\right)}\left(\lambda^{-1}(k)\right)=$ 0 for the remaining $k$, and since the inequality $(p-\nu) p^{q} \leq k<(p-\nu+1) p^{q}, q \in \mathbb{Z}_{+}$is equivalent to $q=\left[\log _{p} \frac{k}{p-\nu}\right]$, it follows that the only nonzero term in the sum $\sum_{q=0}^{\infty}$ has the number $q_{0}:=\left[\log _{p} \frac{k}{p-\nu}\right]$. So

$$
S_{1}=\sum_{\nu=1}^{p-1} p^{-q_{0} / 2} \sum_{n=0}^{p^{q_{0}-1}} c_{q_{0}-j, n+(p-\mu) p^{q_{0}}}^{\nu} \chi\left(\lambda^{-1}(n), D^{-q_{0}} \lambda^{-1}(k)\right) .
$$

By (3) we notice that up to the multiplication by a constant the inner sum in the last expression is the $k$-th term of the discrete Vilenkin-Chrestenson transform of the vector $\left(c_{q_{0}-j, n+(p-\mu) p_{0}^{q}}^{\nu}{ }_{n=0}^{p_{0}^{q}-1}\right.$. Denote this term by $b_{k}^{\nu}$. Finally, for $S_{1}$ we get

$$
S_{1}(x)=\sum_{\nu=1}^{p-1} p^{q_{0} / 2} b_{k}^{\nu} .
$$

Thus, the first sum takes the desired form. To conclude the proof it remains to calculate the following part of the second sum

$$
\begin{gathered}
\int_{G} \overline{\chi\left(k, D^{-j} x\right)} \mathbb{1}_{I_{-j} \oplus \lambda^{-1}\left((p-\mu) p^{j}\right)}(x) d x=p^{j} \int_{G} \overline{\chi(k, x)} \mathbb{1}_{I \oplus \lambda^{-1}(p-\mu)}(x) d x \\
=p^{j} \int_{I} \overline{\chi\left(k, x \ominus \lambda^{-1}(p-\mu)\right)} d x=p^{j} \int_{I} \overline{\chi(k, x)} d x=p^{j} \delta_{k, 0},
\end{gathered}
$$

where $\delta_{0,0}=1$, and $\delta_{k, 0}=0$, if $k \neq 0$.
It is easy to see from (91) that $\min \int_{G}\|t\|_{G}^{2}|F f(t)|^{2} d t=0$ and max $\int_{G}\|t\|_{G}^{2}|F f(t)|^{2} d t=$ $\infty$ under the restriction $\|f\|_{L_{2}(G)}=1$.

Formulas (9) and (10) allow for the following result on estimation of Fourier-Haar coefficients for functions defined on the Vilenkin group.

Corollary 2. Suppose $\|\cdot\|_{G} F f \in L_{2}(G)$, and $f(x)=\sum_{\nu=1}^{p-1} \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}_{+}} c_{j, k}^{\nu} \psi_{j, k}^{\nu}(x)$. Then the series $\sum_{\nu=1}^{p-1} \sum_{j \in \mathbb{Z}} \sum_{k \in \mathbb{Z}_{+}}\left|p^{j} c_{j, k}^{\nu}\right|^{2}$ is convergent.

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## References

[1] E. Breitenberger, Uncertainty measures and uncertainty relations for angle observables, Found. Phys. 15 (1985), 353-364.
[2] B.I. Golubov, Elements of dyadic analysis, [in Russian], Moscow, LKI, 2007.
[3] B.I. Golubov, A.V. Efimov, and V.A. Skvortsov, Walsh series and transforms, English transl.: Kluwer, Dordrecht, 1991.
[4] W. Erb, Uncertainty principles on compact Riemannian manifolds, Appl. Comput. Harmon. Anal., 29 (2010), 182-197.
[5] Yu. A. Farkov, Multiresolution analysis and wavelets on Vilenkin groups, FACTA UNIVERSITATIS (NIS) SER.: ELEC. ENERG. vol. 21, no. 3, December 2008, 309-325.
[6] W. Heisenberg, The actual concept of quantum theoretical kinematics and mechanics, Physikalische Z. 43 (1927), 172.
[7] E. Hewitt, K. A. Ross, Abstract Harmonic Analysis. Springer-Verlag, New York, 1963, 1979.
[8] A. V. Krivoshein, E. A. Lebedeva, Uncertainty Principle for the Cantor Dyadic Group, J. Math. Anal. Appl, 42 (2015), 1231-1242.
[9] W. C. Lang. Orthogonal wavelets on the Cantor dyadic group, SIAM J. Math. Anal. 1996. - Vol. 27. - P. 305-312.
[10] E. Lebedeva, M. Skopina. Walsh and wavelet methods for differential equations on the Cantor group // J. Math. Anal. Appl. - 2015. - Vol. 430. - No. 2. - P. 593-613.
[11] J. F.Price, A. Sitaram, Local uncertainty inequalities for locally compact groups, Trans. of AMS, 3081 (1988), 105-114.
[12] E. Schrödinger, About Heisenberg uncertainty relation, Proc. of The Prussian Acad. of Scien. XIX (1930) 296-303.
[13] F. Schipp, W. R. Wade, P. Simon. Walsh series. An introduction to dyadic harmonic analysis. - Academiai Kiado, Budapest, 1990.


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