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## Recurrence Relations

 for the Fresnel and Similar Integrals

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#### Abstract

The class of functions defined by $\int_{0}^{\infty}[\exp (-c X) d t /$ $\left.(1+Y)(\sqrt{t})^{k}\right]$ where $X$ and $Y$ are either $t$ or $t^{2}$ and $k$ is $\mathbf{- 1 , 0}$, or $\mathbf{1}$ can be evaluated by recurrences for all but small values of the parameter $c$. These recurrences, given here, are more efficient than the usual asymptotic series.

Key Words and Phrases: recurrence relations, Fresnel integral, exponential integral

CR Categories: 5.12


In this paper we give simple recurrence relations that evaluate integrals such as $\int_{0}^{\infty}\left[\exp \left(-c t^{2}\right) d t /(1+t)\right]$ and $\int_{0}^{\infty}\left[\exp (-c t) d t /\left(1+t^{2}\right)\right]$ for values of the argument $2<c<\infty$. These recurrences are simple to program, are remarkably efficient for $\mathrm{c}>5$, and are especially suitable for the small programmable electronic desk calculators that are now appearing where storage for constants is limited. The recurrences provide an alternative to rational approximations (which are not always available) that are derived from asymptotic series via the Q-D algorithm, usually with some kind of economization. (Cody [3] gives some for our sample integral.)

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We emphasize that these recurrences, like asymptotic series, decrease in efficiency as the argument becomes smaller, but they remain viable until $c$ is down to 3 or 2-i.e. well into the range for which a small-argument series is efficient, provided, of course, that any singular behavior at the origin is explicitly removed (see [2, Chap. 10]).

Specifically, the integrals we treat here are of the form
$F(c)=\int_{0}^{\infty}\left[\exp (-c X) d t /(1+Y)(\sqrt{ } \bar{t})^{k}\right]$
where $X$ and $Y$ are either $t$ or $t^{2}$ (in all combinations) and $k$ is $\pm 1$ or zero. The recurrences are all given in Tables I and II. Here we continue our exposition via one example, the integral
$J_{0}(c)=\int_{0}^{\infty}\left[\exp (-c t) d t / \sqrt{t}\left(1+t^{2}\right)\right]$
that occurs as an asymptotic form of the Fresnel integrals (see 7.4.26 and 7.3.5 in [1]).

We define the three sequences of functions
$J_{n}(c)=\int_{0}^{\infty} \frac{e^{-c t}}{\sqrt{t}\left(1+t^{2}\right)}\left[\frac{t^{2}}{1+t^{2}}\right]^{n} d t$

$$
c \geq 0
$$

$I_{n}(c)=\int_{0}^{\infty} \frac{e^{-c t} t}{\sqrt{t}\left(1+t^{2}\right)}\left[\frac{t^{2}}{1+t^{2}}\right]^{n} d t$
$n \geq 0$
$K_{n}(c)=\int_{0}^{\infty} \frac{e^{-c t}}{\sqrt{t}}\left[\frac{t^{2}}{1+t^{2}}\right]^{n} d t$
and note that $J_{0}(c)$ is the integral we seek, $I_{0}(c)$ is essentially the other Fresnel asymptotic form, while $K_{0}(c)$ is elementary, being $\sqrt{\pi / c}$. Further, all three sequences of functions decrease monotonically with increasing $n$. The recurrences
$I_{n-1}=\frac{4 n \cdot I_{n}+2 c \cdot K_{n}}{4 n-1}$
$J_{n-1}=\frac{4 n \cdot J_{n}+2 c \cdot I_{n-1}}{4 n-3}$
$K_{n-1}=K_{n}+J_{n-1}$
may be derived through integration-by-parts with frequent use of the relation
$\frac{t^{2}}{1+t^{2}}=1-\frac{1}{1+t^{2}}$
and the observation that
$\frac{d}{d t}\left[\frac{t^{2}}{1+t^{2}}\right]=\frac{2 t}{\left(1+t^{2}\right)^{2}}$.
The derivation of the first relation is typical:

$$
\begin{aligned}
J_{n} & =\int_{0}^{\infty} \frac{e^{-c t}}{\sqrt{t}\left(1+t^{2}\right)}\left[\frac{t^{2}}{1+t^{2}}\right]^{n-1}\left(1-\frac{1}{1+t^{2}}\right) d t \\
& =J_{n-1}-\frac{1}{2} \int_{0}^{\infty} \frac{e^{-c t}}{t^{3 / 2}}\left[\frac{t^{2}}{1+t^{2}}\right]^{n-1} \frac{2 t d t}{\left(1+t^{2}\right)^{2}}
\end{aligned}
$$

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Table I. Integrals with $\exp (-c t)$ in Their Numerators

*These are the formula numbers in AMS-55 for the integrals.
whence integration-by-parts gives

$$
\begin{aligned}
J_{n}= & J_{n-1}+\frac{1}{2 n} \int_{0}^{\infty} e^{-c t}\left[\frac{t^{2}}{1+t^{2}}\right]^{n}\left(-\frac{c}{t^{3 / 2}}-\frac{3}{2 t^{5 / 2}}\right) d t \\
= & J_{n-1}-\frac{c}{2 n} \int_{0}^{\infty} \frac{e^{-c t} t}{\sqrt{t}\left(1+t^{2}\right)}\left[\frac{t^{2}}{1+t^{2}}\right]^{n-1} d t \\
& -\frac{3}{4 n} \int_{0}^{\infty} \frac{e^{-c t} t}{\sqrt{t}\left(1+t^{2}\right)}\left[\frac{t^{2}}{1+t^{2}}\right]^{n} d t \\
= & J_{n-1}-\frac{c}{2 n} I_{n-1}-\frac{3}{4 n} J_{n-1}
\end{aligned}
$$

so finally
$J_{n}=\left(1-\frac{3}{4 n}\right) J_{n-1}-\frac{c}{2 n} I_{n-1}$.
The system of recurrences is homogeneous, and the familiar scheme of J.C.P. Miller [1, Sec.9.12] may be used in which canonical values of $I_{n}(c)=1, J_{n}(c)=0$, and $K_{n}(c)=0$ are given for some suitably large $n$. The recurrences are then run down to $n=0$ and the value of $I_{0}$ (or $J_{0}$ ) is finally normalized by multiplying by the ratio of $\sqrt{\pi / c}$ to the value of $K_{0}(c)$ that was computed -all values in the computation being erroneous by this common factor. For our example we take $c=5, n=10$ to produce the values in Table III and we see that we have achieved results correct to nearly seven significant figures in $J_{0}$ and five in $I_{0}$. (No intermediate values need be retained. They are given here merely to show the

Table II. Integrals with $\exp \left(-c t^{2}\right)$ in Their Numerators
$I_{0}=\int \frac{e^{-c t^{2}}}{(1+t)} d t$
$J_{1}=\int \frac{e^{-c t_{t}^{2}}}{(1+t)} d t$
Recurrence Normalization

| 1 | $I_{n-1}=I_{n}+\frac{2 c^{\prime}}{n} K_{n}$ | 7.4 .10 |
| :---: | :---: | :---: |
| - | $J_{n-1}=J_{n}+I_{n-1}$ | $J_{0}=\int e^{-c t^{2}} d t=\frac{1}{2} \sqrt{\frac{\pi}{c}}$ |
| 1. | $K_{n-1}=K_{n}+J_{n}$ | $K_{0}=\int e^{-c t^{2}} t d t=\frac{1}{2 c}$ |

$I_{0}=\int \frac{e^{-c t^{2}}}{1+t^{2}} d t$

| 1 |
| :---: |
| $I_{n-1}=\frac{2 n \cdot I_{n}+2 c \cdot J_{n}}{2 n-1}$ |
| $J_{n-1}=J_{n}+I_{n-1}$ |$\quad J_{0}=\int e^{-c t^{2}} d t=\frac{1}{2} \sqrt{\frac{\pi}{c}}$

$\begin{array}{ll}I_{0}=\int_{\sqrt{t}\left(1+t^{2}\right)}^{-c t^{2}}{ }^{2} d t & 1 \\ J_{1}=\int \frac{e^{-c t^{2}} t \sqrt{t} d t}{\left(1+t^{2}\right)} & 0 \quad I_{n-1}=\frac{4 n \cdot I_{n}+4 c \cdot J_{n}}{4 n-3} . \\ J_{n-1}=J_{n}+I_{n-1} & J_{0}=\int \frac{e^{-c t^{2}}}{\sqrt{t}} d t=\frac{\Gamma^{1 / 4}}{2 c^{1 / 4}}\end{array}$
$I_{0}=\int \frac{e^{-c t^{2}} \sqrt{t}}{\left(1+t^{2}\right)} d t$

$I_{0}=\int \frac{e^{-c t^{2}}}{\sqrt{t}(1+t)} d t$

| , | $I_{n-1}=\frac{2 n \cdot I_{n}+4 c \cdot K_{n}}{2 n-1}$ |  |
| :---: | :---: | :---: |
|  | $J_{n-1}=\frac{2 n \cdot J_{n}+4 c \cdot I_{n}}{2 n+1}$ |  |
| 0 | $K_{n-1}=K_{n}+J_{n-1}$ | $K_{0}=\int e^{-c t^{2}} \sqrt{E d t}=\frac{\Gamma c^{3 / 4}}{2 c^{3 / 4}}$ |
| - | $L_{n-1}=L_{n}+K_{n}^{(\text {not } n-1)}$ | $L_{0}=\int e^{-c t^{2}} t \sqrt{t} d t=\frac{\Gamma\left(^{5 / 4}\right)}{2 e^{5 / 4}}$ |

Table III.

| n | $J_{n}$ |  | $I_{n}$ | $\mathrm{K}_{n}$ |
| :---: | :---: | :---: | :---: | :---: |
| 10 | 0 |  | 1.0 | 0 |
| 9 | 0.2772 | 002 | 1.025641 | 0.277200 |
| 8 | 0.6460 | 806 | 1.1341451 | 0.9232809 |
| 7 | 1.2193 | 176 | 1.4685630 | 2.1425985 |
| 6 | 2.2922 | 394 | 2.3165093 | 4.4348380 |
| 5 | 4.6889 | 486 | 4.3454175 | 9.1237866 |
| 4 | 11.0317 | 728 | 9.3761166 | 20.1555954 |
| 3 | 31.6069 | 748 | 23.4382307 | 51.7625342 |
| 2 | 122.8379 | 976 | 72.6258282 | 174.6005318 |
| 1 | 861.4013 | 517 | 332.4302777 | 1036.001883 |
| 0 | 42411.4052 | 3 | 3896.5799 82 | 43447.4071 1 |
| Nomaliz | ed 0.77376 | 43797 | . $071090 \cdot 1885$ |  |
| Correct | 0.77376 | 45665 | .071089873 |  |

growth that is typical in these recurrences.) The larger the starting value of $n$, the more accurate the final results, but the accuracy is also a function of $c$. As a rough guide, $n_{\text {initial }}=150 / \mathrm{c}$ will yield approximately ten significant figures in $I_{0}$ and $J_{0}$. Note that the recurrences involve no loss of precision through subtraction of imprecisely known quantities-all values of $I, J$, and $K$ being positive. Again, this is typical.

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