

An approximative calculation of the fractal structure in self-similar tilings

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Fractal structures emerge from statistical and hierarchical processes in urban development or network evolution. In a class of efficient and robust geographical networks, we derive the size distribution of layered areas, and estimate the fractal dimension by using the distribution without huge computations. This method can be applied to self-similar tilings based on a stochastic process.

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I. INTRODUCTION

Fractal nature is observed in our real-life infrastructures of urban spatial organization and many technological networks. Indeed, the fractal dimensions of land-use and area-perimeter have been measured in real data of urban cities [1][2]. Similar structures with mixing of dense and sparse areas of nodes have been found in router networks [3], air transportation networks [4], and mobile communication networks [5]. These studies show that the spatial distribution of human activities is inhomogeneous and concerning with the population density. Thus, there commonly exists a hierarchical structure based on statistical self-similarity beyond regular mathematical models such as Sierpinski gasket and carpet. In other words, fractal behavior is not limited to objects with a regular morphology, but can be introduced by iterative evolutionary processes of subdivision or growing at different levels of scaling.

Recently, a multi-scale quartered (MSQ) network, which is stochastically constructed by a self-similar tiling, has been proposed [6][7]. The geographical network embedded on a planar space has several advantages [6]: the robustness of connectivity against failures and attacks, the bounded short path lengths, and the decentralized routing algorithm [8] in a distributed manner. Furthermore [7], it is more efficient (economic) with shorter link lengths and more suitable (tolerant) with lower load for avoiding traffic congestion than the state-of-the-art geometric growing networks [9–15] and the spatially preferential attachment models with various topologies ranging from river to scale-free geographical networks. These properties are useful for the future self-organized design of wide-area wireless ad hoc networks.

This paper investigates the size distribution of faces iteratively divided in the generation process of a MSQ network. In particular, we derive an approximative equation for the averaging behavior of random process, and apply it to easily estimate the fractal dimension of such a self-similar tiling.

II. MSQ NETWORK

For generating a MSQ network, the following process is repeated from an initial tiling which consists of same shaped faces. At each time step, a face is chosen, e.g., with a probability proportional to the population in the space of a face for the load balancing of communication requests, or uniformly at random. Then, as shown in Fig. 1, four smaller faces are created in the chosen face, and a planar network is self-organized on a geographical space. Such a fractal structure is also observed in urban road networks [16] [17]. Note that the MSQ network includes a Sierpinski gasket obtained by a special selection when each triangle, except the central one, is hierarchically divided, however its fractal dimension $\log 3 / \log 2 \approx 1.585$ differs to that in the average behavior of the following random selection.

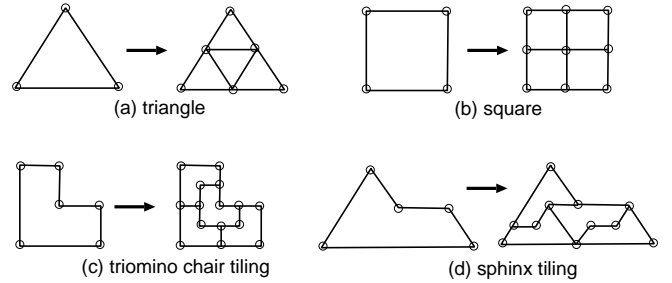


FIG. 1: Basic process of the division.

III. INFINITE STATE MARKOV CHAIN

For simplicity, we treat the uniformly random selection of a face, whose case corresponds to general positions of nodes in the geographical network. In addition, even if a face in the area of high population density tend to be chosen for the subdivision, the assigned population to a node in its territory is asymptotically balanced, then the behavior closes to the case of random selections in course of time. In this setting, the subdivision process makes an infinite state Markov chain as illustrated in Fig. 2. Here, we consider a vector $(n_1(t), n_2(t), \dots, n_l(t) \dots)$ as the state of Markov chain, whose element $n_l(t)$ represents

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the number of faces at time step t on the layer l defined by decreasing order of size. A face on the layer l is chosen with the probability $p_l(t) = n_l(t)/\mathcal{N}(t)$ at the next step $t+1$, then the state is changed from $(\dots, n_l(t), n_{l+1}(t), \dots)$ to $(\dots, n_l(t+1) - 1, n_{l+1}(t+1) + 4, \dots)$, where $\mathcal{N}(t) = \sum_k n_k(t) = \mathcal{N}_0 + 3t$ denotes the total number of faces at t , and \mathcal{N}_0 is the initial number. Note that $p_l(t)$ is time-dependent on a selected path for the decision tree from the top to the bottom in Fig. 2, equivalently on a selection sequence of faces. Thus, the above Markov chain is different from the Galton-Watson type branching process with a time-independent probability for generating offsprings [18].

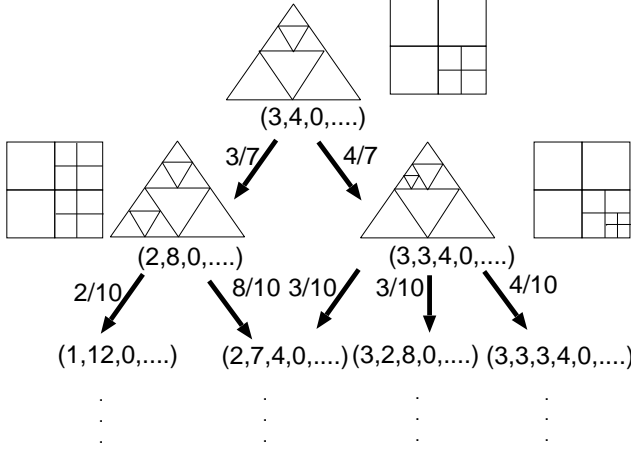


FIG. 2: Iterative division of a face chosen uniformly at random.

We obtain the averaging behavior for

$$\Delta n_l \stackrel{\text{def}}{=} n_l(t+1) - n_l(t). \quad (1)$$

This can be written as

$$\Delta n_l = 4p_{l-1}(t) - p_l(t), \quad (2)$$

since a face on the layer l chosen with the probability $p_l(t)$ is divided into four smaller ones which belong to the layer $l+1$, therefore a face on the layer $l-1$ contributes to increase the number of faces on the layer l . For a large t , by noticing $n_l(t) = \mathcal{N}(t)p_l(t)$ and substituting $\mathcal{N}(t) = \mathcal{N}_0 + 3t \approx 3t$ into the right-hand side of Eq. (1), it is

$$\begin{aligned} \Delta n_l &= 3(t+1)p_l(t+1) - 3tp_l(t), \\ &= 3t[p_l(t+1) - p_l(t)] + 3p_l(t+1). \end{aligned}$$

Using $p_l(t+1) \approx p_l(t)$ because of $t+1 \approx t \gg 1$, Eq. (2) is rewritten to

$$p_l(t+1) - p_l(t) = -\frac{4}{3t}[p_l(t) - p_{l-1}(t)]. \quad (3)$$

We numerically confirm that the average behavior of the Markov chain is well fitting to the solution for Eq.

(3) as shown in Fig. 3. Here, the cross, asterisk, open square, and closed square marks correspond to the average of 100 states reached at each time $t = 10^2, 10^3, 10^4$, and 10^5 , respectively, on the decision tree from the initial state $(3, 4, 0, \dots)$ in Fig. 2. The four types of dashed curves correspond to the distributions calculated by Eq. (3) until these times. The solid curve fitting with the closed square marks is the approximation by a normal distribution. Note that the averaged or calculated number n_l of faces on each layer l at t is normalized as the relative frequency p_l . This bell-shaped distribution means that there exist a mixing of dense and sparse areas with various sizes of faces whose majorities have the intermediate sizes. Figure 4 shows the mean and the variance of the distribution p_l . The upper and lower triangle marks denote the mean and the variance of the distribution obtained for the Markov chain. The solid and dashed lines denote the corresponding results for Eq. (3). In Fig. 5, the circle marks and dotted piecewise linear line denote the skewness of these distributions for the Markov chain and for Eq. (3), respectively. The mean and the variance grow as $O(\log t)$, while the skewness is almost constant around 0.2. Although the p_l is asymptotically a Poisson distribution with a same value for the mean and the variance as mentioned in the Appendix, it is rather fitting with the approximation by Eq. (3) in a finite size.

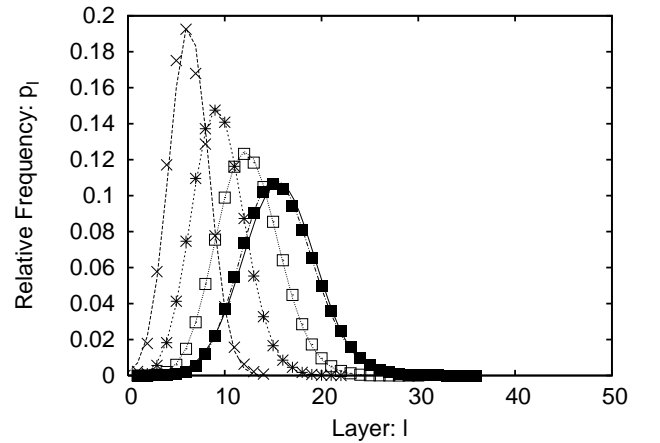


FIG. 3: The distribution of p_l .

On the basis of the numerically obtained $p_l(t)$, the fractal dimension d_f is estimated as 1.2 (a similar value is obtained for a chair or a sphinx tiling [19]) by using a cover-based method in the hierarchically coarse measure of doubling link length on each layer l until $t = 10^8$. Figure 6 shows an example at the state $(1, 11, 4, 0, \dots)$. The number of covered areas (gray rectangles) is

$$N[2] = 2 \times N[1] + 11 + 3 \times (4^2 - (4 \times 1 + 11)),$$

where we set the number of initial coverings $N[1] = 12$ for a square tiling. It depends on the primitive shape; $N[1] = 9$ for a triangle tiling, and $N[1] = 8$ for a chair or a sphinx tiling. In general for a large l , we calculate

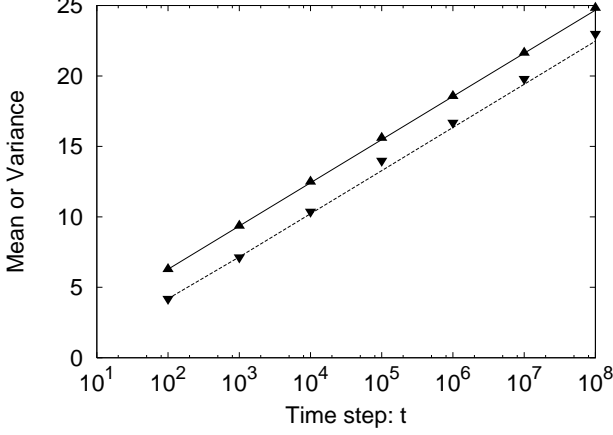


FIG. 4: Mean & Variance of the distribution of p_l .

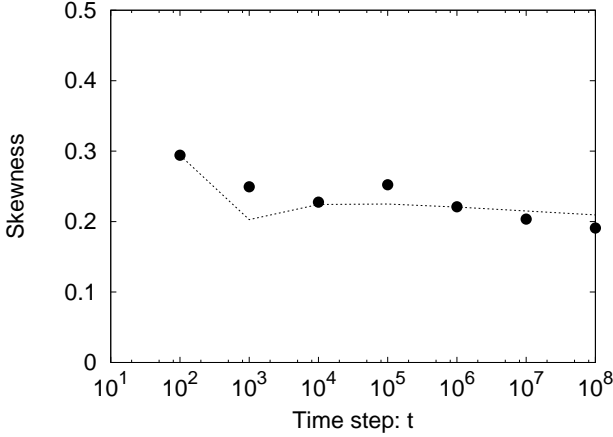


FIG. 5: Skewness of the distribution of p_l .

$d_f = \log N[l] / \log 2^l$ from

$$N[l] = 2 \times N[l-1] + n_l + 3 \times (T_l - S_l), \quad (4)$$

the total number of faces in the completely recursive division $T_l \stackrel{\text{def}}{=} 4^l$, the number of holes without subdivisions $S_l \stackrel{\text{def}}{=} 4 \times S_{l-1} + n_l$, and $S_1 = n_1$. Three terms in the right-hand side of Eq. (4) correspond to thick, thin solid lines, and dashed lines in Fig. 6. Since there is one-to-one correspondence between the divided four faces and the cross-edges in a square at each layer l , e.g., by applying a clockwise mapping rule, n_l represents this effect on the covering. However, the 2nd and 3rd terms are replaced by $3 \times \lceil n_l/4 \rceil$ for a triangle tiling, and by $8 \times \lceil n_l/4 \rceil$ for a chair or a sphinx tiling. We emphasize that our method can be performed by only using the calculated n_l without both averaging and image processings in a box-counting method.

We should remark that the continuous approximation with respect to time and space variables for both sides of Eq. (3)

$$\frac{\partial p_l}{\partial t} = -\frac{4}{3t} \frac{\partial p_l}{\partial l}$$

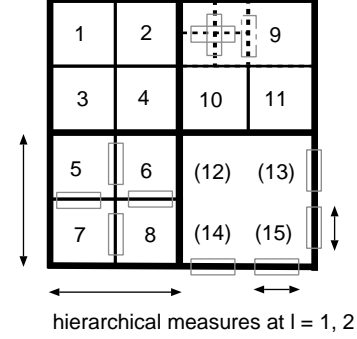


FIG. 6: Illustration for a cover-based method. We count holes by the numbers for n_2 and the parenthesized numbers for n_1 .

gives an incorrect solution, which keeps the shape in the traveling wave

$$p_l(t) = F(l - v \log t),$$

where a function $F()$ is determined by the initial distribution, and $v = 4/3$.

IV. CONCLUSION

Considering the stochastic network construction [6][7] as a Markov chain, the size distribution of the divided faces has been investigated. We have shown, by numerically solving Eq. (3) as the average behavior, that the mean and the variance are proportional to $\log t$, while the skewness is small positive. In a realistic finite size $\mathcal{N} \propto t \leq 10^8$, this approximation of p_l is fitting better than the asymptotical Poisson solution for the interactive particle system (see the Appendix). From the distribution $n_l = \mathcal{N} p_l$, the fractal dimension is estimated as 1.2 without huge computations of averaging and image processings in the conventional box-counting method. The proposed method can be applied to other self-similar tilings based on a stochastic process. Furthermore, it would be interesting to discuss a self-similar modeling which reproduces the fractal structures of urban road [16][17], air-sea transportation [4], communication [5] networks.

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- [1] M. Batty, and P.A. Longley *Fractal Cities: A Geometry of Form and Function*, Academic Press, London, 1994.
 - [2] P. Franlhauser, "The Fractal Approach: A new tool for the spatial analysis of urban agglomerations," *Population: An English Selection*, Vol. 10, pp. 205-240, 1998.
 - [3] S.-H. Yook, H. Jeong, and A.-L. Barabási, "Modeling the Internet's large-scale topology," *PNAS* Vol. 99, No. 21, pp. 13382-13386, 2002.
 - [4] R. Guimerà, S. Mossa, A. Turttschi, and L.A.N. Amaral, "The worldwide air transportation network: Anomalous centrality, community structure, and cities' global roles," *PNAS* Vol. 102, No. 22, pp. 7794-7799, 2005.
 - [5] R. Lambiotte, V.D. Blondel, C. de Kerchove, E. Huens, C. Prieur, Z. Smoreda, and P.V. Dooren, "Geographical dispersal of mobile communication networks," *Physica A* Vol. 387, pp. 5317-5325, 2008.
 - [6] Y. Hayashi, "Evolutionary construction of geographical networks with nearly optimal robustness and efficient routing properties," *Physica A* Vol. 388, pp. 991-998, 2009.
 - [7] Y. Hayashi, and Y. Ono, "Geographical networks stochastically constructed by a self-similar tiling according to population," *Phys. Rev. E* Vol. 82, pp. 016108-1-9, 2010.
 - [8] P. Bose, and P. Morin, "Competitive Online Routing in Geometric Graphs," *Theoretical Computer Science* Vol. 324, No. 2-3, pp. 273-288, 2004.
 - [9] Z. Zhang, S. Zhou, Z. Su, T. Zou, and J. Guan, "Random Sierpinski network with scale-free small-world and modular structure," *Euro. Phys. J. B* Vol. 65, pp. 141-147, 2008.
 - [10] T. Zhou, G. Yan, and B.-H. Wang, "Maximal planar networks with large clustering coefficient and power-law degree distribution," *Phys. Rev. E* Vol. 71, pp. 046141-1-11, 2005.
 - [11] Z. Zhang, and L. Rong, "High dimensional random Apollonian networks," *Physica A* Vol. 364, pp. 610-618, 2006.
 - [12] J.P.K. Doye, and C.P. Massen, "Self-similar disk packings as model spatial scale-free networks," *Phys. Rev. E*, Vol. 71, pp. 016128-1-11, 2005.
 - [13] L. Wang, F. Du, H.P. Dai, and Y.X. Sun, "Random pseudofractal scale-free networks with small-world effect," *Eur. Phys. J. B* Vol. 53, pp. 361-366, 2006.
 - [14] H. D. Rozenfeld, S. Havlin, and D. ben-Avraham, "Fractal and transfractal recursive scale-free nets," *New J. of Phys.* Vol. 6, pp. 175-1-15, 2006.
 - [15] S.N. Dorogovtsev, A.V. Goltsev, and J.F.F. Mendes, "Pseudofractal scale-free web," *Phys. Rev. E* Vol. 65, pp. 066122-1-4, 2002.
 - [16] V. Kalapala, V. Sanwalani, A. Clauset, and C. Moore, "Scale-invariance in road networks," *Phys. Rev. E* Vol. 73, pp. 026130-1-6, 2006.
 - [17] A. Cardillo, S. Scellato, V. Latora, and S. Porta, "Structural properties of planar graphs of urban street patterns," *Phys. Rev. E* Vol. 73, pp. 066107-1-8, 2006.
 - [18] T.M. Liggett, *Stochastic Interacting Systems: Contact, Voter and Exclusion Processes*, Springer, 1999.
 - [19] B. Solomyak, "Dynamics of self-similar tilings," *Ergod. Theory and Dyn. Syst.* Vol. 17, pp. 695-738, 1997.

Appendix A

Let us reformulate the subdivision process for generating a MSQ network as a model in interacting infinite particle systems. Then, according to the notation [18], we consider the generator Ω_H on $Z_+ = \{1, 2, \dots\}$,

$$\Omega_H f(\eta) = \sum_{x \in Z_+} \eta(x) [f(\eta^{x, x+1}) - f(\eta)],$$

where we set $m = 4$ and define the following functions on $\eta \in \{0, 1, 2, \dots\}^{Z_+}$,

$$\eta^{x, x+1}(y) \stackrel{\text{def}}{=} \begin{cases} \eta(x+1) + m, & y = x+1, \\ \eta(x) - 1, & y = x, \\ \eta(y), & y \neq x, x+1. \end{cases}$$

Note that the position x corresponds to the layer l of a chosen face in Sec. 3, therefore the number $\eta(x)$ decreases by 1 at x and increases by m at x due to the contribution at $x-1$. In particular for $f(\eta) = \eta(x)$, we have

$$\frac{dn_x(\tau)}{d\tau} = m \times n_{x-1}(\tau) - n_x(\tau), \quad x \geq 1, \quad (\text{A1})$$

$$\frac{dn_0(\tau)}{d\tau} = -n_0(\tau), \quad (\text{A2})$$

where $n_x(\tau) = E[\eta_\tau(x)]$ is the expectation of the number of particles at position x and at time τ , and the initial configuration is $\eta_0 = (m, 0, 0, \dots)$. We can solve (A1)(A2),

$$n_x(\tau) = m^x \frac{\tau^{x-1}}{(x-1)!} e^{-\tau}, \quad x \geq 1.$$

Also, the expectation of the total number of particles is given by

$$\mathcal{N}(\tau) = m e^{(m-1)\tau}.$$

Therefore, from $p_x(\tau) = n_x(\tau)/\mathcal{N}(\tau)$, we obtain the Poisson distribution with a parameter $m\tau$

$$p_x(\tau) = \frac{(m\tau)^{x-1}}{(x-1)!} e^{-m\tau}, \quad x \geq 1.$$

By the variable transformation between t and τ from the relation $\mathcal{N}(t) = \mathcal{N}(0) + (m-1)t \Leftrightarrow \mathcal{N}(\tau) = m e^{(m-1)\tau}$, the mean and the variance follow $m\tau \propto \log t$. This logarithmic behavior is consistent with the numerical results in Sec. 3.

As another typical approach, if we consider a generating function

$$H^{(t)}(z) \stackrel{\text{def}}{=} \sum_{l=1}^{\infty} p_l(t) z^l,$$

by multiplying z^l to both sides of Eq. (3) and taking the summation $\sum_{l=1}^{\infty}$, we obtain

$$H^{(t+1)}(z) = \left(1 - \frac{4(1-z)}{3t}\right) H^{(t)}(z).$$

Here, we use $\sum_l p_{l-1}(t) z^l = z H^{(t)}(z)$. The frequency $p_l(t)$ is only formally calculated by the recursive differ-

entials

$$p_l(t) = \frac{1}{l!} \frac{\partial^l H^{(t)}(z)}{\partial z^l} \Big|_{z=0}.$$

However, it is practically unsolvable because of involving with a very complicated combinatorial explosion for a large t , equivalently for a large network size $N \propto t$.