Error estimates in periodic homogenization with a non-homogeneous Dirichlet condition.

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Abstract

In this paper we investigate the homogenization problem with a non-homogeneous Dirichlet condition. Our aim is to give error estimates with boundary data in $H^{1/2}(\partial\Omega)$. The tools used are those of the unfolding method in periodic homogenization.

1 Introduction

We consider the following homogenization problem:

$$\phi^{\varepsilon} \in H^1(\Omega), \quad -\operatorname{div}(A_{\varepsilon}\nabla\phi^{\varepsilon}) = f \quad \text{in } \Omega, \quad \phi^{\varepsilon} = g \quad \text{on } \partial\Omega \quad (1.1)$$

where A_{ε} is a periodic matrix satisfying the usual condition of uniform ellipticity and where $f \in L^2(\Omega)$ and $g \in H^{1/2}(\partial \Omega)^1$. We know (see e.g. [4], [10], [13]) that the function ϕ^{ε} weakly converges in $H^1(\Omega)$ towards the solution Φ of the homogenized problem

$$\Phi \in H^1(\Omega), \quad -\operatorname{div}(\mathcal{A}\nabla\Phi) = f \quad \text{in } \Omega, \quad \Phi = g \quad \text{on } \partial\Omega \quad (1.2)$$

where \mathcal{A} is the homogenized matrix (see (4.4) and (4.5)). Using the results in [10] we can give an approximation of ϕ^{ε} belonging to $H^1(\Omega)$ and we easily obtain

$$\phi^{\varepsilon} - \Phi - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_i} \right) \chi_i \left(\frac{\cdot}{\varepsilon} \right) \longrightarrow 0 \quad \text{strongly in} \quad H^1(\Omega)$$

where Q_{ε} is the *scale-splitting operator* (see [10] or Subsection 2.4) and where the χ_i are the correctors (see (4.2)).

One of the aim of this paper is to give error estimates for this homogenization problem. Obviously, if we have $g \in H^{3/2}(\partial\Omega)$ and the appropriate assumptions on

¹The homogenization problem with L^p boundary data is investigated in [3].

the boundary of the domain then we can apply the results in [4], [13], [14], [15], [16] and [22] to deduce error estimates. All of them require that the function Φ belongs at least to $H^2(\Omega)$. Here, the solution Φ of the homogenized problem (1.2) is only in $H^1(\Omega) \cap H^2_{loc}(\Omega)$. In this paper we have to deal with this lack of regularity; this is the main difficulty.

The tools of the unfolding method in periodic homogenization to obtain error estimates (see [14], [15] and [16]) are the projection theorems. This is why we prove two new projection theorems; the Theorems 3.1 and 3.2. Here, both theorems concern the functions $\phi \in H_0^1(\Omega)$ satisfying $\nabla \phi/\rho \in L^2(\Omega; \mathbb{R}^n)$ where $\rho(x)$ is the distance between x and the boundary of Ω . In the first one we give the distance between $\mathcal{T}_{\varepsilon}(\phi)$ (see [10] or Subsection 2.4.1 for the definition of the unfolding operator $\mathcal{T}_{\varepsilon}$) and the space $L^2(\Omega; H_{per}^1(Y))$ in terms of the L^2 norms of ϕ/ρ and $\nabla \phi/\rho$ and obviously ε . In the second one we prove an upper bound for the distance between $\mathcal{T}_{\varepsilon}(\nabla \phi)$ and the space $\nabla H^1(\Omega) \oplus \nabla_y L^2(\Omega; H_{per}^1(Y))$; again in terms of the L^2 norms of ϕ/ρ and $\nabla \phi/\rho$ and $\nabla \phi/\rho$ and ε (see Section 3). This last theorem is partially a consequence of the first one. In this paper we derive the new error estimates from the second projection theorem and those obtained in [16].

Different results are known about the global H^1 error estimate regarding the classical homogenization problem (1.1) (see e.g. [4], [13]). Those with the minimal assumptions are given in [15]; if the solution of the homogenized problem (1.2) belongs to $H^2(\Omega)$ -see Proposition 4.3 in [15]- (respectively $H^{3/2}(\Omega)$; see Theorem 3.3 in [16]) then the H^1 global error is of order $\varepsilon^{1/2}$ (resp. $\varepsilon^{1/4}$) while if this solution belongs to $H^2_{loc}(\Omega) \cap W^{1,p}(\Omega)$ (p > 2) the obtained H^1 global error is smaller and depends on p (see Proposition 4.4 in [15])². Here, with a non-homogeneous Dirichlet condition belonging only to $H^{1/2}(\partial\Omega)$ we do not obtain a global H^1 error estimate. The L^2 global error estimate only requires a boundary of Ω sufficiently smooth (of class $\mathcal{C}^{1,1}$) or a convex open set. Obviously if it is possible to make use of a global H^1 error estimate, the L^2 global error will be better (the reader will be able to compare the Theorem 3.2 in [16] with the Theorem 6.3). The H^1 local error estimate is always linked to the L^2 global error and never needs more assumption (see Theorem 3.2 in [16] or the proof of Theorem 6.1).

The paper is organized as follows. In Section 2 we introduce a few general notations, then we give some reminds³ on lemmas, definitions and results about the unfolding method in periodic homogenization (see [10]), then we prove some new results involving the main operators of this method. Section 3 is devoted to the new projection theorems. In Section 4, we recall the main results on the classical homogenization problem. In Section 5 we introduce an operator which allows to lift the distributions belonging to $H^{-1/2}(\partial\Omega)$ in functions belonging to $L^2(\Omega)$; this lifting operator will play an important role in the case of strongly oscillating boundary data. In Section 6 we derive the error estimates results (Theorems 6.1 and 6.3) with a non-homogenous Dirichlet condition. We

²These propositions or theorem are proved with a Dirichlet condition, with a non-homogenous Dirichlet condition belonging to $H^{3/2}(\partial\Omega)$ the results are obviously the same.

³We want to simplify the reading to a non-familiar reader with the unfolding method

end the paper by investigating a case where the boundary data are strongly oscillating (see Theorem 7.1 in Section 7). A forthcoming paper will be devoted to homogenization problems with other strongly oscillating boundary data.

As general references on the homogenization theory we refer to [1], [4] and [13]. The reader is referred to [10], [12] and [13] for an introduction of the unfolding method in periodic homogenization. The following papers [5], [6], [7], [8], [11], [19], [24] give various applications of the unfolding method in periodic homogenization. As far as the error estimates are concerned, we refer to [2], [4], [14], [15], [16], [20], [22] and [23].

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2 Preliminaries

2.1 Notations

• The space \mathbb{R}^k $(k \ge 1)$ is endowed with the standard basis $(\mathbf{e}_1, \ldots, \mathbf{e}_k)$; the euclidian norm is denoted $|\cdot|$.

• We denote by Ω a bounded domain in \mathbb{R}^n with a Lipschitz boundary.⁴ Let $\rho(x)$ be the distance between $x \in \mathbb{R}^n$ and the boundary of Ω , we set

$$\widetilde{\Omega}_{\gamma} = \left\{ x \in \Omega \mid \rho(x) < \gamma \right\} \qquad \widetilde{\widetilde{\Omega}}_{\gamma} = \left\{ x \in \mathbb{R}^n \mid \rho(x) < \gamma \right\} \qquad \gamma \in \mathbb{R}^{*+}.$$

• There exist constants a, A and γ_0 strictly positive and $M \ge 1$, a finite number N of local euclidian coordinate systems $(O_r; \mathbf{e}_{1r}, \ldots, \mathbf{e}_{nr})$ and mappings $f_r : [-a, a]^{n-1} \longrightarrow \mathbb{R}$, Lipschitz continuous with ratio M, $1 \le r \le N$, such that (see e.g. [17] or [18])

$$\partial\Omega = \bigcup_{r=1}^{N} \left\{ x = x'_{r} + x_{nr} \mathbf{e}_{nr} \in \mathbb{R}^{n} \mid x'_{r} \in \Delta_{a} \text{ and } x_{nr} = f_{r}(x'_{r}) \right\},$$
where $x'_{r} = x_{1r} \mathbf{e}_{1r} + \ldots + x_{n-1r} \mathbf{e}_{n-1r}, \quad \Delta_{a} = \left\{ x'_{r} \mid x_{ir} \in] - a, a[, i \in \{1, \ldots, n-1\} \right\}$

$$\widetilde{\Omega}_{\gamma_{0}} \subset \bigcup_{r=1}^{N} \Omega_{r} \subset \Omega, \qquad \Omega_{r} = \left\{ x \in \mathbb{R}^{n} \mid x'_{r} \in \Delta_{a} \text{ and } f_{r}(x'_{r}) < x_{nr} < f_{r}(x'_{r}) + A \right\}$$

$$\widetilde{\widetilde{\Omega}}_{\gamma_{0}} \subset \bigcup_{r=1}^{N} \left\{ x \in \mathbb{R}^{n} \mid x'_{r} \in \Delta_{a} \text{ and } f_{r}(x'_{r}) - A < x_{nr} < f_{r}(x'_{r}) + A \right\}$$

$$\forall r \in \{1, \ldots, N\}, \quad \forall x \in \Omega_{r} \text{ we have } \frac{1}{2M} (x_{nr} - f_{r}(x'_{r})) \leq \rho(x) \leq x_{nr} - f_{r}(x'_{r}).$$

$$(2.1)$$

⁴In Section 5 and those which follow, we will assume that Ω is a bounded domain of class $\mathcal{C}^{1,1}$ or an open bounded convex set.

• We set

$$Y =]0, 1[^{n}, \qquad \Xi_{\varepsilon} = \left\{ \xi \in \mathbb{Z}^{n} \mid \varepsilon(\xi + Y) \subset \Omega \right\},$$
$$\widehat{\Omega}_{\varepsilon} = \operatorname{interior}\left(\bigcup_{\xi \in \Xi_{\varepsilon}} \varepsilon(\xi + \overline{Y})\right), \qquad \Lambda_{\varepsilon} = \Omega \setminus \widehat{\Omega}_{\varepsilon},$$

where ε is a strictly positive real.

We define

$$\star H^{1}_{\rho}(\Omega) = \left\{ \phi \in L^{2}(\Omega) \mid \rho \nabla \phi \in L^{2}(\Omega; \mathbb{R}^{n}) \right\},$$

$$\star L^{2}_{1/\rho}(\Omega) = \left\{ \phi \in L^{2}(\Omega) \mid \phi/\rho \in L^{2}(\Omega) \right\},$$

$$\star H^{1}_{1/\rho}(\Omega) = \left\{ \phi \in H^{1}_{0}(\Omega) \mid \nabla \phi/\rho \in L^{2}(\Omega; \mathbb{R}^{n}) \right\}.$$
We endow $H^{1}(\Omega)$ (resp. $H^{1}_{+}(\Omega)$) with the norm

We endow $H^1_{\rho}(\Omega)$ (resp. $H^1_{1/\rho}(\Omega)$) with the norm

$$\begin{aligned} \forall \phi \in H^1_{\rho}(\Omega), \quad ||\phi||_{\rho} &= ||\phi||_{L^2(\Omega)} + ||\rho \nabla \phi||_{L^2(\Omega;\mathbb{R}^n)} \\ (\text{ resp. } \forall \phi \in H^1_{1/\rho}(\Omega), \quad ||\phi||_{1/\rho} &= \left\|\nabla \phi / \rho\right\|_{L^2(\Omega;\mathbb{R}^n)}). \end{aligned}$$

Note that if ϕ belongs to $H^1_{\rho}(\Omega)$ then the function $\psi = \rho \phi$ is in $H^1_0(\Omega)$ and vice versa if a function ψ belongs to $H^1_0(\Omega)$ then $\phi = \psi/\rho$ is in $H^1_{\rho}(\Omega)$ since we have (see [9] or [21])

$$\forall \psi \in H_0^1(\Omega), \qquad \left\| \psi/\rho \right\|_{L^2(\Omega)} \le C ||\nabla \psi||_{L^2(\Omega;\mathbb{R}^n)}.$$
(2.2)

Below we recall a classical extension lemma which is proved for example in [15] or which can be proved using the local charts (2.1).

Lemma 2.1. Let Ω be a bounded domain with a Lipschitz boundary, there exist $c_0 \geq 1$ (which depends only on the boundary of Ω) and a linear and continuous extension operator \mathcal{P} from $L^2(\Omega)$ into $L^2(\mathbb{R}^n)$ which also maps $H^1(\Omega)$ into $H^1(\mathbb{R}^n)$ such that

$$\forall \phi \in L^{2}(\Omega), \quad \mathcal{P}(\phi)|_{\Omega} = \phi, \qquad ||\mathcal{P}(\phi)||_{L^{2}(\mathbb{R}^{n})} \leq C||\phi||_{L^{2}(\Omega)}, \\ ||\mathcal{P}(\phi)||_{L^{2}(\widetilde{\Omega}_{\gamma})} \leq C||\phi||_{L^{2}(\widetilde{\Omega}_{c_{0}\gamma})}.$$

$$(2.3)$$

Moreover we have

$$\forall \phi \in H^1(\Omega), \qquad ||\nabla \mathcal{P}(\phi)||_{L^2(\mathbb{R}^n;\mathbb{R}^n)} \le C ||\nabla \phi||_{L^2(\Omega;\mathbb{R}^n)}$$

From now on, if need be, a function ϕ belonging to $L^2(\Omega)$ (resp. $H^1(\Omega)$) will be extended to a function belonging to $L^2(\mathbb{R}^n)$ (resp. $H^1(\mathbb{R}^n)$) using the above lemma. The extension will be still denoted ϕ .

2.2 A characterization of the functions belonging to $H^1_{1/\rho}(\Omega)$

The two first projection theorems (see [15]) regarded the functions belonging to $H_0^1(\Omega)$ while those in [16] regarded the functions in $H^1(\Omega)$. In this paper we prove two

new projection theorems which involve the functions in $H^1_{1/\rho}(\Omega)$; this is why we first give a simple characterization of these functions in the Lemma 2.2 below.

Observe first that if a function ϕ satisfies $\phi/\rho \in H_0^1(\Omega)$ then ϕ belongs to $H_{1/\rho}^1(\Omega)$. The reverse is true.

Lemma 2.2. Let Ω be a bounded domain with a Lipschitz boundary, we have

$$\phi \in H^1_{1/\rho}(\Omega) \iff \phi/\rho \in H^1_0(\Omega).$$

Furthermore there exists a constant which depends only on $\partial\Omega$ such that

$$\forall \phi \in H^{1}_{1/\rho}(\Omega) \qquad \left\| \phi/\rho^{2} \right\|_{L^{2}(\Omega)} + \left\| \phi/\rho \right\|_{H^{1}(\Omega)} \le C ||\phi||_{1/\rho}.$$
(2.4)

Proof. Step 1. Let ϕ be in $H^1(] - a, a[^{n-1} \times]0, A[)$ (a, A > 0) satisfying $\frac{1}{x_n} \nabla \phi(x) \in L^2(] - a, a[^{n-1} \times]0, A[)$ and $\phi(x) = 0$ for a.e. x in $] - a, a[^{n-1} \times \{0\} \cup] - a, a[^{n-1} \times \{A\}$. We have

$$\int_{]-a,a]^{n-1}\times]0,A[}\frac{|\phi(x)|^2}{x_n^4}dx \le \frac{1}{2}\int_{]-a,a]^{n-1}\times]0,A[}\frac{|\nabla\phi(x)|^2}{x_n^2}dx.$$
(2.5)

To prove (2.5), we choose $\eta > 0$ and we integrate by parts $\int_{]-a,a]^{n-1} \times]0,A[} \frac{|\phi(x)|^2}{(\eta + x_n)^4} dx$, then thanks to the identity relation $2bc \leq b^2 + c^2$ we obtain

$$\int_{]-a,a]^{n-1}\times]0,A[} \frac{|\phi(x)|^2}{(\eta+x_n)^4} dx \le \frac{1}{2} \int_{]-a,a]^{n-1}\times]0,A[} \frac{1}{(\eta+x_n)^2} \Big| \frac{\partial\phi}{\partial x_n}(x) \Big|^2 dx$$
$$\le \frac{1}{2} \int_{]-a,a]^{n-1}\times]0,A[} \frac{|\nabla\phi(x)|^2}{x_n^2} dx.$$

Passing to the limit $(\eta \to 0)$ it leads to (2.5). Step 2. Let h be in $W^{1,\infty}(\Omega)$ such that

$$\begin{aligned} h(x) &\in [0, 1], \\ \forall x \in \Omega, \qquad h(x) &= 1 \quad \text{if} \quad \rho(x) \geq \gamma_0, \\ h(x) &= 0 \quad \text{if} \quad \rho(x) \leq \gamma_0/2 \end{aligned}$$

Let ϕ be in $H^1_{1/\rho}(\Omega)$. The function $\phi h/\rho^4$ belongs to $H^1_0(\Omega)$, therefore as a consequence of the Poincaré's inequality we obtain

$$\int_{\Omega} \frac{|\phi(x)h(x)|^2}{\rho(x)^4} dx \leq C \int_{\Omega} \left| \nabla \left(\frac{\phi(x)h(x)}{\rho(x)^4} \right) \right|^2 dx \leq C \int_{\Omega} \left(|\nabla \phi(x)|^2 + |\phi(x)|^2 \right) dx \\
\leq C \int_{\Omega} |\nabla \phi(x)|^2 dx \leq C \int_{\Omega} \frac{|\nabla \phi(x)|^2}{\rho(x)^2} dx.$$
(2.6)

Then using the local chart of Ω_r given by (2.1), the inequality (2.5) and thanks to a simple change of variables we get

$$\int_{\Omega_r} \frac{|\phi(x)(1-h(x))|^2}{\rho(x)^4} dx \le C \int_{\Omega_r} \frac{|\nabla(\phi(x)(1-h(x))|^2}{\rho(x)^2} dx \le C \int_{\Omega_r} \frac{|\nabla\phi(x)|^2 + |\phi(x)|^2}{\rho(x)^2} dx.$$

Since $\phi \in H_0^1(\Omega)$ the function ϕ/ρ belongs to $L^2(\Omega)$ and we have (2.2). Hence, adding these inequalities (r = 1, ..., N) we obtain

$$\int_{\Omega} \frac{|\phi(x)(1-h(x))|^2}{\rho(x)^4} dx \le C \int_{\Omega} \frac{|\nabla\phi(x)|^2}{\rho(x)^2} dx.$$
(2.7)

Finally $\phi/\rho^2 \in L^2(\Omega)$ and (2.6)-(2.7) lead to $\|\phi/\rho^2\|_{L^2(\Omega)} \leq C ||\phi||_{1/\rho}$ and then (2.4). \Box

2.3 Two lemmas

In the Lemma 2.3 we give sharp estimates of a function on the boundary and in a neighborhood of the boundary of Ω . The second estimate in (2.8) is used to obtain the L^2 global error.

Lemma 2.3. Let Ω be a bounded domain with a Lipschitz boundary, there exists $\gamma_0 > 0$ (see Subsection 2.2) such that for any $\gamma \in]0, \gamma_0]$ and for any $\phi \in H^1(\Omega)$ we have

$$\begin{aligned} ||\phi||_{L^{2}(\partial\Omega)} &\leq \frac{C}{\gamma^{1/2}} \left(||\phi||_{L^{2}(\widetilde{\Omega}_{\gamma})} + \gamma||\nabla\phi||_{L^{2}(\widetilde{\Omega}_{\gamma};\mathbb{R}^{n})} \right), \\ ||\phi||_{L^{2}(\widetilde{\Omega}_{\gamma})} &\leq C \left(\gamma^{1/2} ||\phi||_{L^{2}(\partial\Omega)} + \gamma||\nabla\phi||_{L^{2}(\widetilde{\Omega}_{\gamma};\mathbb{R}^{n})} \right). \end{aligned}$$
(2.8)

The constants do not depend on γ .

Proof. Let ψ be in $H^1(] - a, a[^{n-1} \times]0, A[)$. For $\eta \in]0, A[$ we have

$$\begin{aligned} ||\psi||_{L^{2}(]-a,a[^{n-1}\times\{0\})}^{2} &\leq \frac{C}{\eta} ||\psi||_{L^{2}(]-a,a[^{n-1}\times]0,\eta[)}^{2} + C\eta ||\nabla\psi||_{L^{2}(]-a,a[^{n-1}\times]0,\eta[;\mathbb{R}^{n})}^{2}, \\ ||\psi||_{L^{2}(]-a,a[^{n-1}\times]0,\eta[)}^{2} &\leq C\eta ||\psi||_{L^{2}(]-a,a[^{n-1}\times\{0\})}^{2} + C\eta^{2} ||\nabla\psi||_{L^{2}(]-a,a[^{n-1}\times]0,\eta[;\mathbb{R}^{n})}^{2}. \end{aligned}$$

The constants do not depend on η . Now, let ϕ be in $H^1(\Omega)$. We use the above estimates, the local charts of $\widetilde{\Omega}_{\gamma_0}$ given by (2.1) and a simple change of variables to obtain (2.8).

In this second lemma we show that a function in $H_0^1(\Omega)$ can be approached by functions vanishing close to the boundary of Ω . Among other things this lemma is used to give an approximation of ϕ via the scale-splitting operator $\mathcal{Q}_{\varepsilon}$ (see Lemma 2.6) and it is also used in the main projection theorem (Theorem 3.2). **Lemma 2.4.** Let ϕ be in $H_0^1(\Omega)$, there exists $\phi_{\varepsilon} \in H^1(\mathbb{R}^n)$ satisfying

$$\begin{aligned}
\phi_{\varepsilon}(x) &= 0 \quad \text{for a.e. } x \notin \tilde{\Omega}_{6\sqrt{n}\varepsilon}, \\
||\phi - \phi_{\varepsilon}||_{L^{2}(\Omega)} &\leq C\varepsilon ||\nabla\phi||_{L^{2}(\Omega;\mathbb{R}^{n})}, \quad ||\phi_{\varepsilon}||_{H^{1}(\Omega)} \leq C ||\phi||_{H^{1}(\Omega)}.
\end{aligned}$$
(2.9)

Moreover, if $\phi \in H^1_{1/\rho}(\Omega)$ then we have

$$\left\| \left(\phi - \phi_{\varepsilon} \right) / \rho \right\|_{L^{2}(\Omega)} \le C \varepsilon \| \nabla \phi \|_{1/\rho}, \qquad \| \phi_{\varepsilon} \|_{1/\rho} \le C \| \phi \|_{1/\rho}.$$
(2.10)

The constant C is independent of ε .

Proof. Let ϕ be in $H_0^1(\Omega)$. We define ϕ_{ε} by

$$\phi_{\varepsilon}(x) = \begin{cases} \frac{(\rho(x) - 6\sqrt{n\varepsilon})^+}{\rho(x)} \phi(x) & \text{for a. e. } x \in \Omega, \\ 0 & \text{for a. e. } x \in \mathbb{R}^n \setminus \overline{\Omega}. \end{cases}$$

where $\delta^+ = \max\{0, \delta\}$. The above function ϕ_{ε} belongs to $H^1(\mathbb{R}^n)$ and satisfies $\phi_{\varepsilon} = 0$ outside $\widetilde{\Omega}_{6\sqrt{n}\varepsilon}$. Then due to the fact that ϕ/ρ belongs to $L^2(\Omega)$ and verifies $\|\phi/\rho\|_{L^2(\Omega)} \leq C \|\nabla\phi\|_{L^2(\Omega;\mathbb{R}^n)}$ we obtain the estimates in (2.9). If $\phi \in H^1_{1/\rho}(\Omega)$ we use the estimate (2.4) to obtain (2.10).

2.4 Reminds and complements on the unfolding operators

In the sequel, we will make use of some definitions and results from [10] concerning the periodic unfolding method. Below we remind them briefly.

2.4.1 Some reminds

For almost every $x \in \mathbb{R}^n$, there exists an unique element in \mathbb{Z}^n denoted [x] such that

$$x = [x] + \{x\}, \qquad \{x\} \in Y.$$

• The unfolding operator $\mathcal{T}_{\varepsilon}$.

For any $\phi \in L^1(\Omega)$, the function $\mathcal{T}_{\varepsilon}(\phi) \in L^1(\Omega \times Y)$ is given by

$$\mathcal{T}_{\varepsilon}(\phi)(x,y) = \begin{cases} \phi\left(\varepsilon\left[\frac{x}{\varepsilon}\right] + \varepsilon y\right) & \text{for a.e. } (x,y) \in \widehat{\Omega}_{\varepsilon} \times Y, \\ 0 & \text{for a.e. } (x,y) \in \Lambda_{\varepsilon} \times Y. \end{cases}$$
(2.11)

Since $\Lambda_{\varepsilon} \subset \widetilde{\Omega}_{\sqrt{n}\varepsilon}$, using Proposition 2.5 in [10] we get

$$\left|\int_{\Omega}\phi(x)dx - \int_{\Omega\times Y}\mathcal{T}_{\varepsilon}(\phi)(x,y)dxdy\right| \le \int_{\Lambda_{\varepsilon}}|\phi(x)|dx \le ||\phi||_{L^{1}(\widetilde{\Omega}_{\sqrt{n}\varepsilon})}$$
(2.12)

For $\phi \in L^2(\Omega)$ we have

$$||\mathcal{T}_{\varepsilon}(\phi)||_{L^{2}(\Omega)} \leq ||\phi||_{L^{2}(\Omega)}.$$
(2.13)

We also have (see Proposition 2.5 in [10]) for $\phi \in H^1(\Omega)$ (resp. $\psi \in H^1_0(\Omega)$)

$$||\mathcal{T}_{\varepsilon}(\phi) - \phi||_{L^{2}(\widehat{\Omega}_{\varepsilon} \times Y)} \leq C\varepsilon ||\nabla\phi||_{L^{2}(\Omega;\mathbb{R}^{n})}$$

(resp. $||\mathcal{T}_{\varepsilon}(\psi) - \psi||_{L^{2}(\Omega \times Y)} \leq C\varepsilon ||\nabla\psi||_{L^{2}(\Omega;\mathbb{R}^{n})}$). (2.14)

• The local average operator $\mathcal{M}_{\varepsilon}$

For $\phi \in L^1(\mathbb{R}^n)$, the function $\mathcal{M}_{\varepsilon}(\phi) \in L^{\infty}(\mathbb{R}^n)$ is defined by

$$\mathcal{M}_{\varepsilon}(\phi)(x) = \int_{Y} \phi\Big(\varepsilon\Big[\frac{x}{\varepsilon}\Big] + \varepsilon y\Big) dy \quad \text{for a.e. } x \in \mathbb{R}^{n}.$$
(2.15)

The value of $\mathcal{M}_{\varepsilon}(\phi)$ in the cell $\varepsilon(\xi + Y)$ ($\xi \in \mathbb{Z}^n$) will be denoted $\mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi)$. In [10] we proved the following results:

For $\phi \in L^2(\Omega)$ we have

$$||\mathcal{M}_{\varepsilon}(\phi)||_{L^{2}(\Omega)} \leq C||\phi||_{L^{2}(\Omega)}, \qquad ||\mathcal{M}_{\varepsilon}(\phi) - \phi||_{H^{-1}(\Omega)} \leq C\varepsilon||\phi||_{L^{2}(\Omega)}$$
(2.16)

and for $\psi \in H_0^1(\Omega)$ (resp. $\phi \in H^1(\Omega)$) we have

$$\begin{aligned} ||\mathcal{M}_{\varepsilon}(\psi) - \psi||_{L^{2}(\Omega)} &\leq C\varepsilon ||\nabla\psi||_{L^{2}(\Omega;\mathbb{R}^{n})} \\ (\text{resp.} \quad ||\mathcal{M}_{\varepsilon}(\phi) - \phi||_{L^{2}(\widehat{\Omega}_{\varepsilon})} &\leq C\varepsilon ||\nabla\phi||_{L^{2}(\Omega;\mathbb{R}^{n})}). \end{aligned}$$
(2.17)

• The scale-splitting operator $\mathcal{Q}_{\varepsilon}$.

* For $\phi \in L^1(\mathbb{R}^n)$, the function $\mathcal{Q}_{\varepsilon}(\phi) \in W^{1,\infty}(\mathbb{R}^n)$ is given by

$$\mathcal{Q}_{\varepsilon}(\phi)(x) = \sum_{\xi \in \mathbb{Z}^n} \mathcal{M}_{\varepsilon}(\phi)(\varepsilon \xi) H_{\varepsilon,\xi}(x) \quad \text{for a.e. } x \in \mathbb{R}^n,$$

where

$$H_{\varepsilon,\xi}(x) = H\left(\frac{x - \varepsilon\xi}{\varepsilon}\right) \quad \text{with} \\ H(z) = \begin{cases} \left(1 - |z_1|\right) \left(1 - |z_2|\right) \dots \left(1 - |z_n|\right) & \text{if} \quad z \in [-1, 1]^n \\ 0 & \text{if} \quad z \in \mathbb{R}^n \setminus [-1, 1]^n \end{cases}$$

Below, we remind some results about $\mathcal{Q}_{\varepsilon}$ proved in [10] and [16]. \star For $\phi \in L^2(\mathbb{R}^n)$ we have

$$||\mathcal{Q}_{\varepsilon}(\phi)||_{L^{2}(\mathbb{R}^{n})} \leq C||\phi||_{L^{2}(\mathbb{R}^{n})}, \qquad ||\nabla \mathcal{Q}_{\varepsilon}(\phi)||_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} \leq \frac{C}{\varepsilon}||\phi||_{L^{2}(\mathbb{R}^{n})}$$
(2.18)

and

$$\mathcal{Q}_{\varepsilon}(\phi) \longrightarrow \phi$$
 strongly in $L^2(\mathbb{R}^n)$

 \star For $\phi \in H^1(\mathbb{R}^n)$ we have

$$\begin{aligned} ||\nabla \mathcal{Q}_{\varepsilon}(\phi)||_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} &\leq C ||\nabla \phi||_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})}, \\ ||\phi - \mathcal{Q}_{\varepsilon}(\phi)||_{L^{2}(\mathbb{R}^{n})} &\leq C\varepsilon ||\nabla \phi||_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} \end{aligned}$$
(2.19)

and

$$\mathcal{Q}_{\varepsilon}(\phi) \longrightarrow \phi$$
 strongly in $H^1(\mathbb{R}^n)$. (2.20)

* For $\phi \in L^2(\mathbb{R}^n)$ and $\chi \in L^2(Y)$ we have $\mathcal{Q}_{\varepsilon}(\phi)\chi\left(\left\{\frac{\cdot}{\varepsilon}\right\}\right) \in L^2(\mathbb{R}^n), \nabla \mathcal{Q}_{\varepsilon}(\phi)\chi\left(\left\{\frac{\cdot}{\varepsilon}\right\}\right) \in L^2(\mathbb{R}^n)$ and

$$\left\| \mathcal{Q}_{\varepsilon}(\phi) \chi\left(\left\{ \frac{\cdot}{\varepsilon} \right\} \right) \right\|_{L^{2}(\mathbb{R}^{n})} \leq C \|\phi\|_{L^{2}(\mathbb{R}^{n})} \|\chi\|_{L^{2}(Y)},$$

$$\left\| \mathcal{Q}_{\varepsilon}(\phi) \chi\left(\left\{ \frac{\cdot}{\varepsilon} \right\} \right) \right\|_{L^{2}(\widetilde{\Omega}_{\sqrt{n}\varepsilon})} \leq C \|\phi\|_{L^{2}(\widetilde{\widetilde{\Omega}}_{3\sqrt{n}\varepsilon})} \|\chi\|_{L^{2}(Y)}.$$
(2.21)

Moreover, if $\phi \in H^1(\mathbb{R}^n)$ then we have

$$\begin{aligned} \left\| \left(\mathcal{Q}_{\varepsilon}(\phi) - \mathcal{M}_{\varepsilon}(\phi) \right) \chi\left(\left\{\frac{\cdot}{\varepsilon}\right\}\right) \right\|_{L^{2}(\mathbb{R}^{n})} &\leq C\varepsilon \|\nabla\phi\|_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} \|\chi\|_{L^{2}(Y)}, \\ \left\| \nabla\mathcal{Q}_{\varepsilon}(\phi) \chi\left(\left\{\frac{\cdot}{\varepsilon}\right\}\right) \right\|_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} &\leq C \|\nabla\phi\|_{L^{2}(\mathbb{R}^{n};\mathbb{R}^{n})} \|\chi\|_{L^{2}(Y)}, \\ \left\| \nabla\mathcal{Q}_{\varepsilon}(\phi) \chi\left(\left\{\frac{\cdot}{\varepsilon}\right\}\right) \right\|_{L^{2}(\widetilde{\Omega}_{\sqrt{n}\varepsilon};\mathbb{R}^{n})} &\leq C \|\nabla\phi\|_{L^{2}(\widetilde{\widetilde{\Omega}}_{3\sqrt{n}\varepsilon};\mathbb{R}^{n})} \|\chi\|_{L^{2}(Y)}, \end{aligned}$$
(2.22)

2.4.2 Some complements

In this subsection, we extend some results given above to functions belonging to $H^1_{\rho}(\Omega)$. These technical complements intervene in the proofs of the projection theorems and in the Theorem 6.1.

Lemma 2.5. For $\phi \in H^1_{\rho}(\Omega)$ we have

$$\forall i \in \{1, \dots, n\}, \qquad \begin{aligned} &||\rho \big(\mathcal{M}_{\varepsilon}(\phi) - \phi \big)||_{L^{2}(\Omega)} \leq C\varepsilon ||\phi||_{\rho}, \\ &||\rho \big(\phi(\cdot + \varepsilon \mathbf{e}_{i}) - \phi \big)||_{L^{2}(\Omega)} \leq C\varepsilon ||\phi||_{\rho}, \\ &||\rho \big(\mathcal{M}_{\varepsilon}(\phi)(\cdot + \varepsilon \mathbf{e}_{i}) - \mathcal{M}_{\varepsilon}(\phi) \big)||_{L^{2}(\Omega)} \leq C\varepsilon ||\phi||_{\rho}. \end{aligned}$$
(2.23)

For $\phi \in L^2_{1/\rho}(\Omega)$ we have

$$||\mathcal{M}_{\varepsilon}(\phi) - \phi||_{(H^{1}_{\rho}(\Omega))'} \le C\varepsilon ||\phi/\rho||_{L^{2}(\Omega)}.$$
(2.24)

The constants do not depend on ε .

Proof. Step 1. We prove $(2.23)_1$. Let ϕ be in $H^1_{\rho}(\Omega)$ and let $\varepsilon(\xi + Y)$ be a cell included in Ω . Case 1: $\rho(\varepsilon\xi) \ge 2\sqrt{n\varepsilon}$. In this case, observing that

$$1 \le \frac{\max_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}}{\min_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}} \le 3$$

and thanks to the Poincaré-Wirtinger's inequality we obtain

$$\int_{\varepsilon(\xi+Y)} [\rho(x)]^2 |\mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi) - \phi(x)|^2 dx \leq [\max_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}]^2 \int_{\varepsilon(\xi+Y)} |\mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi) - \phi(x)|^2 dx$$
$$\leq [\max_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}]^2 C\varepsilon^2 \int_{\varepsilon(\xi+Y)} |\nabla\phi(x)|^2 dx$$
$$\leq C\varepsilon^2 \int_{\varepsilon(\xi+Y)} [\rho(x)]^2 |\nabla\phi(x)|^2 dx.$$

Case 2: $\rho(\varepsilon\xi) \leq 2\sqrt{n}\varepsilon$. In this case we have

$$\int_{\varepsilon(\xi+Y)} [\rho(x)]^2 |\mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi) - \phi(x)|^2 dx \le C\varepsilon^2 \int_{\varepsilon(\xi+Y)} |\phi(x)|^2 dx$$

The cases 1 and 2 lead to

$$\int_{\widehat{\Omega}_{\varepsilon}} [\rho(x)]^2 |\mathcal{M}_{\varepsilon}(\phi)(x) - \phi(x)|^2 dx \le C\varepsilon^2 \int_{\widehat{\Omega}_{\varepsilon}} \left([\rho(x)]^2 |\nabla\phi(x)|^2 + |\phi(x)|^2 \right) dx.$$
(2.25)

Since $\Lambda_{\varepsilon}\subset \widetilde{\Omega}_{\sqrt{n}\varepsilon}$ and due to Lemma 2.1 we get

$$\int_{\Lambda_{\varepsilon}} [\rho(x)]^2 |\mathcal{M}_{\varepsilon}(\phi)(x) - \phi(x)|^2 dx \le C \varepsilon^2 \int_{\widetilde{\Omega}_{c_0\sqrt{n\varepsilon}}} |\phi(x)|^2 dx$$

which in turn with (2.25) gives $(2.23)_1$. Proceeding in the same way we obtain $(2.23)_2$ and $(2.23)_3$.

Step 2. We prove (2.24).

Let ϕ be in $L^2_{1/\rho}(\Omega)$ and $\psi \in H^1_{\rho}(\Omega)$. We have

$$\int_{\widehat{\Omega}_{\varepsilon}} \left(\mathcal{M}_{\varepsilon}(\phi) - \phi \right) \psi = \int_{\widehat{\Omega}_{\varepsilon}} \left(\mathcal{M}_{\varepsilon}(\psi) - \psi \right) \phi$$

Consequently we obtain

$$\left| \int_{\Omega} \left(\mathcal{M}_{\varepsilon}(\phi) - \phi \right) \psi - \int_{\Omega} \left(\mathcal{M}_{\varepsilon}(\psi) - \psi \right) \phi \right| \leq \int_{\Lambda_{\varepsilon}} \left| \left(\mathcal{M}_{\varepsilon}(\phi) - \phi \right) \psi \right| + \int_{\Lambda_{\varepsilon}} \left| \left(\mathcal{M}_{\varepsilon}(\psi) - \psi \right) \phi \right| \\ \leq C \left(||\phi||_{L^{2}(\Lambda_{\varepsilon})} + ||\mathcal{M}_{\varepsilon}(\phi)||_{L^{2}(\Lambda_{\varepsilon})} \right) ||\psi||_{L^{2}(\Omega)}.$$

The inclusion $\Lambda_{\varepsilon} \subset \widetilde{\Omega}_{\sqrt{n}\varepsilon}$, the fact that $\phi \in L^2_{1/\rho}(\Omega)$ and the estimates $(2.3)_1$ - $(2.23)_1$ lead to

$$\int_{\Omega} \left(\mathcal{M}_{\varepsilon}(\phi) - \phi \right) \psi \leq C \varepsilon ||\phi/\rho||_{L^{2}(\Omega)} ||\psi||_{\rho}.$$

Hence (2.24) is proved.

Lemma 2.6. For $\phi \in H^1_{\rho}(\Omega)$ we have

$$||\rho(\mathcal{Q}_{\varepsilon}(\phi) - \phi)||_{L^{2}(\Omega)} \le C\varepsilon ||\phi||_{\rho}$$
(2.26)

For $\phi \in H^1_{1/\rho}(\Omega)$ and ϕ_{ε} given by Lemma 2.4 we have

$$\begin{aligned} \|\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})\|_{1/\rho} &\leq C \|\phi\|_{1/\rho}, \qquad \left\| \left(\phi - \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon}) \right) / \rho \right\|_{L^{2}(\Omega)} \leq C \varepsilon \|\phi\|_{1/\rho}, \\ \forall \mathbf{i} = i_{1} \mathbf{e}_{1} + \ldots + i_{n} \mathbf{e}_{n}, \qquad (i_{1}, \ldots, i_{n}) \in \{0, 1\}^{n} \\ \left\| \left(\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\cdot + \varepsilon \mathbf{i}) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \right) / \rho \right\|_{L^{2}(\Omega)} \leq C \varepsilon \|\phi\|_{1/\rho}. \end{aligned}$$

$$(2.27)$$

For $\phi \in L^2(\mathbb{R}^n)$ and $\chi \in L^2(Y)$

$$\left\| \left(\mathcal{M}_{\varepsilon}(\rho\phi) - \rho \mathcal{M}_{\varepsilon}(\phi) \right) \chi\left(\left\{ \frac{\cdot}{\varepsilon} \right\} \right) \right\|_{L^{2}(\mathbb{R}^{n})} \leq C\varepsilon \|\phi\|_{L^{2}(\mathbb{R}^{n})} \|\chi\|_{L^{2}(Y)}.$$
(2.28)

For $\phi \in H^1_{\rho}(\Omega)$ and $\chi \in L^2(Y)$

$$\left\| \rho \left(\mathcal{Q}_{\varepsilon}(\phi) - \mathcal{M}_{\varepsilon}(\phi) \right) \chi \left(\left\{ \frac{\cdot}{\varepsilon} \right\} \right) \right\|_{L^{2}(\Omega)} \leq C \varepsilon \|\phi\|_{\rho} \|\chi\|_{L^{2}(Y)},$$

$$\left\| \rho \nabla \mathcal{Q}_{\varepsilon}(\phi) \chi \left(\left\{ \frac{\cdot}{\varepsilon} \right\} \right) \right\|_{L^{2}(\Omega)} \leq C \|\phi\|_{\rho} \|\chi\|_{L^{2}(Y)}.$$

$$(2.29)$$

The constants do not depend on ε .

Proof. Step 1. Let ϕ be in $H^1_{\rho}(\Omega)$. We first prove

$$||\rho (\mathcal{Q}_{\varepsilon}(\phi) - \mathcal{M}_{\varepsilon}(\phi))||_{L^{2}(\Omega)} \leq C\varepsilon ||\phi||_{\rho}.$$
(2.30)

To do that, we proceed as in the proof of $(2.23)_1$. Let $\varepsilon(\xi + Y)$ be a cell included in Ω . Case 1: $\rho(\varepsilon\xi) \ge 3\sqrt{n}\varepsilon$. In this case we have

$$1 \le \frac{\max_{z \in \varepsilon(\xi+Y)}\{\rho(z)\}}{\min_{z \in \varepsilon(\xi+2Y)}\{\rho(z)\}} \le 4 \qquad \text{and} \qquad 1 \le \frac{\max_{z \in \varepsilon(\xi+2Y)}\{\rho(z)\}}{\min_{z \in \varepsilon(\xi+Y)}\{\rho(z)\}} \le \frac{5}{2}$$

By definition of $\mathcal{Q}_{\varepsilon}(\phi)$ we deduce that

$$\begin{split} \int_{\varepsilon(\xi+Y)} [\rho(x)]^2 |\mathcal{Q}_{\varepsilon}(\phi)(x) - \mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi)|^2 dx &\leq [\max_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}]^2 \int_{\varepsilon(\xi+Y)} |\mathcal{Q}_{\varepsilon}(\phi)(x) - \mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi)|^2 dx \\ &\leq [\max_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}]^2 C \varepsilon^2 \int_{\varepsilon(\xi+2Y)} |\nabla\phi(x)|^2 dx \\ &\leq C \varepsilon^2 \int_{\varepsilon(\xi+2Y)} [\rho(x)]^2 |\nabla\phi(x)|^2 dx. \end{split}$$

Case 2: $\rho(\varepsilon\xi) \leq 3\sqrt{n}\varepsilon$. Then again by definition of $\mathcal{Q}_{\varepsilon}(\phi)$ we get

$$\int_{\varepsilon(\xi+Y)} [\rho(x)]^2 |\mathcal{Q}_{\varepsilon}(\phi)(x) - \mathcal{M}_{\varepsilon}(\phi)(\varepsilon\xi)|^2 dx \le C\varepsilon^2 \int_{\varepsilon(\xi+2Y)} |\phi(x)|^2 dx.$$

As a consequence of both cases we get

$$\int_{\widehat{\Omega}_{\varepsilon}} [\rho(x)]^2 |\mathcal{Q}_{\varepsilon}(\phi)(x) - \mathcal{M}_{\varepsilon}(\phi)(x)|^2 dx \le C\varepsilon^2 \int_{\Omega} \left([\rho(x)]^2 |\nabla\phi(x)|^2 + |\phi(x)|^2 \right) dx. \quad (2.31)$$

Furthermore we have

$$\int_{\Lambda_{\varepsilon}} [\rho(x)]^2 |\mathcal{Q}_{\varepsilon}(\phi)(x)|^2 dx \le C\varepsilon^2 \int_{\Lambda_{\varepsilon}} |\mathcal{Q}_{\varepsilon}(\phi)(x)|^2 dx \le C\varepsilon^2 \int_{\Omega} |\phi(x)|^2 dx$$

which with (2.31) lead to (2.30). Then as a consequence of $(2.23)_1$ and (2.30) we get (2.26).

Step 2. We prove $(2.27)_1$.

Let ϕ be in $H^1_{1/\rho}(\Omega)$ and ϕ_{ε} given by Lemma 2.4. Due to the fact that $\phi_{\varepsilon}(x) = 0$ for a.e. $x \in \mathbb{R}^n \setminus \overline{\widetilde{\Omega}}_{6\sqrt{n}\varepsilon}$, hence $\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) = 0$ for every $x \in \Omega$ such that $\rho(x) \leq 4\sqrt{n}\varepsilon$. Again we take a cell $\varepsilon(\xi + Y)$ included in Ω such that $\rho(\varepsilon\xi) \geq 3\sqrt{n}\varepsilon$. The values taken by $\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})$ in the cell $\varepsilon(\xi + Y)$ depend only on the values of ϕ_{ε} in $\varepsilon(\xi + 2Y)$. Then we have

$$\begin{split} &\int_{\varepsilon(\xi+Y)} \frac{1}{[\rho(x)]^2} |\nabla \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x)|^2 dx \leq \frac{C}{[\min_{x \in \varepsilon(\xi+Y)} \{\rho(x)\}]^2} \int_{\varepsilon(\xi+2Y)} |\nabla \phi_{\varepsilon}(x)|^2 dx \\ &\leq C \frac{[\max_{x \in \varepsilon(\xi+2Y)} \{\rho(x)\}]^2}{[\min_{x \in \varepsilon(\xi+Y)} \{\rho(x)\}]^2} \int_{\varepsilon(\xi+2Y)} \frac{1}{[\rho(x)]^2} |\nabla \phi_{\varepsilon}(x)|^2 dx \leq C \int_{\varepsilon(\xi+2Y)} \frac{1}{[\rho(x)]^2} |\nabla \phi_{\varepsilon}(x)|^2 dx. \end{split}$$

Adding all these inequalities gives

$$\int_{\widetilde{\Omega}_{4\sqrt{n\varepsilon}}} \frac{1}{[\rho(x)]^2} |\nabla \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x)|^2 dx \le C \int_{\Omega} \frac{1}{[\rho(x)]^2} |\nabla \phi_{\varepsilon}(x)|^2 dx$$

Since $\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) = 0$ for every $x \in \Omega$ such that $\rho(x) \leq 4\sqrt{n\varepsilon}$, we get $||\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})||_{1/\rho} \leq C||\phi_{\varepsilon}||_{1/\rho}$. We conclude using $(2.10)_2$.

Step 3. Now we prove $(2.27)_2$. Again we consider a cell $\varepsilon(\xi + Y)$ included in Ω such that $\rho(\varepsilon\xi) \geq 3\sqrt{n\varepsilon}$. We have

$$\begin{split} &\int_{\varepsilon(\xi+Y)} \frac{1}{[\rho(x)]^2} |\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) - \phi_{\varepsilon}(x)|^2 dx \leq \frac{C}{[\min_{x \in \varepsilon(\xi+Y)} \{\rho(x)\}]^2} \int_{\varepsilon(\xi+Y)} |\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) - \phi_{\varepsilon}(x)|^2 dx \\ &\leq \frac{C}{[\min_{x \in \varepsilon(\xi+Y)} \{\rho(x)\}]^2} \sum_{\mathbf{i} \in \{0,1\}^n} \int_{\varepsilon(\xi+\mathbf{i}+Y)} |\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon\xi + \varepsilon\mathbf{i}) - \phi_{\varepsilon}(x)|^2 dx \\ &\leq C\varepsilon^2 \frac{[\max_{z \in \varepsilon(\xi+2Y)} \{\rho(z)\}]^2}{[\min_{z \in \varepsilon(\xi+Y)} \{\rho(z)\}]^2} \int_{\varepsilon(\xi+2Y)} \frac{1}{[\rho(x)]^2} |\nabla\phi_{\varepsilon}(x)|^2 dx \leq C\varepsilon^2 \int_{\varepsilon(\xi+2Y)} \frac{1}{[\rho(x)]^2} |\nabla\phi_{\varepsilon}(x)|^2 dx. \end{split}$$

Hence we get

$$\int_{\widetilde{\Omega}_{4\sqrt{n}\varepsilon}} \frac{1}{[\rho(x)]^2} |\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) - \phi_{\varepsilon}(x)|^2 dx \le C\varepsilon^2 \int_{\Omega} \frac{1}{[\rho(x)]^2} |\nabla\phi_{\varepsilon}(x)|^2 dx$$

The above estimate and the fact that $\mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})(x) - \phi_{\varepsilon}(x) = 0$ for a.e. $x \in \Omega$ such that $\rho(x) \leq 4\sqrt{n\varepsilon}$ yield $||(\phi_{\varepsilon} - \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon}))/\rho||_{L^{2}(\Omega)} \leq C\varepsilon ||\phi_{\varepsilon}||_{1/\rho}$. We conclude using both estimates in (2.10).

Proceeding as in the Steps 2 and 3 we obtain $(2.27)_3$, (2.28) and (2.29).

3 Two new projection theorems

Theorem 3.1. Let ϕ be in $H^1_{1/\rho}(\Omega)$. There exists $\widehat{\phi}_{\varepsilon} \in H^1_{per}(Y; L^2(\Omega))$ such that

$$\begin{cases} ||\widehat{\phi}_{\varepsilon}||_{H^{1}(Y;L^{2}(\Omega))} \leq C\{||\phi||_{L^{2}(\Omega)} + \varepsilon||\nabla\phi||_{[L^{2}(\Omega)]^{n}}\}\\ ||\mathcal{T}_{\varepsilon}(\phi) - \widehat{\phi}_{\varepsilon}||_{H^{1}(Y;(H^{1}_{\rho}(\Omega))')} \leq C\varepsilon(||\phi/\rho||_{L^{2}(\Omega)} + \varepsilon||\phi||_{1/\rho}). \end{cases}$$
(3.1)

The constants depend only on n and $\partial \Omega$.

Proof. Here, we proceed as in the proof of Proposition 3.3 in [15]. We first reintroduce the open sets $\widehat{\Omega}_{\varepsilon,i}$ and the "double" unfolding operators $\mathcal{T}_{\varepsilon,i}$. We set

$$\widehat{\Omega}_{\varepsilon,i} = \widehat{\Omega}_{\varepsilon} \cap \left(\widehat{\Omega}_{\varepsilon} - \varepsilon \mathbf{e}_i\right), \qquad K_i = \operatorname{interior}\left(\overline{Y} \cup (\mathbf{e}_i + \overline{Y})\right), \quad i \in \{1, \dots, n\}.$$

The unfolding operator $\mathcal{T}_{\varepsilon,i}$ from $L^2(\Omega)$ into $L^2(\Omega \times K_i)$ is defined by

$$\forall \psi \in L^2(\Omega), \qquad \mathcal{T}_{\varepsilon,i}(\psi)(x,y) = \begin{cases} \psi \left(\varepsilon \left[\frac{x}{\varepsilon} \right]_Y + \varepsilon y \right) & \text{for } x \in \widehat{\Omega}_{\varepsilon,i} \text{ and for a.e. } y \in K_i, \\ 0 & \text{for } x \in \Omega \setminus \overline{\widehat{\Omega}}_{\varepsilon,i} \text{ and for a.e. } y \in K_i. \end{cases}$$

The restriction of $\mathcal{T}_{\varepsilon,i}(\psi)$ to $\widehat{\Omega}_{\varepsilon,i} \times Y$ is equal to $\mathcal{T}_{\varepsilon}(\psi)$.

Step 1. Let us first take $\phi \in L^2_{1/\rho}(\Omega)$. We set $\psi = \frac{1}{\rho}\phi$ and we evaluate the difference $\mathcal{T}_{\varepsilon,i}(\phi)(.,.+\mathbf{e}_i) - \mathcal{T}_{\varepsilon,i}(\phi)$ in $L^2(Y; (H^1_\rho(\Omega))')$. For any $\Psi \in H^1_\rho(\Omega)$ a change of variables gives for a. e. $y \in Y$

$$\int_{\Omega} \mathcal{T}_{\varepsilon,i}(\phi)(x,y+\mathbf{e}_i)\Psi(x)dx = \int_{\widehat{\Omega}_{\varepsilon,i}} \mathcal{T}_{\varepsilon}(\phi)(x+\varepsilon\mathbf{e}_i,y)\Psi(x)dx$$
$$= \int_{\widehat{\Omega}_{\varepsilon,i}+\varepsilon\mathbf{e}_i} \mathcal{T}_{\varepsilon}(\phi)(x,y)\Psi(x-\varepsilon\mathbf{e}_i)dx$$

Then we obtain for a. e. $y \in Y$

$$\begin{split} & \left| \int_{\Omega} \left\{ \mathcal{T}_{\varepsilon,i}(\phi)(.,y+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\phi)(.,y) \right\} \Psi - \int_{\widehat{\Omega}_{\varepsilon,i}} \mathcal{T}_{\varepsilon}(\psi)(.,y) \rho \left\{ \Psi(.-\varepsilon \mathbf{e}_{i}) - \Psi \right\} \right| \\ \leq & \left| \int_{\widehat{\Omega}_{\varepsilon,i}} \mathcal{T}_{\varepsilon}(\psi)(.,y) \left(\mathcal{T}_{\varepsilon}(\rho) - \rho \right) \left\{ \Psi(.-\varepsilon \mathbf{e}_{i}) - \Psi \right\} \right| + C ||\mathcal{T}_{\varepsilon}(\phi)(.,y)||_{L^{2}(\widetilde{\Omega}_{2\sqrt{n}\varepsilon})} ||\Psi||_{L^{2}(\widetilde{\Omega}_{2\sqrt{n}\varepsilon})} \\ \end{aligned}$$

Estimate $(2.23)_2$ leads to

$$||\rho(\Psi(.-\varepsilon \mathbf{e}_i)-\Psi)||_{L^2(\widehat{\Omega}_{\varepsilon,i})} \le C\varepsilon ||\Psi||_{\rho} \qquad \forall i \in \{1,\ldots,n\}.$$

We have

$$||\mathcal{T}_{\varepsilon}(\rho) - \rho||_{L^{\infty}(\Omega)} \le C\varepsilon.$$
(3.2)

The above inequalities imply

$$< \mathcal{T}_{\varepsilon,i}(\phi)(.,y+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\phi)(.,y), \Psi >_{(H^{1}_{\rho}(\Omega))',H^{1}_{\rho}(\Omega)}$$

$$= \int_{\Omega} \{\mathcal{T}_{\varepsilon,i}(\phi)(x,y+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\phi)(x,y)\}\Psi(x)dx$$

$$\le C\varepsilon ||\Psi||_{\rho} ||\mathcal{T}_{\varepsilon}(\psi)(.,y)||_{L^{2}(\Omega)} + C\varepsilon ||\Psi||_{L^{2}(\Omega)} ||\mathcal{T}_{\varepsilon}(\psi)(.,y)||_{L^{2}(\Omega)}$$

$$+ C ||\mathcal{T}_{\varepsilon}(\phi)(.,y)||_{L^{2}(\tilde{\Omega}_{2\sqrt{n}\varepsilon})} ||\Psi||_{L^{2}(\tilde{\Omega}_{2\sqrt{n}\varepsilon})}.$$

Therefore, for a.e. $y \in Y$ we have

$$\left|\left|\mathcal{T}_{\varepsilon,i}(\phi)(.,y+\mathbf{e}_{i})-\mathcal{T}_{\varepsilon,i}(\phi)(.,y)\right|\right|_{(H^{1}_{\rho}(\Omega))'} \leq C\varepsilon \left\|\mathcal{T}_{\varepsilon}(\psi)(.,y)\right\|_{L^{2}(\Omega)} + C\left|\left|\mathcal{T}_{\varepsilon}(\phi)(.,y)\right|\right|_{L^{2}(\widetilde{\Omega}_{2\sqrt{n}\varepsilon})}$$

which leads to the following estimate of the difference between $\mathcal{T}_{\varepsilon,i}(\phi)|_{\Omega \times Y}$ and one of its translated :

$$\begin{aligned} ||\mathcal{T}_{\varepsilon,i}(\phi)(.,.+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\phi)||_{L^{2}(Y;(H^{1}_{\rho}(\Omega))')} &\leq C\varepsilon ||\phi/\rho||_{L^{2}(\Omega)} + C||\phi||_{L^{2}(\widetilde{\Omega}_{2\sqrt{n}\varepsilon})} \\ &\leq C\varepsilon ||\phi/\rho||_{L^{2}(\Omega)}. \end{aligned}$$
(3.3)

The constant depends only on the boundary of Ω .

Step 2. Let $\phi \in H^1_{1/\rho}(\Omega)$. The above estimate (3.3) applied to ϕ and its partial derivatives give

$$||\mathcal{T}_{\varepsilon,i}(\phi)(.,..+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\phi)||_{L^{2}(Y;(H^{1}_{\rho}(\Omega))')} \leq C\varepsilon||\phi/\rho||_{L^{2}(\Omega)}$$
$$||\mathcal{T}_{\varepsilon,i}(\nabla\phi)(.,..+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon,i}(\nabla\phi)||_{[L^{2}(Y;(H^{1}_{\rho}(\Omega))']^{n})} \leq C\varepsilon||\phi||_{1/\rho}.$$

which in turn lead to (we recall that $\nabla_y (\mathcal{T}_{\varepsilon,i}(\phi)) = \varepsilon \mathcal{T}_{\varepsilon,i}(\nabla \phi)$).

$$||\mathcal{T}_{\varepsilon,i}(\phi)(.,.+\mathbf{e}_i) - \mathcal{T}_{\varepsilon,i}(\phi)||_{H^1(Y;(H^1_{\rho}(\Omega))')} \leq C\varepsilon \left(||\phi/\rho||_{L^2(\Omega)} + \varepsilon ||\phi||_{1/\rho}\right).$$

From these inequalities for $i \in \{1, ..., n\}$ we deduce the estimate of the difference of the traces of the function $y \longrightarrow \mathcal{T}_{\varepsilon}(\phi)(., y)$ on the faces $Y_i \doteq \{y \in \overline{Y} \mid y_i = 0\}$ and $\mathbf{e}_i + Y_i$

$$||\mathcal{T}_{\varepsilon}(\phi)(.,.,+\mathbf{e}_{i}) - \mathcal{T}_{\varepsilon}(\phi)||_{H^{1/2}(Y_{i};(H^{1}_{\rho}(\Omega))')} \leq C\varepsilon \left(||\phi/\rho||_{L^{2}(\Omega)} + \varepsilon ||\phi||_{1/\rho}\right).$$
(3.4)

These estimates $(i \in \{1, \ldots, n\})$ give a measure of the periodic defect of the function $y \longrightarrow \mathcal{T}_{\varepsilon}(\phi)(., y)$ (see [15]).

Then we decompose $\mathcal{T}_{\varepsilon}(\phi)$ into the sum of an element belonging to $H^1_{per}(Y; L^2(\Omega))$ and one to $(H^1(Y; L^2(\Omega)))^{\perp}$ (the orthogonal of $H^1_{per}(Y; L^2(\Omega))$ in $H^1(Y; L^2(\Omega))$, see [15])

$$\mathcal{T}_{\varepsilon}(\phi) = \widehat{\phi}_{\varepsilon} + \overline{\phi}_{\varepsilon}, \qquad \widehat{\phi}_{\varepsilon} \in H^{1}_{per}(Y; L^{2}(\Omega)), \qquad \overline{\phi}_{\varepsilon} \in \left(H^{1}(Y; L^{2}(\Omega))\right)^{\perp}.$$
(3.5)

The function $y \longrightarrow \mathcal{T}_{\varepsilon}(\phi)(., y)$ takes its values in a finite dimensional space,

$$\overline{\phi}_{\varepsilon}(.,..) = \sum_{\xi \in \Xi_{\varepsilon}} \overline{\phi}_{\varepsilon,\xi}(..) \chi_{\varepsilon,\xi}(.)$$

where $\chi_{\varepsilon,\xi}(.)$ is the characteristic function of the cell $\varepsilon(\xi + Y)$ and where $\overline{\phi}_{\varepsilon,\xi}(..) \in (H^1(Y))^{\perp}$ (the orthogonal of $H^1_{per}(Y)$ in $H^1(Y)$, see [15]). The decomposition (3.5) is the same in $H^1(Y; (H^1_{\rho}(\Omega))')$ and we have

$$||\widehat{\phi}_{\varepsilon}||^{2}_{H^{1}(Y;L^{2}(\Omega))} + ||\overline{\phi}_{\varepsilon}||^{2}_{H^{1}(Y;L^{2}(\Omega))} = ||\mathcal{T}_{\varepsilon}(\phi)||^{2}_{H^{1}(Y;L^{2}(\Omega))} \leq C\{||\phi||_{L^{2}(\Omega)} + \varepsilon ||\nabla\phi||_{[L^{2}(\Omega)]^{n}}\}^{2}.$$

It gives the first inequality in (3.1) and the estimate of $\overline{\phi}_{\varepsilon}$ in $H^1(Y; L^2(\Omega))$. From Theorem 2.2 in [15] and (3.4) we obtain a finer estimate of $\overline{\phi}_{\varepsilon}$ in $H^1(Y; (H^1_{\rho}(\Omega))')$

$$||\phi_{\varepsilon}||_{H^{1}(Y;(H^{1}_{\rho}(\Omega))')} \leq C\varepsilon(||\phi/\rho||_{L^{2}(\Omega)} + \varepsilon||\phi||_{1/\rho})$$

It is the second inequality in (3.1).

Theorem 3.2. For $\phi \in H^1_{1/\rho}(\Omega)$, there exists $\widehat{\phi}_{\varepsilon} \in H^1_{per}(Y; L^2(\Omega))$ such that

$$\begin{aligned} ||\widehat{\phi}_{\varepsilon}||_{H^{1}(Y;L^{2}(\Omega))} &\leq C||\nabla\phi||_{[L^{2}(\Omega)]^{n}}, \\ ||\mathcal{T}_{\varepsilon}(\nabla\phi) - \nabla\phi - \nabla_{y}\widehat{\phi}_{\varepsilon}||_{[L^{2}(Y;(H^{1}_{\rho}(\Omega))')]^{n}} &\leq C\varepsilon||\phi||_{1/\rho}. \end{aligned}$$
(3.6)

The constants depend only on $\partial\Omega$.

Proof. Let ϕ be in $H^1_{1/\rho}(\Omega)$ and $\psi = \phi/\rho \in H^1_0(\Omega)$. The function ϕ is extended by 0 outside of Ω . We decompose ϕ as

$$\phi = \Phi + \varepsilon \underline{\phi}$$
, where $\Phi = \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon})$ and $\underline{\phi} = \frac{1}{\varepsilon} \Big(\phi - \mathcal{Q}_{\varepsilon}(\phi_{\varepsilon}) \Big)$

where ϕ_{ε} is given by Lemma 2.4. We have Φ and $\phi \in H_0^1(\Omega)$ and due to (2.27) we get the following estimates:

$$||\Phi||_{1/\rho} + \varepsilon ||\underline{\phi}||_{1/\rho} + ||\underline{\phi}/\rho||_{L^2(\Omega)} \le C ||\phi||_{1/\rho}.$$

$$(3.7)$$

The projection Theorem 3.1 applied to $\underline{\phi} \in H^1_{1/\rho}(\Omega)$ gives an element $\widehat{\phi}_{\varepsilon}$ in $H^1_{per}(Y; L^2(\Omega))$ such that

$$\begin{aligned} ||\phi_{\varepsilon}||_{H^{1}(Y;L^{2}(\Omega))} &\leq C ||\phi||_{1/\rho}, \\ ||\mathcal{T}_{\varepsilon}(\underline{\phi}) - \widehat{\phi}_{\varepsilon}||_{H^{1}(Y;(H^{1}_{\rho}(\Omega))')} &\leq C\varepsilon ||\phi||_{1/\rho}. \end{aligned}$$
(3.8)

Now we evaluate $||\mathcal{T}_{\varepsilon}(\nabla\Phi) - \nabla\Phi||_{[L^{2}(Y;(H^{1}_{\rho}(\Omega))')]^{n}}$.

From (2.24), $(2.27)_1$ and (3.7) we get

$$\|\nabla\Phi - \mathcal{M}_{\varepsilon}(\nabla\Phi)\|_{(H^{1}_{\rho}(\Omega;\mathbb{R}^{n}))'} \leq C\varepsilon \|\phi\|_{1/\rho}.$$
(3.9)

We set

$$H^{(1)}(z) = \begin{cases} (1 - |z_2|)(1 - |z_3|) \dots (1 - |z_n|) & \text{if } z = (z_1, z_2, \dots, z_n) \in [-1, 1]^n, \\ 0 & \text{if } z \in \mathbb{R}^n \setminus [-1, 1]^n. \end{cases}$$
$$\mathbf{I} = \left\{ \mathbf{i} \mid \mathbf{i} = i_2 \mathbf{e}_2 + \dots + i_n \mathbf{e}_n, \quad (i_2, \dots, i_n) \in \{0, 1\}^{n-1} \right\}$$

For $\xi \in \mathbb{Z}^n$ and for every $(x, y) \in \varepsilon(\xi + Y) \times Y$ we have

$$\mathcal{T}_{\varepsilon} \Big(\frac{\partial \Phi}{\partial x_1} \Big)(x, y) = \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \big(\varepsilon(\xi + \mathbf{e}_1 + \mathbf{i})\big) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \big(\varepsilon(\xi + \mathbf{i})\big)}{\varepsilon} H^{(1)}(y - \mathbf{i})$$
$$\mathcal{M}_{\varepsilon} \Big(\frac{\partial \Phi}{\partial x_1} \Big)(\varepsilon\xi) = \frac{1}{2^{n-1}} \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \big(\varepsilon(\xi + \mathbf{e}_1 + \mathbf{i})\big) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \big(\varepsilon(\xi + \mathbf{i})\big)}{\varepsilon}.$$

Now, let us take $\psi \in H^1_{\rho}(\Omega)$. We recall that $\phi_{\varepsilon}(x) = 0$ for a.e. $x \in \mathbb{R}^n \setminus \overline{\widetilde{\Omega}}_{6\sqrt{n}\varepsilon}$, hence $\Phi(x) = 0$ for $x \in \mathbb{R}^n \setminus \overline{\widetilde{\Omega}}_{3\sqrt{n}\varepsilon}$; as a first consequence $\mathcal{M}_{\varepsilon}\left(\frac{\partial \Phi}{\partial x_1}\right) = 0$ in Λ_{ε} . For $y \in Y$ we have

$$<\mathcal{T}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)(.,y)-\mathcal{M}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big),\psi>_{(H^{1}_{\rho}(\Omega))',H^{1}_{\rho}(\Omega)}=\int_{\Omega}\Big\{\mathcal{T}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)(x,y)-\mathcal{M}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)(x)\Big\}\psi(x)dx\\=\int_{\widehat{\Omega}_{\varepsilon}}\Big\{\mathcal{T}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)(x,y)-\mathcal{M}_{\varepsilon}\Big(\frac{\partial\Phi}{\partial x_{1}}\Big)(x)\Big\}\mathcal{M}_{\varepsilon}(\psi)(x)dx.$$

Besides we have

$$\begin{split} \int_{\widehat{\Omega}_{\varepsilon}} \mathcal{M}_{\varepsilon} \Big(\frac{\partial \Phi}{\partial x_{1}} \Big)(x) \mathcal{M}_{\varepsilon}(\psi)(x) dx &= \varepsilon^{n} \sum_{\xi \in \mathbb{Z}^{n}} \mathcal{M}_{\varepsilon} \Big(\frac{\partial \Phi}{\partial x_{1}} \Big)(\varepsilon\xi) \mathcal{M}_{\varepsilon}(\psi)(\varepsilon\xi) \\ &= \frac{\varepsilon^{n}}{2^{n-1}} \sum_{\xi \in \mathbb{Z}^{n}} \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \big(\varepsilon(\xi + \mathbf{e}_{1} + \mathbf{i})\big) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \big(\varepsilon(\xi + \mathbf{i})\big)}{\varepsilon} \mathcal{M}_{\varepsilon}(\psi)(\varepsilon\xi) \\ &= \frac{\varepsilon^{n}}{2^{n-1}} \sum_{\xi \in \mathbb{Z}^{n}} \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon}(\psi) \big(\varepsilon(\xi - \mathbf{e}_{1})\big) - \mathcal{M}_{\varepsilon}(\psi) \big(\varepsilon\xi\big)}{\varepsilon} \mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon(\xi + \mathbf{i})) \end{split}$$

and

$$\begin{split} &\int_{\widehat{\Omega}_{\varepsilon}} \mathcal{T}_{\varepsilon} \Big(\frac{\partial \Phi}{\partial x_{1}} \Big)(x, y) \mathcal{M}_{\varepsilon}(\psi)(x) dx \\ = &\varepsilon^{n} \sum_{\xi \in \mathbb{Z}^{n}} \sum_{\mathbf{i} \in \mathbf{I}} \Big[\frac{\mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \big(\varepsilon(\xi + \mathbf{e}_{1} + \mathbf{i})\big) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \big(\varepsilon(\xi + \mathbf{i})\big)}{\varepsilon} \Big] H^{(1)}(y - \mathbf{i}) \mathcal{M}_{\varepsilon}(\psi)(\varepsilon\xi) \\ = &\varepsilon^{n} \sum_{\xi \in \mathbb{Z}^{n}} \sum_{\mathbf{i} \in \mathbf{I}} \frac{\mathcal{M}_{\varepsilon}(\psi) \big(\varepsilon(\xi - \mathbf{e}_{1})\big) - \mathcal{M}_{\varepsilon}(\psi) \big(\varepsilon\xi)\big)}{\varepsilon} H^{(1)}(y - \mathbf{i}) \mathcal{M}_{\varepsilon}(\phi_{\varepsilon}) \big(\varepsilon(\xi + \mathbf{i})\big) \end{split}$$

Due to the fact that $\phi_{\varepsilon}(x) = 0$ for a.e. $x \in \mathbb{R}^n \setminus \overline{\widetilde{\Omega}}_{6\sqrt{n}\varepsilon}$, in the above summations we only take the ξ 's belonging to Ξ_{ε} and satisfying $\rho(\varepsilon\xi) \geq 3\sqrt{n}\varepsilon$. Hence

$$<\mathcal{T}_{\varepsilon}\left(\frac{\partial\Phi}{\partial x_{1}}\right)(.,y)-\mathcal{M}_{\varepsilon}\left(\frac{\partial\Phi}{\partial x_{1}}\right),\psi>_{(H^{1}_{\rho}(\Omega))',H^{1}_{\rho}(\Omega)}$$
$$=\varepsilon^{n}\sum_{\xi\in\mathbb{Z}^{n}}\frac{\mathcal{M}_{\varepsilon}(\psi)\left(\varepsilon(\xi-\mathbf{e}_{1})\right)-\mathcal{M}_{\varepsilon}(\psi)\left(\varepsilon\xi\right)\right)}{\varepsilon}\sum_{\mathbf{i}\in\mathbf{I}}\left[H^{(1)}(y-\mathbf{i})-\frac{1}{2^{n-1}}\right]\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon(\xi+\mathbf{i})).$$

Thanks to the identity relation $\sum_{\mathbf{i}\in\mathbf{I}} \left[H^{(1)}(y-\mathbf{i}) - \frac{1}{2^{n-1}}\right] = 0$ we obtain that

$$\Big|\sum_{\mathbf{i}\in\mathbf{I}} \Big[H^{(1)}(y-\mathbf{i}) - \frac{1}{2^{n-1}}\Big]\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon(\xi+\mathbf{i}))\Big| \leq \sum_{\mathbf{i}\in\mathbf{I}} \Big|\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon(\xi+\mathbf{i})) - \mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon\xi)\Big|.$$

Taking into account the last equality and inequality above we deduce that

$$<\mathcal{T}_{\varepsilon}\left(\frac{\partial\Phi}{\partial x_{1}}\right)(.,y)-\mathcal{M}_{\varepsilon}\left(\frac{\partial\Phi}{\partial x_{1}}\right),\psi>_{(H^{1}_{\rho}(\Omega))',H^{1}_{\rho}(\Omega)}$$
$$=\varepsilon^{n}\sum_{\xi\in\mathbb{Z}^{n}}\sum_{\mathbf{i}\in\mathbf{I}}\left|\frac{\mathcal{M}_{\varepsilon}(\psi)\left(\varepsilon(\xi-\mathbf{e}_{1})\right)-\mathcal{M}_{\varepsilon}(\psi)\left(\varepsilon\xi\right)\right)}{\varepsilon}\left|\left|\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon(\xi+\mathbf{i}))-\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\varepsilon\xi)\right|\right|$$
$$=\frac{1}{\varepsilon}\sum_{\mathbf{i}\in\mathbf{I}}\int_{\Omega}\left|\mathcal{M}_{\varepsilon}(\psi)(\cdot-\varepsilon\mathbf{e}_{1})-\mathcal{M}_{\varepsilon}(\psi)\right|\left|\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\cdot+\varepsilon\mathbf{i})-\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})\right|$$
$$\leq\frac{C}{\varepsilon}\sum_{\mathbf{i}\in\mathbf{I}}\left\|\rho\left(\mathcal{M}_{\varepsilon}(\psi)(\cdot-\varepsilon\mathbf{e}_{1})-\mathcal{M}_{\varepsilon}(\psi)\right)\right\|_{L^{2}(\Omega)}\left\|\frac{1}{\rho}\left(\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})(\cdot+\varepsilon\mathbf{i})-\mathcal{M}_{\varepsilon}(\phi_{\varepsilon})\right)\right\|_{L^{2}(\Omega)}.$$

Due to $(2.23)_3$ and $(2.27)_3$ we finally get

$$< \mathcal{T}_{\varepsilon} \Big(\frac{\partial \Phi}{\partial x_1} \Big) (., y) - \mathcal{M}_{\varepsilon} \Big(\frac{\partial \Phi}{\partial x_1} \Big), \psi >_{(H^1_{\rho}(\Omega))', H^1_{\rho}(\Omega)} \le C \varepsilon ||\phi_{\varepsilon}||_{1/\rho} ||\psi||_{\rho}.$$

It leads to

$$\left\| \mathcal{T}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_1} \right) - \mathcal{M}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_1} \right) \right\|_{L^{\infty}(Y; (H^1_{\rho}(\Omega))')} \le C \varepsilon ||\phi_{\varepsilon}||_{1/\rho}.$$
(3.10)

Besides we have

$$\int_{\Omega} \frac{\partial \underline{\phi}}{\partial x_1}(x)\psi(x)dx = -\int_{\Omega} \underline{\phi}(x)\frac{\partial \psi}{\partial x_1}(x)dx \le C||\underline{\phi}/\rho||_{L^2(\Omega)}||\psi||_{\rho} \le C||\phi||_{1/\rho}||\psi||_{\rho}.$$

Hence $\left\|\varepsilon \frac{\partial \phi}{\partial x_1}\right\|_{(H^1_{\rho}(\Omega;\mathbb{R}^n))'} \leq C\varepsilon ||\phi||_{1/\rho}$. This last estimate with (2.10)₂, (3.9) and (3.10) yield

$$\left\|\mathcal{T}_{\varepsilon}\left(\frac{\partial\Phi}{\partial x_{1}}\right)-\frac{\partial\phi}{\partial x_{1}}\right\|_{L^{\infty}(Y;(H^{1}_{\rho}(\Omega))')}\leq C\varepsilon||\phi_{\varepsilon}||_{1/\rho}.$$

In the same way we prove the estimates for the partial derivatives of Φ with respect to $x_i, i \in \{2, ..., n\}$. Hence we get $\|\mathcal{T}_{\varepsilon}(\nabla \Phi) - \nabla \phi\|_{[L^{\infty}(Y;(H^1_{\rho}(\Omega))')]^n} \leq C\varepsilon ||\phi_{\varepsilon}||_{1/\rho}$. Then thanks to (3.8) the second estimate in (3.6) is proved.

4 Reminds about the classical periodic homogenization problem

We consider the homogenization problem

$$\phi^{\varepsilon} \in H_0^1(\Omega), \qquad \int_{\Omega} A_{\varepsilon}(x) \nabla \phi^{\varepsilon}(x) \nabla \psi(x) dx = \int_{\Omega} f(x) \psi(x) dx, \qquad \forall \psi \in H_0^1(\Omega), \quad (4.1)$$

where

• $A_{\varepsilon}(x) = A\left(\left\{\frac{x}{\varepsilon}\right\}\right)$ for a.e. $x \in \Omega$, where A is a square matrix belonging to $L^{\infty}(Y; \mathbb{R}^{n \times n})$ and satisfying the condition of uniform ellipticity $c|\xi|^2 \leq A(y)\xi \cdot \xi$ for a.e. $y \in Y$, with c a strictly positive constant,

•
$$f \in L^2(\Omega)$$

We showed in [10] that

$$\mathcal{T}_{\varepsilon}(\nabla\phi^{\varepsilon}) \longrightarrow \nabla\Phi + \nabla_y\widehat{\phi} \quad \text{strongly in} \quad L^2(\Omega \times Y; \mathbb{R}^n)$$

where $(\Phi, \hat{\phi}) \in H^1_0(\Omega) \times L^2(\Omega; H^1_{per}(Y))$ is the solution of the problem of unfolding homogenization

$$\begin{aligned} \forall (\Psi, \widehat{\psi}) &\in H^1_0(\Omega) \times L^2(\Omega; H^1_{per}(Y)) \\ &\int_{\Omega} \int_Y A(y) \big\{ \nabla \Phi(x) + \nabla_y \widehat{\phi}(x, y) \big\} \big\{ \nabla \Psi(x) + \nabla_y \widehat{\psi}(x, y) \big\} dx dy = \int_{\Omega} f(x) \Psi(x) dx. \end{aligned}$$

The correctors χ_i , $i \in \{1, \ldots, n\}$, are the solutions of the variational problems

$$\chi_i \in H^1_{per}(Y), \qquad \int_Y \chi_i = 0,$$

$$\int_Y A(y) \nabla_y (\chi_i(y) + y_i) \nabla_y \psi(y) dy = 0, \qquad \forall \psi \in H^1_{per}(Y).$$
(4.2)

They allow to express $\widehat{\phi}$ in terms of the partial derivatives of Φ

$$\widehat{\phi} = \sum_{i=1}^{n} \frac{\partial \Phi}{\partial x_i} \chi_i \tag{4.3}$$

and to give the homogenized problem satisfied by Φ

$$\Phi \in H_0^1(\Omega), \qquad \int_{\Omega} \mathcal{A} \nabla \Phi(x) \nabla \Psi(x) dx = \int_{\Omega} f(x) \Psi(x) dx, \qquad \forall \Psi \in H_0^1(\Omega)$$
(4.4)

where (see [10])

$$\mathcal{A}_{ij} = \sum_{k,l=1}^{n} \int_{Y} a_{kl}(y) \frac{\partial(y_j + \chi_j(y))}{\partial y_l} \frac{\partial(y_i + \chi_i(y))}{\partial y_k} dy.$$
(4.5)

5 An operator from $H^{-1/2}(\partial\Omega)$ into $L^2(\Omega)$

From now on, Ω is a bounded domain with a $\mathcal{C}^{1,1}$ boundary or an open bounded convex set.

In this section we first introduce a lifting operator \mathbf{T} (defined by (5.1)) from $H^{1/2}(\partial\Omega)$ into $H^1(\Omega)$. This operator and the estimate (5.2) are in fact sufficient to obtain the error estimates with a non-homogeneous Dirichlet condition (Theorem 6.3); one of the aim of this paper. Then we extend this operator. The extension of \mathbf{T} from $H^{-1/2}(\partial\Omega)$ into $H^1_{\rho}(\Omega)$ is essential in order to get a sharper estimate (6.3) than (6.2)₁. In Theorem 7.1 we give an application based on (6.3), in this theorem we investigate a first case of strongly oscillating boundary data.

Let g be in $H^{1/2}(\partial\Omega)$, there exists one $\phi_g \in H^1(\Omega)$ such that

$$\operatorname{div}(\mathcal{A}\nabla\phi_g) = 0 \quad \text{in} \quad \Omega, \quad \phi_g = g \quad \text{on} \quad \partial\Omega \tag{5.1}$$

where \mathcal{A} is the matrix given by (4.5). We have

$$||\phi_g||_{H^1(\Omega)} \le C||g||_{H^{1/2}(\partial\Omega)}.$$
 (5.2)

We denote by **T** the operator from $H^{1/2}(\partial\Omega)$ into $H^1(\Omega)$ which associates to $g \in H^{1/2}(\partial\Omega)$ the function $\phi_g \in H^1(\Omega)$.

Now, let (ψ, Ψ) be a couple in $[\mathcal{C}^{\infty}(\overline{\Omega})]^2$, integrating by parts over Ω gives

$$\int_{\Omega} \mathcal{A}\nabla\psi(x)\nabla\Psi(x)dx = -\int_{\Omega}\psi(x)\mathrm{div}(\mathcal{A}^{T}\nabla\Psi)(x)dx + \int_{\partial\Omega}\psi(x)(\mathcal{A}^{T}\nabla\Psi)(x)dx \cdot \nu(x)d\sigma.$$

The space $\mathcal{C}^{\infty}(\overline{\Omega})$ being dense in $H^1(\Omega)$ and $H^2(\Omega)$, hence the above equality holds true for any $\psi \in H^1(\Omega)$ and any $\Psi \in H^2(\Omega)$. Hence, for $\Psi \in H^1_0(\Omega) \cap H^2(\Omega)$ and ϕ_g defined by (5.1) we get

$$\int_{\Omega} \phi_g(x) \operatorname{div}(\mathcal{A}^T \nabla \Psi)(x) dx = \int_{\partial \Omega} g(x) \left(\mathcal{A}^T \nabla \Psi \right)(x) \cdot \nu(x) d\sigma.$$
(5.3)

Under the assumption on Ω the function $\Psi(g)$ defined by

$$\Psi(g) \in H_0^1(\Omega), \quad \operatorname{div}(\mathcal{A}^T \nabla \Psi(g)) = \phi_g \quad \text{in} \quad \Omega$$

belongs to $H^1_0(\Omega) \cap H^2(\Omega)$ and satisfies

$$||\Psi(g)||_{H^2(\Omega)} \le C ||\phi_g||_{L^2(\Omega)}.$$

Taking $\Psi = \Psi(g)$ in the above equality (5.3) we obtain

$$\int_{\Omega} |\phi_g(x)|^2 dx = \int_{\partial\Omega} g(x) \left(\mathcal{A}^T \nabla \Psi(g)(x) \right) \cdot \nu(x) d\sigma \le ||g||_{H^{-1/2}(\partial\Omega)} ||(\mathcal{A}^T \nabla \Psi(g)) \cdot \nu||_{H^{1/2}(\partial\Omega)} \le C ||g||_{H^{-1/2}(\partial\Omega)} ||\Psi(g)||_{H^2(\Omega)}.$$

This leads to

$$||\phi_g||_{L^2(\Omega)} \le C||g||_{H^{-1/2}(\partial\Omega)}.$$
 (5.4)

Due to (5.4), the operator **T** admits an extension (still denoted **T**) from $H^{-1/2}(\partial\Omega)$ into $L^2(\Omega)$ and we have

$$\forall g \in H^{-1/2}(\partial\Omega), \qquad ||\mathbf{T}(g)||_{L^2(\Omega)} \le C||g||_{H^{-1/2}(\partial\Omega)}$$

For $g \in H^{-1/2}(\partial \Omega)$, we also denote $\phi_g = \mathbf{T}(g)$. This function is the "very weak" solution of the problem

$$\phi_g \in L^2(\Omega), \quad \operatorname{div}(\mathcal{A}\nabla\phi_g) = 0 \quad \text{in } \Omega, \quad \phi_g = g \quad \text{on } \partial\Omega$$

or the solution of the following:

$$\begin{split} \phi_g &\in L^2(\Omega), \\ \int_{\Omega} \phi_g(x) \operatorname{div}(\mathcal{A}^T \nabla \psi(x)) dx = \langle g, (\mathcal{A}^T \nabla \psi) \cdot \nu \rangle_{H^{-1/2}(\partial\Omega), H^{1/2}(\partial\Omega)}, \end{split}$$
(5.5)
$$\forall \psi \in H^1_0(\Omega) \cap H^2(\Omega). \end{split}$$

Lemma 5.1. The operator **T** is a bicontinuous linear operator from $H^{-1/2}(\partial \Omega)$ onto

$$\mathbf{H} = \Big\{ \phi \in L^2(\Omega) \mid \operatorname{div}(\mathcal{A}\nabla\phi) = 0 \quad \operatorname{in} \quad \Omega \Big\}.$$

There exists a constant $C \ge 1$ such that

$$\forall g \in H^{-1/2}(\partial\Omega), \qquad \frac{1}{C} ||g||_{H^{-1/2}(\partial\Omega)} \le ||\mathbf{T}(g)||_{L^2(\Omega)} \le C ||g||_{H^{-1/2}(\partial\Omega)}. \tag{5.6}$$

Proof. Let ϕ be in **H** we are going to prove that there exists an element $g \in H^{-1/2}(\partial\Omega)$ such that $\mathbf{T}(g) = \phi$. To do that, we consider a continuous linear lifting operator **R** from $H^{1/2}(\partial\Omega)$ into $H^1_0(\Omega) \cap H^2(\Omega)$ satisfying for any $h \in H^{1/2}(\partial\Omega)$

$$\mathbf{R}(h) \in H_0^1(\Omega) \cap H^2(\Omega),$$

$$\mathcal{A}^T \nabla \mathbf{R}(h)_{|\partial\Omega} \cdot \nu = h \quad \text{on} \quad \partial\Omega.$$

$$||\mathbf{R}(h)||_{H^2(\Omega)} \leq C||h||_{H^{1/2}(\partial\Omega)}.$$

The map $h \mapsto \int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \mathbf{R}(h))$ is a continuous linear form defined over $H^{1/2}(\partial \Omega)$. Thus, there exists $g \in H^{-1/2}(\partial \Omega)$ such that

$$\int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \mathbf{R}(h)) = \langle g, h \rangle_{H^{-1/2}(\partial\Omega), H^{1/2}(\partial\Omega)} .$$
(5.7)

Since $\phi \in \mathbf{H}$, we deduce that for any $\psi \in \mathcal{C}_0^{\infty}(\Omega)$ we have $\int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \psi) = 0$. Therefore, for any $\psi \in H_0^2(\Omega)$ we have $\int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \psi) = 0$. Taking into account (5.7) we get

$$\int_{\Omega} \phi \operatorname{div}(\mathcal{A}^T \nabla \psi) = \langle g, (\mathcal{A}^T \nabla \psi) \cdot \nu \rangle_{H^{-1/2}(\partial\Omega), H^{1/2}(\partial\Omega)}, \qquad \forall \psi \in H^1_0(\Omega) \cap H^2(\Omega).$$

It yields $\phi = \phi_g$ and then (5.6).

Remark 5.2. It is well known (see e.g. [18]) that every function $\phi \in \mathbf{H}$ also belongs to $H^1_{\rho}(\Omega)$ and verifies

$$||\phi||_{\rho} \le C||\phi||_{L^2(\Omega)}.$$
 (5.8)

6 Error estimates with a non-homogeneous Dirichlet condition

Theorem 6.1. Let $(\phi^{\varepsilon})_{\varepsilon>0}$ be a sequence of functions belonging to $H^1(\Omega)$ such that

$$div(A_{\varepsilon}\nabla\phi^{\varepsilon}) = 0 \qquad in \quad \Omega.$$
(6.1)

Setting $g_{\varepsilon} = \phi_{|\partial\Omega}^{\varepsilon}$ and $\phi_{g_{\varepsilon}} = \mathbf{T}(g_{\varepsilon}) \in H^1(\Omega)$, there exists $\varepsilon_0 > 0$ such that for every $\varepsilon \leq \varepsilon_0$ we have

$$\begin{aligned} ||\phi^{\varepsilon}||_{H^{1}(\Omega)} &\leq C||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}, \qquad ||\phi^{\varepsilon} - \phi_{g_{\varepsilon}}||_{L^{2}(\Omega)} \leq C\varepsilon^{1/2}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}, \\ \left\|\rho\left(\nabla\phi^{\varepsilon} - \nabla\phi_{g_{\varepsilon}} - \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon}\left(\frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}}\right)\nabla_{y}\chi_{i}\left(\frac{\cdot}{\varepsilon}\right)\right)\right\|_{L^{2}(\Omega;\mathbb{R}^{n})} \leq C\varepsilon^{1/2}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}. \end{aligned}$$
(6.2)

Moreover we have

$$||\phi^{\varepsilon}||_{\rho} \leq C\left(\varepsilon^{1/2}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} + ||g_{\varepsilon}||_{H^{-1/2}(\partial\Omega)}\right).$$
(6.3)

The χ_i 's are the correctors introduced in Section 4 and **T** is the operator defined in Section 5.

Proof. Step 1. We prove the first estimate in (6.2). From Section 5 we get

$$||\phi_{g_{\varepsilon}}||_{H^{1}(\Omega)} \leq C||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} \qquad ||\phi_{g_{\varepsilon}}||_{\rho} \leq C||g_{\varepsilon}||_{H^{-1/2}(\partial\Omega)}.$$
(6.4)

We write (6.1) in the following weak form:

$$\phi^{\varepsilon} = \check{\phi}_{\varepsilon} + \phi_{g_{\varepsilon}}, \quad \check{\phi}_{\varepsilon} \in H^{1}_{0}(\Omega)$$

$$\int_{\Omega} A_{\varepsilon} \nabla \check{\phi}_{\varepsilon} \nabla v = -\int_{\Omega} A_{\varepsilon} \nabla \phi_{g_{\varepsilon}} \nabla v \qquad \forall v \in H^{1}_{0}(\Omega).$$
(6.5)

The solution $\check{\phi}_{\varepsilon}$ of the above variational problem satisfies

$$||\dot{\phi}_{\varepsilon}||_{H^1(\Omega)} \le C ||\nabla \phi_{g_{\varepsilon}}||_{L^2(\Omega;\mathbb{R}^n)}$$

Hence, from $(6.4)_1$ and the above estimate we get the first inequality in (6.2). Step 2. We prove the second estimate in (6.2). For every test function $v \in H_0^1(\Omega)$ we have

$$\int_{\Omega} A_{\varepsilon} \nabla \phi^{\varepsilon} \nabla v = 0.$$
(6.6)

Now, in order to obtain the L^2 error estimate we proceed as in the proof of the Theorem 3.2 in [16]. We first recall that for any $\phi \in H^1(\Omega)$ we have (see Lemma 2.3) for every $\varepsilon \leq \varepsilon_0 \doteq \gamma_0/3\sqrt{n}$

$$||\phi||_{L^2(\widetilde{\Omega}_{3c_0\sqrt{n}\varepsilon})} \le C\varepsilon^{1/2} ||\phi||_{H^1(\Omega)}.$$

Let U be a test function belonging to $H_0^1(\Omega) \cap H^2(\Omega)$. The above estimate yields

$$||\nabla U||_{L^2(\widetilde{\Omega}_{3c_0\sqrt{n\varepsilon}};\mathbb{R}^n)} \le C\varepsilon^{1/2}||U||_{H^2(\Omega)}$$
(6.7)

which in turn with (2.12)-(2.13)- $(2.14)_1$ and $(6.2)_1$ -(6.6) lead to

$$\left|\int_{\Omega \times Y} A(y) \mathcal{T}_{\varepsilon}(\nabla \phi^{\varepsilon})(x, y) \nabla U(x) dx dy\right| \le C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)} ||U||_{H^{2}(\Omega)}.$$
 (6.8)

The Theorem 2.3 in [16] gives an element $\widehat{\phi}_{\varepsilon} \in L^2(\Omega; H^1_{per}(Y))$ such that

$$\begin{aligned} ||\mathcal{T}(\nabla\phi^{\varepsilon}) - \nabla\phi^{\varepsilon} - \nabla_{y}\widehat{\phi}_{\varepsilon}||_{[L^{2}(Y;(H^{1}(\Omega))')]^{n}} &\leq C\varepsilon^{1/2}||\nabla\phi^{\varepsilon}||_{L^{2}(\Omega;\mathbb{R}^{n})} \\ &\leq C\varepsilon^{1/2}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}. \end{aligned}$$
(6.9)

The above inequalities (6.8) and (6.9) yield

$$\left|\int_{\Omega \times Y} A\left(\nabla \phi^{\varepsilon} + \nabla_{y} \widehat{\phi}_{\varepsilon}\right) \nabla U\right| \le C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} ||U||_{H^{2}(\Omega)}.$$
(6.10)

We set

$$\forall x \in \mathbb{R}^n, \qquad \rho_{\varepsilon}(x) = \inf \left\{ 1, \frac{\rho(x)}{\varepsilon} \right\}.$$

Now, we take $\overline{\chi} \in H^1_{per}(Y)$ and we consider the test function $u_{\varepsilon} \in H^1_0(\Omega)$ defined for a.e. $x \in \Omega$ by

$$u_{\varepsilon}(x) = \varepsilon \rho_{\varepsilon}(x) \mathcal{Q}_{\varepsilon} \left(\frac{\partial U}{\partial x_i}\right)(x) \overline{\chi} \left(\frac{x}{\varepsilon}\right).$$

Due to $(2.21)_2$ and (6.7) we get

$$\left\| \mathcal{Q}_{\varepsilon} \left(\frac{\partial U}{\partial x_{i}} \right) \nabla_{y} \overline{\chi} \left(\frac{\cdot}{\varepsilon} \right) \right\|_{L^{2}(\widetilde{\Omega}_{\sqrt{n}\varepsilon};\mathbb{R}^{n})} \leq C \varepsilon^{1/2} ||U||_{H^{2}(\Omega)} ||\overline{\chi}||_{H^{1}(Y)}$$
(6.11)

Then by a straightforward calculation and thanks to $(2.21)_2$ - $(2.22)_2$ and (6.7)-(6.11) we obtain

$$\left\|\nabla u_{\varepsilon} - \mathcal{Q}_{\varepsilon}\left(\frac{\partial U}{\partial x_{i}}\right)\nabla_{y}\overline{\chi}\left(\frac{\cdot}{\varepsilon}\right)\right\|_{L^{2}(\Omega;\mathbb{R}^{n})} \leq C\varepsilon^{1/2}||U||_{H^{2}(\Omega)}||\overline{\chi}||_{H^{1}(Y)}$$

which in turn with again (6.11) give

$$\|\nabla u_{\varepsilon}\|_{L^{2}(\widetilde{\Omega}_{\sqrt{n}\varepsilon};\mathbb{R}^{n})} \leq C\varepsilon^{1/2}||U||_{H^{2}(\Omega)}||\overline{\chi}||_{H^{1}(Y)}$$
(6.12)

and then with $(2.22)_1$ they lead to

$$\left\|\nabla u_{\varepsilon} - \mathcal{M}_{\varepsilon}\left(\frac{\partial U}{\partial x_{i}}\right)\nabla_{y}\overline{\chi}\left(\frac{\cdot}{\varepsilon}\right)\right\|_{L^{2}(\Omega;\mathbb{R}^{n})} \leq C\varepsilon^{1/2}||U||_{H^{2}(\Omega)}||\overline{\chi}||_{H^{1}(Y)}$$

In (6.6) we replace ∇u_{ε} with $\mathcal{M}_{\varepsilon}\left(\frac{\partial U}{\partial x_{i}}\right)\nabla_{y}\overline{\chi}\left(\frac{\cdot}{\varepsilon}\right)$; we continue using (2.12)-(2.13) and (6.2)₁-(6.12) to obtain

$$\left|\int_{\Omega\times Y} A(y)\mathcal{T}_{\varepsilon}(\nabla\phi^{\varepsilon})(x,y)\mathcal{M}_{\varepsilon}\left(\frac{\partial U}{\partial x_{i}}\right)(x)\nabla_{y}\overline{\chi}(y)dxdy\right| \leq C\varepsilon^{1/2}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}||U||_{H^{2}(\Omega)}||\overline{\chi}||_{H^{1}(Y)}$$

which with $(2.17)_2$ and then (6.9) give

$$\left|\int_{\Omega\times Y} A(y) \left(\nabla\phi^{\varepsilon}(x) + \nabla_{y}\widehat{\phi}_{\varepsilon}(x,y)\right) \frac{\partial U}{\partial x_{i}}(x)\nabla_{y}\overline{\chi}(y)dx\,dy\right| \leq C\varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} ||U||_{H^{2}(\Omega)} ||\overline{\chi}||_{H^{1}(Y)}$$

$$\tag{6.13}$$

As in [16] we introduce the adjoint correctors $\overline{\chi}_i \in H^1_{per}(Y)$, $i \in \{1, \ldots, n\}$, defined by

$$\int_{Y} A(y) \nabla_{y} \psi(y) \nabla_{y}(\overline{\chi}_{i}(y) + y_{i}) dy = 0 \qquad \forall \psi \in H^{1}_{per}(Y).$$
(6.14)

From (6.13) we get

$$\left|\int_{\Omega\times Y} A\left(\nabla\phi^{\varepsilon} + \nabla_{y}\widehat{\phi}_{\varepsilon}\right)\nabla_{y}\left(\sum_{i=1}^{n}\frac{\partial U}{\partial x_{i}}\overline{\chi}_{i}\right)\right| \leq C\varepsilon^{1/2}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}||U||_{H^{2}(\Omega)}$$

and from the definition (4.2) of the correctors χ_i we have

$$\int_{\Omega \times Y} A\Big(\nabla \phi^{\varepsilon} + \sum_{i=1}^{n} \frac{\partial \phi^{\varepsilon}}{\partial x_{i}} \nabla_{y} \chi_{i}\Big) \nabla_{y}\Big(\sum_{j=1}^{n} \frac{\partial U}{\partial x_{j}} \overline{\chi}_{j}\Big) = 0.$$

Thus

$$\left|\int_{\Omega\times Y} A\nabla_y \left(\widehat{\phi}_{\varepsilon} - \sum_{i=1}^n \frac{\partial \phi^{\varepsilon}}{\partial x_i} \chi_i\right) \nabla_y \left(\sum_{j=1}^n \frac{\partial U}{\partial x_j} \overline{\chi}_j\right)\right| \le C\varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} ||U||_{H^2(\Omega)}$$

and thanks to (6.14) we obtain

$$\left|\int_{\Omega\times Y} A\nabla_y \left(\widehat{\phi}_{\varepsilon} - \sum_{i=1}^n \frac{\partial \phi^{\varepsilon}}{\partial x_i} \chi_i\right) \nabla U\right| \le C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} ||U||_{H^2(\Omega)}.$$

The above estimate, (6.10) and the expression (4.5) of the matrix \mathcal{A} yield

$$\left|\int_{\Omega} \mathcal{A}\nabla\phi^{\varepsilon}\nabla U\right| \leq C\varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} ||U||_{H^{2}(\Omega)}.$$

Finally, since we have $\int_{\Omega} \mathcal{A} \nabla \phi_{g_{\varepsilon}} \nabla v = 0$ for any $v \in H_0^1(\Omega)$, we deduce that

$$\forall U \in H_0^1(\Omega) \cap H^2(\Omega), \qquad \left| \int_{\Omega} \mathcal{A} \nabla (\phi^{\varepsilon} - \phi_{g_{\varepsilon}}) \nabla U \right| \le C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)} ||U||_{H^2(\Omega)}.$$

Now, let $U_{\varepsilon} \in H_0^1(\Omega)$ be the solution of the following variational problem:

$$\int_{\Omega} \mathcal{A} \nabla v \nabla U_{\varepsilon} = \int_{\Omega} v(\phi^{\varepsilon} - \phi_{g_{\varepsilon}}), \qquad \forall v \in H_0^1(\Omega).$$

Under the assumption on the boundary of Ω , we know that U_{ε} belongs to $H_0^1(\Omega) \cap H^2(\Omega)$ and satisfies $||U_{\varepsilon}||_{H^2(\Omega)} \leq C ||\phi^{\varepsilon} - \phi_{g_{\varepsilon}}||_{L^2(\Omega)}$ (the constant do not depend on ε). Therefore, the second estimate in (6.2) is proved.

Step 3. We prove the third estimate in (6.2) and (6.3). The partial derivative $\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_i}$ satisfies

$$\operatorname{div}\left(\mathcal{A}\nabla\left(\frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}}\right)\right) = 0 \quad \text{in} \quad \Omega, \quad \frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}} \in L^{2}(\Omega)$$

Thus, from Remark 5.8 and estimate $(6.4)_2$ we get

$$\left\|\rho\nabla\left(\frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}}\right)\right\|_{L^{2}(\Omega;\mathbb{R}^{n})} \leq C\left\|\frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}}\right\|_{L^{2}(\Omega)} \leq C||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$
(6.15)

Now, let U be in $H_0^1(\Omega)$, the function ρU belongs to $H_{1/\rho}^1(\Omega)$. Applying the Theorem 3.2 with the function ρU , there exists $\widehat{u}_{\varepsilon} \in L^2(\Omega; H_{per}^1(Y))$ such that

$$||\mathcal{T}_{\varepsilon}(\nabla(\rho U)) - \nabla(\rho U) - \nabla_{y} \widehat{u}_{\varepsilon}||_{L^{2}(Y;(H^{1}_{\rho}(\Omega;\mathbb{R}^{n}))')} \leq C\varepsilon ||\rho U||_{H^{1}_{1/\rho}(\Omega)} \leq C\varepsilon ||U||_{H^{1}(\Omega)}.$$
 (6.16)

The above estimates (6.15) and (6.16) lead to

$$\left|\int_{\Omega\times Y} A\Big(\nabla\phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y}\chi_{i}\Big)\Big(\mathcal{T}_{\varepsilon}\big(\nabla(\rho U)\big) - \nabla(\rho U) - \nabla_{y}\widehat{u}_{\varepsilon}\Big)\right| \leq C\varepsilon||U||_{H^{1}(\Omega)}||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}$$

By definition of the correctors χ_i we have

$$\int_{\Omega \times Y} A \Big(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y} \chi_{i} \Big) \nabla_{y} \widehat{u}_{\varepsilon} = 0.$$

Besides, from the definitions of the function $\phi_{g_{\varepsilon}}$ and the homogenized matrix \mathcal{A} we have

$$0 = \int_{\Omega} \mathcal{A} \nabla \phi_{g_{\varepsilon}} \nabla(\rho U) = \int_{\Omega \times Y} A \Big(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y} \chi_{i} \Big) \nabla(\rho U).$$

The above inequality and equalities yield

$$\left|\int_{\Omega\times Y} A\left(\nabla\phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y}\chi_{i}\right) \mathcal{T}_{\varepsilon}\left(\nabla(\rho U)\right)\right| \leq C\varepsilon ||\nabla U||_{L^{2}(\Omega;\mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$
 (6.17)

We have

$$\nabla(\rho U) = \rho \Big(\nabla U + \nabla \rho \frac{U}{\rho} \Big).$$

Then since $U/\rho \in L^2(\Omega)$ and $||U/\rho||_{L^2(\Omega)} \leq C ||\nabla U||_{L^2(\Omega;\mathbb{R}^n)}$ and due to (3.2) we get

$$\left\|\mathcal{T}_{\varepsilon}\left(\nabla(\rho U)\right) - \rho \mathcal{T}_{\varepsilon}\left(\nabla U + \nabla \rho \frac{U}{\rho}\right)\right\|_{L^{2}(\Omega;\mathbb{R}^{n})} \leq C\varepsilon \left\|\nabla U + \nabla \rho \frac{U}{\rho}\right\|_{L^{2}(\Omega;\mathbb{R}^{n})} \leq C\varepsilon ||U||_{H^{1}(\Omega)}.$$

From (6.17) and the above inequalities we deduce that

$$\left|\int_{\Omega\times Y} A\left(\rho\nabla\phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \rho \frac{\partial\phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y}\chi_{i}\right) \mathcal{T}_{\varepsilon}\left(\nabla U + \nabla\rho \frac{U}{\rho}\right)\right| \leq C\varepsilon ||\nabla U||_{L^{2}(\Omega;\mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$

We recall that $\rho \nabla \phi_{g_{\varepsilon}} \in H_0^1(\Omega; \mathbb{R}^n)$, hence from $(2.14)_2$, $(2.17)_1$ and (6.15) we get

$$\left| \int_{\Omega \times Y} A\left(\rho \nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \rho \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \nabla_{y} \chi_{i}\right) \mathcal{T}_{\varepsilon}\left(\nabla U + \nabla \rho \frac{U}{\rho}\right) - \int_{\Omega \times Y} A\left(\mathcal{T}_{\varepsilon}(\rho \nabla \phi_{g_{\varepsilon}}) + \sum_{i=1}^{n} \mathcal{M}_{\varepsilon}\left(\rho \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}}\right) \nabla_{y} \chi_{i}\right) \mathcal{T}_{\varepsilon}\left(\nabla U + \nabla \rho \frac{U}{\rho}\right) \right| \leq C\varepsilon ||\nabla U||_{L^{2}(\Omega;\mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}$$

Then transforming by inverse unfolding we obtain

$$\left|\int_{\widehat{\Omega}_{\varepsilon}} A_{\varepsilon} \left(\rho \nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \mathcal{M}_{\varepsilon} \left(\rho \frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}}\right) \nabla_{y} \chi_{i} \left(\frac{\cdot}{\varepsilon}\right) \right) \left(\nabla U + \nabla \rho \frac{U}{\rho}\right) \right| \leq C\varepsilon ||\nabla U||_{L^{2}(\Omega;\mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$

Now, thanks to (2.28) and (6.15) we get

$$\left|\int_{\Omega} A_{\varepsilon} \rho \left(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \mathcal{M}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \nabla_{y} \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \right) \left(\nabla U + \nabla \rho \frac{U}{\rho} \right) \right| \leq C \varepsilon ||\nabla U||_{L^{2}(\Omega; \mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$

Then using $(2.29)_1$ it leads to

$$\left|\int_{\Omega} A_{\varepsilon} \left(\nabla \phi_{g_{\varepsilon}} + \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \nabla_{y} \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \right) \nabla(\rho U) \right| \leq C \varepsilon ||\nabla U||_{L^{2}(\Omega; \mathbb{R}^{n})} ||g_{\varepsilon}||_{H^{1/2}(\partial \Omega)}.$$

We recall that $\int_{\Omega} A_{\varepsilon} \nabla \phi^{\varepsilon} \nabla (\rho U) = 0$. We choose $U = \rho \left(\phi^{\varepsilon} - \phi_{g_{\varepsilon}} - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \right)$

which belongs to $H_0^1(\Omega)$. Due to the second estimate in (6.2), the third one in (6.2) follows immediately.

The estimate (6.3) is the consequence of $(2.29)_2$, $(6.2)_2$, $(6.2)_3$, $(6.4)_2$ and (6.15).

Corollary 6.2. Let $(\phi^{\varepsilon})_{\varepsilon>0}$ be a sequence of functions belonging to $H^1(\Omega)$ and satisfying (6.1). We set $g_{\varepsilon} = \phi^{\varepsilon}_{|\partial\Omega}$, if we have

$$g_{\varepsilon} \rightharpoonup g \quad weakly \ in \quad H^{1/2}(\partial \Omega)$$

then we obtain

$$\phi^{\varepsilon} \rightharpoonup \phi_g \quad weakly \ in \qquad H^1(\Omega),$$

$$\phi^{\varepsilon} - \phi_g - \varepsilon \sum_{i=1}^n \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_g}{\partial x_i}\right) \chi_i \left(\frac{\cdot}{\varepsilon}\right) \longrightarrow 0 \quad strongly \ in \qquad H^1_{\rho}(\Omega). \tag{6.18}$$

Moreover, if

$$g_{\varepsilon} \longrightarrow g \quad strongly \ in \quad H^{1/2}(\partial\Omega)$$
 (6.19)

then we have

$$\phi^{\varepsilon} - \phi_g - \varepsilon \sum_{i=1}^n \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_g}{\partial x_i} \right) \chi_i \left(\frac{\cdot}{\varepsilon} \right) \longrightarrow 0 \quad strongly \ in \quad H^1(\Omega). \tag{6.20}$$

Proof. Thanks to $(6.2)_1$ the sequence $(\phi^{\varepsilon})_{\varepsilon>0}$ is uniformly bounded in $H^1(\Omega)$. Then due to Lemma 5.1 and Remark 5.8 we get

$$||\phi_g - \phi_{g_{\varepsilon}}||_{\rho} \le C||g - g_{\varepsilon}||_{H^{-1/2}(\partial\Omega)}$$

which with $(6.2)_2$ (resp. $(6.2)_3$) give the convergence $(6.18)_1$ (resp. $(6.18)_2$). Under the assumption (6.19), we use (5.2) and we proceed as in the proof of the Theorem 6.1 of [10] in order to obtain the strong convergence (6.20). **Theorem 6.3.** Let ϕ^{ε} be the solution of the following homogenization problem:

 $-\operatorname{div}(A_{\varepsilon}\nabla\phi^{\varepsilon})=f\quad in\quad \Omega,\qquad \phi^{\varepsilon}=g\quad on\quad \partial\Omega$

where $f \in L^2(\Omega)$ and $g \in H^{1/2}(\partial \Omega)$. We have

$$\begin{split} \|\phi^{\varepsilon} - \Phi\|_{L^{2}(\Omega)} &\leq C\left\{\varepsilon\|\|f\|_{L^{2}(\Omega)} + \varepsilon^{1/2}\|g\|_{H^{1/2}(\partial\Omega)}\right\},\\ \left\|\rho\left(\nabla\phi^{\varepsilon} - \nabla\Phi - \sum_{i=1}^{n}\mathcal{Q}_{\varepsilon}\left(\frac{\partial\Phi}{\partial x_{i}}\right)\nabla_{y}\chi_{i}\left(\frac{\cdot}{\varepsilon}\right)\right)\right\|_{L^{2}(\Omega;\mathbb{R}^{n})} &\leq C\left\{\varepsilon\|\|f\|_{L^{2}(\Omega)} + \varepsilon^{1/2}\|g\|_{H^{1/2}(\partial\Omega)}\right\}. \end{split}$$

where Φ is the solution of the homogenized problem

 $-div(\mathcal{A}\nabla\Phi) = f$ in Ω , $\Phi = g$ on $\partial\Omega$.

Moreover we have

$$\phi^{\varepsilon} - \Phi - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \Phi}{\partial x_{i}} \right) \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \longrightarrow 0 \quad strongly \ in \quad H^{1}(\Omega). \tag{6.21}$$

Proof. Let $\widetilde{\phi}^{\varepsilon}$ be the solution of the homogenization problem

$$\widetilde{\phi}^{\varepsilon} \in H_0^1(\Omega), \quad -\operatorname{div}(A_{\varepsilon}\nabla\widetilde{\phi}^{\varepsilon}) = f \quad \text{in} \quad \Omega$$

and $\widetilde{\Phi}$ the solution of the homogenized problem

$$\widetilde{\Phi} \in H_0^1(\Omega), \quad -\operatorname{div}(\mathcal{A}\nabla\widetilde{\Phi}) = f \quad \text{in } \Omega.$$

The Theorem 3.2 in [16] gives the following estimate:

$$\left\|\widetilde{\phi}^{\varepsilon} - \widetilde{\Phi}\right\|_{L^{2}(\Omega)} + \left\|\rho \nabla \left(\widetilde{\phi}^{\varepsilon} - \widetilde{\Phi} - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \widetilde{\Phi}}{\partial x_{i}}\right) \chi_{i}\left(\frac{\cdot}{\varepsilon}\right)\right)\right\|_{L^{2}(\Omega;\mathbb{R}^{n})} \leq C\varepsilon \left\|f\right\|_{L^{2}(\Omega)} \quad (6.22)$$

while the Theorem 4.1 in [15] gives

$$\left\|\widetilde{\phi}^{\varepsilon} - \widetilde{\Phi} - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \widetilde{\Phi}}{\partial x_{i}}\right) \chi_{i} \left(\frac{\cdot}{\varepsilon}\right) \right\|_{H^{1}(\Omega)} \leq C \varepsilon^{1/2} ||f||_{L^{2}(\Omega)}.$$
(6.23)

The function $\phi^{\varepsilon} - \widetilde{\phi}^{\varepsilon}$ satisfies

 $\operatorname{div} \left(A_{\varepsilon} \nabla (\phi^{\varepsilon} - \widetilde{\phi}^{\varepsilon}) \right) = 0 \quad \text{in} \quad \Omega, \quad \phi^{\varepsilon} - \widetilde{\phi}^{\varepsilon} = g \quad \text{on} \quad \partial \Omega.$

Thanks to the inequalities (6.2) and (6.22) we deduce the estimates of the theorem. The strong convergence (6.21) is a consequence of (6.23) and the strong convergence (6.20) after having observed that $\Phi - \tilde{\Phi} = \phi_q$.

Remark 6.4. In Theorem 6.3, if $g \in H^{3/2}(\partial\Omega)$ then in the estimates therein, we can replace $\varepsilon^{1/2}||g||_{H^{1/2}(\partial\Omega)}$ with $\varepsilon||g||_{H^{3/2}(\partial\Omega)}$. Moreover we have the following H^1 -global error estimate:

$$\left\|\phi^{\varepsilon} - \Phi - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon}\left(\frac{\partial \Phi}{\partial x_{i}}\right) \chi_{i}\left(\frac{\cdot}{\varepsilon}\right)\right\|_{H^{1}(\Omega)} \leq C\varepsilon^{1/2}\left\{||f||_{L^{2}(\Omega)} + ||g||_{H^{3/2}(\partial\Omega)}\right\}.$$

7 A first result with strongly oscillating boundary data

In this section we consider the solution ϕ^{ε} of the homogenization problem

$$\operatorname{div}(A_{\varepsilon}\nabla\phi^{\varepsilon}) = 0 \quad \text{in} \quad \Omega$$

$$\phi^{\varepsilon} = g_{\varepsilon} \quad \text{on} \quad \partial\Omega$$
(7.1)

where $g_{\varepsilon} \in H^{1/2}(\partial \Omega)$. As a consequence of the Theorem 6.1 we obtain the following result:

Theorem 7.1. Let ϕ^{ε} be the solution of the problem (7.1). If we have

$$g_{\varepsilon} \rightharpoonup g \quad weakly \ in \quad H^{-1/2}(\partial \Omega)$$

and

$$\varepsilon^{1/2}g_{\varepsilon} \longrightarrow 0 \quad strongly \ in \quad H^{1/2}(\partial\Omega)$$

$$(7.2)$$

then

$$\phi^{\varepsilon} \rightharpoonup \phi_g \quad weakly \ in \quad H^1_{\rho}(\Omega).$$
 (7.3)

Furthermore, if we have

$$g_{\varepsilon} \longrightarrow g \quad strongly \ in \quad H^{-1/2}(\partial \Omega)$$

then

$$\phi^{\varepsilon} - \phi_g - \varepsilon \sum_{i=1}^n \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_i} \right) \chi_i \left(\frac{\cdot}{\varepsilon} \right) \longrightarrow 0 \quad strongly \ in \quad H^1_{\rho}(\Omega). \tag{7.4}$$

Proof. Due to (6.3) the sequence $(\phi^{\varepsilon})_{\varepsilon>0}$ is uniformly bounded in $H^1_{\rho}(\Omega)$. From the estimates $(6.2)_3$ and $(6.4)_2$ we get

$$\left\|\phi^{\varepsilon} - \phi_{g_{\varepsilon}} - \varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}}\right) \chi_{i} \left(\frac{\cdot}{\varepsilon}\right) \right\|_{H^{1}_{\rho}(\Omega)} \leq C \varepsilon^{1/2} ||g_{\varepsilon}||_{H^{1/2}(\partial\Omega)}.$$

Then using the variational problem (5.5) and estimate $(6.4)_2$ we obtain

$$\phi_{g_{\varepsilon}} \rightharpoonup \phi_g$$
 weakly in $H^1_{\rho}(\Omega)$.

Since the sequence $\varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \chi_{i} \left(\frac{\cdot}{\varepsilon} \right)$ is uniformly bounded in $H^{1}_{\rho}(\Omega)$ and strongly converges to 0 in $L^{2}(\Omega)$, we have $\varepsilon \sum_{i=1}^{n} \mathcal{Q}_{\varepsilon} \left(\frac{\partial \phi_{g_{\varepsilon}}}{\partial x_{i}} \right) \chi_{i} \left(\frac{\cdot}{\varepsilon} \right) \to 0$ weakly in $H^{1}_{\rho}(\Omega)$. Therefore the weak convergence (7.3) is proved.

In the case $g_{\varepsilon} \longrightarrow g$ strongly in $H^{-1/2}(\partial \Omega)$, the estimates (5.4) and (5.8) lead to

$$||\phi_{g_{\varepsilon}} - \phi_{g}||_{H^{1}_{\rho}(\Omega)} \leq C||g_{\varepsilon} - g||_{H^{-1/2}(\partial\Omega)}$$

Hence with $(2.29)_2$ they yield (7.4).

In a forthcoming paper we will show that in both cases (weak or strong convergence of the sequence $(g_{\varepsilon})_{\varepsilon>0}$ towards g in $H^{-1/2}(\partial\Omega)$) the assumption (7.2) is essential in order to obtain at least (7.3).

References

- G. Allaire. Homogenization and two-scale convergence. SIAM J. Math. Anal., 23, 1992, 1482-1518.
- [2] G. Allaire and M. Amar. Boundary layer tails in periodic homogenization, ESAIM: Control, Optimization and Calc. of Variations, 4, 1999, 209-243.
- [3] M. Avellaneda and F.-H. Lin. Homogenization of elliptic problems with L^p boundary data, Appl. Math. Optim., 15, 1987, 93-107.
- [4] A. Bensoussan, J.-L.Lions and G.Papanicolaou. Asymptotic Analysis for Periodic Structures. North Holland, Amsterdam, 1978.
- [5] D. Blanchard, A. Gaudiello and G. Griso. Junction of a periodic family of elastic rods with a thin plate. II, J. Math. Pures Appl., 88, 2007, 149-190.
- [6] D. Blanchard, G. Griso. Microscopic effects in the homogenization of the junction of rods and a thin plate Asymptot. Anal., 56, 1, 2008, 1-36.
- [7] A. Blasselle and G. Griso. Mechanical modeling of the skin, Asymptot. Anal., 74, 3-4, 2011, 167-198.
- [8] A. Bossavit, G. Griso, and B. Miara. Modeling of periodic electro-magnetic structures. Bianisotropic materials with memory effects, J.M.P.A., 84, 2005, 819-850.
- [9] H. Brezis. Analyse fonctionnelle. Théorie et applications, Masson, Paris, 1983.
- [10] D. Cioranescu, A. Damlamian and G. Griso. The periodic unfolding method in homogenization, SIAM J. of Math. Anal., 40 (4), 2008, 1585-1620.
- [11] D. Cioranescu, A. Damlamian, G. Griso and D. Onofrei. The periodic unfolding method for perforated domains and Neumann sieve models, J. Math. Pures Appl., 89, 2008, 248-277.
- [12] D. Cioranescu, A. Damlamian, P. Donato, G. Griso and R. Zaki. The periodic unfolding method in domains with holes, SIAM J. of Math. Anal., 44 (2), 2012, 718-760.
- [13] D. Cioranescu and P. Donato. An Introduction to Homogenization. Oxford Lecture Series in Mathematics and its Applications. Oxford University Press, 1999.

- [14] G. Griso. Estimation d'erreur et éclatement en homogénéisation périodique. C. R. Acad. Sci. Paris, Ser. I 335, 2002, 333-336.
- [15] G. Griso. Error estimate and unfolding for periodic homogenization, Asymptot. Anal., 40, 2004, 269-286.
- [16] G. Griso. Interior error estimates for periodic homogenization, Analysis and Applications, 4, 2006, 61-79.
- [17] G. Griso. Asymptotic behavior of structures made of plates, Analysis Appl., 3 (4), 2005, 325-356.
- [18] G. Griso. Decomposition of displacements of thin structures. J. Math. Pures Appl., 89, 2008, 199-233.
- [19] G. Griso, E. Rohan. On the homogenization of a diffusion-deformation problem in strongly heterogeneous media. Ricerche di Matematica, 56 (2), 2007, 161-188.
- [20] S. Kesavan. Homogenization of elliptic eigenvalue problems: Part 1. Appl. Math. Optim., 5 (1), 1979, 153-167.
- [21] J.L. Lions and E. Magenes. Problèmes aux limites non homogènes, Dunod, Paris, 1968.
- [22] S. Moskow and M. Vogelius. First-order corrections to the homogenised eigenvalues of a periodic composite medium. A convergence proof. Proceedings of the Royal Society of Edimburgh, 127A, 1997, 1263-1299.
- [23] D. Onofrei, B. Vernescu. Error estimates for periodic homogenization with nonsmooth coefficients, Asymptot. Anal., 54, 2007, 103-123.
- [24] O. Ouchetto, S. Zouhdi, A. Bossavit, G. Griso, B. Miara, A. Razek. Homogenization of structured electromagnetic materials and metamaterials. Journal of materials processing technology, 181 (1), 2007, 225-229.